

Moebius Algorithm[†]

for

Domain Wall and GapDW Fermions

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with

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[†] **Brower, Neff and Orginos, Lattice 2004; hep-lat/0409118**

Outline

- Moebius Domain Wall algorithm
- Smaller L_s at fixed m_{res}
 - Quenched $\beta=6.0$ $16^3 \times 32$ lattices
 - DW lattices RBC $16^3 \times 32$ and $24^4 \times 64$
 - Gapped (quenched) lattices
- Elegant formalism*
 - Exact map to overlap at finite L_s
 - Vector and Axial currents.
- Code for BG/L, BG/P,
 - Andrew Polchinsky's Invert plus RHMC
 - Application Toolbox: Techni-color/SUSY/High T(?)

*RCB, Hartmut Neff and Kostas Orginos [hep-lat/0808XXX](#)

Choosing the “best” Lattice Action?

$$Z[U] = \int \mathcal{D}\bar{\psi}\mathcal{D}\psi e^{\beta\text{Tr}[U_P] + S_{\text{improved}}[U] + \bar{\psi}(D_{\text{ov}}[U] + m_f)\psi}$$

*Single Plaquette
Wilson Gauge
(ultra-local)*

*Symansik, tadpole,
Iwasaki, DBW2
(ultra-local), or
Gapped (local ?)*

*Shamir, Borici,
Moebius:
Rational approx for
DomainWall/Pauli-
Villars (local ?)*

Many choices & interact with each other !

Need to implement Overlap Operator (aka Ginsparg Wilson Relation)

$$D_{ov}(m) = \frac{1+m}{2} + \frac{1-m}{2} \gamma_5 \epsilon_L [\gamma_5 D^{kernel}(M_5)]$$

Two steps

ALGORITHM:

Choose Rational approx:

$$e_L[\mathbf{x}] \simeq \mathbf{x}/|\mathbf{x}|$$

PHYSICS:

Choose 4-d "kernel":

$$H_5 \equiv e_5 D(M_5)$$

G-W error operator: $\Delta_L[H] = \frac{1}{4}(1 - \epsilon_L^2[H])$

$$\gamma_5 D_{ov}(0) + D_{ov}(0) \gamma_5 - 2D_{ov}(0) \gamma_5 D_{ov}(0) = 2\gamma_5 \Delta_L$$

Chiral violation for Overlap Action (Kikukawa & Noguchi hep-lat/9902022)

$$(1-m)^{-1} \delta(\bar{\psi} D_{ov}(m) \psi) = m_q \bar{\psi} (\gamma_5 + \hat{\gamma}_5) \psi + 2\bar{\psi} \gamma_5 \Delta_L \psi$$

$$m_q = m/(1-m) \quad \text{and} \quad \hat{\gamma}_5 = \gamma_5(1-2D_0) = -\epsilon_L[H]$$

NOTE : $\frac{D_{ov}(m)}{1-m} = D_{ov}(0) + m_q$ AND $D_{ov}(0) = \frac{1}{2} + \frac{1}{2} \gamma_5 \epsilon_L[H]$

Shamir vs Borici kernels

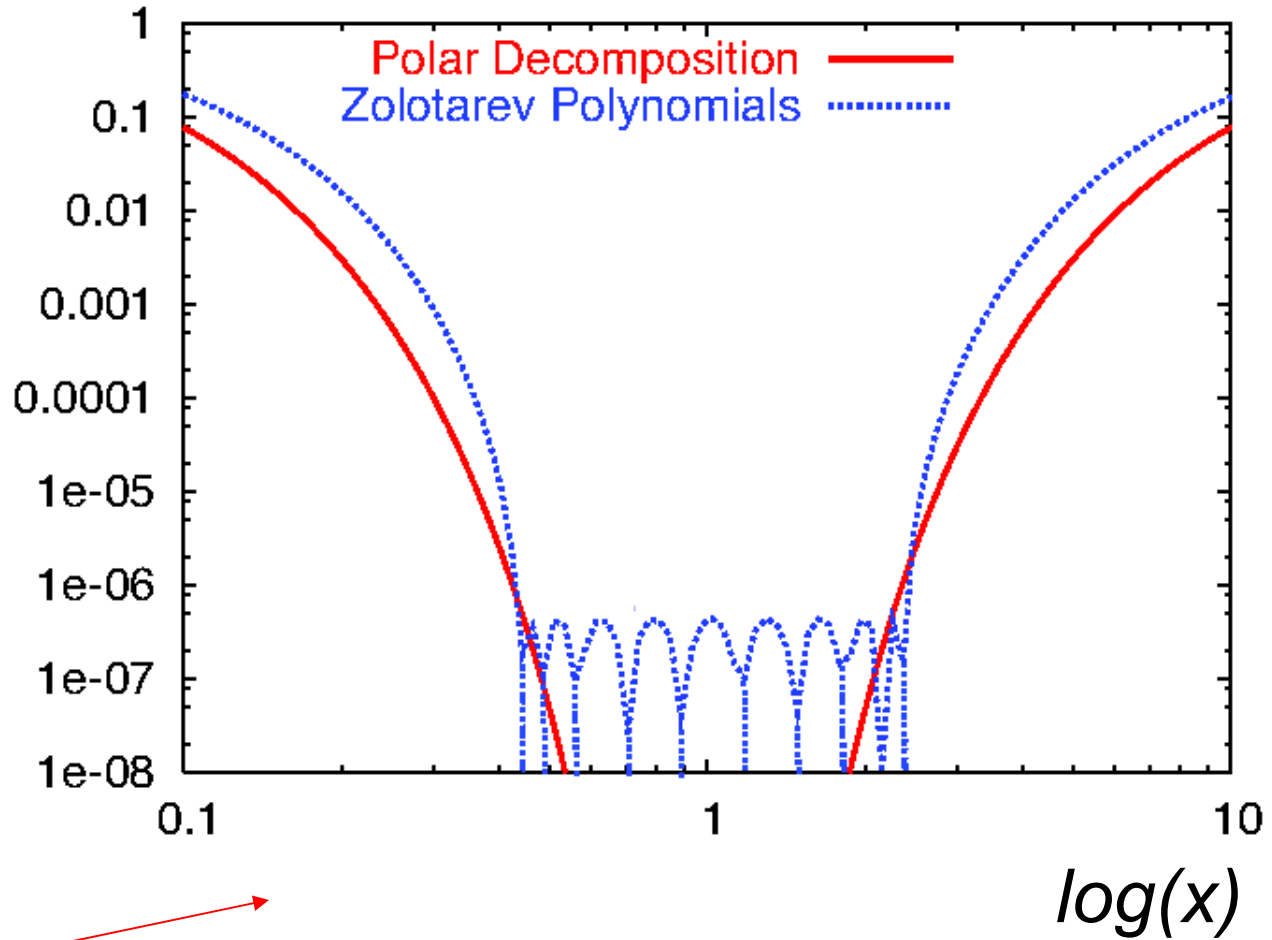
$$\text{Shamir: } D(M_5) = \frac{a_5 D_w(M_5)}{2 + a_5 D_w(M_5)}$$

$$\text{Borici: } D(M_5) = a_5 D_w(M_5)$$

$$\begin{aligned} D_w(M_5)_{xy} &= (4 + M_5)\delta_{x,y} \\ &- \frac{1}{2}(1 - \gamma_\mu)U_\mu(x)\delta_{x+\mu,y} - \frac{1}{2}(1 + \gamma_\mu)U_\mu^\dagger(y)\delta_{x,y+\mu} \end{aligned}$$

Shamir Polar ($L_s = 16$) vs Zolotarev ($L_s = 8$) approx.

W_1 $L[x]$



Can we re-scale window: $\log(\alpha x) = \log(x) + c$?

Moebius Generalization

$$D_{\text{Moebius}}(M_5) = \frac{(b_5 + c_5)D_w(M_5)}{2 + (b_5 - c_5)D_w(M_5)}$$
$$= \alpha D_{\text{Shamir}}(M_5) = \alpha \frac{a_5 D_w(M_5)}{2 + a_5 D_w(M_5)}$$

Parameters: scale: $\alpha = b_5 + c_5$, $a_5 = b_5 - c_5$ and M_5

Since $\epsilon[\alpha x] = \epsilon[x]$, **Moebius is “just” a new “algorithm”**

DW action (with s-dependent terms[†])

$L_s \times L_s$ DW Matrix:

$$D_{s',s}^{DW} = \begin{bmatrix} D_+^{(1)} & D_-^{(1)} P_- & 0 & \dots & -m D_-^{(1)} P_+ \\ D_-^{(2)} P_+ & D_+^{(2)} & D_-^{(2)} P_- & \dots & 0 \\ 0 & D_-^{(3)} P_+ & D_+^{(3)} & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots \\ -m D_-^{(L_s)} P_- & 0 & 0 & \dots & D_+^{(L_s)} \end{bmatrix}$$

