

Challenges in Hadronic Form-Factor Calculations



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In collaboration with

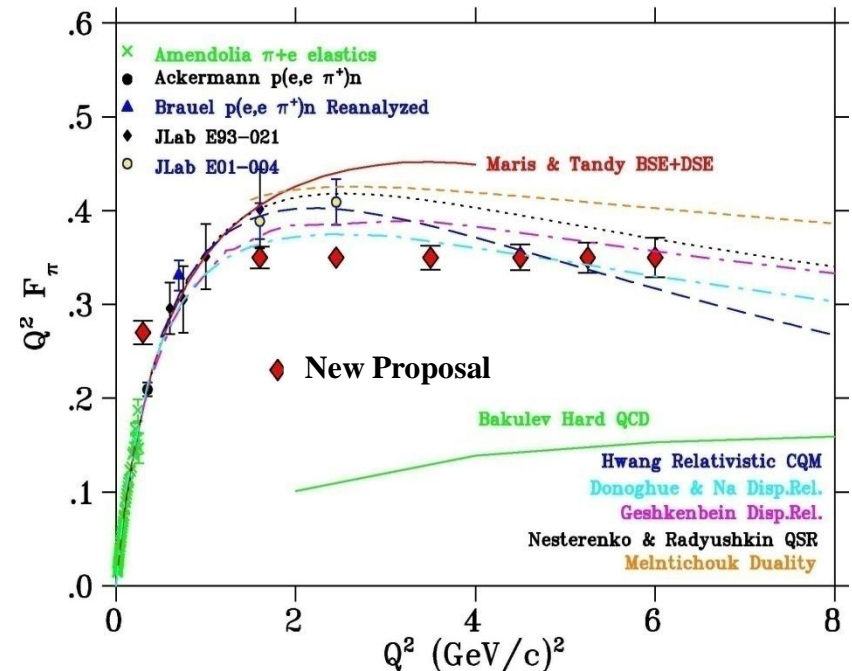
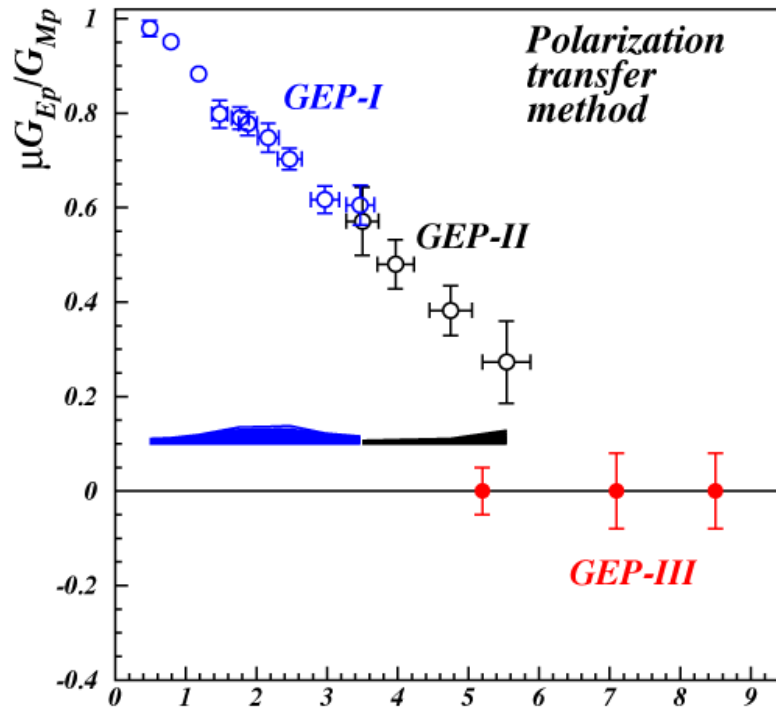
Saul Cohen, Robert Edwards, and David Richards (JLab)



