

Wilson Chiral Perturbation Theory for twisted mass QCD at NLO

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Lattice Simulation

- Machine power can simulate QCD at small mass region.
- ETMC's simulations reach $m_q \sim a^2 \Lambda_{\text{QCD}}^3$.

ChPT for twisted mass QCD

- regime : $m_q \sim a \Lambda_{\text{QCD}}^2$
- $O(a^2)$ effects are NLO.

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ChPT for twisted mass QCD

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We need Wilson ChPT at small mass regime, $m_q \sim a^2 \Lambda_{\text{QCD}}$.

Previous Study

- LO: $O(m, a)$
- NLO: LO^2

Our Study

- LO: $O(m, a)$, sub LO(SLO): $O(a^2, am)$
- NLO: LO^2 , $LO \cdot SLO \sim O(a^2 m, am^2)$, $SLO^2 \sim O(a^4, a^3 m, a^2 m^2)$

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- non-trivial phase structure appears at LO.
cf. S.R.Sharpe & R.L.Singleton ('98), S.R.Sharpe & J.M.S.Wu ('04)

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- Vacuum expectation has divergence at 1-loop.

Talk Plan

- 1 Introduction
 - power counting
- 2 Leading Order
 - gap equation and phase structure
 - pion mass
- 3 NLO: 1-loop and renormalization
 - vertexes and 1-loop diagram
 - Local Counter Term
 - renormalization
- 4 Maximal Twist
- 5 Summary

LO Lagrangian

$$\begin{aligned} \mathcal{L}_{\text{LO}} = & \frac{f^2}{4} \langle \partial_\mu \Sigma \partial_\mu \Sigma^\dagger \rangle - \frac{f^2}{4} \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle - \frac{f^2}{4} \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle \\ & + W_{45} \langle \partial_\mu \Sigma \partial_\mu \Sigma^\dagger \rangle \langle (\Sigma - \Sigma_0) \hat{a}^\dagger + \hat{a} (\Sigma - \Sigma_0)^\dagger \rangle \\ & - W_{68} \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle - W'_{68} \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle^2 \end{aligned}$$

- $\Sigma = \Sigma_0^{1/2} \Sigma_{ph} \Sigma_0^{1/2}$, $\Sigma_0 = \exp[i\tau^3 \phi]$, $\Sigma_{ph} = \exp[i\tau^a \pi_a / f]$
- ϕ : vacuum expectation, π : pion fields, τ^a : Pauli Matrix
- $M = 2B_0(m + i\tau^3 \mu)$, $\hat{a} = 2W_0 a$

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- $\tilde{M} = 2B_0(\tilde{m} + i\tau^3 \mu)$, $\hat{a} = 2W_0 a$
- $2B_0 \tilde{m} = 2B_0 m + 2W_0 a$
- $\tilde{W}'_{68} = W'_{68} - W_{68}$

Gap Equation

Gap equation, $d\mathcal{L}_{\text{LO}}^{(0)}/d\phi|_{\phi=\phi_0} = 0$

$$2B_0\tilde{m}\sin\phi_0 - (c_2a^2 - \tilde{c}_2a2B_0\tilde{m})\sin 2\phi_0 = 2B_0\mu\cos\phi_0 + \tilde{c}_2a2B_0\mu\cos 2\phi_0$$

The low energy constants,

- $c_2 = -64\tilde{W}'_{68}W_0^2/f^2, \tilde{c}_2 = 32W_{68}W_0/f^2$

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Non-trivial phase structure appears.

Phase Structure

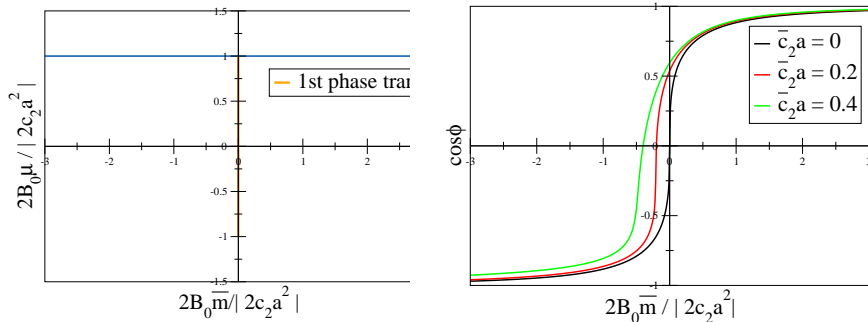


Figure: $c_2 < 0$, $2B_0\mu/|2c_2a^2| = 1$

- This phase structure is similar to previous study at NLO.
cf. S.R.Sharpe & J.M.S.Wo('04)

