

Topological susceptibility from lattice QCD



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Quantum Chromodynamics (QCD)

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}(i\not{D} - m_q)\psi - \frac{1}{4}G_{\mu\nu}G^{\mu\nu}$$

- parameters are the **quark masses** m_q and the **dimensionless gauge coupling**,
- in the chiral limit $m_q \rightarrow 0$, a scale is generated through **dimensional transmutation**,
- all dimensionful quantities can be expressed in units of **one characteristic scale**, e.g. the proton mass,

There is another mass generating mechanism:

- spontaneous breaking of chiral symmetry produces eight (or three) massless Goldstone bosons,
- while the flavour-singlet meson η' (or η_2) becomes massive.
- Mass is generated through **topological charge fluctuations**.
- This non-perturbative effect is mediated through disconnected contributions.

- For N_f degenerate fermions of mass m_q we have

$$N_f Q = m_q \sum_x \bar{\psi} \gamma_5 \psi = \sum_x \frac{m_q \gamma_5}{D + m_q}$$

in the limit as $m_q \rightarrow 0$.

- For twisted mass fermions we have in the twisted (lattice) basis $Z_P \bar{\psi} \gamma_5 \psi \rightarrow i Z_S \bar{\chi} \tau_3 \chi$ and hence

$$N_f \rho(x) = i \frac{Z_S}{Z_P} \mu_q \bar{\chi} \tau_3 \chi = \frac{Z_S}{Z_P} \frac{i \mu_q \tau_3}{D_W + i \mu_q \tau_3 \gamma_5}$$

- Note that $Z_S/Z_P = Z_\mu Z_S$ since $Z_\mu Z_P = 1$.
- The topological charge is the sum of that density over the whole lattice.

- Density is the same (up to a proportionality factor) that we use for the *disconnected correlator* of the flavour-singlet η' .
- If we define $\mu_q d(t) = \sum_{\mathbf{x}} \rho(\mathbf{x})$ at time t , then the disconnected correlator is

$$D(t) = (Z_P/Z_S)^2 \sum_{t'} d(t')d(t+t')/V$$

- The topological susceptibility is

$$\chi = \frac{1}{V} \sum_{x,y} \rho(x)\rho(y) = \frac{\mu_q^2}{V} \sum_{t,t'} d(t)d(t')$$

and hence

$$\chi = (Z_S/Z_P)^2 \mu_q^2 \sum_t D(t).$$

- Vice versa, the disconnected η' correlator can be calculated from the (gluonic) charge density correlator.

- The full η' correlator has connected (C) and disconnected (D) components:

$$C_{\text{tot}}(t) = C(t) - 2D(t) \sim e^{-m(\eta')t}$$

- At small quark mass the connected piece is dominated by the pion

$$C(t) = g_\pi^2 e^{-m(\pi)t} / 2m(\pi),$$

- If the η' is finite in the chiral limit, we have that

$$D(t) = (C(t) - C_{\text{tot}}(t)) / 2$$

is dominated by the ground state pion in $C(t)$.

- Summing this over t gives the top. susceptibility

$$\chi = \frac{1}{2} (Z_S / Z_P)^2 \mu_q^2 g_\pi^2 / m(\pi)^2$$

- For twisted mass pions, at maximal twist and finite a , the connected contributions to π^+ and π^0 are different, however,