$$D_+^{(s)} = b_5(s) D_w(M) + 1, \quad D_-^{(s)} = c_5(s) D_w(M) - 1$$

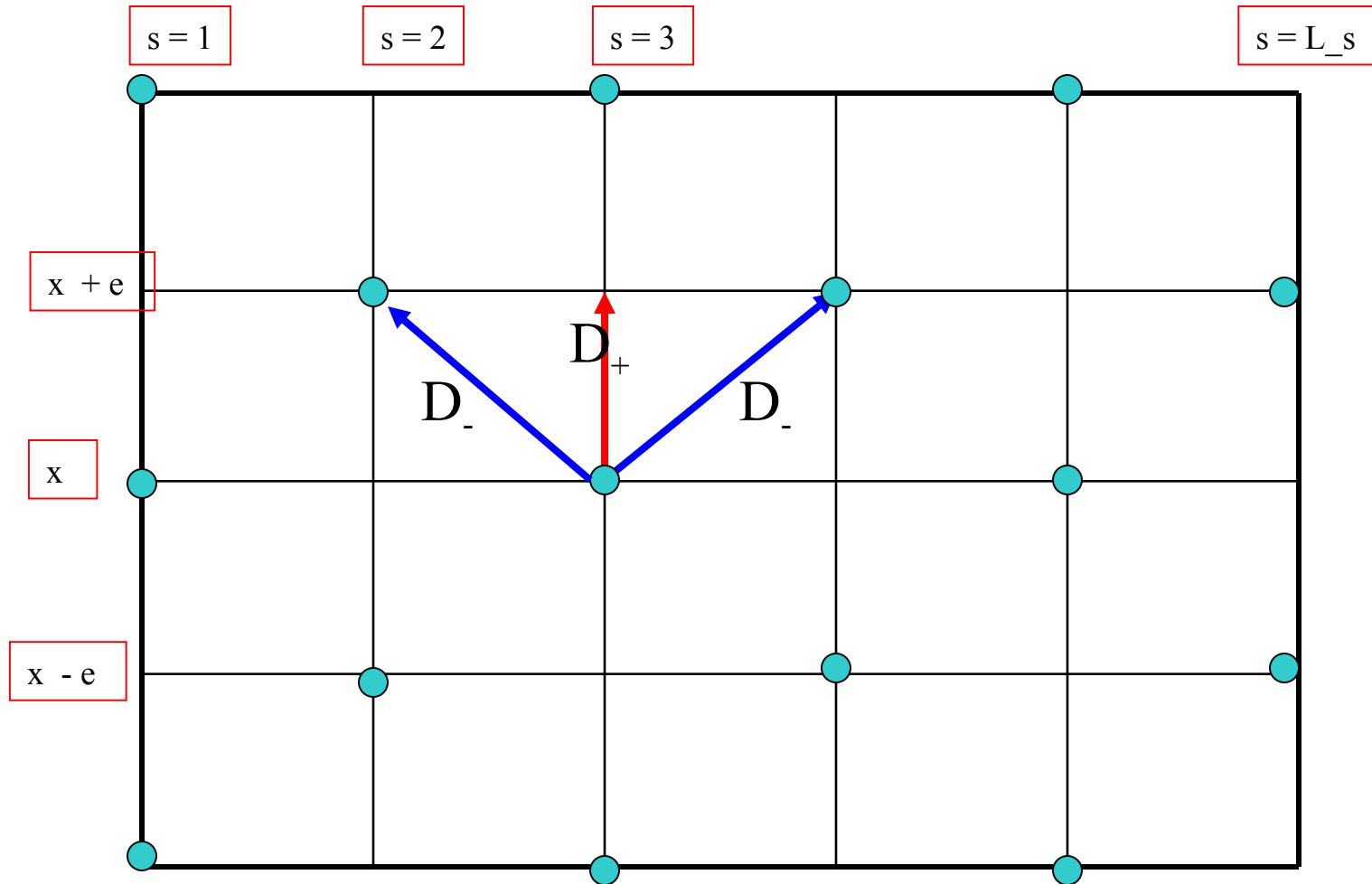
[†] Include for

Zolotarev (Chiu): $\omega(s) = b_5(s) + c_5(s)$ $a_5 = e b_5(s) - c_5(s)$

Fluctuation 5-d fields (like AdS/QCD)

Domain wall filter (Bar, Narayanan, Neuberger, Witzel)

Mobius generalization of Shamir/Borici

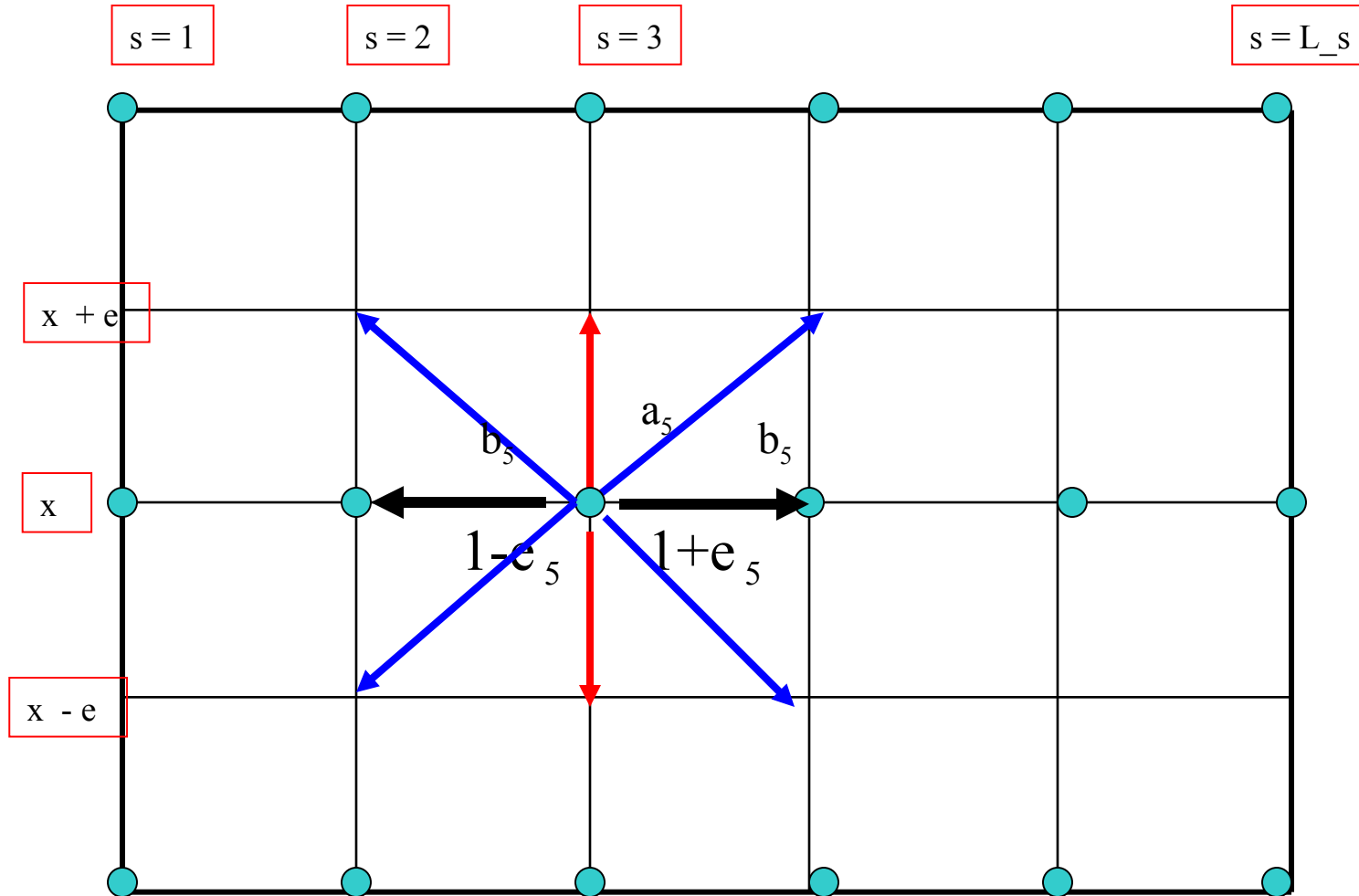


Shamir: $b_5 = a_5, c_5 = 0$

Borici: $b_5 = c_5 = a_5$

$$D_+ = b_5 D_w(M) + 1, \quad D_- = c_5 D_w(M) - 1$$

Modified Even/Odd \rightarrow 4-d Checkerboard



Code in Chroma (R. Edwards) and QOP (A. Polchinsky)

Even/Odd Partition of Matrix

$$D_w(M) = \begin{bmatrix} I_{ee} & D_{eo}^{DW'} \\ D_{oe}^{DW'} & I_{oo} \end{bmatrix}$$

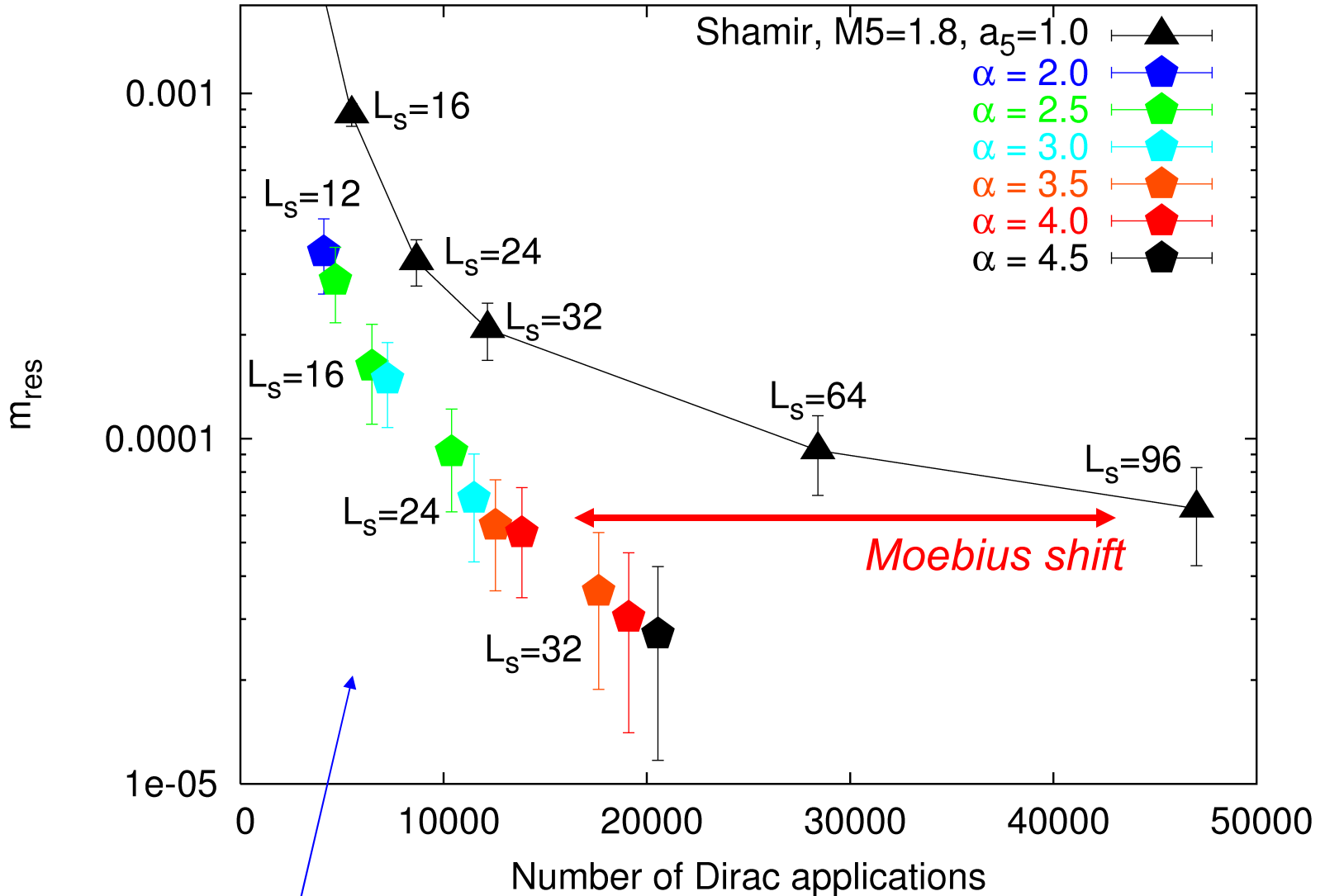
The Schur decomposition

$$D_w(M) = \begin{bmatrix} 1 & 0 \\ D_{oe}^{DW'} I_{ee}^{-1} & 1 \end{bmatrix} \begin{bmatrix} I_{ee} & 0 \\ 0 & I_{oo} - D_{oe}^{DW'} I_{ee}^{-1} D_{eo}^{DW'} \end{bmatrix} \begin{bmatrix} 1 & I_{ee}^{-1} D_{eo}^{DW'} \\ 0 & 1 \end{bmatrix}$$

$$D_{preconditioned}^{DW} = 1 - I_{oo}^{-1} D_{oe}^{DW} I_{ee}^{-1} D_{eo}^{DW}$$

