

# Topological susceptibility in the SU(3) random vortex world-surface model

M. Engelhardt

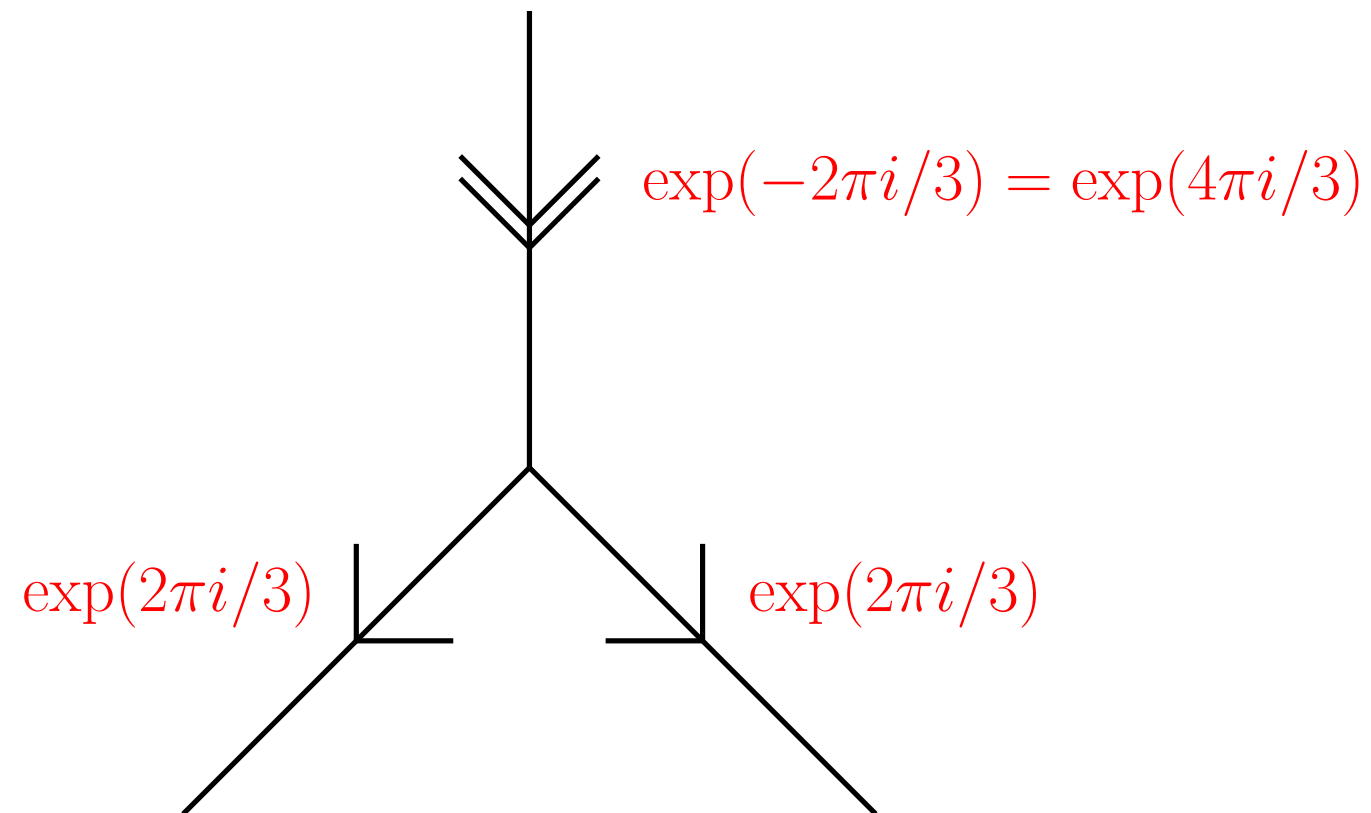
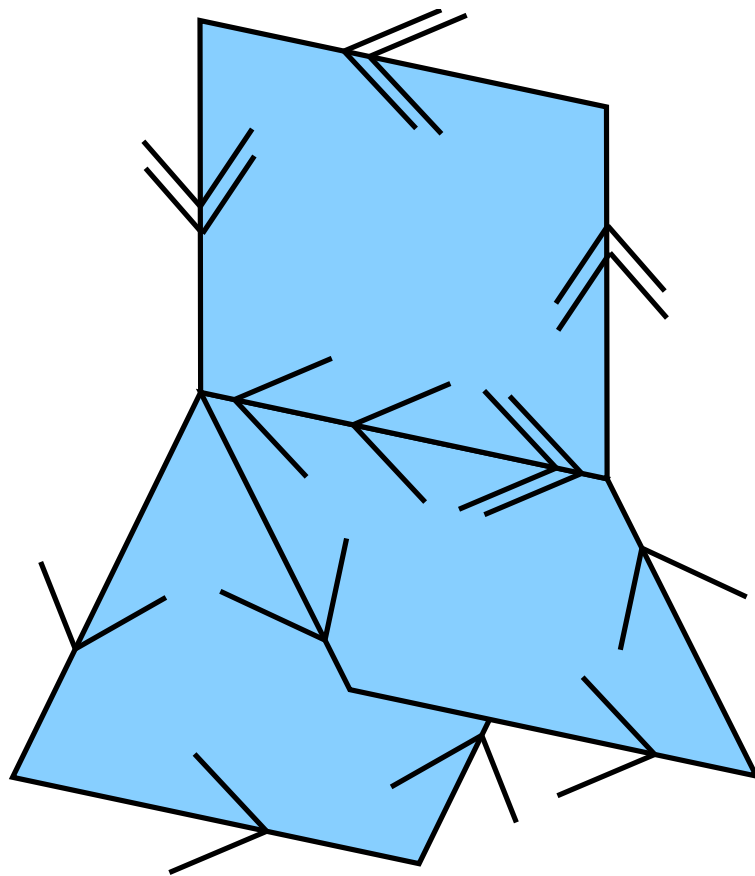
New Mexico State University

