From 12 GeV to EIC: Imaging

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HELMHOLTZ RESEARCH FOR GRAND CHALLENGES



Introduction	GPDs	TMDs	Wigner fcts	Summary
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Parton imaging: why? what? how? where?

scope of this talk

- overarching physics motivation concrete examples of physics accessible with parton imaging
- explain some underlying physical principles
- some comments on theory status
- some comments on specific issues at JLab@12 GeV and at EIC
- not a comprehensive overview

more information in later talks at this meeting

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Physics questions

outstanding question in QCD: relation between

 $\begin{array}{ll} q, \bar{q}, g \text{ in } \mathcal{L}_{\mathsf{QCD}} & \text{constituent quarks} \\ \text{current quarks, } m_{u,d} \ll 100 \, \mathrm{MeV} & \text{effective masses } m \sim 300 \, \mathrm{MeV} \\ \text{manifest at small distances} & \text{spectroscopy, quark models} \\ \text{hadrons seen at high resolution} & "p = uud, n = ddu, \dots" \end{array}$

 simplest idea: at low resolution p = uud gluons and sea quarks from perturbative parton splitting Parisi, Petronzio '76, Novikov et al '77; Glück, Reya '77 is incompatible with measured parton densities Glück, Reva, Vogt '90 ff.; MD talk at POETIC 2018

other scenarios

- proton = quarks and antiquarks, gluons from evolution have \bar{q} e.g. from meson cloud Thomas '83 or in quark-soliton model Diakokov et al '86
- proton = UUD with "composite valence quarks" U, D containing antiquarks and gluons Altarelli et al '74

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Physics questions

outstanding question in QCD: relation between

 q, \bar{q}, q in \mathcal{L}_{QCD} current quarks, $m_{u,d} \ll 100 \,\mathrm{MeV}$ effective masses $m \sim 300 \,\mathrm{MeV}$ manifest at small distances hadrons seen at high resolution "p = uud, n = ddu, ..."

constituent quarks spectroscopy, quark models

general idea:

seek qualitatively new and quantitatively precise information on q, \bar{q}, q in hadrons at high resolution

several lines of study:

- partons in "valence region" $x \sim 1/3$
- partons at large $x \gg 1/3 \rightsquigarrow JLab@12 \text{ GeV}$
- sea guarks and gluons → EIC
- · spin and orbital angular momentum of partons
- distribution of partons in three dimensions ~> imaging

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Imaging: longitudinal vs. transverse directions

hard processes single out (at least) one spatial direction



holds both in collision c.m. and in target rest frame

- different roles played by longitud. and transv. directions
 lose manifest 3dim rotation symmetry in target rest frame
- usual parton densities: longitudinal information aims: achieve high precision, details of flavor structure, q vs. q
 , polarisation, nuclear effects
- transverse structure: much less well known in first instance aim to see general trends/patterns but may require high-precision to expose subtle effects
- new d.o.f.: orbital angular momentum (classically: $\vec{L} = \vec{r} \times \vec{p}$)

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Imaging basics: transverse momentum and transverse position

 variables related by 2d Fourier transforms, e.g. for quark field operator (in QFT) or wave function (in QM)

$$\phi(\mathbf{k}) = \int d^2 \mathbf{z} \ e^{i\mathbf{z}\mathbf{k}} \psi(\mathbf{z})$$

display only transverse variables (boldface), omit longitudinal ones

- fully relativistic: can localise only in 2 dimensions in 3d can only localise object within its Compton wavelength
- at level of squared amplitudes/probabilities/density matrices have

$$\begin{split} \overline{\phi}(\boldsymbol{k}) \, \phi(\boldsymbol{l}) &= \int d^2 \boldsymbol{y} \, d^2 \boldsymbol{z} \, e^{-i(\boldsymbol{y}\boldsymbol{k}-\boldsymbol{z}\boldsymbol{l})} \, \overline{\psi}(\boldsymbol{y}) \, \psi(\boldsymbol{z}) \\ \boldsymbol{y} \boldsymbol{k} - \boldsymbol{z} \boldsymbol{l} &= \frac{1}{2} (\boldsymbol{y} + \boldsymbol{z}) (\boldsymbol{k} - \boldsymbol{l}) + \frac{1}{2} (\boldsymbol{y} - \boldsymbol{z}) (\boldsymbol{k} + \boldsymbol{l}) \end{split}$$

'average' transv. momentum \leftrightarrow position difference transv. momentum transfer \leftrightarrow 'average' position

- 'average' transv. mom. and position not Fourier conjugate nor are parton distributions in transv. mom. or in impact parameter
- ▶ integrate over all transv. mom. ~> all fields at transv. position zero

