

Search for Chiral Fermion Actions on Non-Orthogonal Lattices

Presented by Michael I. Buchoff

In collaboration with Paulo Bedaque, Brian Tiburzi, and
Andre Walker-Loud

P.F. Bedaque, M.I.B, B.C. Tiburzi, and A. Walker-Loud (Phys. Lett. B 662, 449-455), 2008

P.F. Bedaque, M.I.B, B.C. Tiburzi, and A. Walker-Loud (Phys. Rev. D 78, 017502), 2008

Introduction

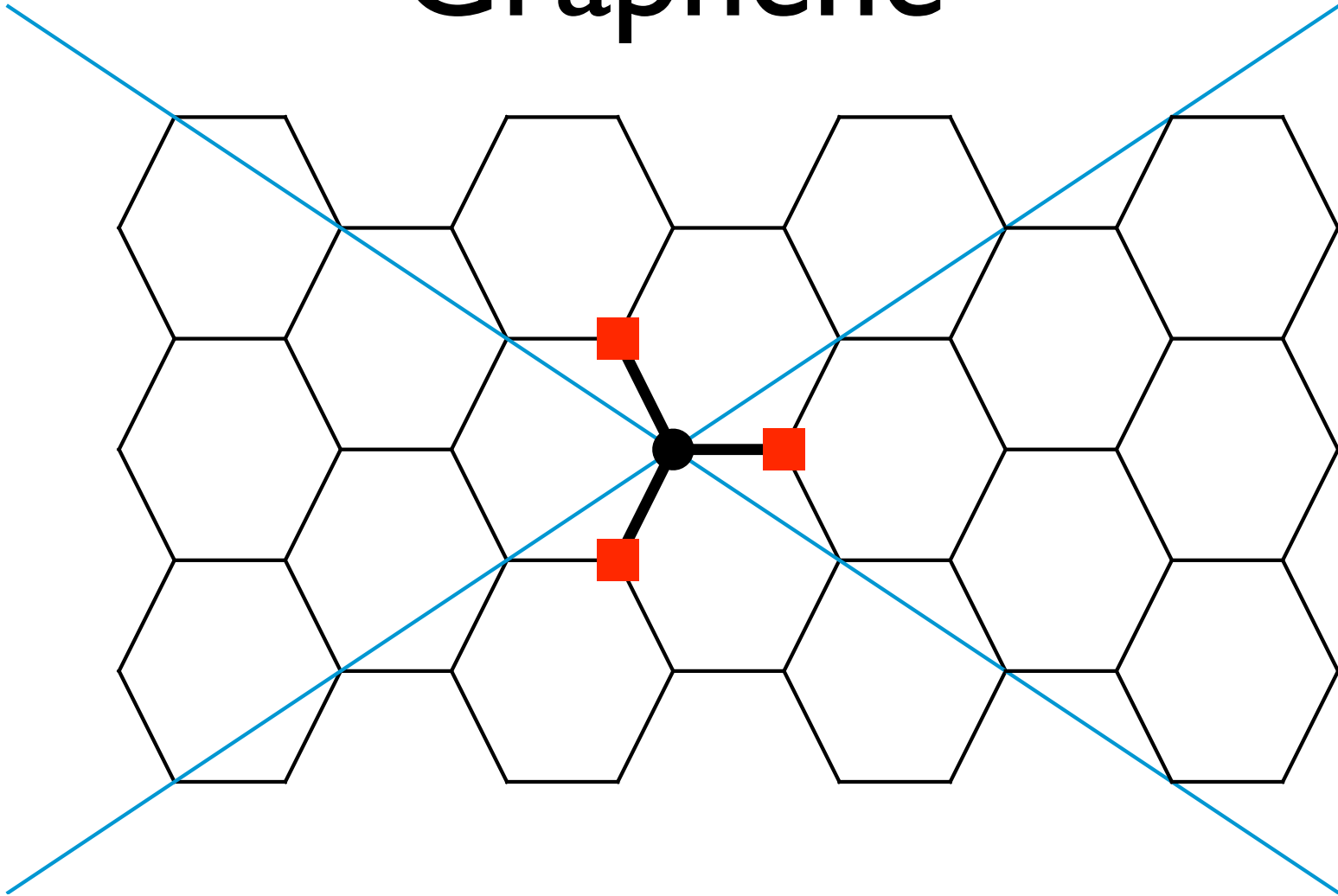
- Why use non-orthogonal lattices (“graphene”, “hyperdiamond”, etc.)?
- Complications with proposed “graphene” lattice actions
- Advantages and disadvantages of other non-orthogonal lattice actions