Kostas Orginos

Higher- Q^2 Form Factor

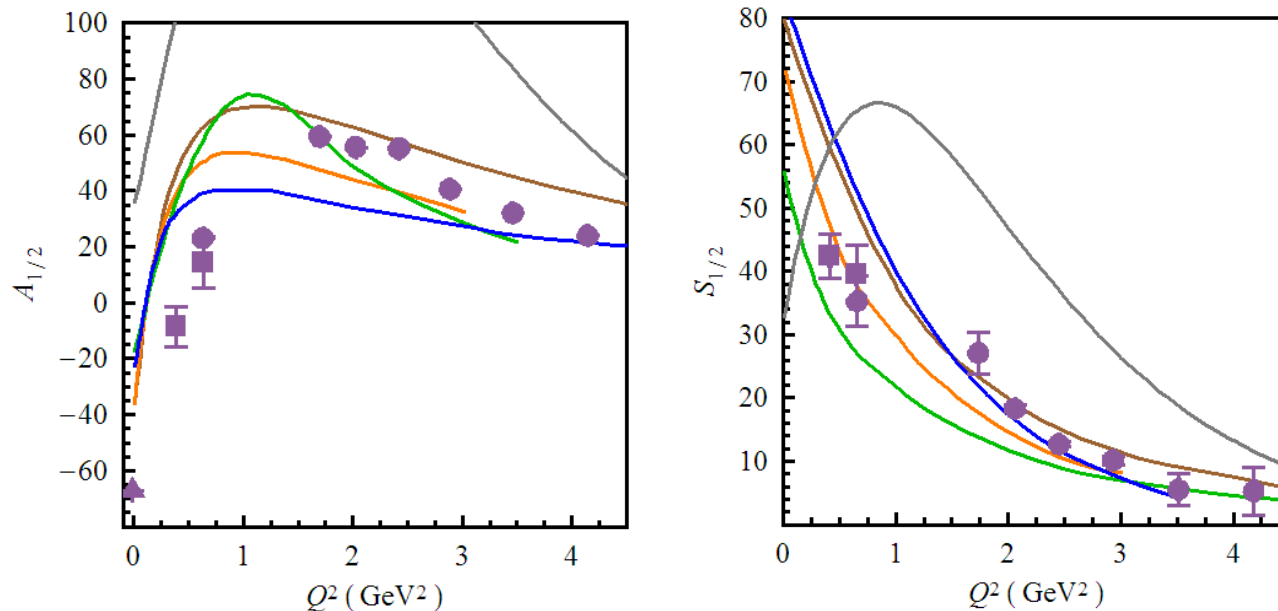
- Higher- Q^2 data will help us to understand hadrons and challenge QCD-based models
- Experiments are looking...
(JLab 12-GeV upgrade will provide more promising data)



- Want to know: does G_E^p cross zero at large Q^2 ? If so, where?
What's the pion form factor's behavior?

$N-P_{11}$ Form Factor

- ◆ Experiments at Jefferson Laboratory (**CLAS**), MIT-Bates, LEGS, Mainz, Bonn, GRAAL, and Spring-8
- ◆ **Helicity amplitudes** are measured (in $10^{-3} \text{ GeV}^{-1/2}$ units)
- ◆ Many models disagree (a selection are shown below)

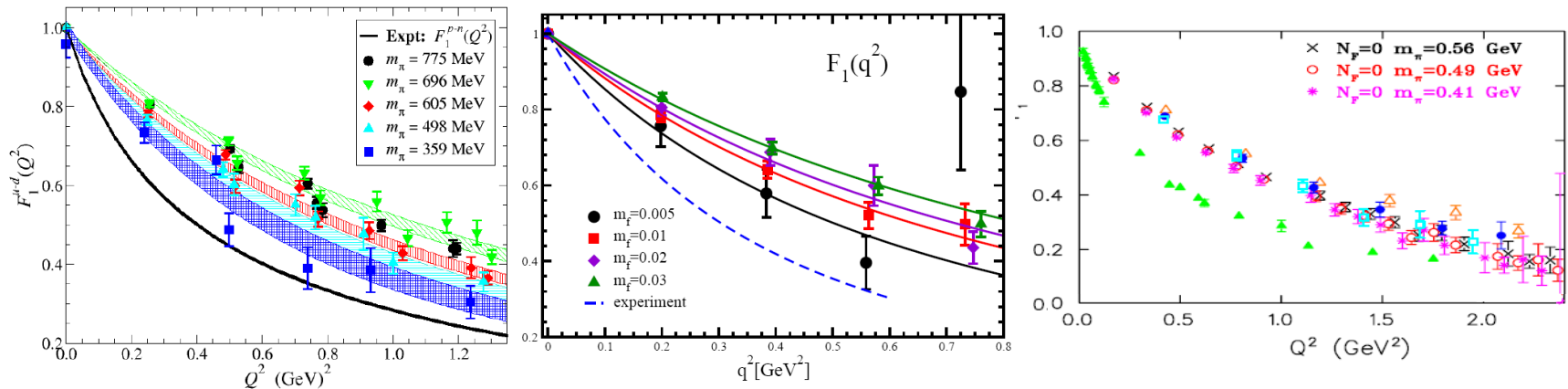


- ◆ If the Roper is the first radially excited state of nucleon, this is the data to compare with