The $O(\pi^2)$ lagrangian

$$\mathcal{L}_{LO}^{(2)} = \frac{1}{2} \partial_\mu \pi_a \partial_\mu \pi_a + \frac{(m_\pi^a)^2}{2} \pi_a^2$$

The pion mass term

- $(m_\pi^{1,2})^2 = 2B_0 m' - 2c_2 a^2 \cos^2 \phi_0 + 2\tilde{c}_2 a(2B_0 m') \cos \phi_0$
- $(m_\pi^3)^2 = (m_\pi^{1,2})^2 + (\Delta m_\pi^3)^2$
- $(\Delta m_\pi^3)^2 = 2c_2 a^2 \sin^2 \phi_0 + 2\tilde{c}_2 a(2B_0 \mu') \sin \phi_0$
- The short-hand notation for quark mass term,

$$\begin{pmatrix} m' \\ \mu' \end{pmatrix} = \begin{pmatrix} \cos \phi_0 & \sin \phi_0 \\ -\sin \phi_0 & \cos \phi_0 \end{pmatrix} \begin{pmatrix} \tilde{m} \\ \mu \end{pmatrix}$$

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Pion mass splitting.

$O(\pi^3)$ and $O(\pi^4)$ Lagrangian

$$\mathcal{L}_{\text{LO}}^{(3)} = \frac{\lambda_{3p}}{2f} \pi_3 (\partial_\mu \pi_a)^2 + \frac{\lambda_3}{2f} \pi^2 \pi_3$$

$$\mathcal{L}_{\text{LO}}^{(4)} = \frac{\lambda_{4p}}{6f^2} (\pi_a \partial_\mu \pi_a)^2 + \frac{\lambda'_{4p}}{6f^2} \pi^2 (\partial_\mu \pi_a)^2 + \frac{\lambda_4}{6f^2} (\pi^2)^2 + \frac{\lambda'_4}{6f^2} \pi^2 (\pi_3)^2$$

These vertex factor λ is given by,

$$\lambda_{3p} = -c_0 a \sin \phi_0, \quad \lambda_{4p} = 1, \quad \lambda'_{4p} = - \left(1 + \frac{3}{2} c_0 a \cos \phi_0 \right),$$

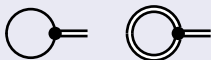
$$\lambda_3 = c_2 a^2 \sin 2\phi_0 - \tilde{c}_2 a (2B_0 m' \sin \phi_0 - 2B_0 \mu' \cos \phi_0),$$

$$\lambda_4 = -\frac{1}{4} 2B_0 m' + 2c_2 a^2 \cos^2 \phi_0 - 2\tilde{c}_2 a 2B_0 m' \cos \phi_0,$$

$$\lambda'_4 = -2c_2 a^2 \sin^2 \phi_0 - 2\tilde{c}_2 a 2B_0 \mu' \sin \phi_0$$

1-Loop Diagram for Vacuum Expectation

tadpole diagram



- — : $\pi_{1,2}$
- = : π_3

1-Loop Diagram for Pion Mass

2-points function 1



2-points function 2



● — : $\pi_{1,2}$

● == : π_3

1-Loop Diagram for Pion Mass

2-points function 1 $\rightarrow LO^2$



2-points function 2 $\rightarrow SLO^2$



• — : $\pi_{1,2}$

• == : π_3

Local Counter Term

$$O(p^2 m), O(m^2) : \# = 2$$

$$\langle \Sigma M^\dagger + M \Sigma^\dagger \rangle \langle L_{\mu\mu} \rangle, \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle^2$$

$$O(ap^2 m), O(am^2) : \# = 6$$

$$\langle L_{\mu\mu} \rangle \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle, \langle L_{\mu\mu} \rangle \langle \hat{a} M^\dagger + M \hat{a}^\dagger \rangle, \\ \langle \partial_\mu \Sigma \hat{a}^\dagger + \hat{a} \partial_\mu \Sigma^\dagger \rangle \langle \partial_\mu \Sigma M^\dagger + M \partial_\mu \Sigma^\dagger \rangle, \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle^2, \\ \langle M M^\dagger \rangle \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle, \langle \hat{a} M^\dagger + M \hat{a}^\dagger \rangle \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle$$