Brief Intro on Graphene

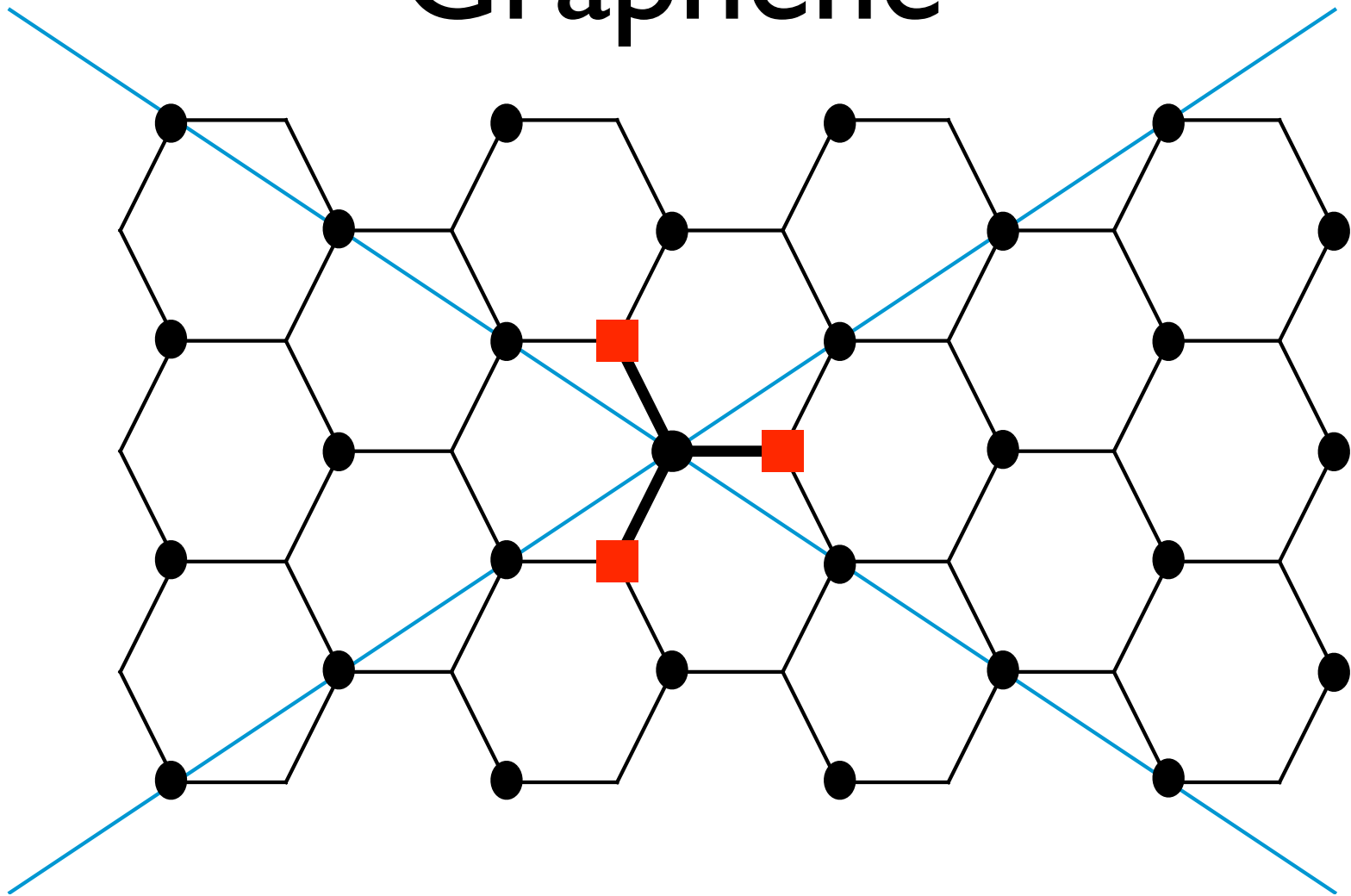
- Two dimensional honeycomb (hexagonal) lattice
- Leads to Dirac fermions in the massless limit
 - This effect lead to great interest throughout condensed matter community

