

# The hadron spectrum in full QCD: Setup and parameter selection

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# Outline

## Why the spectrum?

## Action & Algorithm

Action

Locality

Algorithm & Stability

## Simulation parameters

Parameter selection

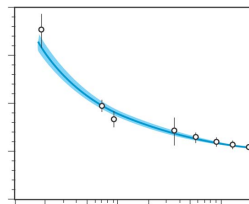
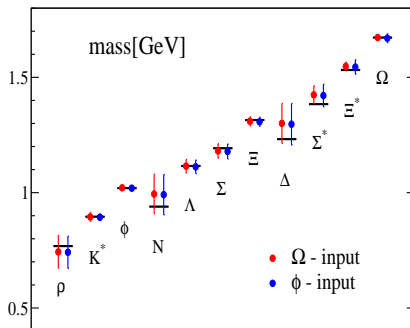
Simulation points

## Analysis

Analysis details

# Status

- ▶ Asymptotic freedom: good agreement between theory and experiment



- ▶ Good evidence that QCD describes the strong interaction in the non-perturbative domain (e.g. CP-PACS '07,  $N_f=2+1$ ,  $210\text{MeV} \leq M_\pi \leq 730\text{MeV}$ ,  $a \simeq 0.087\text{ fm}$ ,  $L \lesssim 2.8\text{ fm}$ ,  $M_\pi L \simeq 2.9$ )
- ▶ However, systematic errors **not** yet under control

# Controlling systematic errors

- ▶ Include  $u$ ,  $d$  and  $s$  quarks into the simulation with an action, whose universality class is QCD.
- ▶ Use large volumes ( $M_\pi L \gtrsim 4$ ) to guarantee small finite-size effects and at least one simulation at a significantly larger volume to confirm the smallness of these effects.
- ▶ Use controlled interpolations and extrapolations of the results to physical  $m_{ud}$  and  $m_s$
- ▶ Use controlled extrapolations to the continuum limit, requiring that the calculations be performed at no less than three values of the lattice spacing.

## Choice of action

- Explicit form of our action (thin links  $U_{n,\mu}$ , smeared links  $V_{n,\mu}$ )

$$\begin{aligned}
 S &= S_G^{\text{Sym}} + S_F^{\text{SW}} \\
 S_G^{\text{Sym}} &= \beta \left[ \frac{c_0}{3} \sum_{\text{plaq}} \text{Re Tr} (1 - U_{\text{plaq}}) + \frac{c_1}{3} \sum_{\text{rect}} \text{Re Tr} (1 - U_{\text{rect}}) \right] \\
 S_F^{\text{SW}} &= S_F^{\text{W}}[V] - \frac{c_{\text{SW}}}{4} \sum_n \sum_{\mu,\nu} \bar{\psi}_n \sigma_{\mu\nu} F_{\mu\nu,n}[V] \psi_n,
 \end{aligned}$$

- with  $c_{\text{SW}} = 1$ ,  $c_1 = -1/12$ ,  $c_0 = 1 - 8c_1 = 5/3$ . Smearing:

$$\begin{aligned}
 V^{(n+1)} &= e^{\rho S^{(n)}} U^{(n)}, \\
 S^{(n)} &= \frac{1}{2} (\Gamma^{(n)} V^{(n)\dagger} - V^{(n)} \Gamma^{(n)\dagger}) - \frac{1}{6} \text{Re Tr} (\Gamma^{(n)} V^{(n)\dagger} - V^{(n)} \Gamma^{(n)\dagger}) \\
 \Gamma_{n,\mu}^{(n)} &= \sum_{\nu \neq \mu} V_{n,\nu}^{(n)} V_{n+\nu,\mu}^{(n)} V_{n+\mu,\nu}^{(n)\dagger}
 \end{aligned}$$

# Locality properties of the action

- ▶ locality in position space:

$|D(x, y)| < \text{const } e^{-\lambda|x-y|}$  with  $\lambda = O(a^{-1})$  for all couplings.