$$O(a^2 p^2), O(a^2 m) : \# = 6$$

$$\langle L_{\mu\mu} \rangle \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle^2, \langle L_{\mu\mu} \rangle \langle \hat{a} \hat{a}^\dagger \rangle, \langle \partial_\mu \Sigma \hat{a}^\dagger + \hat{a} \partial_\mu \Sigma^\dagger \rangle^2, \\ \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle^2 \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle, \langle \hat{a} \hat{a}^\dagger \rangle \langle \Sigma M^\dagger + M \Sigma^\dagger \rangle, \\ \langle \hat{a} M^\dagger + M \hat{a}^\dagger \rangle \langle \Sigma \hat{a}^\dagger + \hat{a} \Sigma^\dagger \rangle$$

where $L_{\mu\nu} = \partial_\mu \Sigma \partial_\nu \Sigma^\dagger$.

of divergence terms

- vacuum expectation : 6
- pion mass term : 13

- # of counter terms (14) < # of divergences terms (19)
- Can not renormalize

of independent divergence terms

- vacuum expectation : 6 \rightarrow 2
- pion mass term : 13 \rightarrow 5

- # of counter terms (14) $>$ # of divergences terms (7)
- mass term reducing by gap eq
 - $2B_0\mu' = -c_2 a^2 \sin 2\phi_0 + \tilde{c}_2 a(2B_0 m' \sin \phi_0 - 2B_0\mu' \cos \phi_0)$
 - $(2B_0 m')(2B_0\mu') = -c_2 a^2(2B_0 m') \sin 2\phi_0 + \tilde{c}_2 a(2B_0 m')^2 \sin \phi_0 + H.O.$
 - $(2B_0\mu')^2 = a(2B_0 m')(2B_0\mu') = a(2B_0\mu')^2 = a^2(2B_0\mu') = H.O.$
- Can renormalize

(twisted) PCAC mass

$$m_{\text{PCAC}} = \frac{\sum_{\vec{x}} \langle \partial_0 A_0^a(\vec{x}, t) P^a(0) \rangle}{2 \sum_{\vec{x}} \langle P^a(\vec{x}, t) P^a(0) \rangle} = 0, \quad a = 1, 2$$

- (twisted) PCAC mass at LO

$$m_{\text{PCAC}} = \frac{(m_\pi^a)_{\text{LO}}^2 \cos \phi_0}{2B_0(1 + \tilde{c}_2 a \cos \phi_0)}$$

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- maximal twist condition

$$\cos \phi_0 = 0$$

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- maximal twist condition

$$\cos \phi_0 = 0$$

- untwisted quark mass

$$2B_0 \tilde{m} = -2B_0 \mu \tilde{c}_2 a$$

pion mass at maximal twist

- $(m_\pi^{1,2})_{\text{LO}}^2 = 2B_0\mu$
- $(m_\pi^3)_{\text{LO}}^2 = 2B_0\mu + 2c_2a^2 - 2(\tilde{c}_2a)^2 2B_0\mu$
- $(m_\pi^{1,2})_{\text{NLO}}^2 = (m_\pi^{1,2})_{\text{LO}}^2 \{1 + C(\mu_{\text{ChPT}}) + \mathcal{C}L_\pi^{1,2} + \mathcal{C}'L_\pi^3\}$

The coefficients is

- $C = -16(m_\pi^a)_{\text{LO}}^2(2L_{68} + L_{45})/f^2 + a^2 X_2/f^2$
- $\mathcal{C} = 0$
- $\mathcal{C}' = \frac{3}{4}$
- $L_\pi^a = \frac{(m_\pi^a)^2}{16\pi^2 f^2} \log\left(\frac{m_\pi^a}{\mu_{\text{ChPT}}}\right)^2$

Summary

- Construct WChPT for twisted mass QCD at the regime, $m_q \sim a^2$.
- It appears the non-trivial phase structure and pion mass splitting at LO.
- At NLO, a vacuum expectation has divergence, but it can be renormalized.
- Construct Pion mass at NLO.

Working Problems

- Decay constant at NLO
- Phase structure at NLO
- Maximal twist condition at NLO