$$g_{\pi^0} = \frac{Z_S}{Z_P} g_{\pi^+}.$$

- From the vector Ward identity we have

$$g_{\pi^+} = m(\pi^+)^2 f_{\pi^+} / (2\mu_q),$$

and hence

$$\chi = f_{\pi^+}^2 m(\pi^+)^4 / 8m(\pi_{\text{conn}}^0)^2.$$

- In the continuum limit $m(\pi^+) = m(\pi_{\text{conn}}^0)$, so

$$\chi = f_{\pi^+}^2 m(\pi^+)^2 / 8 \quad .$$

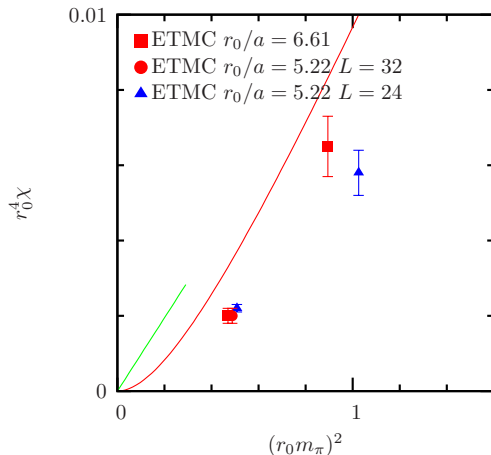
- In practice, connected pion dominance is not fulfilled:

$$\sum_t 2D(t) = 0.65 \sum_t C(t),$$

so chiral formula overestimates χ , e.g.,

$$\begin{array}{ll} r_0^4 \chi = 0.0030 & \text{from chiral formula,} \\ r_0^4 \chi = 0.0020 & \text{from direct evaluation} \end{array}$$

at lowest quark mass at $\beta = 4.05$.



Red line from TM formula for $\beta = 4.05$, green line is continuum χ_{PT} .

- An astonishing property of the charge correlator is

$$\langle q(0)q(\mathbf{x}) \rangle < 0 \quad \text{for } \mathbf{x} > 0$$

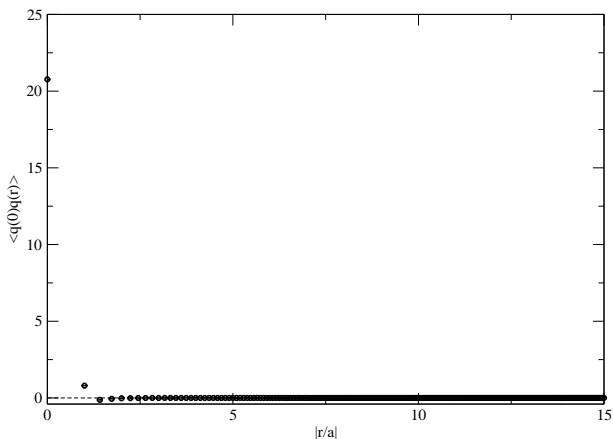
due to reflection positivity, however,

$$\chi = \sum_{\mathbf{x}} \langle q(0)q(\mathbf{x}) \rangle > 0.$$

- Correlator $\langle q(0)q(\mathbf{x}) \rangle$ contains contact term at $\mathbf{x} = 0$ which is (partially) canceled at $\mathbf{x} > 0$.
- Define field theoretic topological charge density via

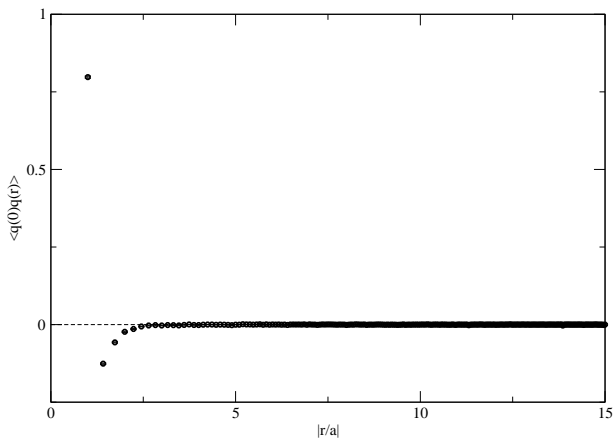
$$q(\mathbf{x}) = \frac{1}{16\pi^2} \text{tr} \left[F_{\mu\nu}(\mathbf{x}) \tilde{F}_{\mu\nu}(\mathbf{x}) \right].$$

- Topological charge correlator from gauge fields:

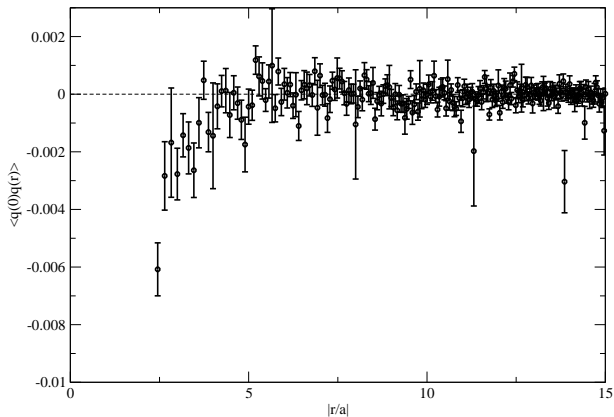


⇒ contact term, smeared out over one lattice spacing.