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fully relativistic: can localise only in 2 dimensions Soper '72; Burkardt '02 in 3d can only localise object within its Compton wavelength

at level of squared amplitudes/probabilities/density matrices have

$$\begin{split} \overline{\phi}(\boldsymbol{k}) \, \phi(\boldsymbol{l}) &= \int d^2 \boldsymbol{y} \, d^2 \boldsymbol{z} \, e^{-i(\boldsymbol{y}\boldsymbol{k}-\boldsymbol{z}\boldsymbol{l})} \, \overline{\psi}(\boldsymbol{y}) \, \psi(\boldsymbol{z}) \\ \boldsymbol{y} \boldsymbol{k} - \boldsymbol{z} \boldsymbol{l} &= \frac{1}{2} (\boldsymbol{y} + \boldsymbol{z}) (\boldsymbol{k} - \boldsymbol{l}) + \frac{1}{2} (\boldsymbol{y} - \boldsymbol{z}) (\boldsymbol{k} + \boldsymbol{l}) \end{split}$$

'average' transv. momentum \leftrightarrow position difference transv. momentum transfer \leftrightarrow 'average' position

▶ full information: Wigner phase space distributions $W(x, \mathbf{k}, \mathbf{b})$ give probabilities $\int d^2 \mathbf{k} W = f(x, \mathbf{b})$ and $\int d^2 \mathbf{b} W = f(x, \mathbf{k})$

in context of parton distributions: Belitsky, Ji, Yuan '03

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Factorisation: exclusive or inclusive

- separate dynamics into "hard probe" and "structure of target" nontrivial, requires hard scale Q², comes with corrections ~ (Λ/Q)^k
- ► exclusive processes: nucleon in final state, longitudinal and transverse momentum transfer ~→ transverse position of struck parton
 - PDF factorization
 - ★ inclusive processes
 - ★ p_T ~ hardest scale or unmeasured



- TMD factorization
 - 🛨 inclusive
 - ★ p_T « hardest scale



- GPD factorization
 - ★ exclusive processes
 - 🖈 non-forward kinematics



- small-x factorization
 - ★ inclusive or exclusive
 - ★ unintegrated gluon dist's



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Factorisation: exclusive or inclusive

- separate dynamics into "hard probe" and "structure of target" nontrivial, requires hard scale Q², comes with corrections ~ (Λ/Q)^k
- inclusive processes: parton distribution in squared amplitude, no mom. transfer on hadron, but can have transv. parton momentum
 - PDF factorization
 - ★ inclusive processes
 - ★ p_T ~ hardest scale or unmeasured



- TMD factorization
 - 🛨 inclusive
 - ★ p_T « hardest scale



- GPD factorization
 - ★ exclusive processes
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Longitudinal momentum: moments and sum rules

Iongitudinal parton mom. fraction \leftrightarrow

 $z^{-}=rac{1}{\sqrt{2}}\left(z^{0}-z^{3}
ight)~\sim$ long. position in frame where hadron moves fast

parton distributions defined via

$$f(x) \sim \int dz^- e^{ixp^+z^-} \bar{\psi}\left(-\frac{1}{2}z^-\right) \dots \psi\left(\frac{1}{2}z^-\right)$$

display only longitudinal variables (boldface), omit transverse ones

Iowest moments:

- n = 1, unpolarised quarks → vector current long. pol. quarks → axial vector current
- n = 2, unpolarised quarks or gluons \rightsquigarrow energy-momentum tensor
 - → momentum, angular momentum (Ji's sum rule '96) pressure and shear forces Polyakov '03, talk by F X Girod

note: gravitons couple directly to energy-momentum tensor, but cannot tell a quark from a gluon

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Parton correlation functions and their descendants



- graph for unpolarised partons, with polarisation even more structure
- subtleties of resolution scale dependence are understood, not shown here

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Parton correlation functions and their descendants



▶ several of these functions can be computed in lattice QCD \rightarrow talk by M Constantinou

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Exclusive processes and GPDs: transverse position

▶ DVCS and meson production ~→ generalised parton distrib's



similar theory as for usual parton densities have factorisation proofs, evolution in resolution scale Q

longit. mom. transfer \rightsquigarrow two parton mom. fractions $x \pm \xi$

- at LO in α_s measure ${\sf GPD}(x,\xi=x,{\bf \Delta})$
- in general x "smeared" around ξ
- separate dep'ce on x and ξ from scaling violations in Q^2
 - need largest possible Q^2 range
- ▶ imaging: measure △ and Fourier transform to b

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Exclusive processes and GPDs: transverse position

▶ DVCS and meson production ~→ generalised parton distrib's



▶ '1st stage' imaging: amplitude $\rightarrow_{\text{FT}} \text{GPD}(x, \xi = x, b)$

$$x \qquad x' = 0$$

no probability interpretation, but $\boldsymbol{b} =$ well defined transverse distance