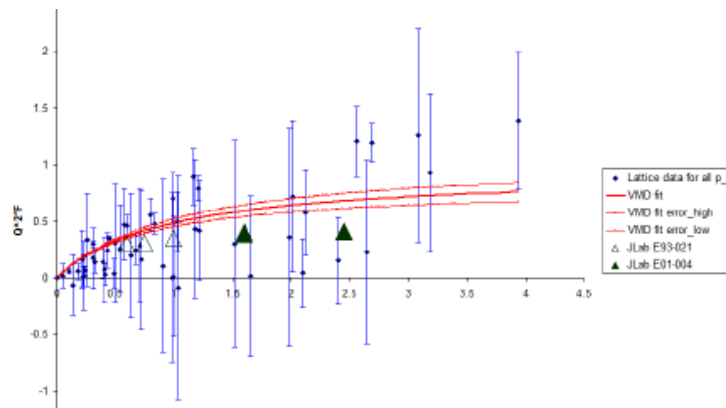
Focusing on Two Challenges

◆ Large- Q^2 calculations

◆ Typical Q^2 range for nucleon form factors is $< 3.0 \text{ GeV}^2$



◆ Higher- Q^2 calculations suffer from poor noise-to-signal ratios

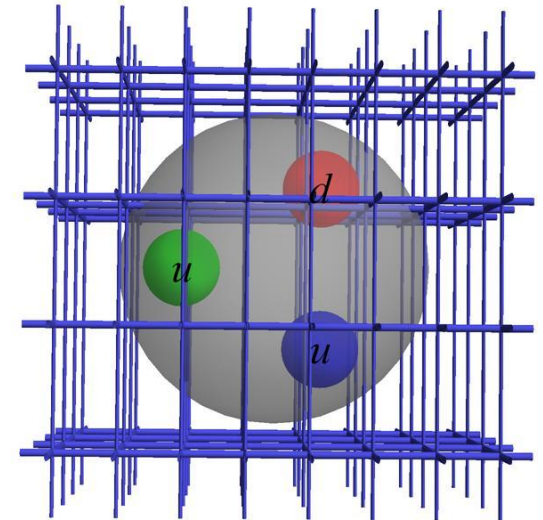


Focusing on Two Challenges

- ◆ Large- Q^2 calculations
 - ◆ Typical Q^2 range for nucleon form factors is $< 3.0 \text{ GeV}^2$
 - ◆ Higher- Q^2 calculations suffer from poor noise-to-signal ratios
- ◆ Radially excited transition form factor calculations
 - ◆ A different way to understand the properties of excited hadrons
 - ◆ Euclidean space: signal falls exponentially with time
 - ◆ Help from anisotropic lattices

Lattice Setup

- ◆ “Quenched” for exploratory study
 - ◆ $16^3 \times 64$ **anisotropic** lattice, $\xi = 3$
 - ◆ Better resolution for excited states
 - ◆ Wilson gauge action + clover fermion action
 - ◆ $a_t^{-1} \approx 6$ GeV and $a_s \approx 0.125$ fm ($L < 2$ fm)
 - ◆ $m_\pi \approx 720$ (480 and 1100) MeV
 - ◆ 200 configurations



Anisotropic lattice

- ◆ Baryon interpolating field

$$J_\alpha(\vec{p}, t) = \sum_{\vec{x}, a, b, c} e^{i\vec{p} \cdot \vec{x}} \epsilon^{abc} [u_a^T(y_1, t) C \gamma_5 d_b(y_2, t)] u_{c, \alpha}(y_3, t) \phi(y_1 - x) \phi(y_2 - x) \phi(y_3 - x)$$

- ◆ No disconnected contributions
- ◆ Multiple Gaussian smearings

Form Factors

- ◆ The form factors are buried in the amplitudes

$$\begin{aligned} & \Gamma_{\mu,AB}^{(3),T}(t_i, t, t_f, \vec{p}_i, \vec{p}_f) \\ &= a^3 \sum_n \sum_{n'} \frac{1}{Z_j} \frac{Z_{n',B}(p_f) Z_{n,A}(p_i)}{4E'_n(\vec{p}_f) E_n(\vec{p}_i)} e^{-(t_f-t)E'_n(\vec{p}_f)} e^{-(t-t_i)E_n(\vec{p}_i)} \\ & \times \sum_{s,s'} T_{\alpha\beta} u_{n'}(\vec{p}_f, s')_{\beta} \langle N_{n'}(\vec{p}_f, s') | j_{\mu}(0) | N_n(\vec{p}_i, s) \rangle \bar{u}_n(\vec{p}_i, s)_{\alpha} \end{aligned}$$

- ◆ Nucleon form factor ($n = n' = 0$)

$$\langle N | V_{\mu} | N \rangle(q) = \bar{u}_N(p') \left[\gamma_{\mu} F_1(q^2) + \sigma_{\mu\nu} q_{\nu} \frac{F_2(q^2)}{2m} \right] u_N(p) e^{-iq \cdot x}$$

- ◆ Nucleon-Roper form factor ($n = 0, n' = 1$ or $n = 1, n' = 0$)

$$\langle N_2 | V_{\mu} | N_1 \rangle_{\mu}(q) = \bar{u}_{N_2}(p') \left[F_1(q^2) \left(\gamma_{\mu} - \frac{q_{\mu}}{q^2} \not{q} \right) + \sigma_{\mu\nu} q_{\nu} \frac{F_2(q^2)}{M_{N_1} + M_{N_2}} \right] u_{N_1}(p) e^{-iq \cdot x}$$

- ◆ Need best possible input from two-point correlators

Variational Method

◆ Generalized eigenvalue problem:

[C. Michael, Nucl. Phys. B 259, 58 (1985)]

[M. Lüscher and U. Wolff, Nucl. Phys. B 339, 222 (1990)]

◆ Construct the matrix

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t)^\dagger \mathcal{O}_j(0) | 0 \rangle$$

◆ Solve for the generalized eigensystem of

$$C(t_0)^{-1/2} C(t) C(t_0)^{-1/2} v = \lambda(t, t_0) v$$

with eigenvalues

$$\lambda_n(t, t_0) = e^{-(t-t_0)E_n} (1 + \mathcal{O}(e^{-|\delta E|(t-t_0)}))$$

Now the original correlator matrix can be approximated by

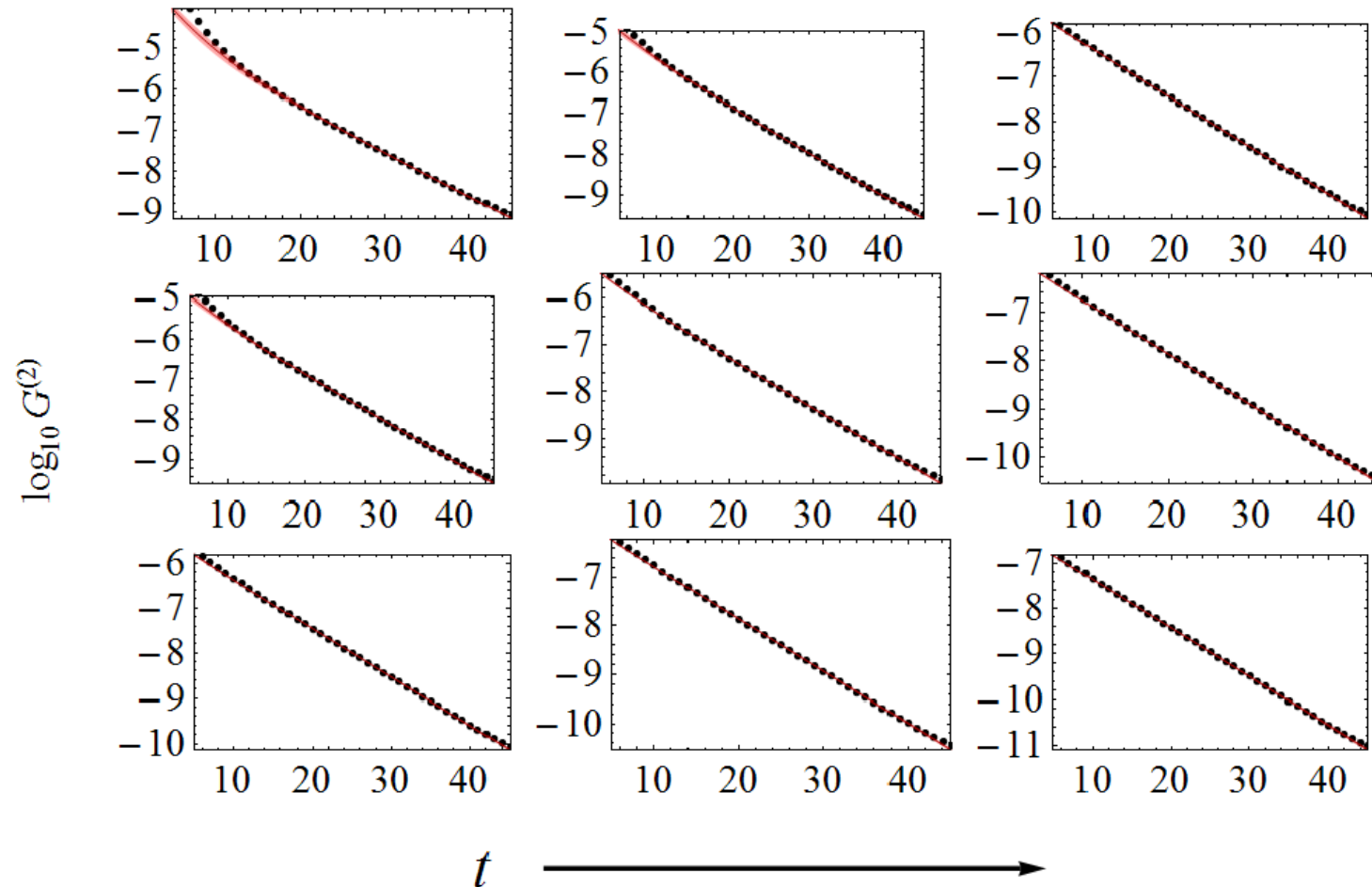
$$C_{ij} = \sum_{n=1}^r (C(t_0)^{1/2} v_n^*)_i (v_n C(t_0)^{1/2})_j \lambda_n(t, t_0) = \sum_n \frac{E_n + m}{2E_n} Z_{i,n} Z_{j,n} e^{-E_n t}$$