## Random vortex world-surface model:

- SU(2)
  - Confined, deconfined phases, second order deconfinement transition
  - Topological susceptibility
  - Quenched chiral condensate
- SU(3)
  - Confined, deconfined phases, weakly first order deconfinement transition
  - Baryonic Y law
  - Topological susceptibility
- SU(4), Sp(2)
  - Confined, deconfined phases, first order deconfinement transition
  - Increasing complexity of effective dynamics

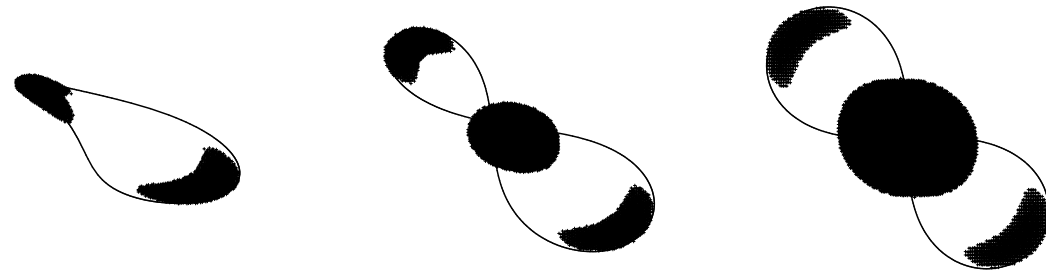
## SU(3) center vortices

- Two-dimensional closed world-surfaces of chromomagnetic flux
- Flux quantized such as to contribute a center phase  $\exp(\pm 2\pi i/3)$  to any Wilson loop to which they are linked (orientation of flux determines sign)
- Branching possible (as opposed to SU(2)):