- **Shamir:** 4-d & 5-d Even/Odd give ~ 2.7 speed up.
- **Borici/Moebius:** 4-d Even/Odd gives ~ 2.7 speed up
- **Solves explicitly one axis:** So probably 4-d is better than 5-d ?
- **Balint:** 3-d beats 4-d for asymmetric clover lattices

Pure Gauge: $16^3 \times 32$ @ $\beta = 6$ & $m_\pi = 0.44$



Optimal. $\alpha \simeq 1 + L_s/8$

Pure gauge $\beta = 6.0$ is not that smooth?

QuickTime[®] and a
decompressor
are needed to see this picture.

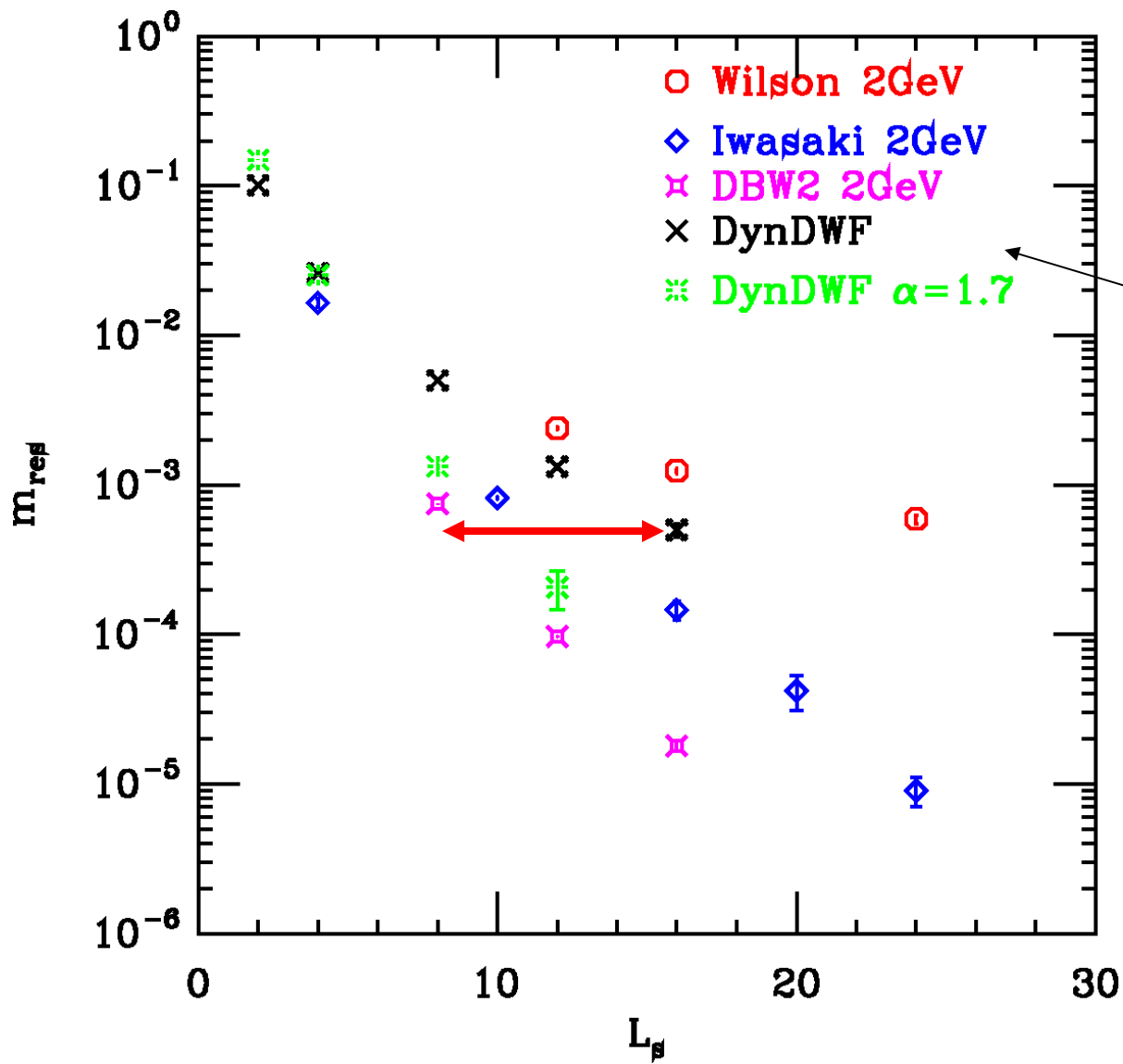
$-m_0$

$-M_5$

Vranas: for 20 config, 20 lowest e.v. of $\gamma_5 D_{Wil}(m_0)$

Early DW test for Moebius

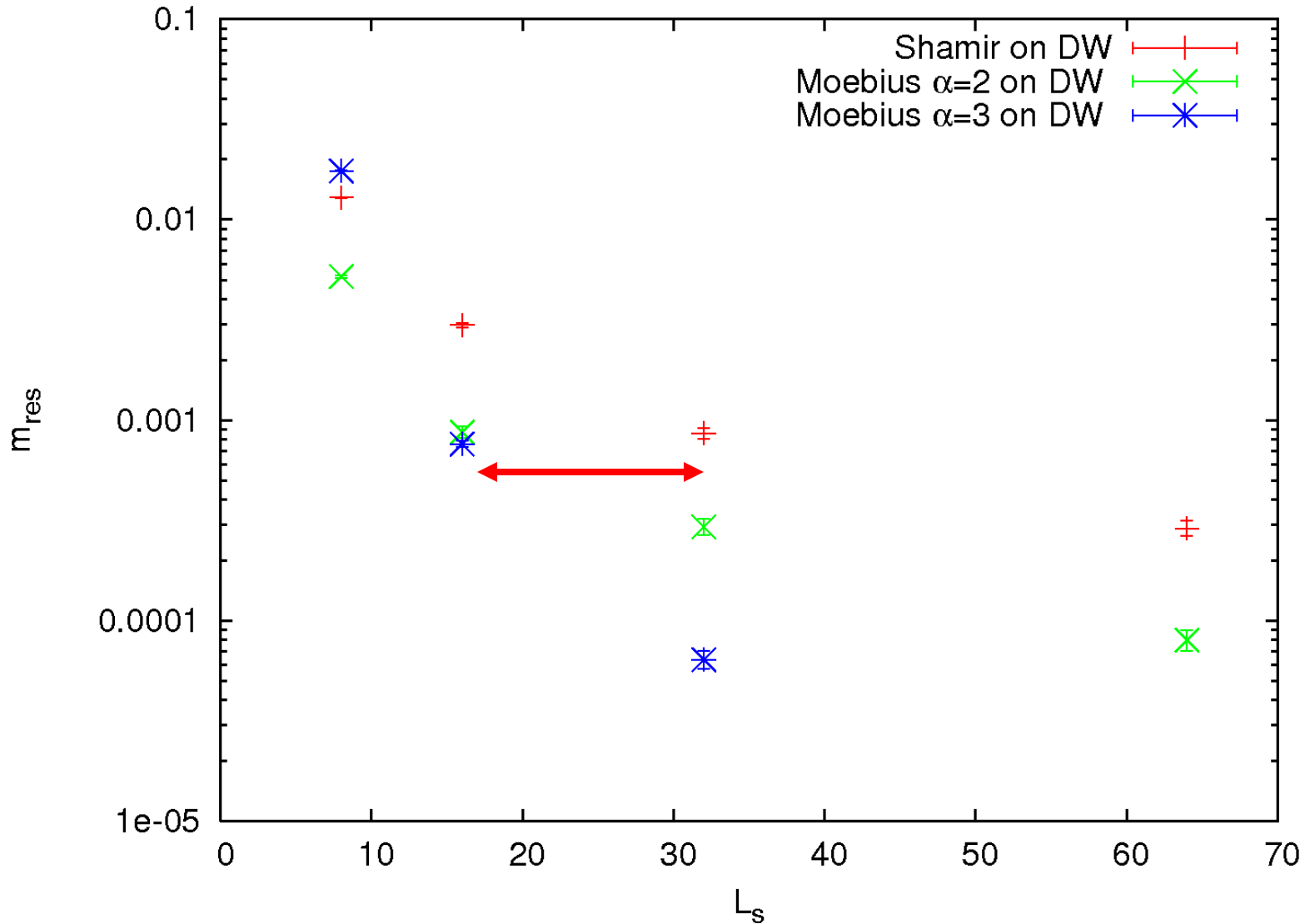
(Sept 23, ILFT04 Shuzenji, Japan)



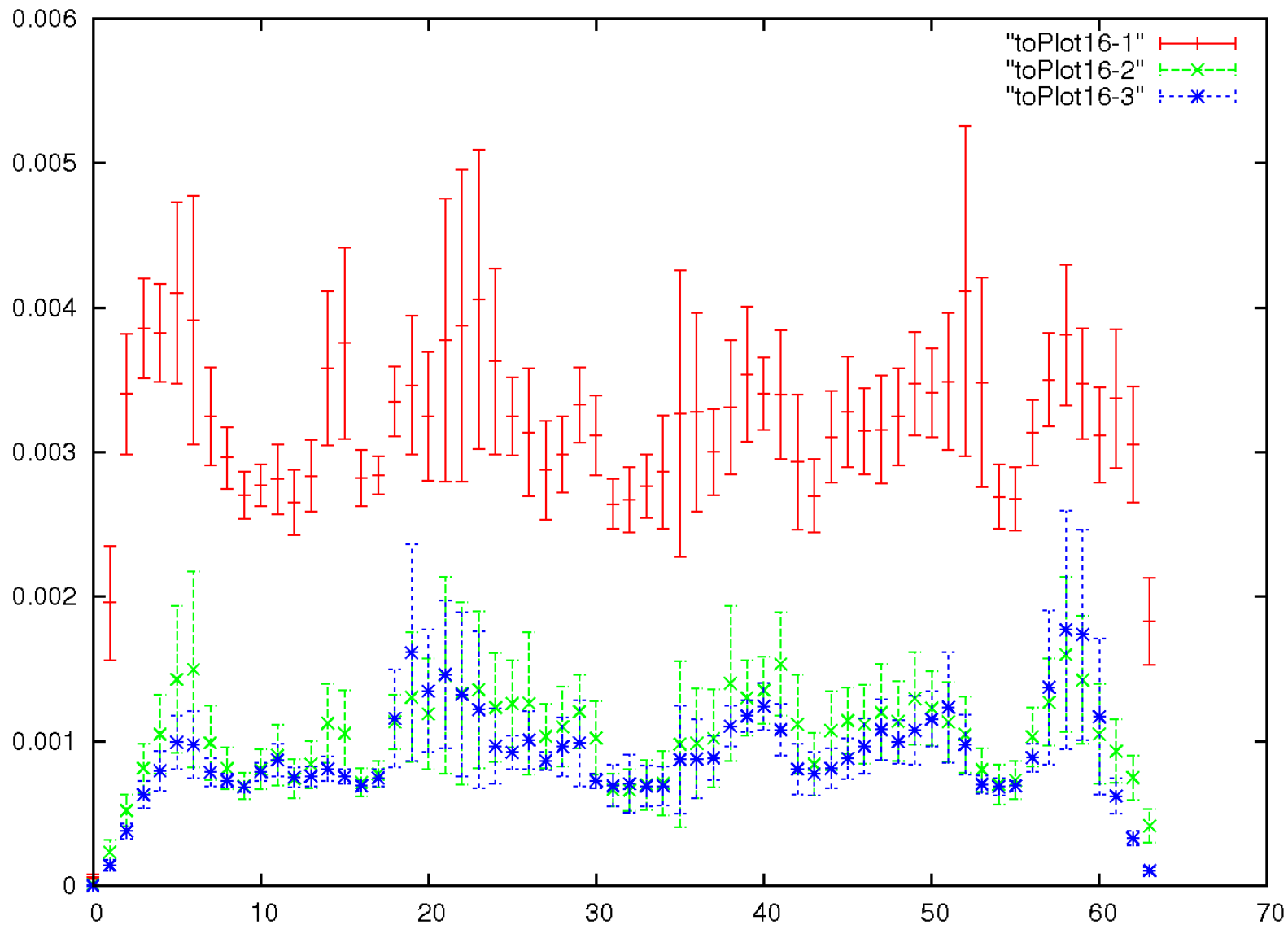
RBC DWF $16^3 \times 32$
 $m_f = 0.20$
 $a = 1.7\text{GeV}$
 $M_5 = 1.8$

Domain Wall Lattices (BNL archive)

($24^3 \times 64 L_s = 16$, 2+1 Iwaski $\beta = 2.13$ $m_s = 0.4$ $m_l/m_s = 1/4$)



$m_{\text{res}}(t)$ for $L_s = 16$ on one lattice



Gapped Fermions (Vranas :hep-lat/0606014v2)

$$\text{Det}[D^\dagger(M_5) D(M_5)]$$

factor in path integral opens a “gap” in $H_5 = \gamma_5 D_{wil}(M_5)$

QuickTime[®] and a
decompressor
are needed to see this picture.

M_5

M_5 $-m_0$

Ten smallest magnitude eigenvalues of $\gamma_5 D_W(m_0)$ vs. m_0 on 20 independent configurations . 0-flavor and 2-flavor β values correspond to the same lattice spacing $a^{-1} = 1.4$ Gev.

Gap gives exponentially **local** “effective” gauge action (just like overlap actions)

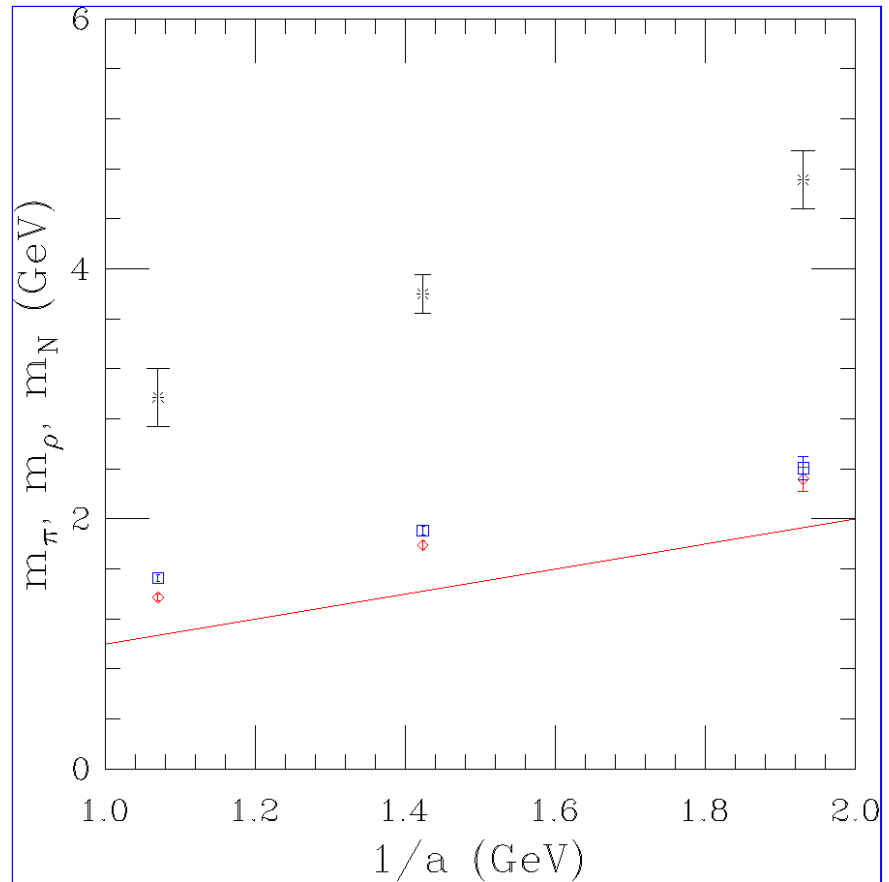
$$S_{gap}[U] = \log[\int \psi \bar{\psi} e^{-\bar{\psi} H_5(M_5) \psi}] = Tr[\log D^\dagger(M_5) D(M_5)]$$

Axial -Axial correlator measures range for gap action!

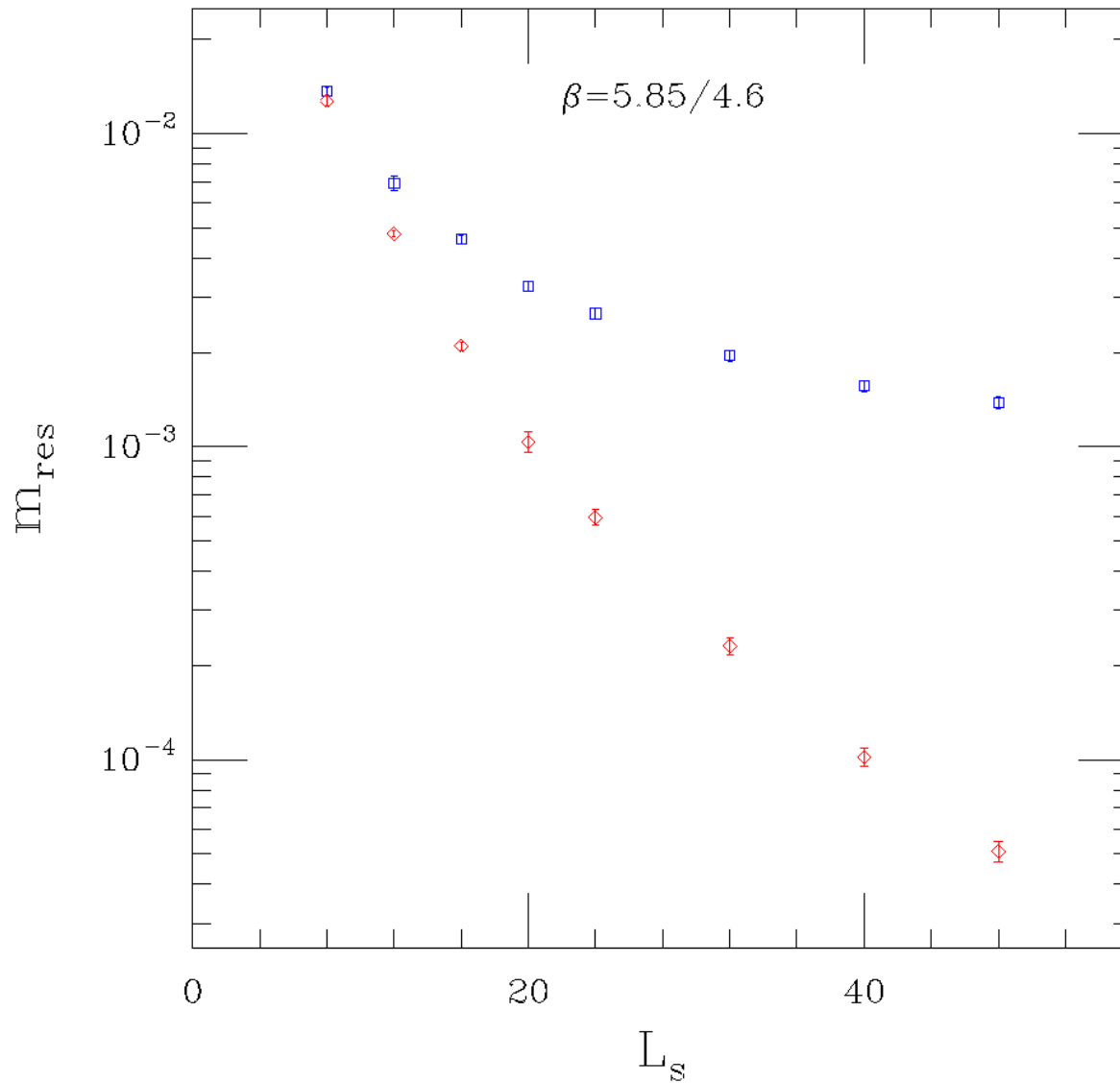
$$\delta_{A_\mu(x)} \delta_{A_\nu(0)} S_{gap}[U]$$

$$= \langle \bar{\psi}(x) \gamma_5 V_\mu(x) \psi(x) \bar{\psi}(0) \gamma_5 V_\nu(0) \psi \rangle$$

$$M_5 = -1.9$$

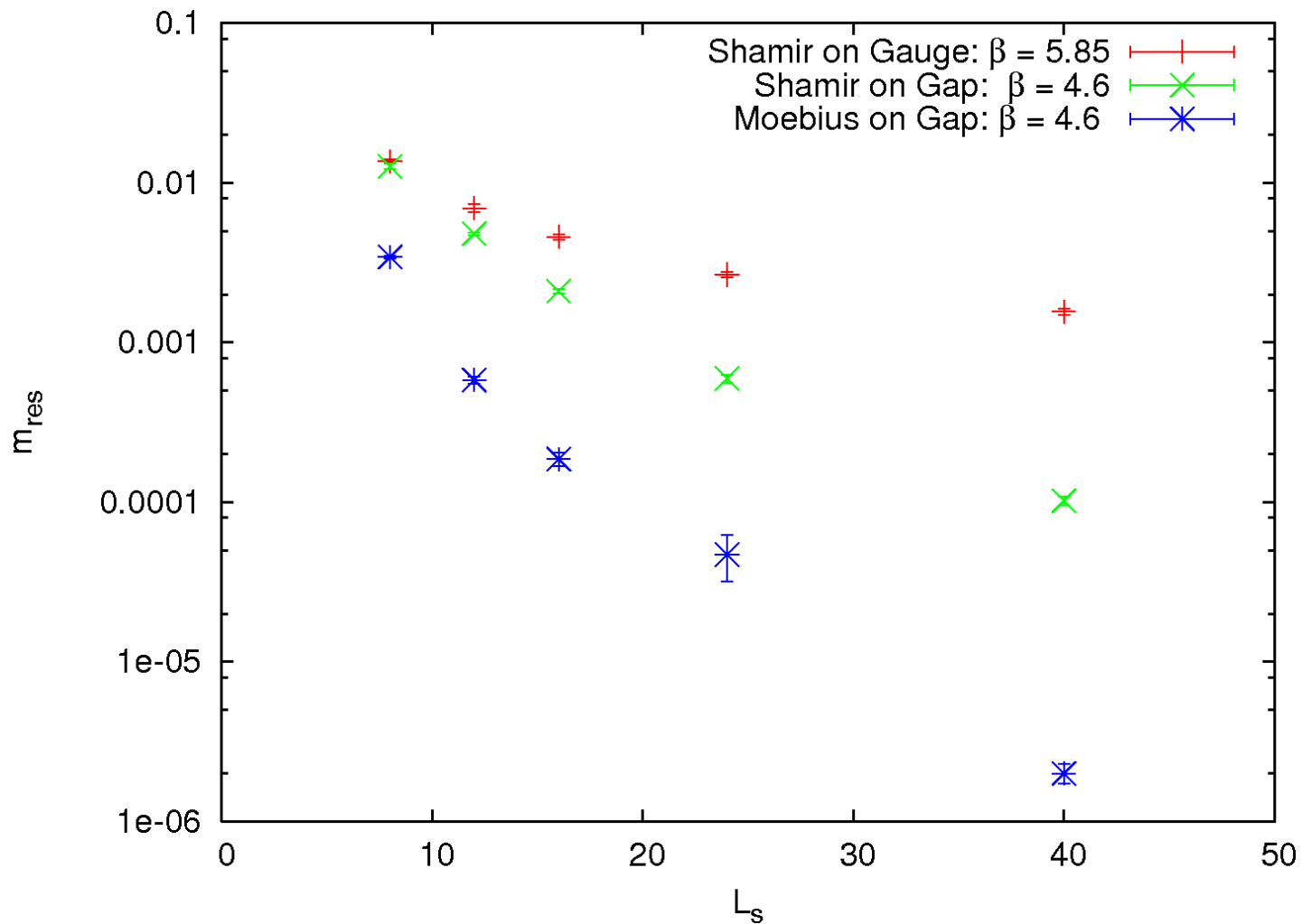


Shamir on Gauge ($\beta = 5.85$) vs Gapped ($\beta = 4.6$) Lattices

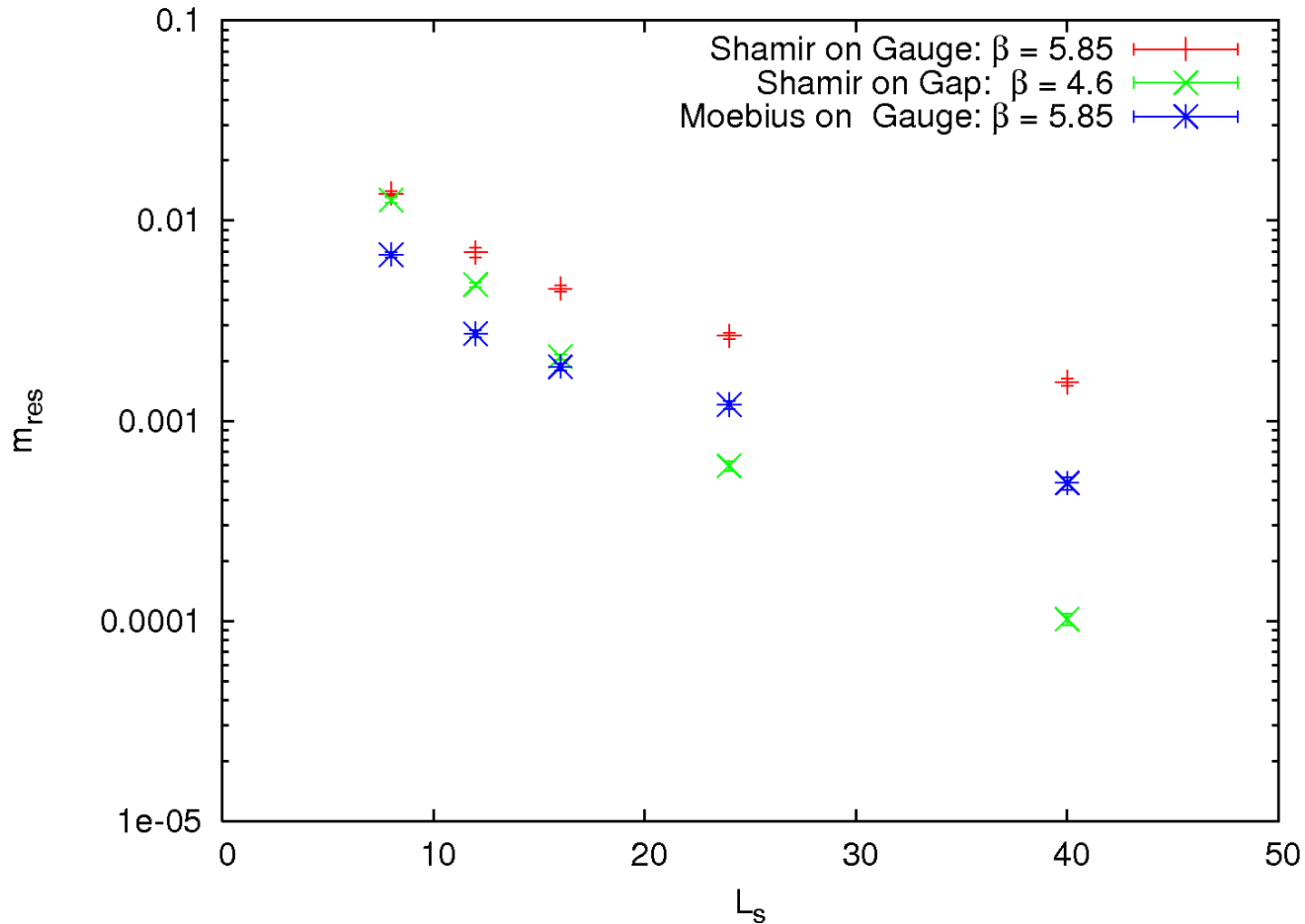


† Pavlos Vranas arXiv:hep-lat/0606014

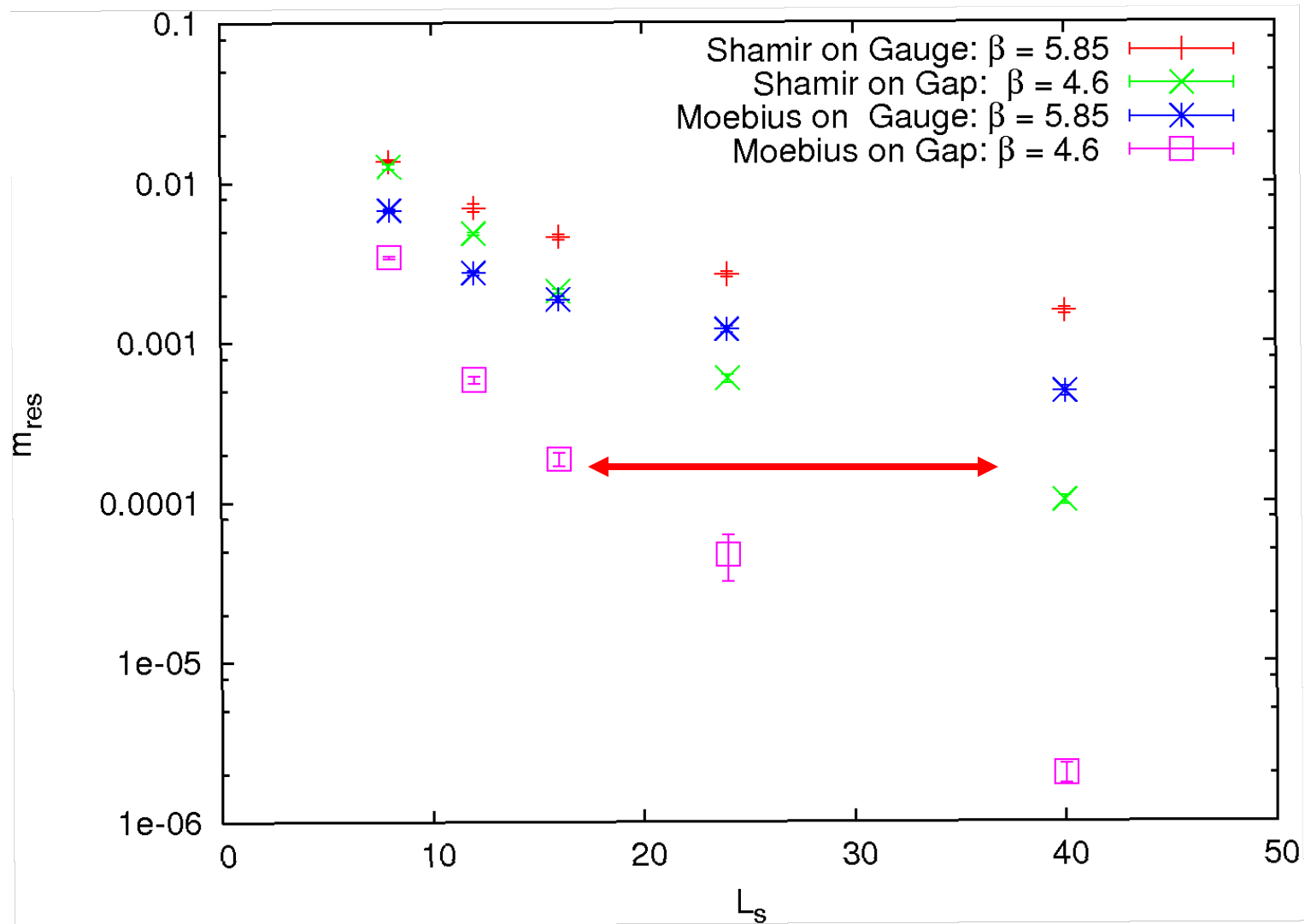
Gapped Shamir vs Moebius Lattices



Gauge Shamir vs Moebius Lattices



Combined Plot for Gapped Lattices



Generalized D_5 Hermiticity and All That

$$D_{DW} = D_- \times \widetilde{D}_{DW} \equiv$$

$$\begin{bmatrix} D_-^{(1)} & 0 & 0 & 0 \\ 0 & D_-^{(2)} & 0 & 0 \\ 0 & 0 & D_-^{(3)} & 0 \\ 0 & 0 & 0 & D_-^{(4)} \end{bmatrix} \times \begin{bmatrix} D_+^{(1)}/D_-^{(1)} & P_- & 0 & -mP_+ \\ P_+ & D_+^{(2)}/D_-^{(2)} & P_- & 0 \\ 0 & P_+ & D_+^{(3)}/D_-^{(3)} & P_- \\ -mP_- & 0 & P_+ & D_+^{(4)}/D_-^{(4)} \end{bmatrix}$$

Symmetry $D_-^{(s)} = D_-^{(L_s+1-s)}$, allows $\mathcal{R}(-D_-)\gamma_5$ to acts like "gamma 5"

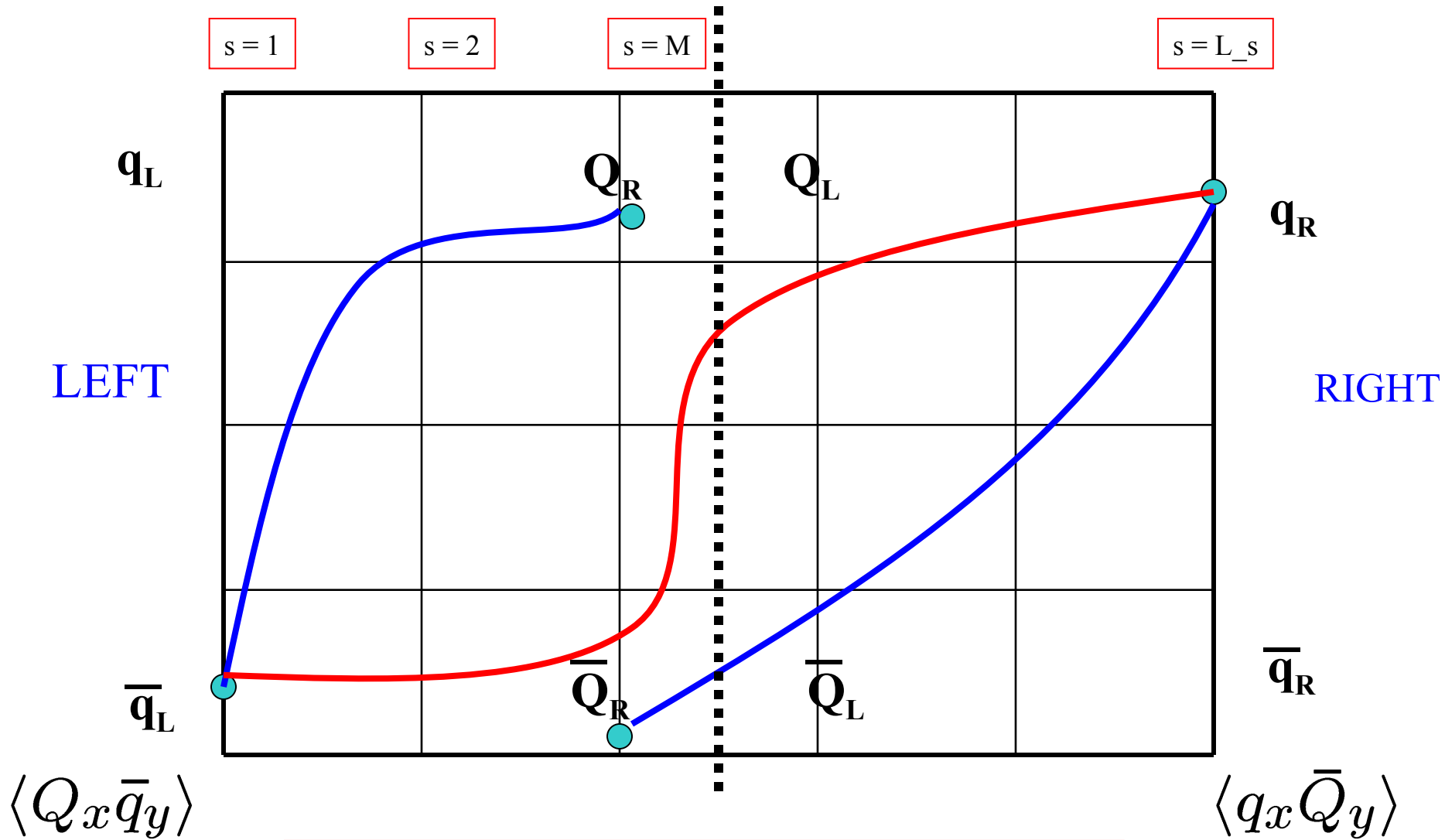
Chiral boundaries

$$q = P_- [\Psi]_1 + P_+ [\Psi]_{L_s}$$

re-defined

$$\bar{q} = [\bar{\Psi} D^{DW}(1)]_1 P_+ + [\bar{\Psi} D^{DW}(1)]_{L_s} P_-$$

Split Screen Correlators: 5-d Vector \Rightarrow 4-d Axial

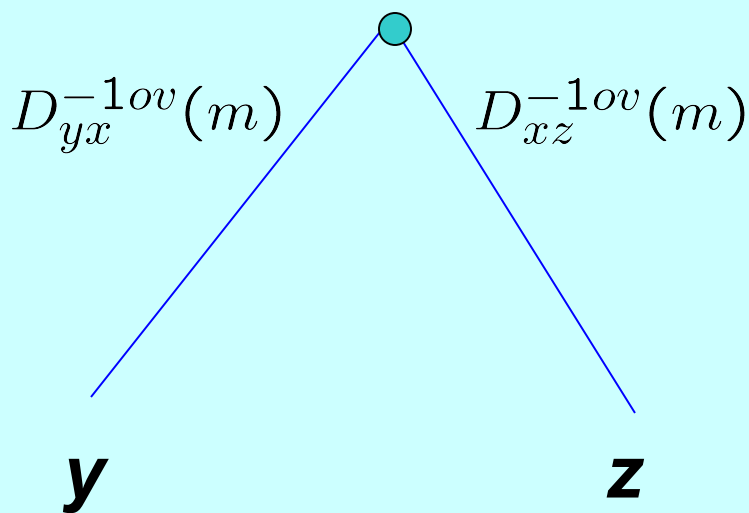


$$\langle q_x \bar{q}_y \rangle_{DW} = [D_{ov}^{-1}(m)]_{xy}$$

DW/overlap map for all correlators

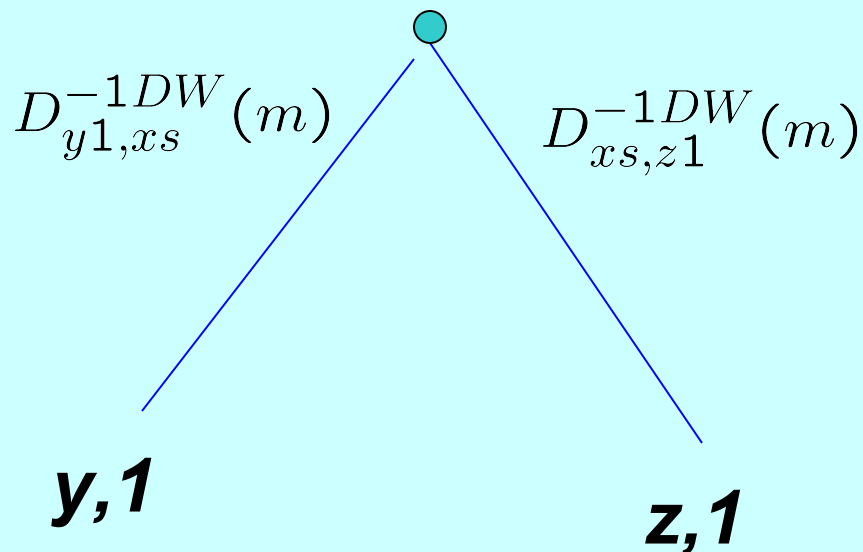
$$\langle J_\mu^{ov}(x) \psi_y \bar{\psi}_z \rangle_{ov} = \langle J_\mu^{DW}(x) q_y \bar{q}_z \rangle_{DW}$$

$$J_\mu^{(a)ov}(x)$$



\equiv

$$J_\mu^{(a)DW}(x, s)$$



Nice Definition of Overlap Axial Current:

$$\langle J_\mu^{ov}(x) \psi_y \bar{\psi}_z \rangle_{ov} \equiv \langle J_\mu^{DW}(x) q_y \bar{q}_z \rangle_{DW}$$

$$\Delta_{-\mu} J_\mu^{a,DW}(x) = 2m \tilde{q}_x \lambda^a \gamma_5 q_x + 2\tilde{Q}_x \gamma_5 \lambda^a Q_x + \text{P-V}$$

implies $\Delta_{-\mu} J_\mu^{a,ov}(x) = m_q \bar{\psi} \lambda^a (\gamma_5 + \hat{\gamma}_5) \psi + 2\bar{\psi} \gamma_5 \rho_L(x) \lambda^a \psi$

$$\rho_{L_s}^{zy}(x) = \Delta_{zx}^L \Delta_{xy}^R = \left[\frac{T_1^{-1} \dots T_{L_s/2}^{-1}}{1 + \mathbf{T}^{-L_s}} \right]_{zx} \left[\frac{T_{L_s/2+1}^{-1} \dots T_{L_s}^{-1}}{1 + \mathbf{T}^{-L_s}} \right]_{xy}$$

where

$$\mathbf{T}^{L_s} = T_{L_s} \cdots T_2 T_1 \quad , \quad \sum_x \rho_{L_s}(x) = \Delta_{L_s}$$

Note: Anomaly comes from P-V term at the boundary!

*CG Convergence: Effects of precondition & increasing α
at fixed m_{res} (need more study but preliminary)*

Dirac applications on DW $24^3 \times 64$ Lattices

L_s	Shamir, $\alpha = 1$	Moebius, $\alpha = 2$	Moebius, $\alpha = 3$
8	165 \pm 46	205 \pm 72	123 \pm 36
16	196 \pm 77	220 \pm 92	132 \pm 55
32	204 \pm 88	224 \pm 96	131 \pm 56
64	206 \pm 91	225 \pm 99	

*10-15% cost for Moebius at $\alpha=2$ (as seen on Gauge lattices)
but $\alpha=3$ is cheaper?*

*Real benchmarks by direct comparison of RHMC codes
will be performed on the BG/L in Fall 2008*

Conclusions

- " *Moebius algorithm: $L_s \Rightarrow L_s/2$ or more at fixed m_{res}*
- " *Gap + Moebius gives independent reduction in L_s*
- " *Promising to avoid (very) large L_s considered necessary for*
 - Technicolor
 - Finite T
 - SUSY, etc
- " *Code requires 4-d red/black precondition.*
 - Identical number of Dslash(M_5) operations per CG step
 - RHMC can use same building blocks CG and Force
- " *Stay tuned for more detailed benchmarks on*
 - Precondition comparisons
 - RHMC performance.

Phen. model of m_{res} dependence on e & L_s

$$m_{res} = \sum_{\lambda} w(\lambda) \Delta_L(\lambda) \quad w(\lambda) = \frac{\langle \lambda | G_{ov} G_{ov}^{\dagger} | \lambda \rangle}{\sum_{\lambda} \langle \lambda | G_{ov} G_{ov}^{\dagger} | \lambda \rangle} \geq 0$$

(D) has negligible dependence on D and L_s

$$m_{res} \simeq \int dn(\lambda) w(\lambda) \Delta_L(\alpha\lambda)$$

$$\Delta_L(\lambda) = \langle \lambda | \Delta_L(H) | \lambda \rangle = \frac{4}{2 + [\frac{1+\lambda}{1-\lambda}]^{-L} + [\frac{1+\lambda}{1-\lambda}]^L} \geq 0$$

$\rightarrow e^{-L \log[(1+\lambda)/(1-\lambda)]}$ for $O(L^{-1}) < \lambda < O(L)$

(Parameterize and fit m_{res} data)

Measuring the Operator D_{Ls}

(use Plateau region away from sources)

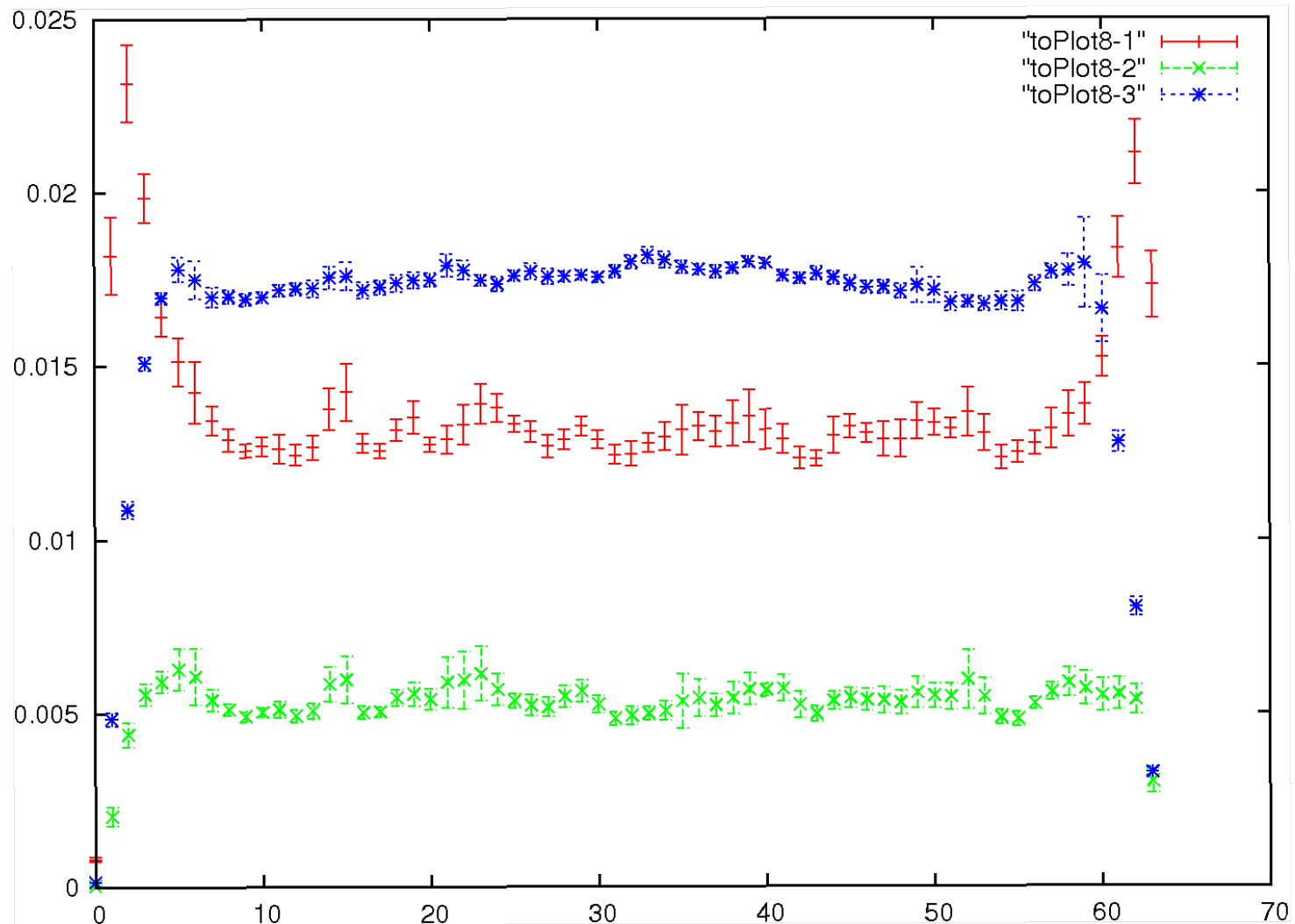
$$m_{res}(t) \equiv \frac{\sum_x \langle \tilde{Q}_{t,x} \gamma_5 Q_{t,x} \tilde{q}_0 \gamma_5 q_0 \rangle_c}{\sum_x \langle \bar{q}_{t,x} \gamma_5 q_{t,x} \bar{q}_0 \gamma_5 q_0 \rangle_c}$$

Theoretical m_{res} : Sum over $t \rightarrow$ Measure Matrix element of e_L operator

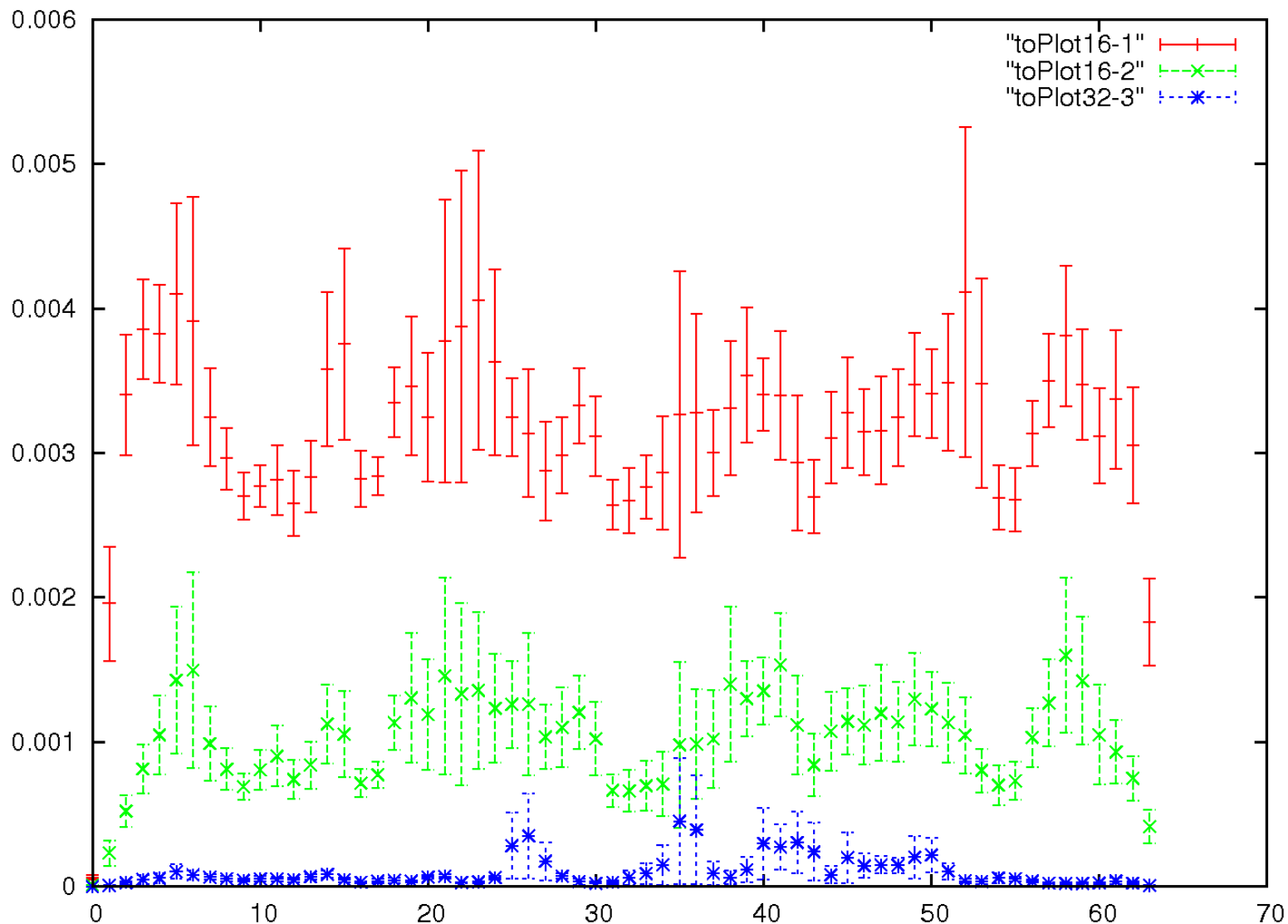
$$m_{res} \equiv \frac{Tr[\Delta_L(H) D_{ov}^{-1} D_{ov}^{\dagger-1}]}{Tr[D_{ov}^{-1} D_{ov}^{\dagger-1}]} = \sum_{\lambda} w(\lambda) \Delta_L(\lambda)$$

$|D\rangle$ in the Eigen basis of $H = D_5 D(-M)$

$m_{\text{res}}(t)$ for $L_s = 8$ on one lattice



$m_{\text{res}}(t)$ for $L_s = 32$ on one lattice



Edwards & Heller use “Standard” UDL decomposition

$$' D^{DW} \mathcal{P} = U D L(m) '$$

$$D_{DW}(m) \mathcal{P} = [\gamma_5 Q_- U Q_-^{-1}] \quad Q_- \begin{bmatrix} D_4(m) & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad [L(m)]$$

Step #1: Prepare the Pivots by Permute Columns

$$\mathcal{P} = \begin{bmatrix} P_- & P_+ & 0 & 0 \\ 0 & P_- & P_+ & 0 \\ 0 & 0 & P_- & P_+ \\ P_+ & 0 & 0 & P_- \end{bmatrix}$$

Step #2: Do Gaussian Elimination to get U matrix

$$U = \begin{bmatrix} 1 & -T_1^{-1} & -T_1^{-1}T_2^{-1} & -T_1^{-1}T_2^{-1}T_3^{-1} \\ 0 & 1 & -T_2^{-1} & -T_2^{-1}T_3^{-1} \\ 0 & 0 & 1 & -T_3^{-1} \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

Step #3 Back substitution to get L matrix

$$L(m) = \begin{bmatrix} -1 & 0 & 0 & 0 \\ -T_2^{-1}T_3^{-1}T_4^{-1}c_+ & 1 & 0 & 0 \\ -T_3^{-1}T_4^{-1}c_+ & 0 & 1 & 0 \\ -T_4^{-1}c_+ & 0 & 0 & 1 \end{bmatrix}$$

where

$$Q_-^{-1} = \text{Diag}[(Q_-^{(1)})^{-1}(Q_-^{(2)})^{-1}(Q_-^{(3)})^{-1}(Q_-^{(4)})^{-1}]$$

$$Q_-^{(s)} = \gamma_5[D_-^{(s)}P_+ + D_+^{(s)}P_-] \quad c_- = P_- - mP_+$$

$$Q_+^{(s)} = \gamma_5[D_+^{(s)}P_+ + D_-^{(s)}P_-] \quad c_+ = P_+ - mP_-$$

LUD \Rightarrow

$$\mathcal{P}^\dagger \frac{1}{D_{DW}(1)} D_{DW}(m) \mathcal{P} = \begin{bmatrix} D_{ov}(m) & 0 & 0 & \cdots & 0 \\ -(1-m)\Delta_2^R & 1 & 0 & \cdots & 0 \\ -(1-m)\Delta_3^R & 0 & 1 & \cdots & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ -(1-m)\Delta_{L_s}^R & 0 & 0 & \cdots & 1 \end{bmatrix}$$

$$\Delta_{s+1}^R = \frac{T_{s+1}^{-1} T_{s+2}^{-1} \cdots T_{L_s}^{-1}}{1 + \mathbf{T}^{-L_s}}, \quad \mathbf{T}^{L_s} = T_{L_s} \cdots T_2 T_1$$

$$T_s = \frac{1 - H_s}{1 + H_s}, \quad H_s = \gamma_5 D_{Moebius}^{(s)}(M_5)$$

$$D_{ov}(m) = \frac{1+m}{2} + \frac{1-m}{2} \gamma_5 \frac{T^{-L} - 1}{T^{-L} + 1}$$

DW/Overlap Equivalence:

$$\langle \mathcal{O}(q, \bar{q}) \rangle_{DW} = \langle \mathcal{O}(\psi, \bar{\psi}) \rangle_{ov}$$

where $q = [\mathcal{P}^\dagger \Psi]_1$, $\bar{q} = [\bar{\Psi} D_{DW}(1) \mathcal{P}]_1$

$$\Rightarrow \langle q_y \bar{q}_x \rangle \equiv [\dots]_{x1,y1} = D_{ov}^{-1}(m)_{xy} \equiv \langle \psi_y \bar{\psi}_x \rangle$$

$$\mathcal{P}^\dagger \frac{1}{D_{DW}(m)} D_{DW}(1) \mathcal{P} = \begin{bmatrix} D_{ov}^{-1}(m) & 0 & 0 & \dots & 0 \\ (1-m)\Delta_2^R D_{ov}^{-1}(m) & 1 & 0 & \dots & 0 \\ (1-m)\Delta_3^R X_3 D_{ov}^{-1}(m) & 0 & 1 & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots \\ (1-m)\Delta_{L_s}^R D_{ov}^{-1}(m) & 0 & 0 & \dots & 1 \end{bmatrix}$$

note: Standard approach

$$\bar{q} = [\bar{\Psi} (-D_-^{(1)}) \mathcal{R} \mathcal{P}]_1 \Rightarrow \langle q \bar{q} \rangle = \frac{1}{1-m} [D_{ov}^{-1}(m) - 1]$$

Bulk to Boundary Propagators

$$\langle Q_s \tilde{q} \rangle = \frac{T_{s+1}^{-1} \cdots T_{L_s}^{-1}}{1 + \mathbf{T}^{-L_s}} D_{ov}^{-1}(m)$$

$$\langle q \tilde{Q}_s \rangle = D_{ov}^{-1}(m) \gamma_5 \frac{1}{1 + \mathbf{T}^{-L_s}} [T_1^{-1} \cdots T_s^{-1}] \gamma_5$$

where $s = L_s/2$ plane for m calculation

$$Q_s = P_- \Psi_{s+1} + P_+ \Psi_s$$

$$\tilde{Q}_s = -\bar{\Psi}_{s+1} D_-^{(s+1)} P_+ - \bar{\Psi}_s D_-^{(s)} P_-$$

†

(See Kikukawa and Noguchi, hep-lat/99902022)

Derivation:

$$m_{res} \equiv \frac{\sum_x \langle \tilde{Q}_x \gamma_5 Q_x \tilde{q}_0 \gamma_5 q_0 \rangle_c}{\sum_x \langle \bar{q}_x \gamma_5 q_x \bar{q}_0 \gamma_5 q_0 \rangle_c} \Rightarrow$$

$$\frac{\sum_x Tr[\langle Q_x \tilde{q}_0 \rangle \gamma_5 \langle q_0 \tilde{Q}_x \rangle \gamma_5]}{\sum_x Tr[\langle q_x \bar{q}_0 \rangle \gamma_5 \langle q_0 \bar{q}_x \rangle \gamma_5]} = \frac{\sum_x Tr[\Delta_{xy}^R \langle q_y \bar{q}_0 \rangle \langle q_z \bar{q}_0 \rangle^\dagger \Delta_{zx}^L]}{\sum_x Tr[\langle q_x \bar{q}_0 \rangle \langle q_x \bar{q}_0 \rangle^\dagger]}$$

$$= \frac{\sum_{zy} \langle \bar{q}_z \gamma_5 \Delta_{zy} q_y \bar{q}_0 \gamma_5 q_0 \rangle}{\sum_x \langle \bar{q}_x \gamma_5 q_x \bar{q}_0 \gamma_5 q_0 \rangle} = \frac{Tr[\Delta_L D_{ov}^{-1} D_{ov}^{\dagger-1}]}{Tr[D_{ov}^{-1} D_{ov}^{\dagger-1}]} \simeq \frac{\langle 0 | \bar{q} \gamma_5 \Delta q | \pi \rangle}{\langle 0 | \bar{q} \gamma_5 q | \pi \rangle}$$