Our case:  $D(x, y) = 0$  as soon as  $|x - y| > 1$

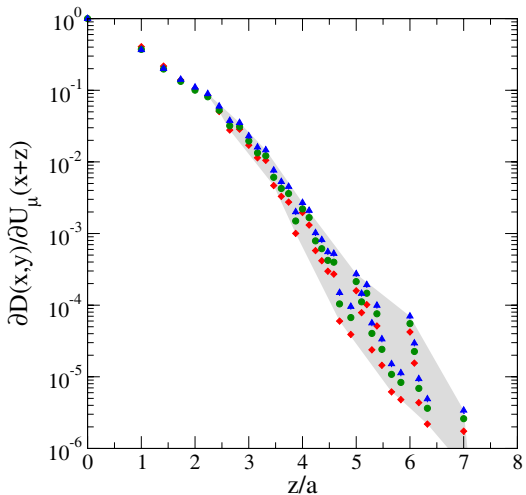
(despite 6 smearings).

- ▶ locality in space of gauge fields:

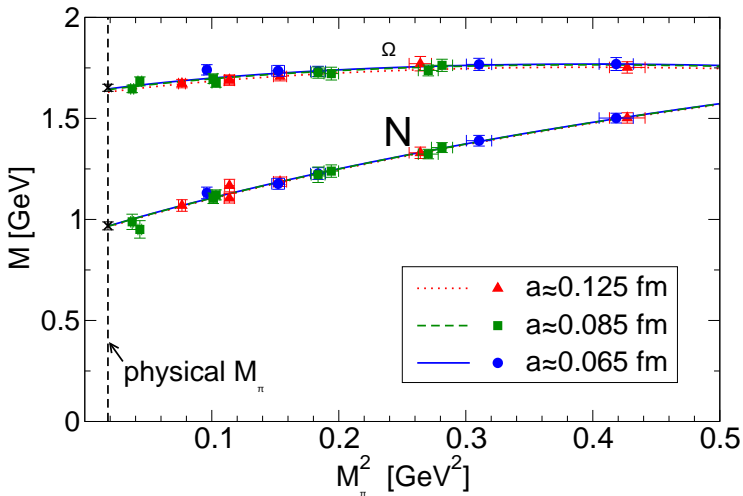
$|\delta D(x, y) / \delta A(z)| < \text{const } e^{-\lambda|(x+y)/2-z|}$  with  $\lambda = O(a^{-1})$  for all

couplings.

# Locality properties of the action



# Scaling of this action





# Dynamical fermion Algorithm

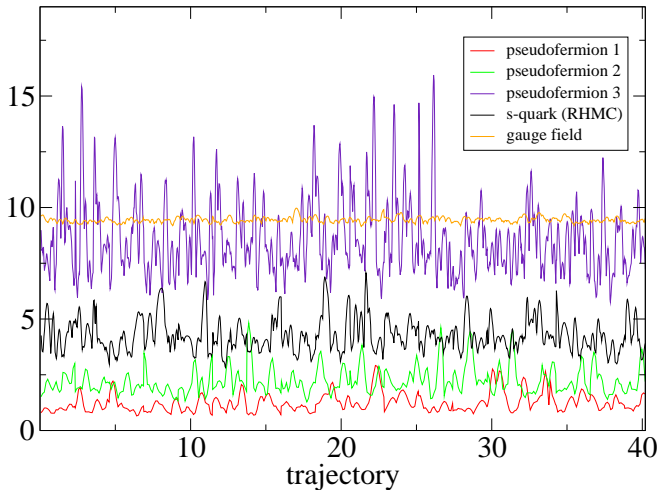
## ▶ Action

- ▶ Clover tree level improved Wilson fermions
- ▶ Symanzik improved gauge action
- ▶ Stout links

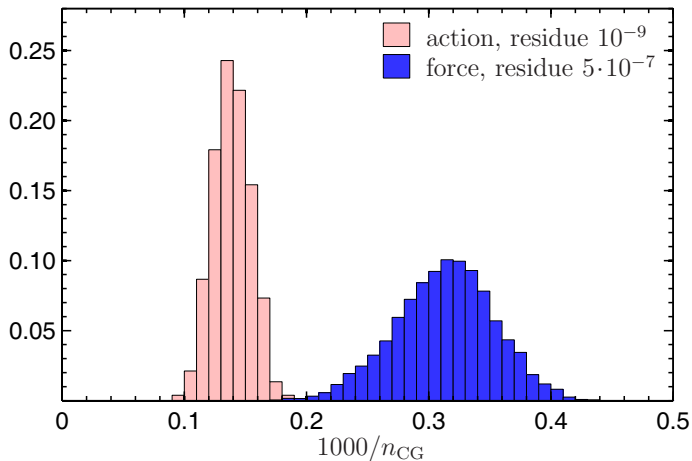
## ▶ Algorithm

- ▶ Rational HMC for strange quark
- ▶ Mass preconditioning (“Hasenbusch trick”)
- ▶ Multi scale integration scheme (“Sexton-Weingarten”)
- ▶ Omelyan integrator (“non equidistant leap frog”), increasing integration precision
- ▶ Mixed precision inverters