Graphene

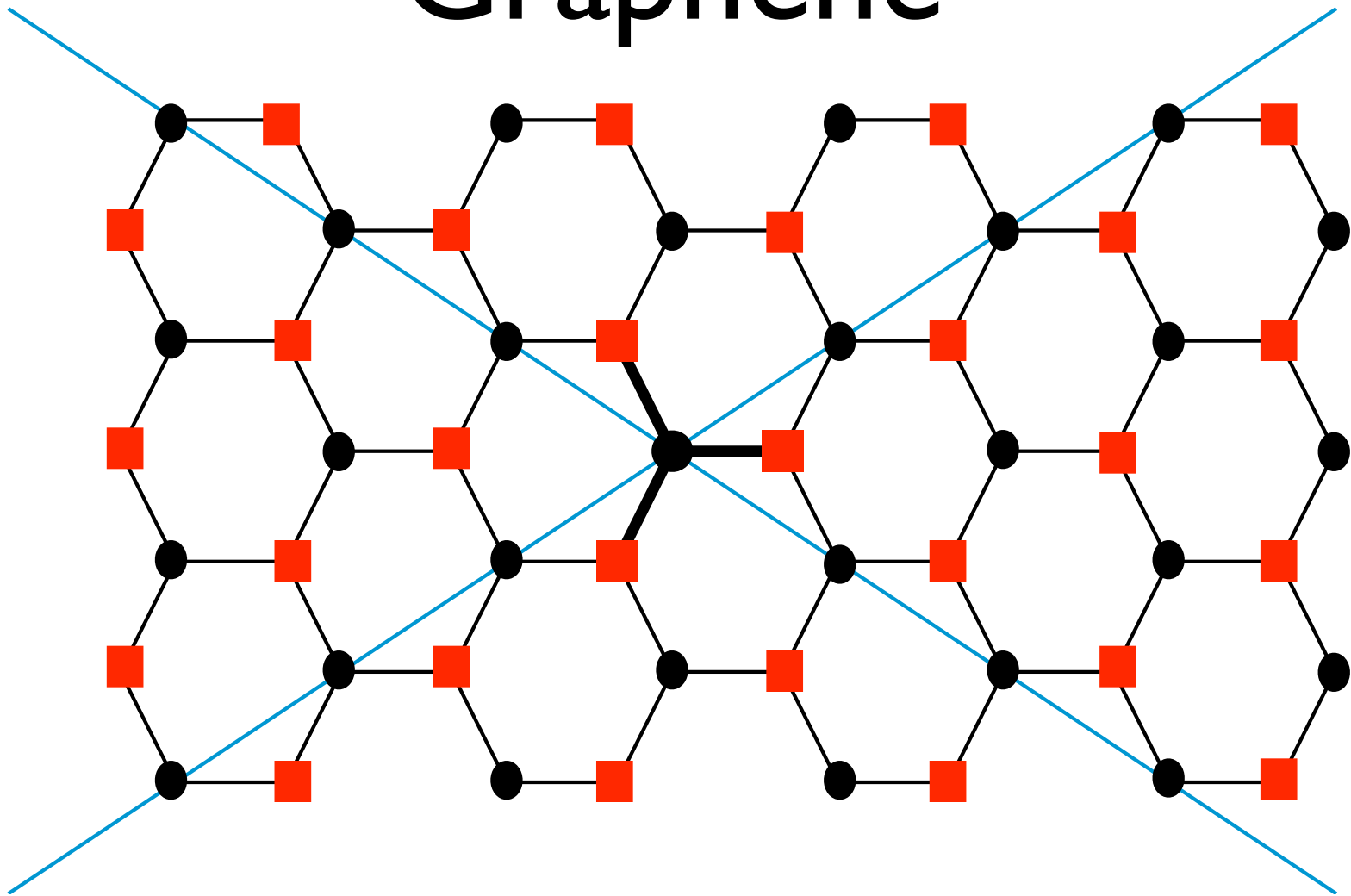
Graphene



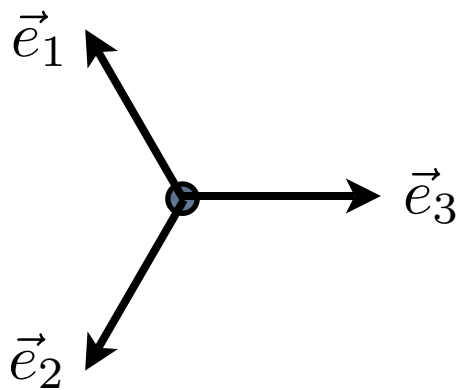
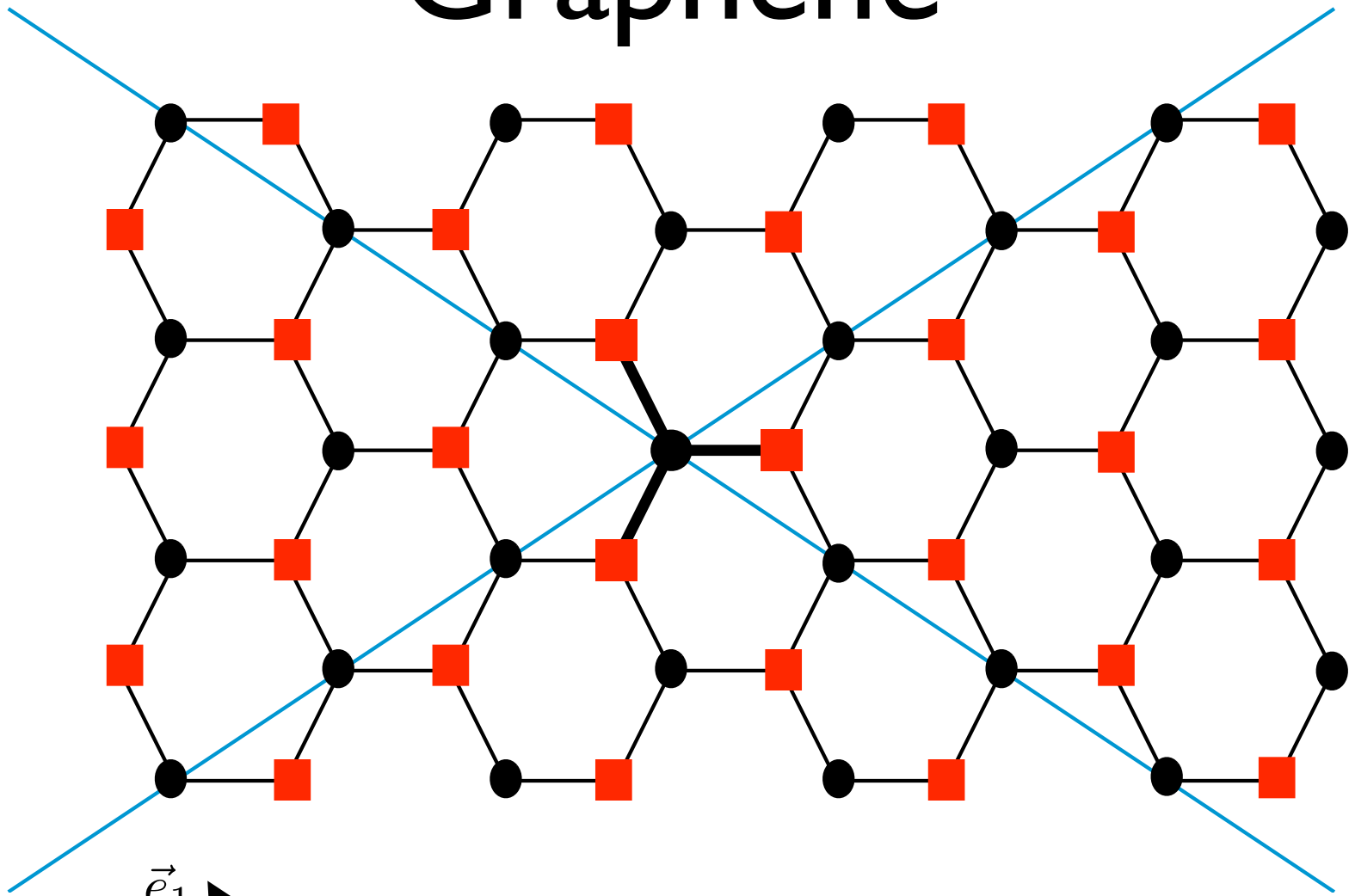
Graphene