◆ Three smearings (i, j) are chosen for this work

◆ 2nd excited state is contaminated by remaining states

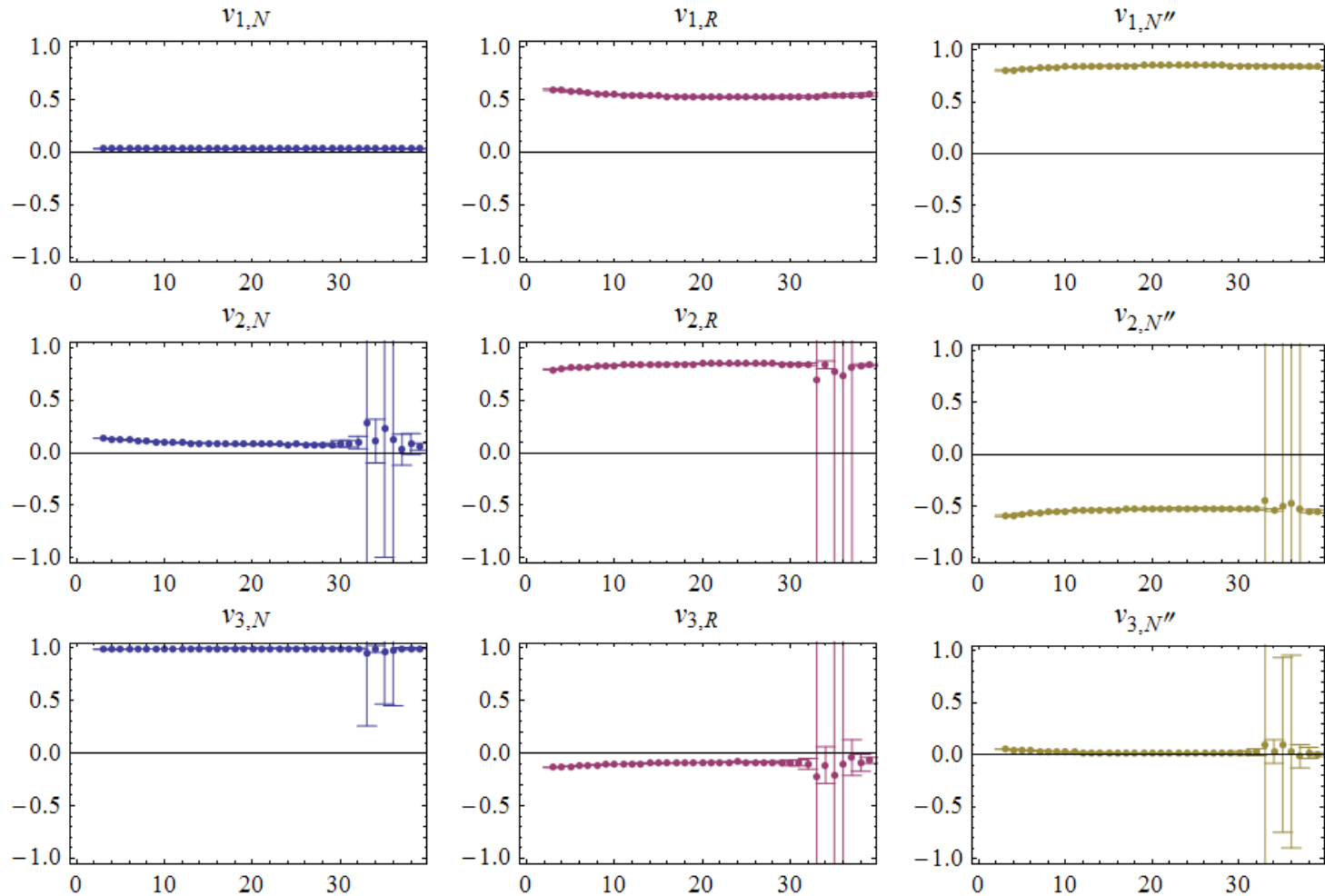
Variational Method

- ◆ Reconstruct two-point correlators from Z and λ



Variational Method

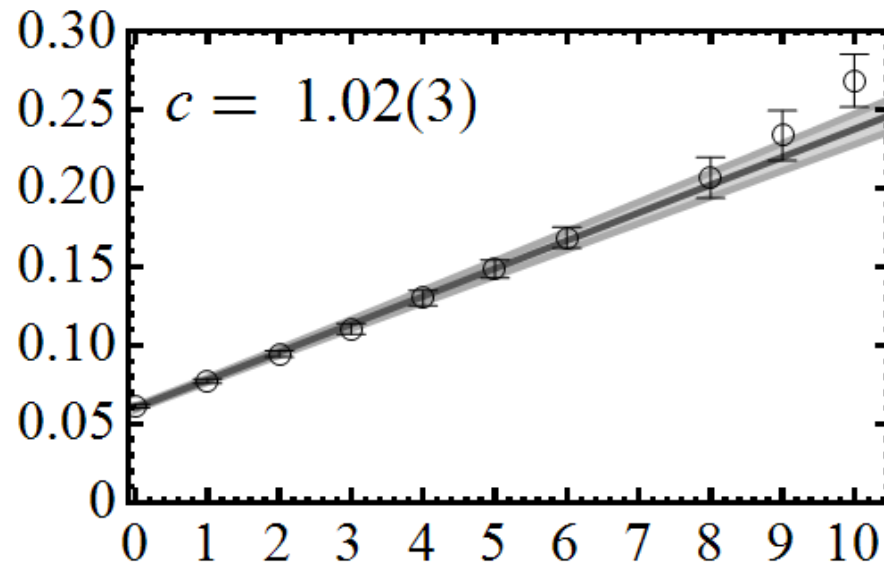
- ◆ Eigenvectors (at $p = 0$) show overlap of smearings with states