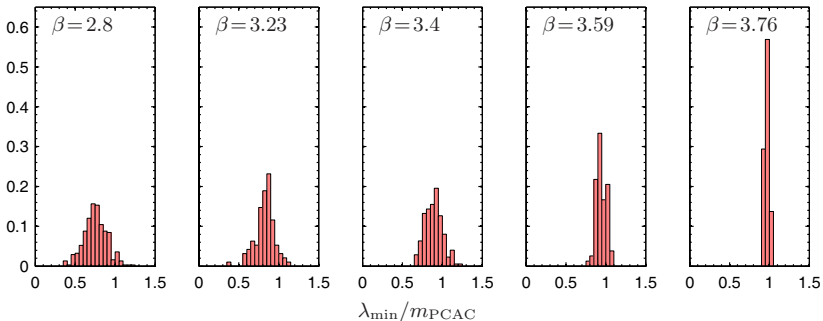
# Fermionic force history



# Inverse iteration count distribution

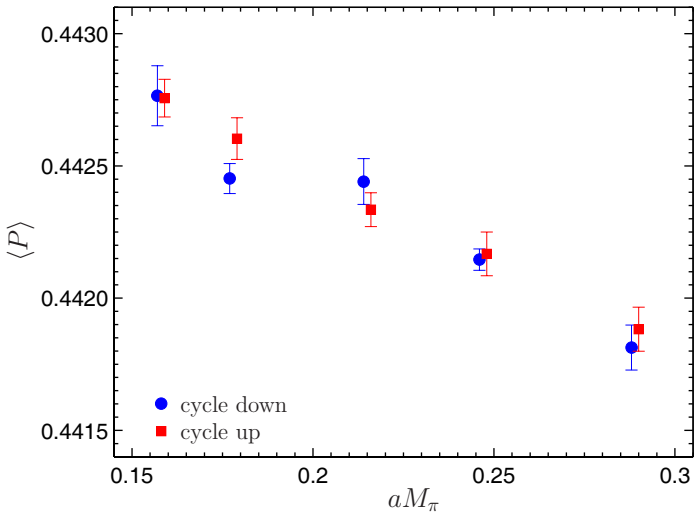


# $\lambda_{min}^{-1}$ distribution



→ Simulations

# “Thermal cycle”

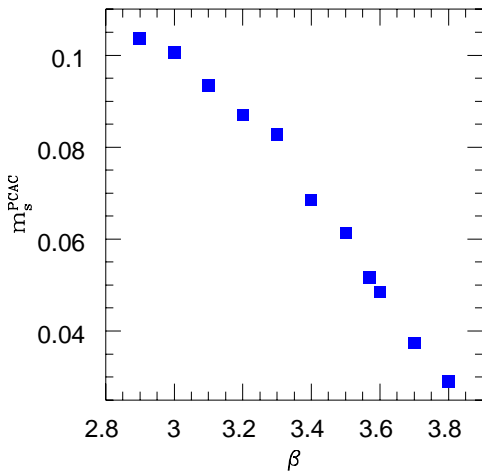


# Setting the strange quark mass

- ▶ We use a  $N_f = 3$  simulation to set the strange quark mass
- ▶ For each beta, search for the  $\kappa$  where