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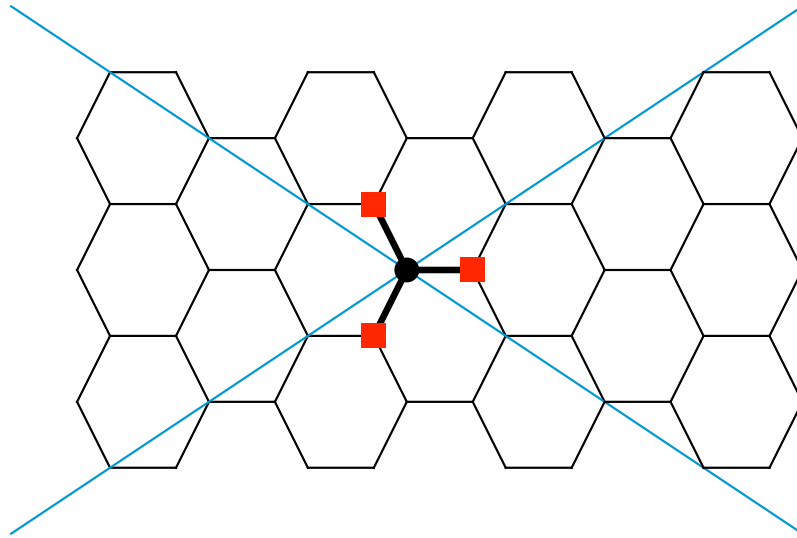


Graphene

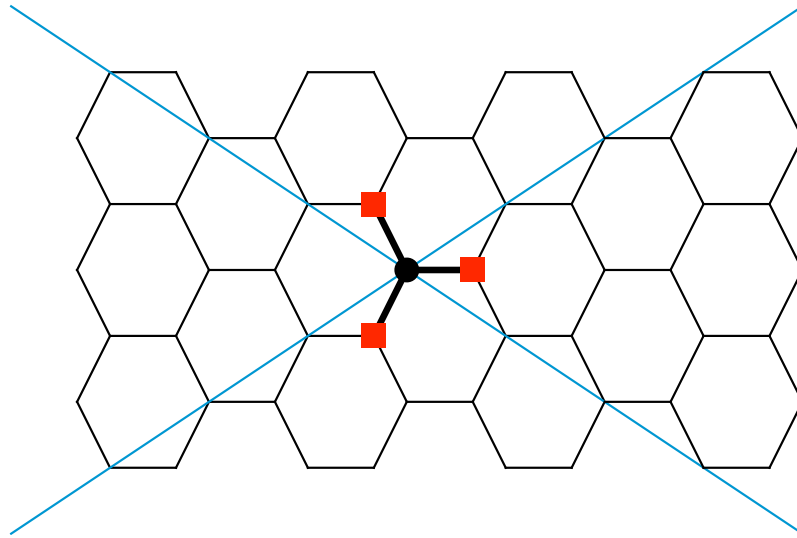


$$\sum_{\alpha} \mathbf{e}^{\alpha} = 0 \quad \text{for } \alpha = 1, 2, 3$$

Graphene (cont.)

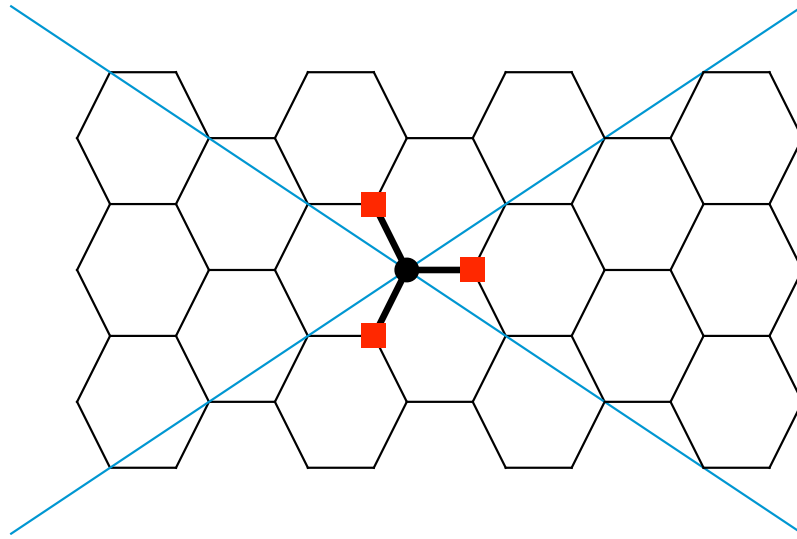


Graphene (cont.)



- 1+1 dimensional graphene lattices in the massless limit lead to chiral Dirac fermions.

Graphene (cont.)



- $1+1$ dimensional graphene lattices in the massless limit lead to chiral Dirac fermions.
- Question: Can this construction be extended to a four dimensional lattice gauge theory?