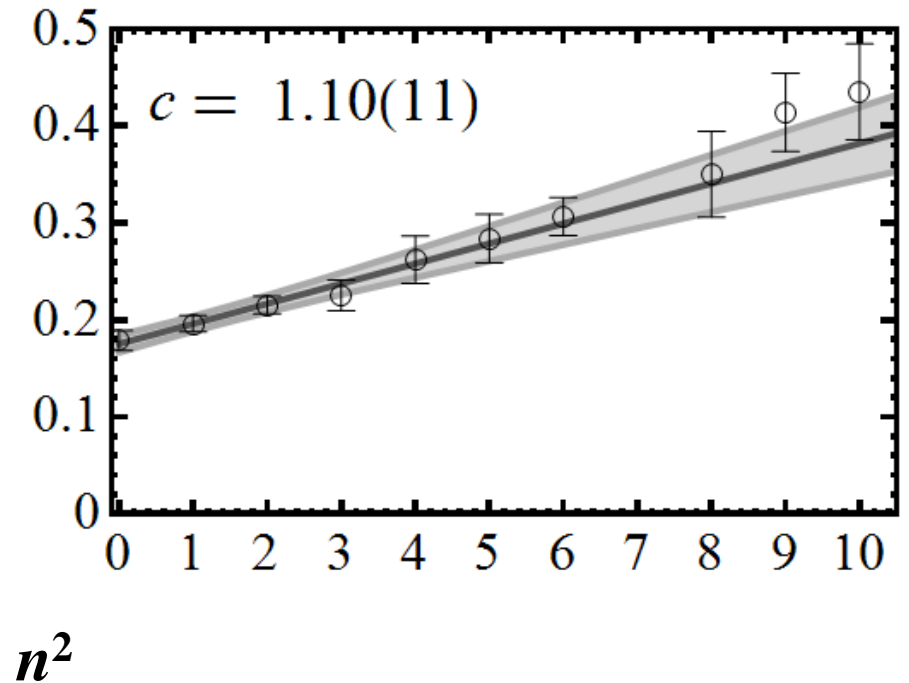
Dispersion Relation

◆ Example: $m_\pi = 720$ MeV

$a_t^2 E^2$ **Nucleon**



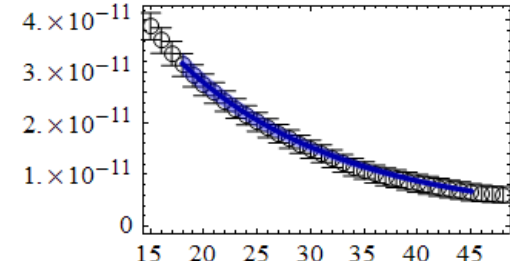
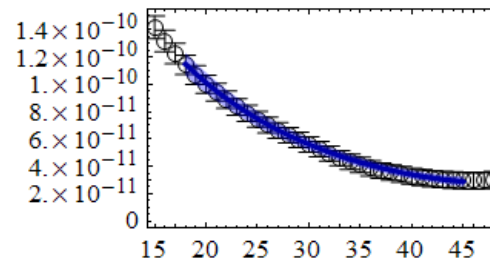
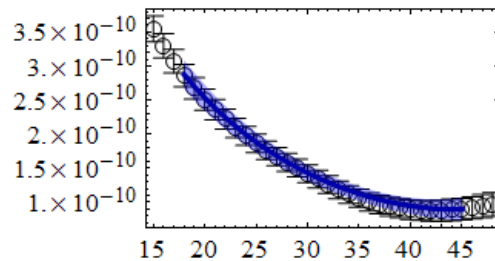
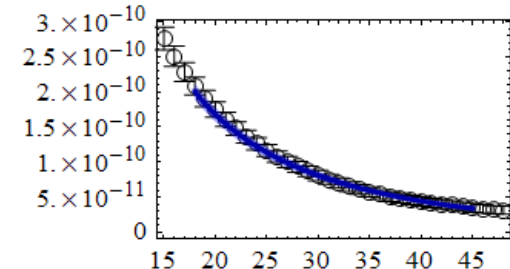
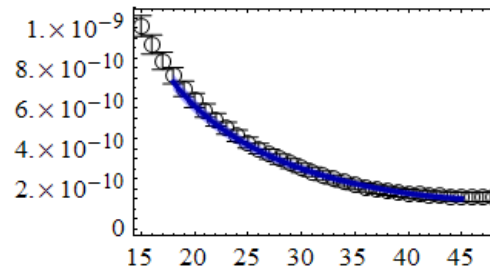
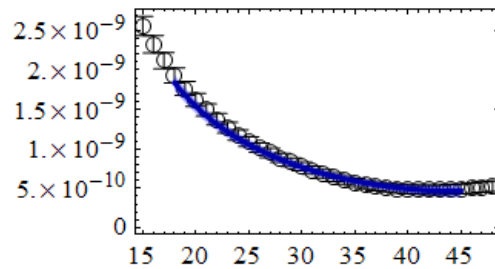