$$\frac{m_{ps}}{m_V} = \frac{\sqrt{2m_K^2 - m_\pi^2}}{m_\phi}$$

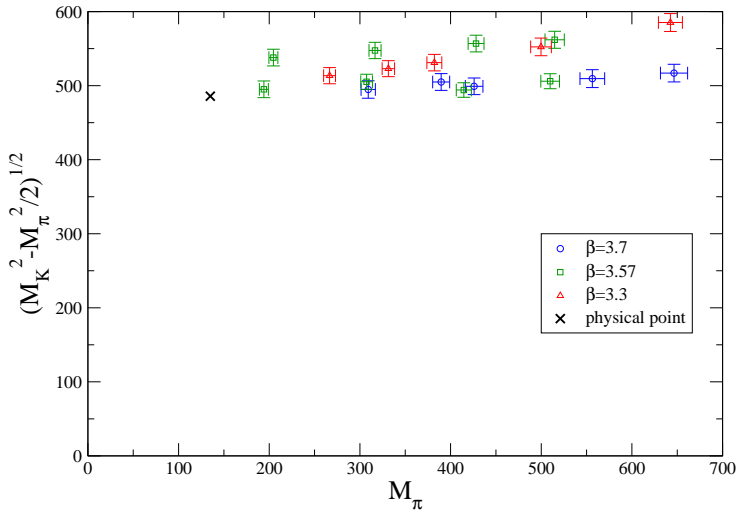
- ▶ We determined the  $\beta$  dependency in the range ( $\beta = 2.9 \dots 3.8$ )



# Simulation points

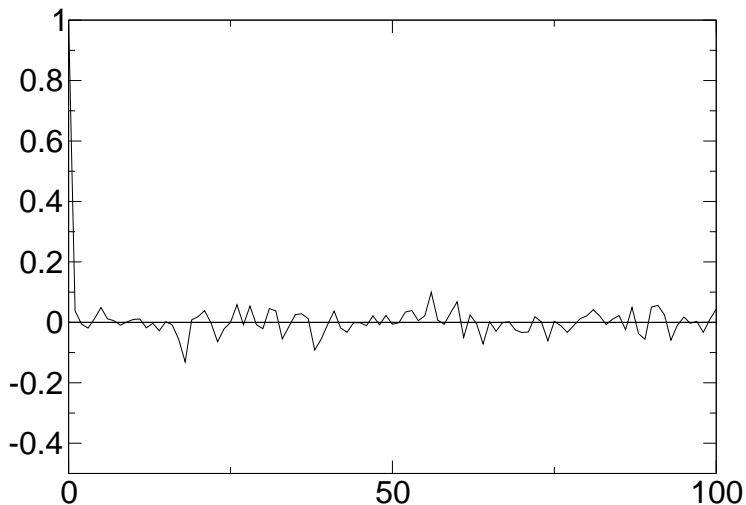
$\beta$	$am_{ud}$	$M_\pi$ [GeV]	$am_s$	$L^3 \times T$	# traj.
3.3	-0.0960	.55	-0.057	$16^3 \times 32$	10000
	-0.1100	.45	-0.057	$16^3, 32^3 \times 32$	1450,1800
	-0.1200	.36	-0.057	$16^3 \times 64$	4500
	-0.1233	.32	-0.057	$16^3, 24^3, 32^3 \times 64$	5000,2000,1300
	-0.1265	.26	-0.057	$24^3 \times 64$	2100
3.57	-0.0318	.46,.48	0.0, -0.01	$24^3 \times 64$	3300
	-0.0380	.39,.40	0.0, -0.01	$24^3 \times 64$	2900
	-0.0440	.31,.32	0.0, -0.007	$32^3 \times 64$	3000
	-0.0483	.19,.21	0.0, -0.007	$48^3 \times 64$	1500
3.7	-0.007	.58	0.0	$32^3 \times 96$	1100
	-0.013	.50	0.0	$32^3 \times 96$	1450
	-0.020	.40	0.0	$32^3 \times 96$	2050
	-0.022	.36	0.0	$32^3 \times 96$	1350
	-0.025	.29	0.0	$40^3 \times 96$	1450