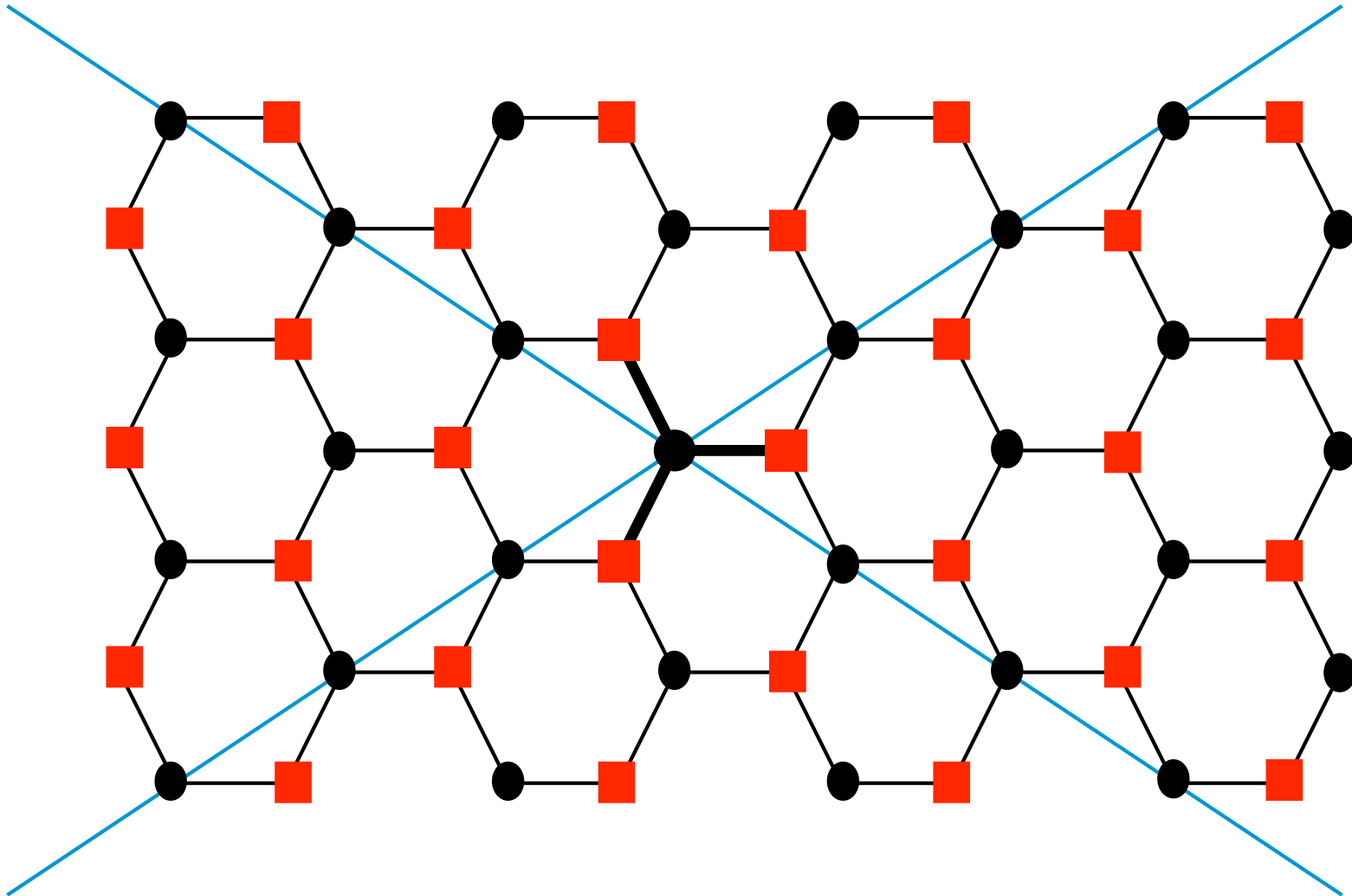
Graphene-Inspired Lattice

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- Method recently proposed by Creutz (JHEP 0804:017,2008)