- Zoom in to see it become negative:

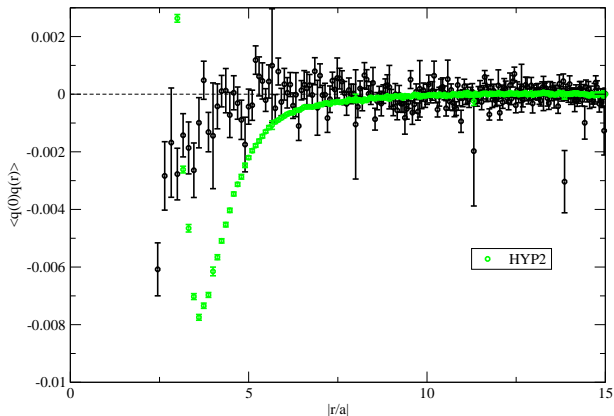


- Sum over space gives the susceptibility, but



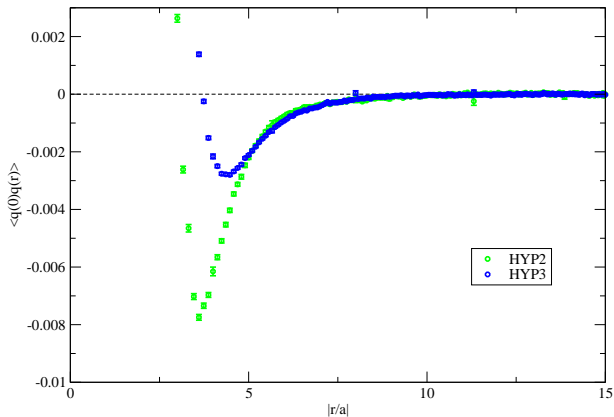
⇒ very noisy due to contributions at large r .

- Smoothing gauge field to suppress ultraviolet fluctuations:



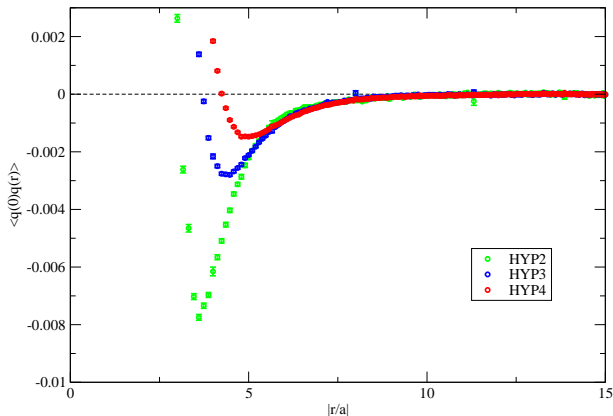
⇒ signal is very much improved.

- More smoothing improves signal further:



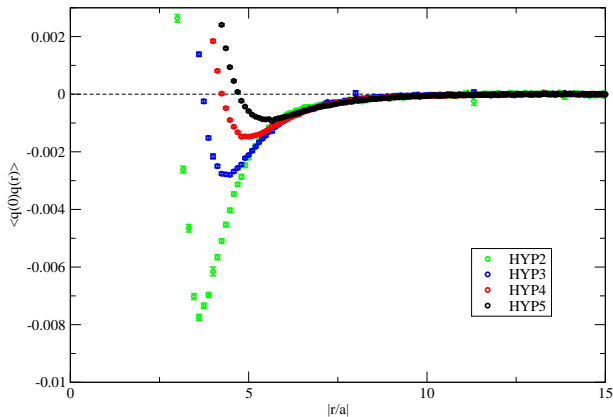
⇒ contact term is smeared over larger range.