# “Landscape” plot

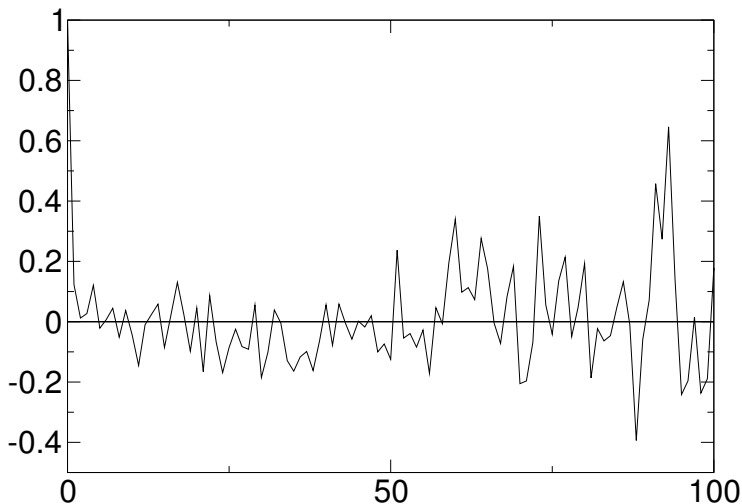




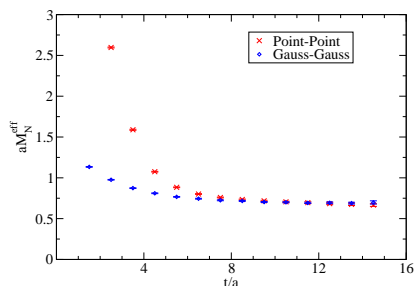
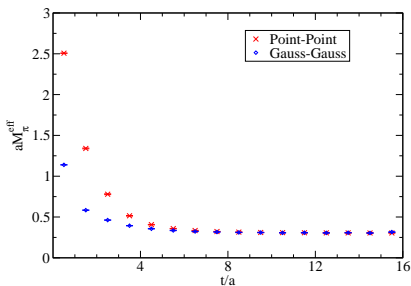
# Nucleon autocorrelation ( $M_\pi = 550$ MeV, $\beta = 3.3$ )



# Pion autocorrelation ( $M_\pi = 190 \text{ MeV}$ , $\beta = 3.57$ )

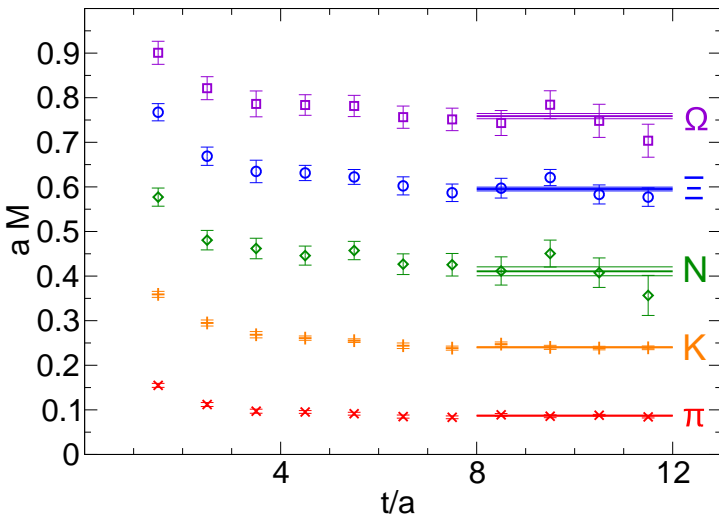


# Sources



- ▶ Gaussian sources  $r = 0.32$  fm
- ▶ Coulomb gauge
- ▶ Gauss-Gauss less contaminated by excited states

# Effective masses and correlated fits



# Setting the lattice spacing via hadron mass

The particle selected should have a mass

1. that is experimentally well known
2. that is independent of light quark mass  $\rightarrow$  large strange content
3. that can be simulated with small statistical errors  $\rightarrow$  octet better suited than decuplet

All points cannot be fulfilled simultaneously, but

- ▶  $\Xi$ : largest strange content of the octet, but still dependent on  $ud$  mass
- ▶  $\Omega$ : member of the decuplet, but largest strange content of particles included in analysis

# Conclusions

- ▶ To control systematic errors in spectrum calculations, we
  1. include dynamical  $N_f = 2 + 1$  degrees of freedom,
  2. use an action in the universality class of QCD,
  3. use lattices with  $M_\pi L \gtrsim 4$ ,
  4. use controlled chiral extrapolations,
  5. and include ensembles at three different  $\beta$  values.
- ▶ Our simulation algorithm is save even at  $M_\pi < 200$  MeV.
- ▶ Our action has a very good scaling behavior.

# Result