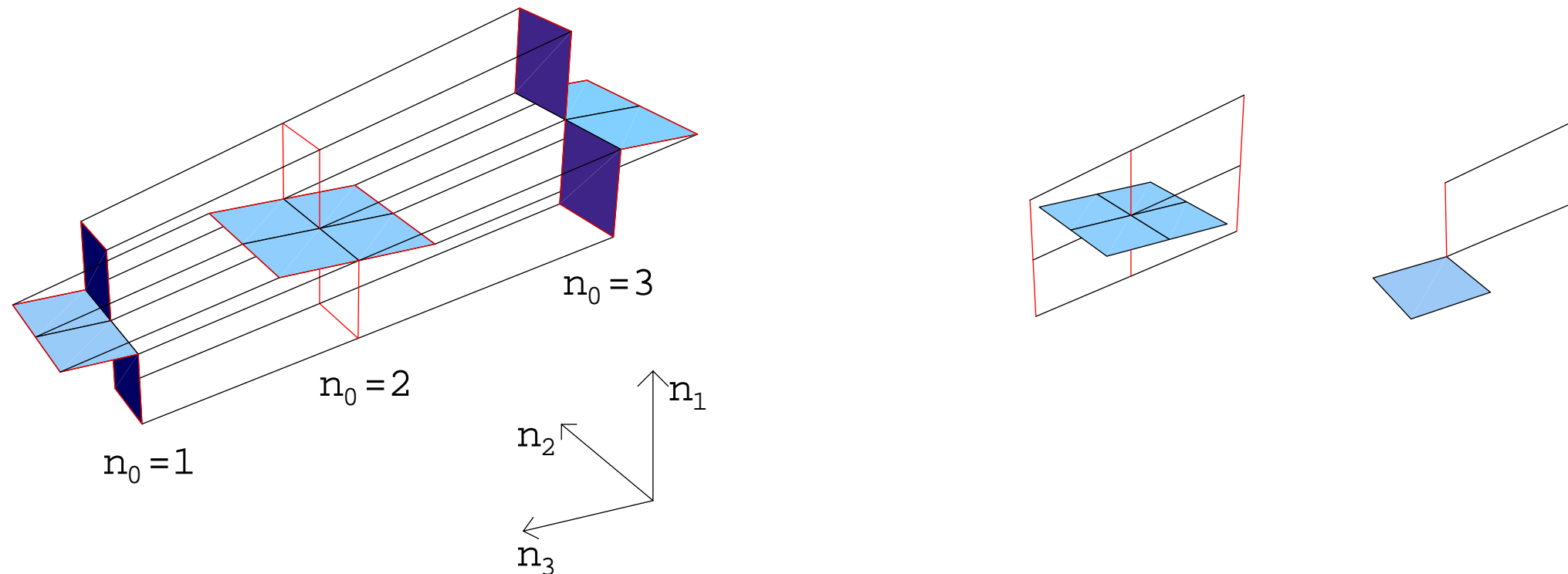
## Vortex topological charge

- Topological charge carried by vortex surface writhe and intersection points



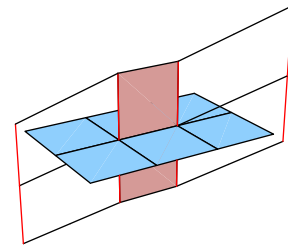
$$Q = \frac{1}{32\pi^2} \int d^4x \epsilon_{\mu\nu\lambda\tau} \text{Tr} F_{\mu\nu} F_{\lambda\tau}$$

- Orientation of surfaces enters via sign of  $F$  and generic surfaces are non-orientable – in Abelian gauge, surface patches of varying orientation separated by monopoles
- On lattice, topological charge concentrated at sites:



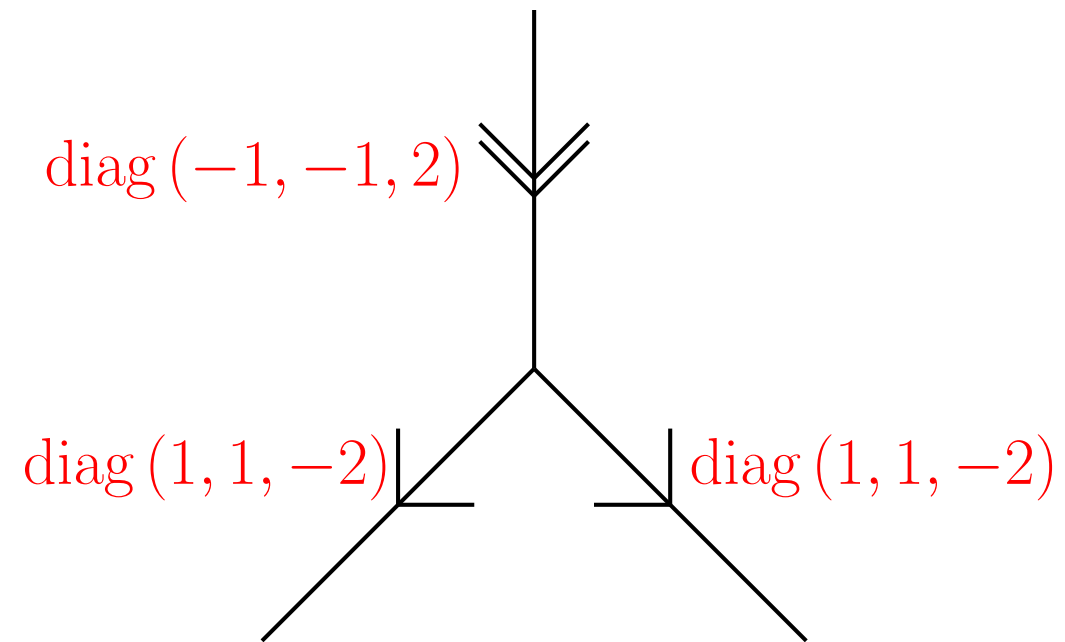
## Resolving ambiguities in lattice surfaces