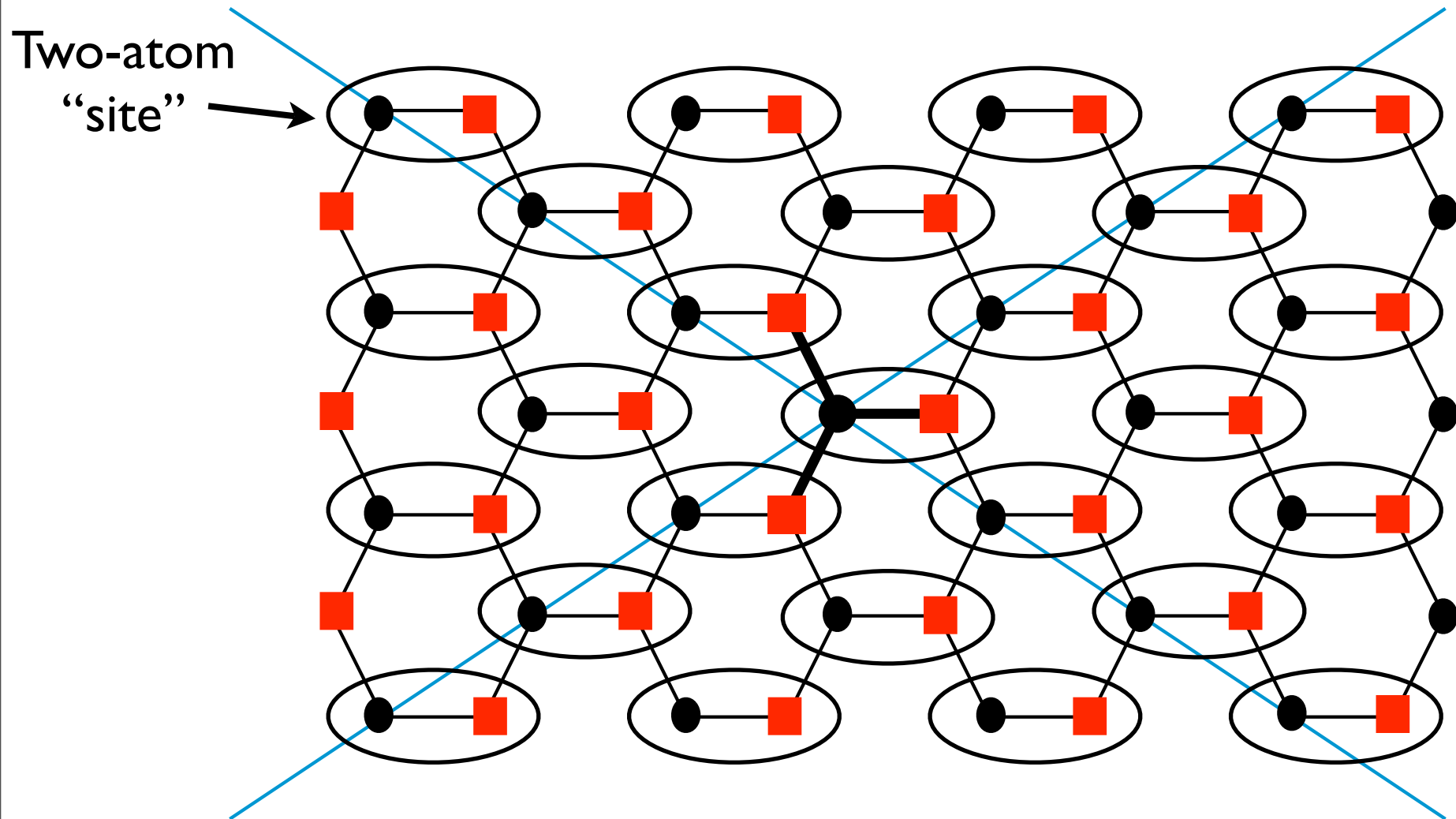
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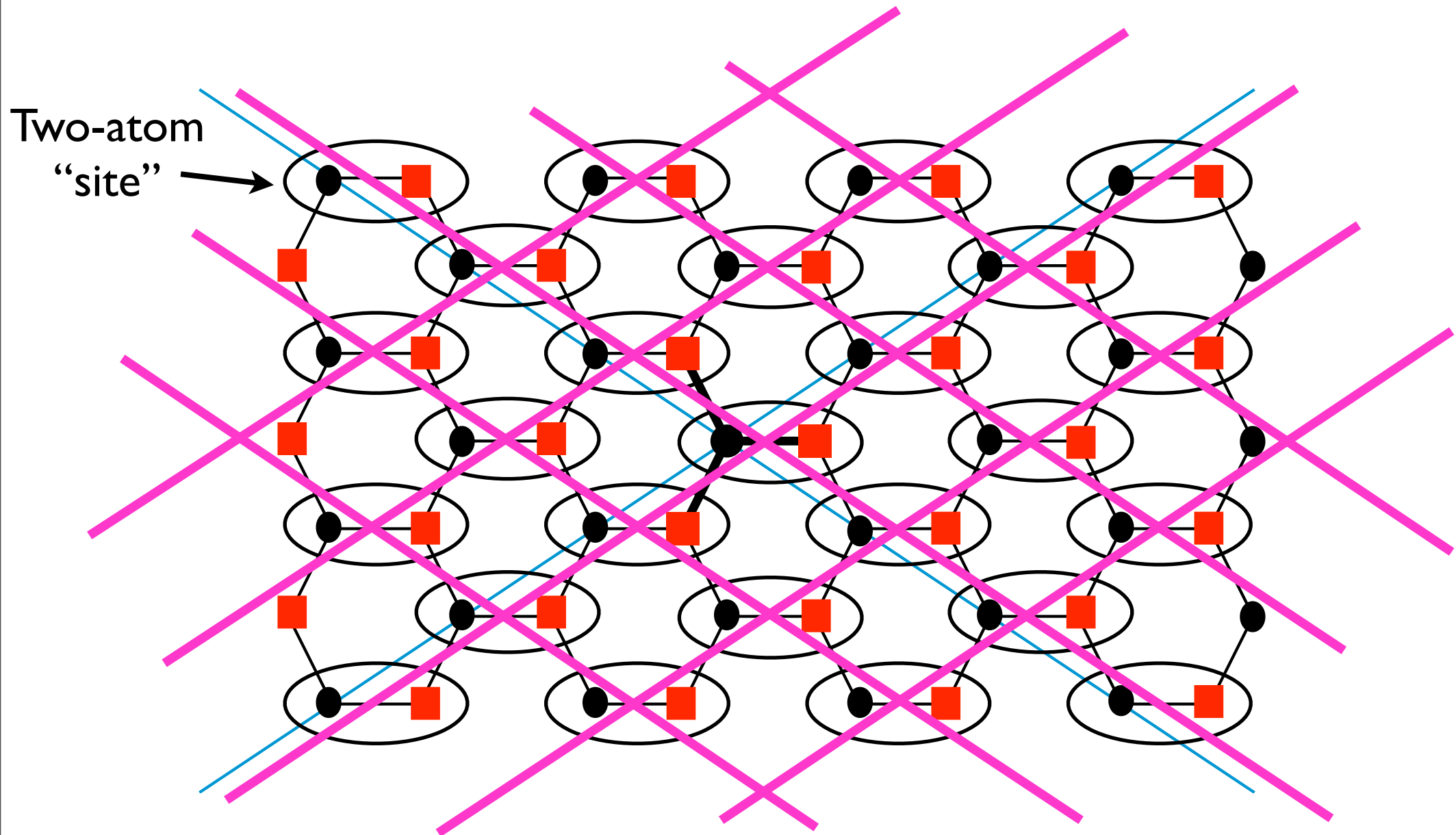
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Borici-Creutz Action

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$$S_{BC} = \frac{1}{2} \sum_x \left[\sum_{\mu} (\bar{\psi}_{x-\mu} \mathbf{e}^{\mu} \cdot \Gamma \psi_x - \bar{\psi}_{x+\mu} \mathbf{e}^{\mu} \cdot \Gamma^{\dagger} \psi_x) + \bar{\psi}_x \mathbf{e}^5 \cdot \Gamma \psi_x - \bar{\psi}_x \mathbf{e}^5 \cdot \Gamma^{\dagger} \psi_x \right]$$

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$$\Gamma_{\mu} = (\vec{\gamma}, i\gamma_4)$$

$$\mu = 1, 2, 3, 4$$

$$\mathbf{e}^1 = (1, 1, 1, B)$$

$$\mathbf{e}^2 = (1, -1, -1, B)$$

$$\mathbf{e}^3 = (-1, -1, 1, B)$$

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$$\mathbf{e}^5 = -(0, 0, 0, 4BC)$$

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Momentum Space:

$$S_{BC} = \int_p \bar{\psi}_p \left[i \sum_{\mu} \left(\sin(p_{\mu}) \vec{\mathbf{e}}^{\mu} \cdot \vec{\gamma} + B\gamma_4 (\cos(p_{\mu}) - C) \right) \right] \psi_p$$

Borici-Creutz (cont.)