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Exclusive processes and GPDs: transverse position

▶ DVCS and meson production ~→ generalised parton distrib's



► '2nd stage': $GPD(x, \xi = x, b) \rightarrow GPD(x, \xi = 0, b)$

- density interpretation: $GPD(x, \xi = 0, b) = f(x, b)$
- access only via $lpha_s$ effects $\rightsquigarrow Q^2$ dependence
- extrapolation to $\xi = 0$ depends on theoretical assumptions

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Correlations between x and b

Introduction lattice calculations (moments ∫ dx xⁿ f(x, b_T) with n = 0, 1, 2) find significant correlation between b_T and x average x in moments ~ 0.2 to 0.4 ↔ size of "valence" configurations

direct measurement in scattering experiments? ~> JLab@12 GeV

• at large b prediction from chiral dynamics $f(x,b) \sim e^{-\kappa b}/b$ with $\kappa \sim 2m_{\pi} = (0.7 \,\text{fm})^{-1}$

Strikman, Weiss '09

sets in for $x \leq m_{\pi}/m_p$ requires precise measurement at low Δ first glimpse from COMPASS DVCS full potential at EIC



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Correlations between x and b

• at small x < 0.01 find $\langle b^2 \rangle \propto \text{const} + \alpha' \log \frac{1}{x}$





• for gluons $lpha'\sim 0.15\,{
m GeV^{-2}}$ from HERA J/Ψ prod'n

much smaller than in soft hadronic procs.

• for valence quarks $(q - \bar{q})$ get $\alpha' \sim 0.9 \, {\rm GeV}^{-2}$ in GPD models fitted to e.m. form factors MD et al '04; Guidal et al '04; MD, Kroll '13

in line with meson Regge trajectories

direct measurement in scattering experiments?

value for sea quarks? cross talk with gluons?

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Spin and orbital angular momentum

- GPD $E \leftrightarrow$ nucleon helicity flip $\langle \downarrow | \mathcal{O} | \uparrow \rangle$
 - →→ interference between wave fcts. with L^z and $L^z \pm 1$ no direct relation with $\langle L^z \rangle$, but indicator of large L^z
- ▶ helicity flip ↔ transverse polarisation asymmetry parton dist's in proton polarised along x are shifted along y:

$$f^{\uparrow}(x, \mathbf{b}) = f(x, b^2) - \frac{b^y}{m} \frac{\partial}{\partial b^2} e(x, b^2)$$

 $e(x, b^2) =$ Fourier transform of $E(x, \xi = 0, \Delta_T)$

▶ connection to orbital angular momentum via *b* × *p* ▶ shift known to be large for valence combinations *u* − *ū*, *d* − *d* from sum rule connecting with magnetic moments of *p* and *n* unknown for sea guarks and gluons

Burkardt '02, '05; Burkardt and Schnell '05

E key part of Ji's angular momentum sum rule

$$2J^{q} = \int dx \, x[q(x) + \bar{q}(x)] + \int dx \, x[e^{q}(x) + e^{\bar{q}}(x)]$$

and its analogue for gluons



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Exclusive processes

 \rightarrow talk by C Hyde

Compton scattering: deeply virtual (DVCS) or timelike (TCS)

$$ep \to ep\gamma \text{ or } \gamma p \to \ell^+ \ell^- p$$

• best theory control: NNLO, twist three, corr's in m^2/Q^2 and t/Q^2

Müller et al., Braun and Manashov

- interference with Bethe-Heitler process (calculable)
 - \rightarrow phase of Compton amplitude





meson production

- many channels, separation of quark flavors and gluons
- theory more involved: meson wave fct. NLO and 1/Q corrections can be large
- strong indications that need $Q^2\gtrsim {\rm several}~10\,{\rm GeV}^2$ to reach factorisation regime

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JLab@12 GeV and EIC

- different x ranges \rightsquigarrow different configurations in nucleon
- special opportunity at 12 GeV: very rare processes double DVCS
 - $\gamma^*p \to \gamma^*p$ measured in $ep \to e + \mu^+\mu^- + p$
 - two independent photon virtualities
 → more detailed access to longitudinal kinematics
 → disentangle x and ξ in GPDs
- special opportunities at EIC
 - high $Q^2 \rightsquigarrow$ greater lever arm for scale evolution, cleaner theory for meson production
 - production of J/Ψ and $\Upsilon \; \rightsquigarrow$ selects gluons in target diminished thy. uncertainty from meson wave function

rate estimates cf. e.g. talk by S Joosten at QCD Evolution 2018

at 12 GeV J/Ψ prod'n close to threshold \rightsquigarrow dynamics not described by GPDs