- Resolve intersection lines  
lattices  $(1/3, 1/9)$

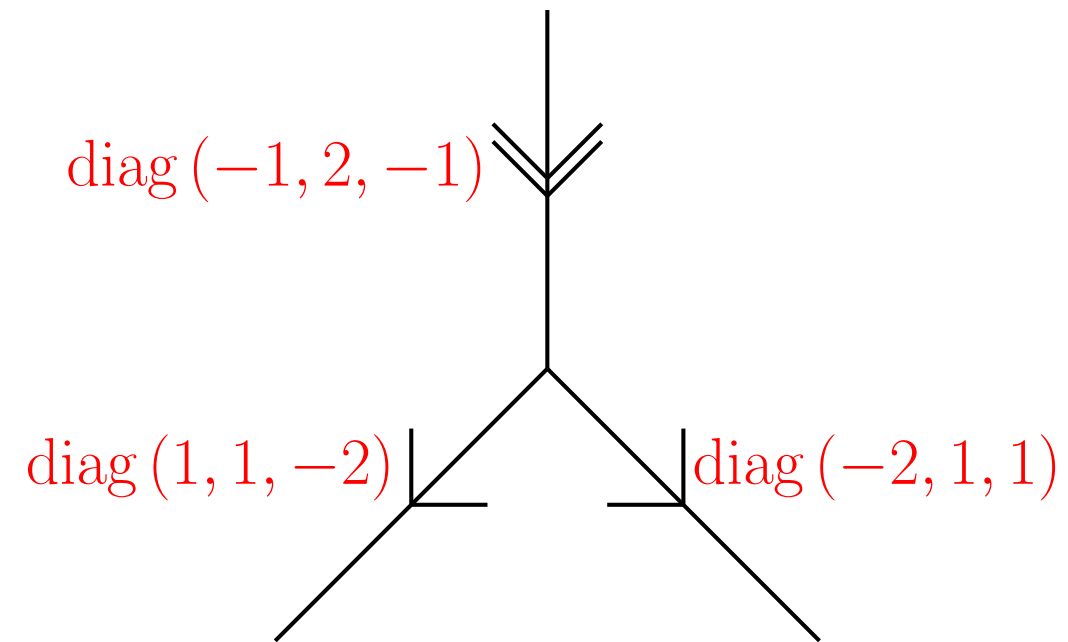


via slight deformations of surfaces on finer

- Edges of patches coinciding with sites carrying topological charge – use suitable color structure at branchings!



Minimal description:  
 $T = \pm \text{diag}(1, 1, -2)$



Non-minimal description:  
 $T = \pm \text{diag}(1, 1, -2),$   
 $\pm \text{diag}(1, -2, 1),$   
 $\pm \text{diag}(-2, 1, 1)$