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Chiral Symmetry

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Chiral Symmetry

Minimal Doubling (two Poles): If $B \neq 0$ and $0 < C < 1$

$$p_{\mu}^{(1)} = \tilde{p}(1, 1, 1, 1) \quad p_{\mu}^{(2)} = -\tilde{p}(1, 1, 1, 1) \quad \cos(\tilde{p}) = C$$

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Bad News:

Broken Discrete Symmetries:

Parity: $\psi(\vec{p}, p_4) \rightarrow \gamma_4 \psi(-\vec{p}, p_4)$

Charge Conjugation: $\psi(\vec{p}, p_4) \rightarrow C \bar{\psi}^T(\vec{p}, p_4)$

Time Reversal: $\psi(\vec{p}, p_4) \rightarrow \gamma_5 \gamma_4 \psi(\vec{p}, -p_4)$

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- Terms will be generated unless additional symmetry of the action prevents them.

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4D: If \mathbb{Z}_5 (or A_5 or S_5) symmetric

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- Does Borici-Creutz have the minimal \mathbb{Z}_5 symmetry?

Permutation Symmetry?

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$$\begin{aligned}
 S_{BC} = \frac{1}{2} \sum_x \left[\sum_{\mu} (\bar{\phi}_{x-\mu} \Sigma \cdot \mathbf{e}^{\mu} \chi_x - \bar{\chi}_{x+\mu} \Sigma \cdot \mathbf{e}^{\mu} \phi_x) \right. \\
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 \end{aligned}$$

$$\psi_x = \begin{pmatrix} \phi_x \\ \chi_x \end{pmatrix}$$

$$\bar{\psi}_p = (\bar{\phi}_p, \bar{\chi}_p)$$

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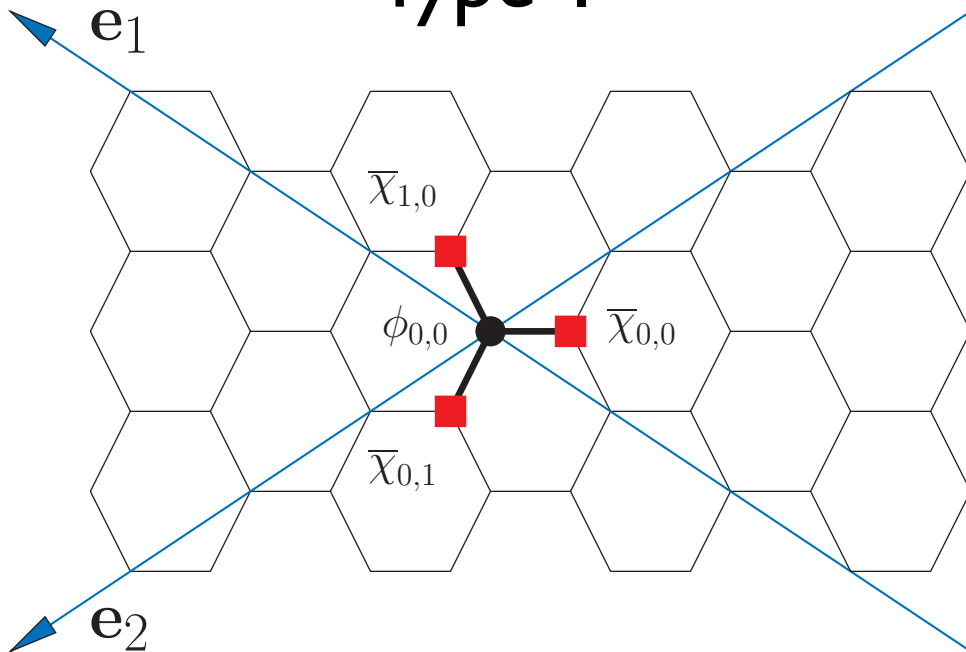
$$\bar{\psi}_p = (\bar{\phi}_p, \bar{\chi}_p)$$

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For $B = 1/\sqrt{5}$
and $C = 1$

Type I



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Type 1

Type 2

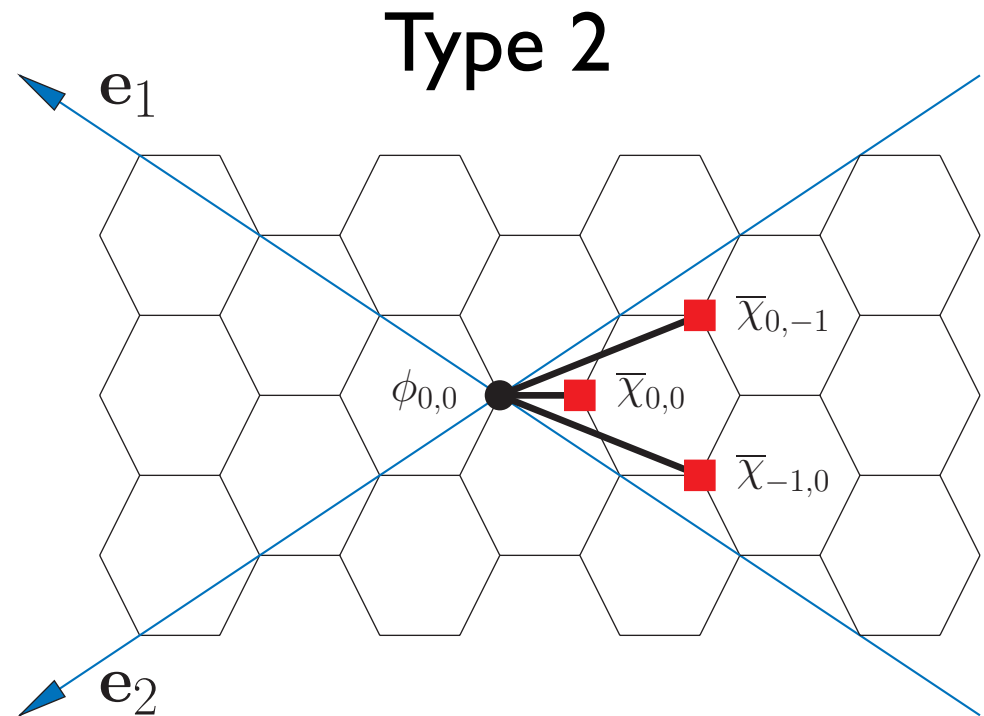
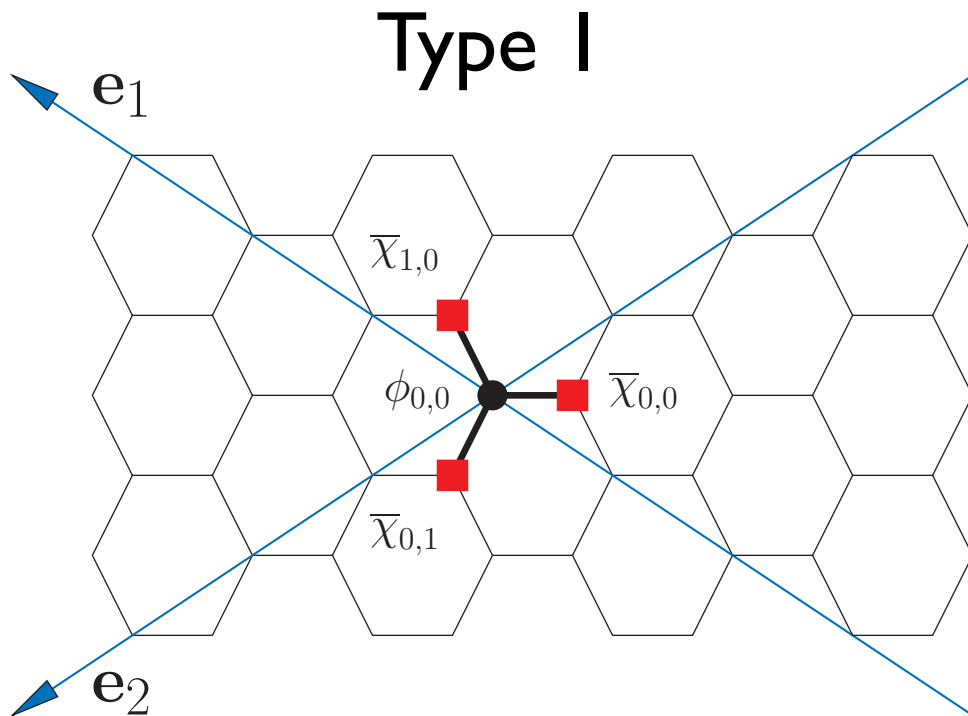
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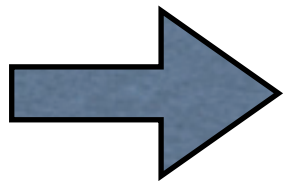
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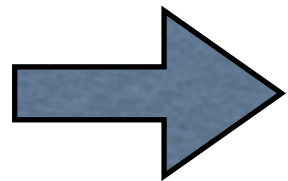


Action around pole $p_{\mu} \simeq 0$: $\bar{\psi}(i\vec{\gamma} \cdot \vec{k} + \gamma_5 \gamma_4 k_4)\psi$

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“Mutilated” Fermions

“Hyperdiamond”

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$$S = \sum_x \left[\sum_{\mu=1}^4 \left(\bar{\phi}_{x-\mu} \sigma \cdot \mathbf{e}^\mu \chi_x - \bar{\chi}_{x+\mu} \bar{\sigma} \cdot \mathbf{e}^\mu \phi_x \right) + \bar{\phi}_x \sigma \cdot \mathbf{e}^5 \chi_x - \bar{\chi}_x \bar{\sigma} \cdot \mathbf{e}^5 \phi_x \right]$$

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$$\mathbf{e}^1 = \frac{1}{4} (\sqrt{5}, \sqrt{5}, \sqrt{5}, 1)$$

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$$\sum_{\alpha} \mathbf{e}^{\alpha} = 0$$

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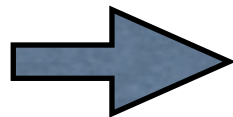
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Has A_5 symmetry



No relevant operators

Hyperdiamond (cont.)

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- Hyperdiamond action in momentum space:

$$S = \int_p \bar{\psi}_p \left[i \sum_{\mu} \sin(p_{\mu}) \mathbf{e}^{\mu} \cdot \boldsymbol{\gamma} - \left(\sum_{\mu} \cos(p_{\mu}) \mathbf{e}^{\mu} + \mathbf{e}^5 \right) \cdot \boldsymbol{\gamma} \gamma_5 \right] \psi_p$$

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Ex:

$$p_1 = -p_2 = -p_3 = p_4 = \cos^{-1}(-2/3)$$

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- Additional broken symmetries can (and often will) lead to relevant and marginal operators from radiative corrections.
- An intricate balance of symmetry is needed for chiral symmetry, minimal doubling, and no relevant operators.
 - To this point, no non-orthogonal action has accomplished this balance.