- More smoothing improves signal further:



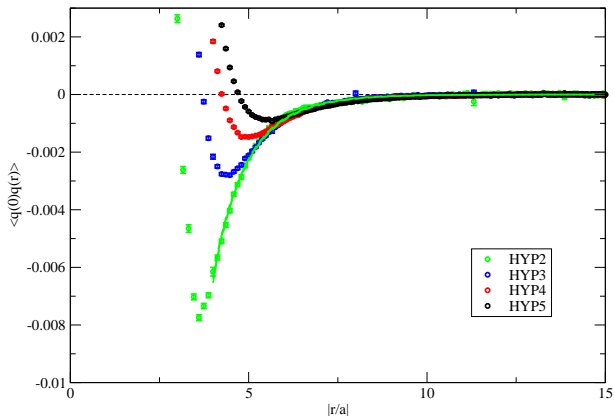
⇒ autocorrelation is essentially absent (self-averaging).

- Signal at large r remains unaffected by smearing:



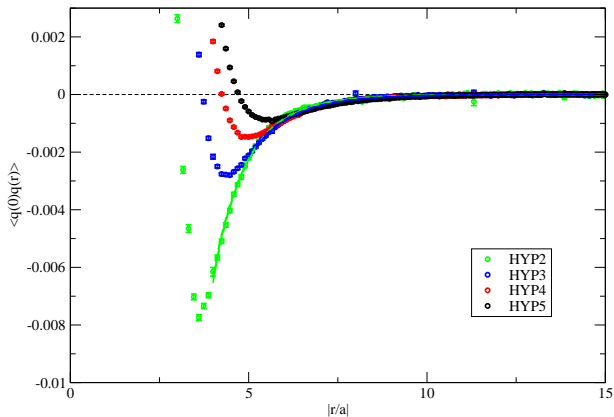
\Rightarrow this is just the disconnected part of the η' -correlator.

- We can fit this by the scalar propagator $m/r \cdot K_1(mr)$:



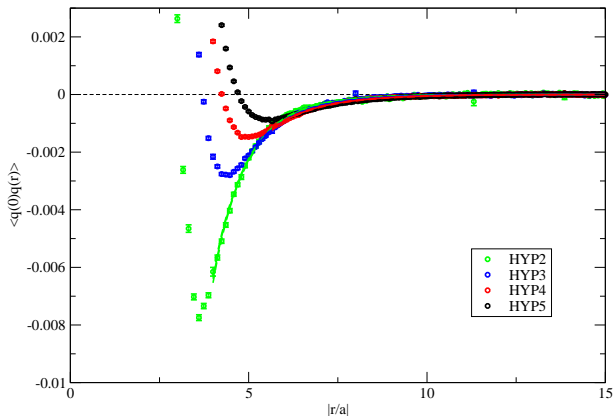
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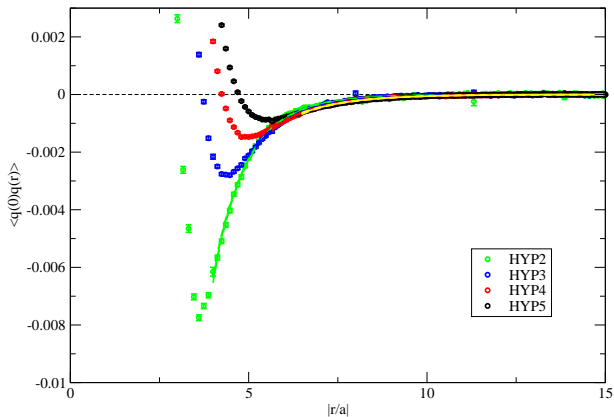
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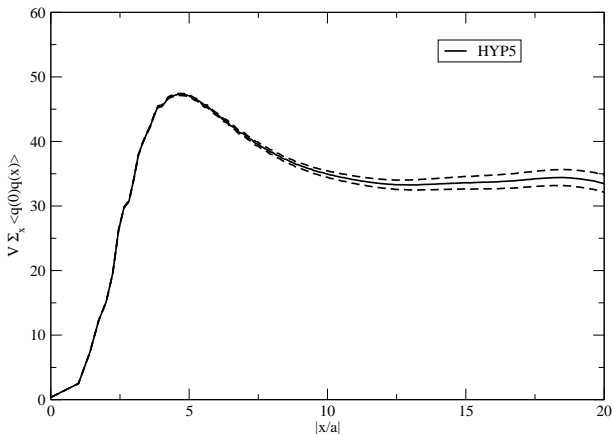
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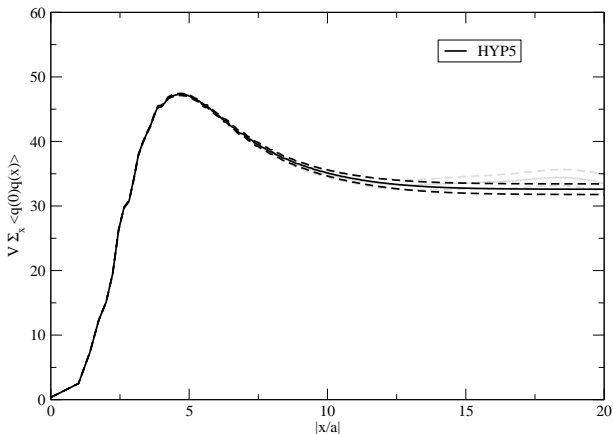
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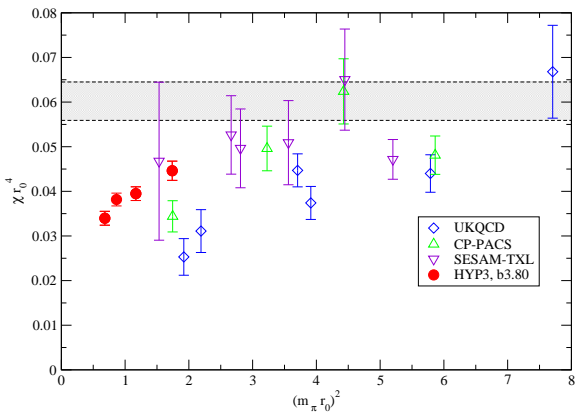


\Rightarrow noise at large r distorts the signal.

- Replacing the long distance with the fitted correlator:

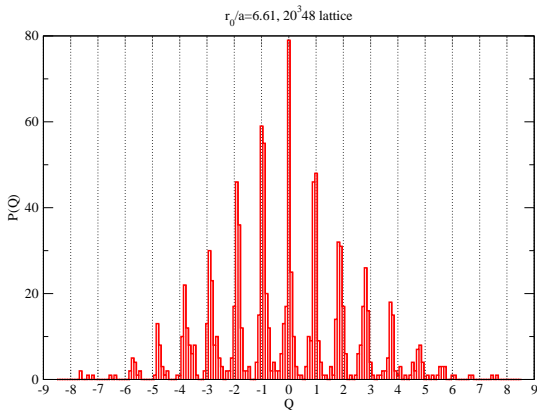


⇒ signal is stabilised and more reliable.



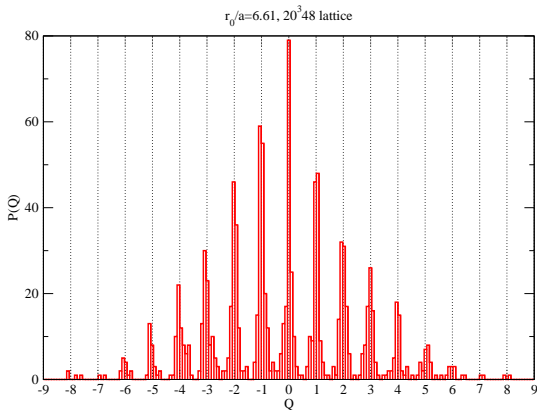
⇒ needs additive and multiplicative renormalisation...