Special opportunity wherever can be realised: positron beam DVCS beam charge asymmetry → clean access to Compton amplitude via interference with Bethe-Heitler

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$Transverse momentum dependent distributions \quad \rightarrow {\sf talk \ by \ H \ Gao}$

- theoretical description of transv. momenta in final state:
 - for small transv. momenta (or transv. momentum differences) described by transv. mom. dependent distributions etc.
 - if large then generate perturbatively = hard radiation
 - graphs for SIDIS $ep \rightarrow eh + X$:



analogous for Drell-Yan process $pp \to \ell^+\ell^- + X$

no sharp boundary between "intrinsic" and "radiative" regimes but transition between the two interesting and practically relevant basic theory well understood, open questions remain

cf. e.g. Bacchetta et al '08; Collins et al '16; Gamberg et al '17

 radiative corrections known at NLO, in parts also NNLO partial understanding of 1/Q corrections

cf. e.g. talk by A Vladimirov at QCD Evolution 2018

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The importance of gluons "accompanying" a parton



- in general, coloured objects are surrounded by gluons profound consequence of gauge invariance technically implemented in Wilson lines
- k_T dep't distributions can be time reversal odd
 e.g. Sivers function: unpol. quarks in proton pol. along x axis:

$$f^{X}(x, \mathbf{k}_{T}) = f(x, k_{T}^{2}) + \frac{k^{y}}{M} f_{1T}^{\perp}(x, k_{T}^{2})$$

Sivers fct. has opposite sign when gluons couple after quark scatters (SIDIS) or before quark annihilates (DY) would be zero if gluons were absent

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fragment'n fct's: similar dynamics, with important differences

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Orbital angular momentum again



 Sivers fct. ↔ proton helicity flip
 → interference of config's with L^z and L'_z = L^z ± 1 another indicator of L^z

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Orbital angular momentum again





- chromodynamic lensing: transverse shift in b space (described by E)
 → transverse shift in k_T (described by f[⊥]_{1T})
 - generated by gluon exchange, opposite signs for SIDIS and DY
 - no calculation in full QCD (is highly nonperturbative) but seen in model calculations

should test experimentally for different x and different parton species

both E and f[⊥]_{1T} exist for quarks and gluons could become sizeable at small x by parton splitting, provided that are not small at low scale/low k_T

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JLab@12 GeV and EIC

- ▶ different x ranges \rightsquigarrow different configurations in nucleon
- ► challenge at 12 GeV: factorisation for SIDIS $ep \rightarrow eh + X$ requires large invariant hadronic mass: $m_X \gg$ resonance masses



$$m_X^2 = Q^2 \frac{x_B}{1 - x_B} (1 - z) - \frac{1}{z} P_{h\perp}^2$$

+hadron mass terms

detailed analysis: Boglione et al '17

at EIC: several benefits from larger phase space:

- larger $Q^2 \rightsquigarrow$ higher theory precision can study scale evolution
- large range of hadron transv. mom. $P_{h\perp}$
 - \rightsquigarrow study transition between "intrinsic" and "radiative" regimes

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Wigner functions

- relatively young field, many aspects remain to be worked out
- theory tool to relate GPDs and TMDs, exhibit unifying aspects work by Metz et al; Hatta et al; Lorcé and Pasquini; Rajan et al

▶ direct measurement possible? proposal: exclusive electroproduction of dijets $ep \rightarrow e + jet 1 + jet 2 + p$

Hatta, Xiao, Yuan '16

- access to generalised TMDs (x, q, ξ, Δ)
- open questions regarding theory quite sure that need γ*, not photoproduction and experimental realisation
- requires EIC at high energy

also discussed in small-x framework and in collinear factorisation (GPDs)



plot: Y Hatta at POETIC 2018

Altinoluk et al '15; Boussarie et al '16 Braun, Ivanov '05

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Summary

 exclusive and semi-inclusive procs. with measured transv. momenta
 study trv. parton position and momentum in quantitative, theoretically controlled ways

 relates to deep concepts and questions in QCD some concrete, others more generic

- interplay of pert. and nonpert. phenomena
 - radiatively generated vs. nonpert. sea, flavour and spin strct.
 - transition from small to large k_T
- spatial distribution of partons in hadron \leftrightarrow confinement
- role of π fluctuations \leftrightarrow chiral dynamics
- spin-orbit correlations (k_T or b_T vs. polarisation)
 → orbital angular momentum
- dynamics of gluons that accompany any coloured particle
 → gauge symmetry, Wilson lines

have common theory framework to interpret meas'ts at 12 GeV and at EIC pointed out specific differences, but see coherent science programme for the two facilities