$$W = (1/3) \text{Tr} \exp(2\pi i T/3)$$

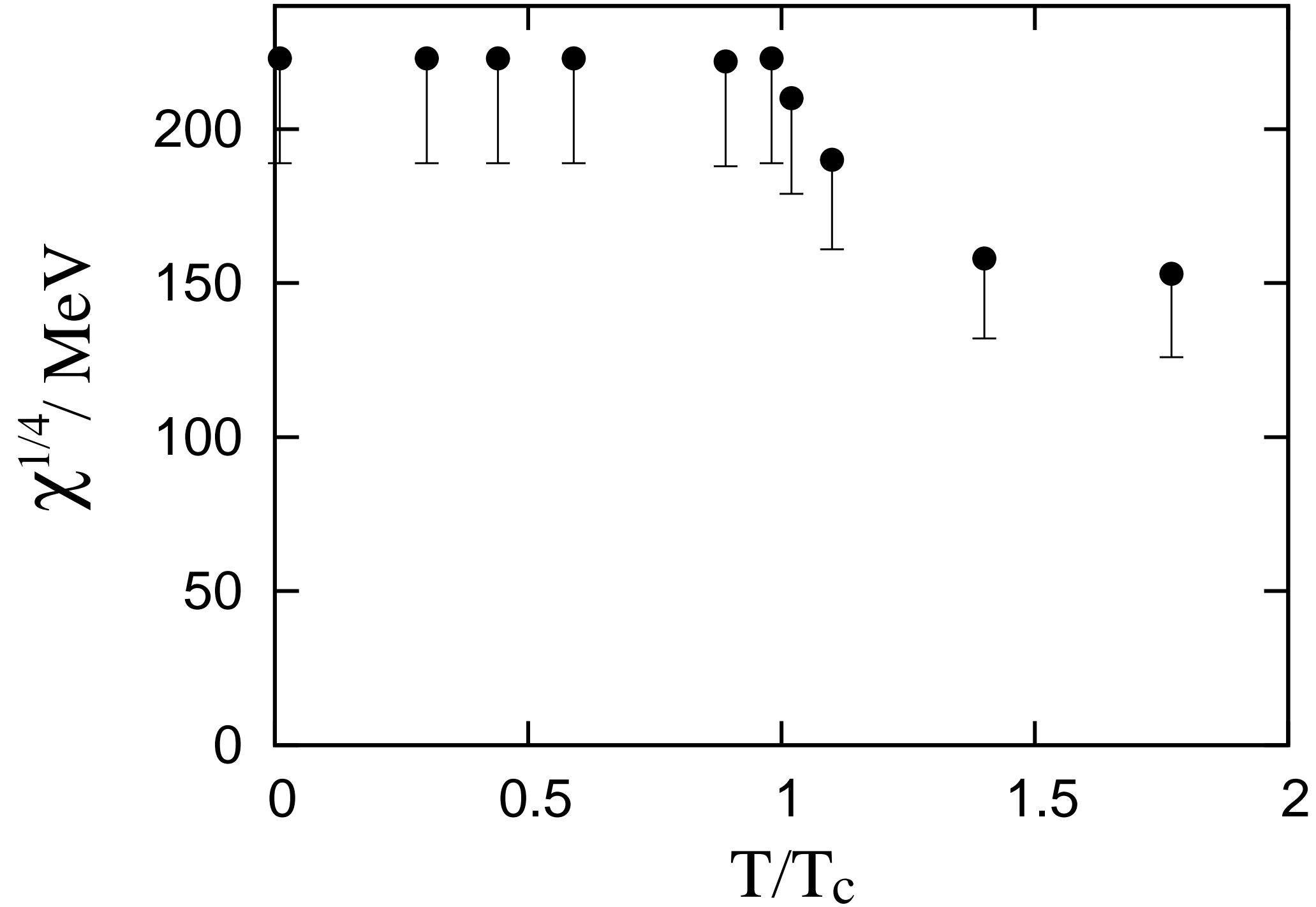
## SU(3) random vortex world-surface model

- Vortex world-surfaces composed of elementary squares on a hypercubic lattice
- Fixed lattice spacing  $\longleftrightarrow$  vortex thickness
- Vortex world-surface curvature action

$$S = c \times \text{[Diagram of a square with a thickened top edge representing a vortex world-surface element]}$$

- Dimensionless coupling constant  $c$  tuned to reproduce ratio  $T_c/\sqrt{\sigma} = 0.63$  from SU(3) Yang-Mills theory  $\longrightarrow c = 0.21$
- Scale determined by setting  $\sigma = (440 \text{ MeV})^2 \longrightarrow$  lattice spacing  $a = 0.39 \text{ fm}$

# Topological susceptibility



## Conclusions

- Substantial systematic uncertainty of topological charge measurement engendered by the necessity to remove ambiguities introduced through the hypercubic description
- Taking into account this uncertainty, topological susceptibility predicted by the SU(3) random vortex world-surface model is compatible with SU(3) Yang-Mills theory
- Alternative representation of the vortex surfaces (e.g., random triangulations) would remove this uncertainty