- Renormalisation constants from charge distribution:



⇒ multiplicative renormalisation from shift of peaks, additive renormalisation from width.

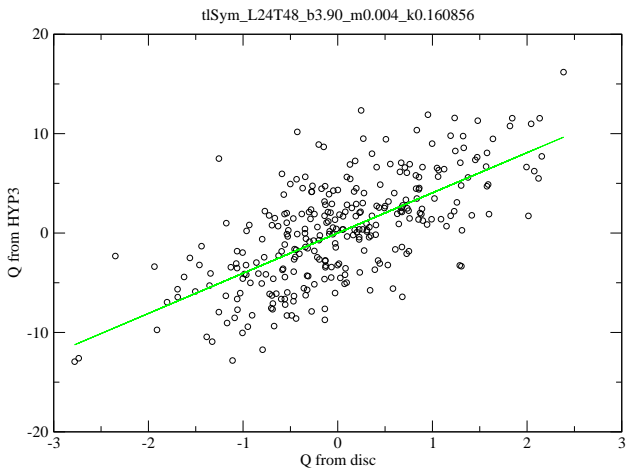
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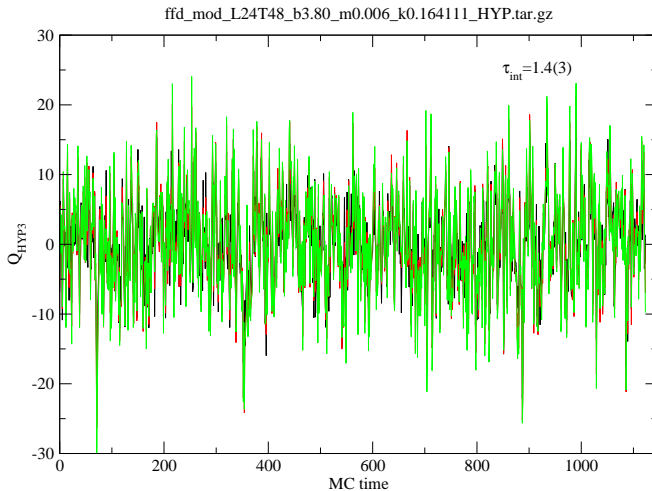


⇒ multiplicative renormalisation from shift of peaks, additive renormalisation from width.

- We have elaborated on the connection of the topological susceptibility and the disconnected η' -correlator.
- We find that χ_t is indeed suppressed towards the chiral limit and vanishes at $m_q \rightarrow 0$.
- Gluonic definition of the charge correlator contains information on the disconnected η' -correlator.
- Gluonic point-to-point correlators can be useful to cheaply extract (fermionic) physics.

Determination of gluonic χ_t is underway, but needs multiplicative and additive renormalisation.





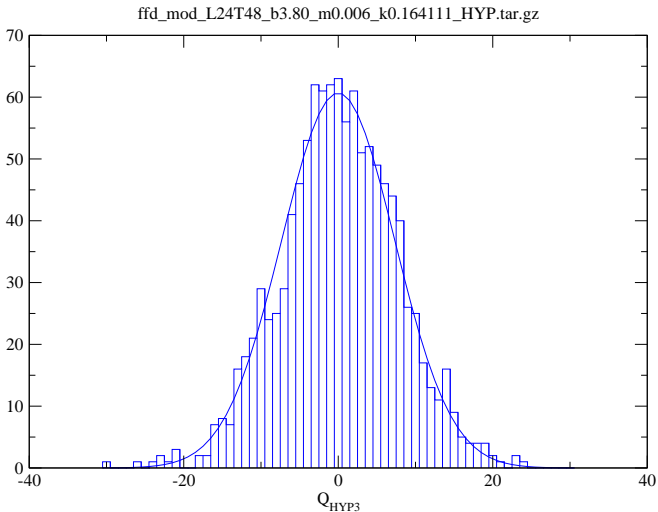


Illustration on 12^4 at $\beta = 6.0$ with HYP3 smearing:

