

# Nucleon Wave Function from Lattice QCD

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July 2008





# ■ NUCLEON PUZZLE

What we really want:

Nucleon wave function

What we really get:

Lowest moments of the Nucleon  
Distribution Amplitude (NDA)

- ★ What are NDA's
- ★ How to calculate their moments on the lattice
- ★ Possible forms of the NDA's
- ★ From NDA to experiment



# ■ SIMPLIFYING THE PROBLEM

## “What is the Nucleon Wave Function”

$$\Psi_{\text{BS}}(x_i) = \langle 0 | T [q(x_1, k_{1,\perp}) q(x_2, k_{2,\perp}) q(x_3, k_{3,\perp})] | p \rangle$$

- $x_i$  – longitudinal momentum fractions of partons
- $k_{i,\perp}$  – transverse momenta of quarks
- $|p\rangle$  – baryon state with momentum  $p$  (e.g., Proton).

## “Distribution Amplitude(s)”

$$\Phi(x_i, \mu) = Z(\mu) \int^{|\mathbf{k}_\perp| \leq \mu} [d\mathbf{k}_{i,\perp}] \Psi_{\text{BS}}(x_i, \mathbf{k}_\perp)$$

Further simplifications:

- ↔ Twist expansion
- ↔ Expansion in local operators (moments of DA's)



# MOMENTS

## “Expansion in **local** operators”

$$\begin{aligned} \langle 0 | \epsilon^{abc} [i^l D^{\lambda_1} \dots D^{\lambda_l} q_\alpha^a(0)] (C\gamma^\rho)_{\alpha\beta} [i^m D^{\mu_1} \dots D^{\mu_m} q_\beta^b(0)] [i^n D^{\nu_1} \dots D^{\nu_n} (\gamma_5 q^c(0))_\gamma] | p \rangle \\ = - f_V V^{lmn} p^\rho p^{\lambda_1} \dots p^{\lambda_l} p^{\mu_1} \dots p^{\mu_m} p^{\nu_1} \dots p^{\nu_n} N_\gamma(p) \end{aligned}$$

analogous for all Distribution Amplitudes

Moments of DA's:

$$V^{lmn} = \frac{1}{f_V} \int_0^1 [dx] x_1^l x_2^m x_3^n V(x_1, x_2, x_3), \quad [dx] = dx_1 dx_2 dx_3 \delta(1 - \sum_{i=1}^3 x_i)$$

$\delta(1 - \sum_{i=1}^3 x_i)$  - Momentum conservation

$f_V$  - DA normalization constant



# ■ "PROBLEM REDUCTION FOR NUCLEON"

Two *up*-quarks, Isospin symmetry in the proton

↪ Only **one** Distribution amplitude:

$$\varphi^{lmn} = V^{lmn} - A^{lmn} \quad \text{or} \quad \phi^{lmn} = \frac{1}{3}(V^{lmn} - A^{lmn} + 2T^{lnm})$$

$\phi^{lmn}$  is preferable on the lattice. On the other hand

$$|p, \uparrow\rangle = \int_0^1 [dx] \frac{\varphi(x_i)}{\sqrt{96x_1x_2x_3}} |u^\uparrow(x_1) [u^\downarrow(x_2)d^\uparrow(x_3) - d^\downarrow(x_2)u^\uparrow(x_3)]\rangle$$

↪  $\varphi^{100}$ ,  $\varphi^{010}$  und  $\varphi^{001}$  - momentum fractions of quarks in nucleon.

## Momentum conservation

$$\phi^{lmn} = \phi^{(l+1)mn} + \phi^{l(m+1)n} + \phi^{lm(n+1)} \quad \phi^{000} \equiv 1$$



# FORM FACTORS AND DISTRIBUTION AMPLITUDES

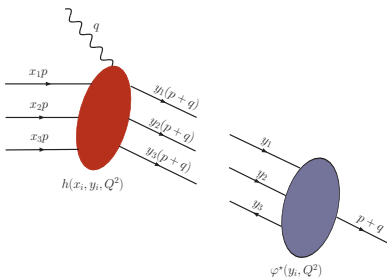
Factorization for (very) large  $Q^2$

exclusive processes, hard scattering

Generalization of Form Factors.

Different form factors are related to each other through (one) distribution amplitude.

$$\varphi(x_i, Q^2 \rightarrow \infty) = 120 x_1 x_2 x_3 f_N$$



$$G_M(Q^2) = \int_0^1 [dx] \int_0^1 [dy] \varphi^*(y_i, Q^2) h(x_i, y_i, Q^2) \varphi(x_i, Q^2)$$



# ■ OPERATOR CHOICE

Reduced symmetry due to discretization  $\leftrightarrow$   
 ☹ additional operator mixing

## Construction of proper operators required

- $\leftrightarrow$  Irreducibly transforming operators have “minimal” mixing
- $\leftrightarrow$  Construct such operators from DA's operators

$$\text{Isospin } 1/2: \mathcal{O}_{1,I=1/2} = \mathcal{O}_1(u_\alpha u_\beta d_\gamma) + i\gamma_4 \mathcal{O}_1(u_\alpha d_\beta u_\gamma)$$

## Correlator ( $t \rightarrow \infty$ )

$$\langle \mathcal{O}_{1,I=1/2}(t) \bar{N}(0) \rangle = f_N \phi^{lmn} (ip_1 \gamma_1 - ip_2 \gamma_2) (ip_3 \gamma_3 - \gamma_4 E(\vec{p})) N(\vec{p}) e^{-E(\vec{p})t}$$

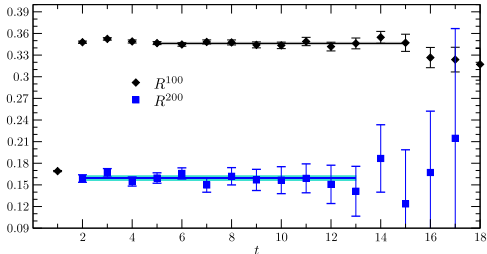
$\leftrightarrow \vec{p} \neq 0 \leftrightarrow$  “good” Mixing



# ■ CONSTRAINED ANALYSIS

## DA correlator ratios

$$R^{lmn} = \frac{\phi^{lmn}}{S_{(l+m+n)}}$$



$$S_1 = \phi^{100} + \phi^{010} + \phi^{001}$$

$$S_2 = \phi^{200} + \phi^{020} + \phi^{002}$$

$$+ 2(\phi^{011} + \phi^{101} + \phi^{110})$$

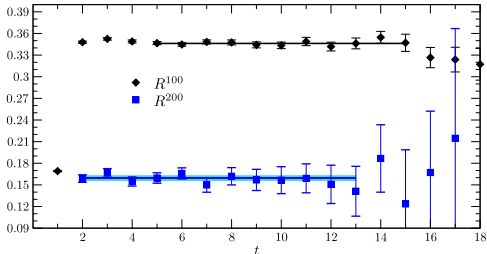




# ■ CONSTRAINED ANALYSIS

## DA correlator ratios

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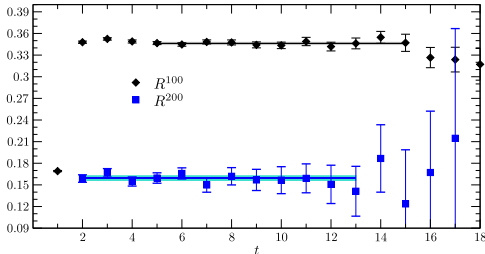
Problem: normalization  
of  $\phi^{lmn}$  lost



# ■ CONSTRAINED ANALYSIS

## DA correlator ratios

$$R^{lmn} = \frac{\phi^{lmn}}{S_{(l+m+n)}}$$



Solution:

$$S_1 = \phi^{100} + \phi^{010} + \phi^{001} \equiv 1$$

Restore the normalization from the constraint

$$\sum_{l+m+n=1} \phi^{lmn} = \frac{1}{\sum_{l+m+n=1} Z_{lmn} R^{lmn}}$$

$Z^{lmn}$ : Renormalization constants



# ■ JUST NUMBERS

Results from QCDSF/DIK configurations at  $\beta = 5.40$

Comparison with estimates at 1 GeV

	Asympt.	QCD-SR	BK	BLW	Latt. $N$	Latt. $N^*(1535)$
$\varphi^{100}$	$\frac{1}{3} \approx 0.333$	0.560(60)	0.38	0.415	0.3999(13)(126)	0.4765(18)(92)
$\varphi^{010}$	$\frac{1}{3} \approx 0.333$	0.192(12)	0.31	0.285	0.2986(22)(111)	0.2523(28)(111)
$\varphi^{001}$	$\frac{1}{3} \approx 0.333$	0.229(29)	0.31	0.300	0.3015(9)(18)	0.2712(10)(29)
$\varphi^{200}$	$\frac{1}{7} \approx 0.143$	0.350(70)	0.18*	0.212	0.1792(26)(157)	0.2289(40)(71)
$\varphi^{020}$	$\frac{1}{7} \approx 0.143$	0.084(19)	0.13*	0.123	0.1459(66)(63)	0.0839(80)(264)
$\varphi^{002}$	$\frac{1}{7} \approx 0.143$	0.109(19)	0.13*	0.132	0.1354(42)(270)	0.1005(38)(176)
$\varphi^{011}$	$\frac{2}{21} \approx 0.095$	-0.030(30)	0.08*	0.053	0.0491(54)(351)	0.0457(58)(200)
$\varphi^{101}$	$\frac{2}{21} \approx 0.095$	0.102(12)	0.10*	0.097	0.1171(21)(66)	0.1250(29)(44)
$\varphi^{110}$	$\frac{2}{21} \approx 0.095$	0.090(10)	0.10*	0.093	0.1037(34)(266)	0.1227(50)(116)



# ■ FROM THE MOMENTS TO THE DA

NDA is not completely determined by the (lowest) moments.

↪ Additional input required!



# FROM THE MOMENTS TO THE DA

NDA is not completely determined by the (lowest) moments.

↪ Additional input required!

## From renormalization group

$$\begin{aligned} \varphi(x_1, x_2, x_3, \mu) = & 120x_1x_2x_3f_N(\mu_0)L^{\frac{2}{3\beta_0}} \left\{ 1 + h_{10}(\mu_0)(x_1 - 2x_2 + x_3)L^{\frac{8}{3\beta_0}} \right. \\ & + h_{11}(\mu_0)(x_1 - x_3)L^{\frac{20}{9\beta_0}} \\ & + h_{20}(\mu_0) \left[ 1 + 7(x_2 - 2x_1x_3 - 2x_2^2) \right] L^{\frac{14}{3\beta_0}} \\ & + h_{21}(\mu_0) (1 - 4x_2)(x_1 - x_3) L^{\frac{40}{9\beta_0}} \\ & \left. + h_{22}(\mu_0) \left[ 3 - 9x_2 + 8x_2^2 - 12x_1x_3 \right] L^{\frac{32}{9\beta_0}} \right\} \end{aligned}$$

$h_{ij}$  determined by  $\phi^{lmn}$

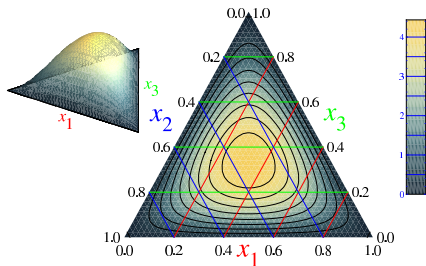
$$L = \frac{\alpha_s(\mu)}{\alpha_s(\mu_0)}$$

$$\beta_0 = 11 - \frac{2}{3}n_f$$

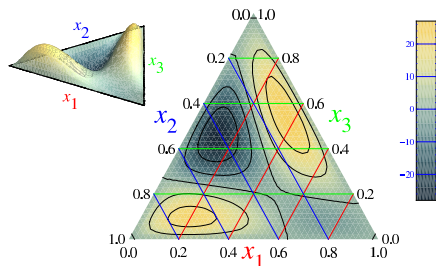


# ■ NEXT<sup>2</sup>-TO-LEADING CONFORMAL SPIN

## Asymptotic form

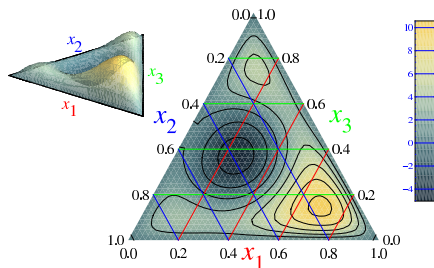


## Nucleon model ( $\phi^{101}$ , $\phi^{002}$ , $\phi^{200}$ )

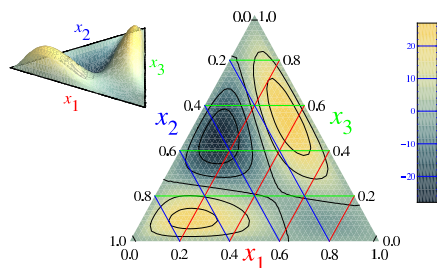


# ■ NEXT<sup>2</sup>-TO-LEADING CONFORMAL SPIN

## QCD-SR inspired form

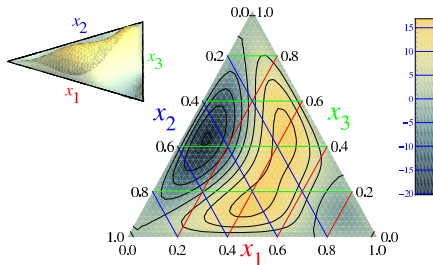


## Nucleon model ( $\phi^{101}$ , $\phi^{002}$ , $\phi^{200}$ )

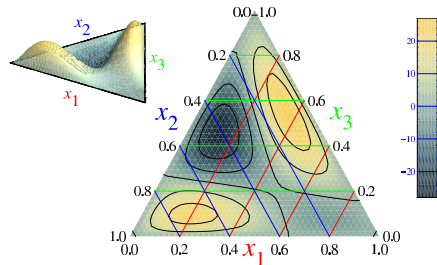


# ■ NEXT<sup>2</sup>-TO-LEADING CONFORMAL SPIN

$N^*(1535)$  model ( $\phi^{101}$ ,  $\phi^{002}$ ,  $\phi^{200}$ )



Nucleon model ( $\phi^{101}$ ,  $\phi^{002}$ ,  $\phi^{200}$ )





# ■ CONCLUSIONS

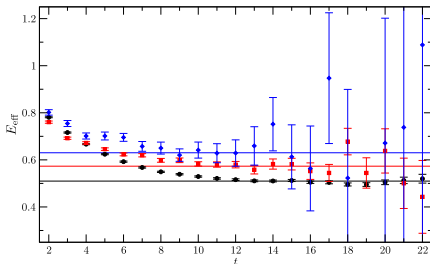
- ▶ Lattice QCD at the moment the method of choice for DA
- ▶ Accuracy improved by explicitly exploiting momentum conservation.
- ▶ Asymmetries are less pronounced compared to QCD-SR
- ▶ Symmetry in  $u^\downarrow$  and  $d^\uparrow$  (diquark)
- ▶ In good agreement with phenomenological values



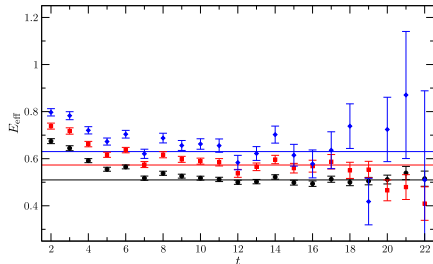
# THE USUAL APPROACH

## Effective energies for required correlators

### DA Correlators



### Nucleon correlators



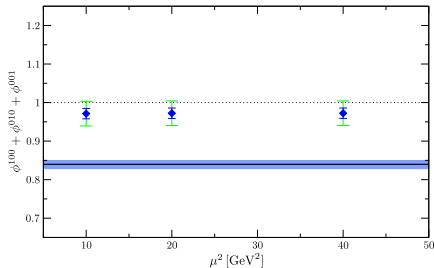
Nucleon correlator adds noise but does not really contain any additional information



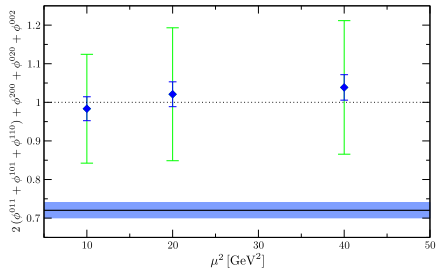
# QUALITY OF THE RESULTS

## Sum of the moments

$$\phi^{100} + \phi^{010} + \phi^{001}$$



$$\begin{aligned} &\phi^{200} + \phi^{020} + \phi^{002} \\ &+ 2(\phi^{011} + \phi^{101} + \phi^{110}) \end{aligned}$$



Errors are too large for the interesting quantities (asymmetries).

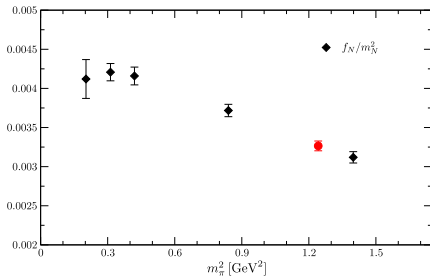


# CHIRAL EXTRAPOLATION I

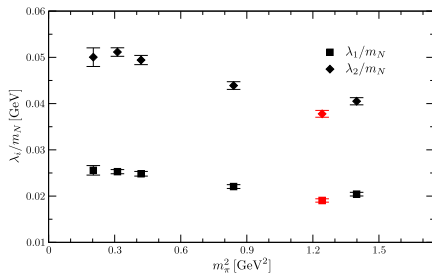
No results from  $\chi$ PT

## Normalization constants

### Leading twist



### Next-to-leading twist

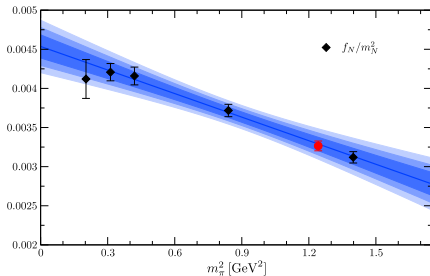


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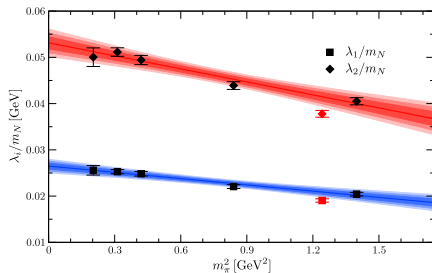
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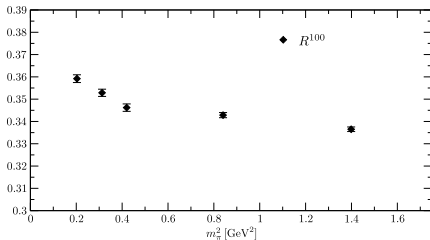


# CHIRAL EXTRAPOLATION II

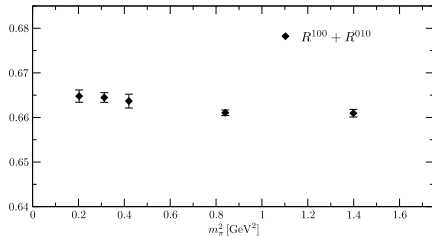
$R^{100}$  and  $R^{010}$  not really linear  
No strong deviation for other moments

## First moments

$R^{100}$



$R^{100} + R^{010}$

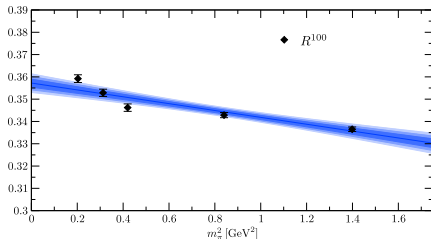


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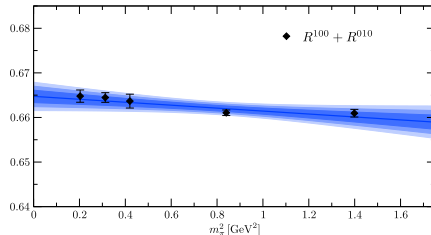
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## First moments

$R^{100}$  (linear)



$R^{100} + R^{010}$  (linear)

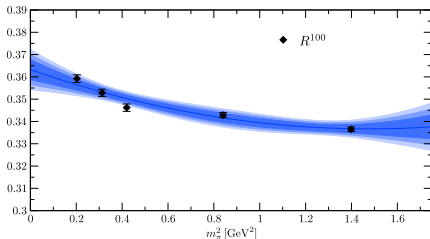


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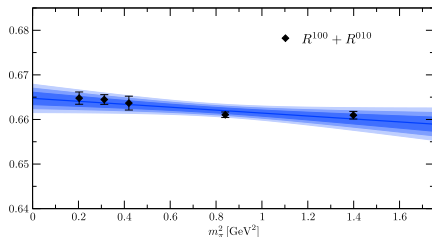
$R^{100}$  and  $R^{010}$  not really linear  
No strong deviation for other moments

## First moments

$R^{100}$  (quadratic)



$R^{100} + R^{010}$  (linear)



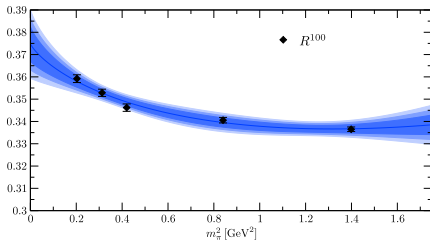


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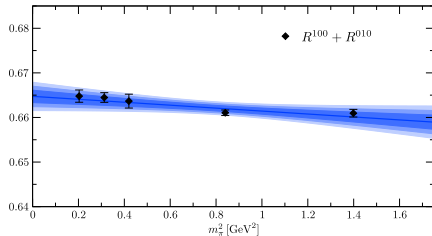
$R^{100}$  and  $R^{010}$  not really linear  
No strong deviation for other moments

## First moments

$R^{100}$  (“chiral log”)



$R^{100} + R^{010}$  (linear)

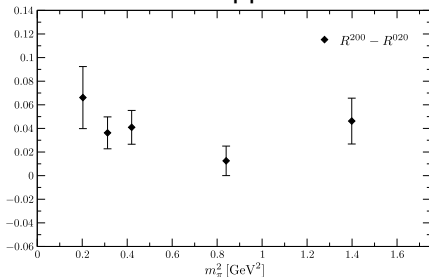


# CHIRAL EXTRAPOLATION III

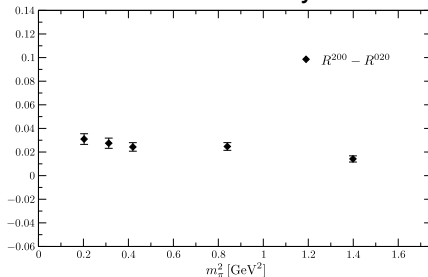
Asymmetries are the important quantities

For the second moments

Standard approach



Constrained analysis

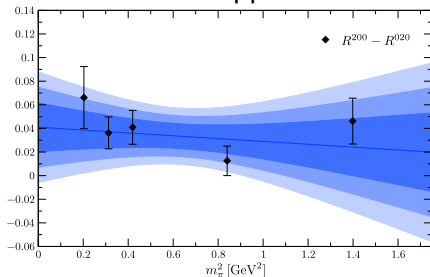


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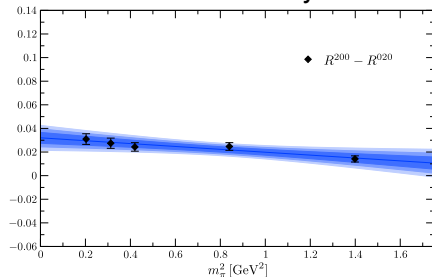
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# THANK YOU