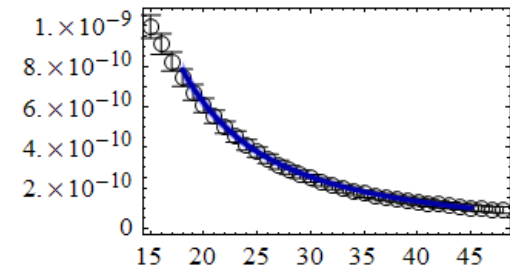
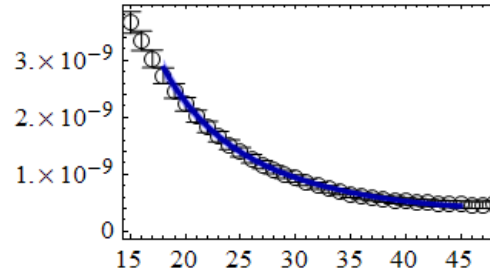
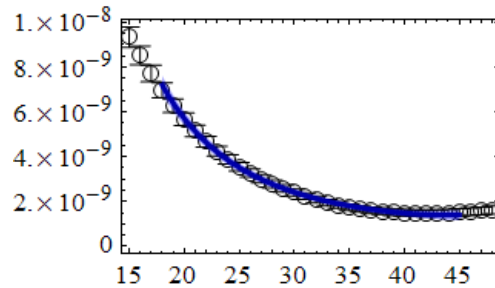
Roper



n^2

Three-Point Fitting

◆ Example: $P_f = \{0,0,0\}$, $P_i = \{0,1,1\}$, V_4



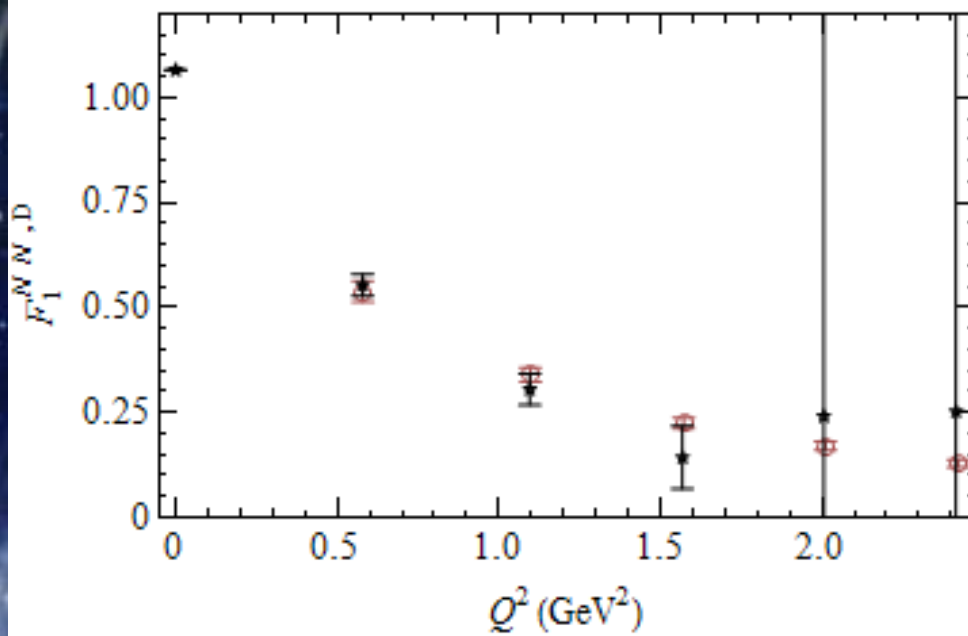
Comparison with Ratio Method

- ◆ Compare calculation with “ratio” approach:

$$R_\mu = \frac{\Gamma_{\mu,GG}^{(3),\mathcal{P}}(t_i, t, t_f, \vec{p}_i, \vec{p}_f)}{\Gamma_{GG}^{(2)}(t_i, t_f, \vec{p}_f)} \times \left(\frac{\Gamma_{LG}^{(2)}(t, t_f, \vec{p}_i) \Gamma_{GG}^{(2)}(t_i, t, \vec{p}_f) \Gamma_{LG}^{(2)}(t_i, t_f, \vec{p}_f)}{\Gamma_{LG}^{(2)}(t, t_f, \vec{p}_f) \Gamma_{GG}^{(2)}(t_i, t, \vec{p}_i) \Gamma_{LG}^{(2)}(t_i, t_f, \vec{p}_i)} \right)^{\frac{1}{2}}$$

and SVD solutions

- ◆ Results



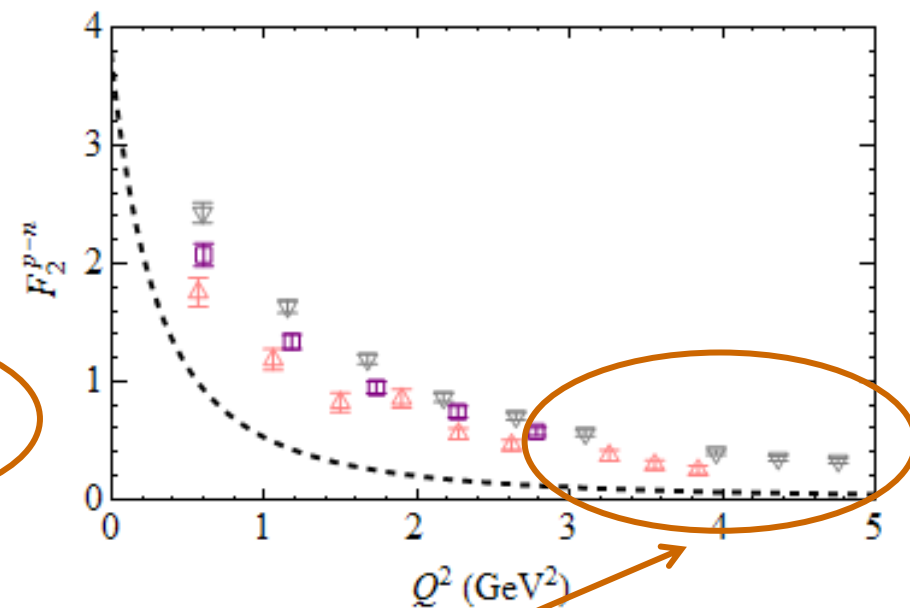
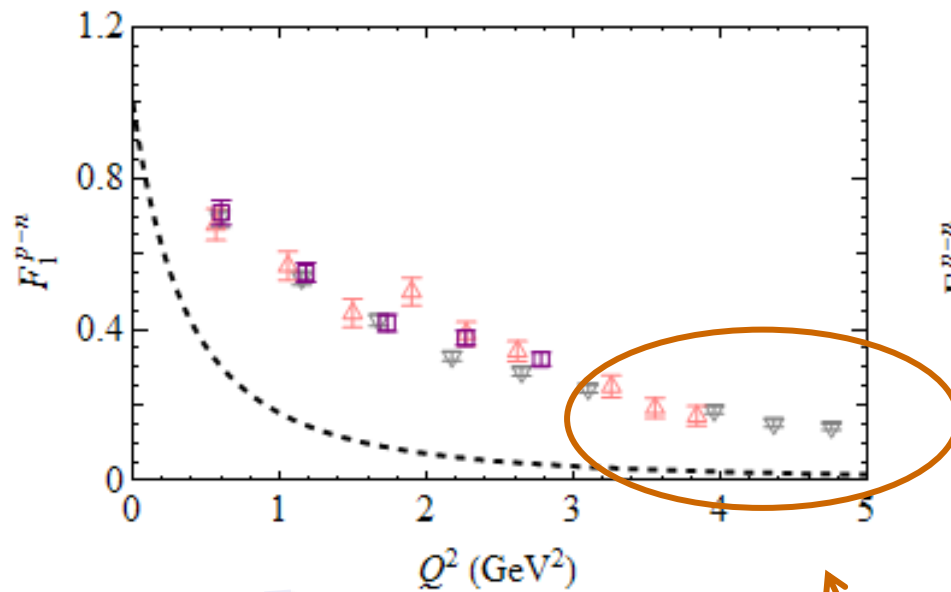
- ◆ Consistent with conventional ratio approach
- ◆ Smaller errorbar at larger Q^2
- ◆ Normally smearings are tuned to eliminate higher-energy contribution

Nucleon Form Factors

- ◆ Pion masses around 480, 720 and 1100 MeV

Isovector F_1

Isovector F_2



Preliminary

Clean signal for momentum region of $3 < Q^2 < 5 \text{ GeV}^2$

- ◆ Working on larger- Q^2 regions with non-zero p_f
- ◆ $N_f = 3$ and 2+1 data are on the way

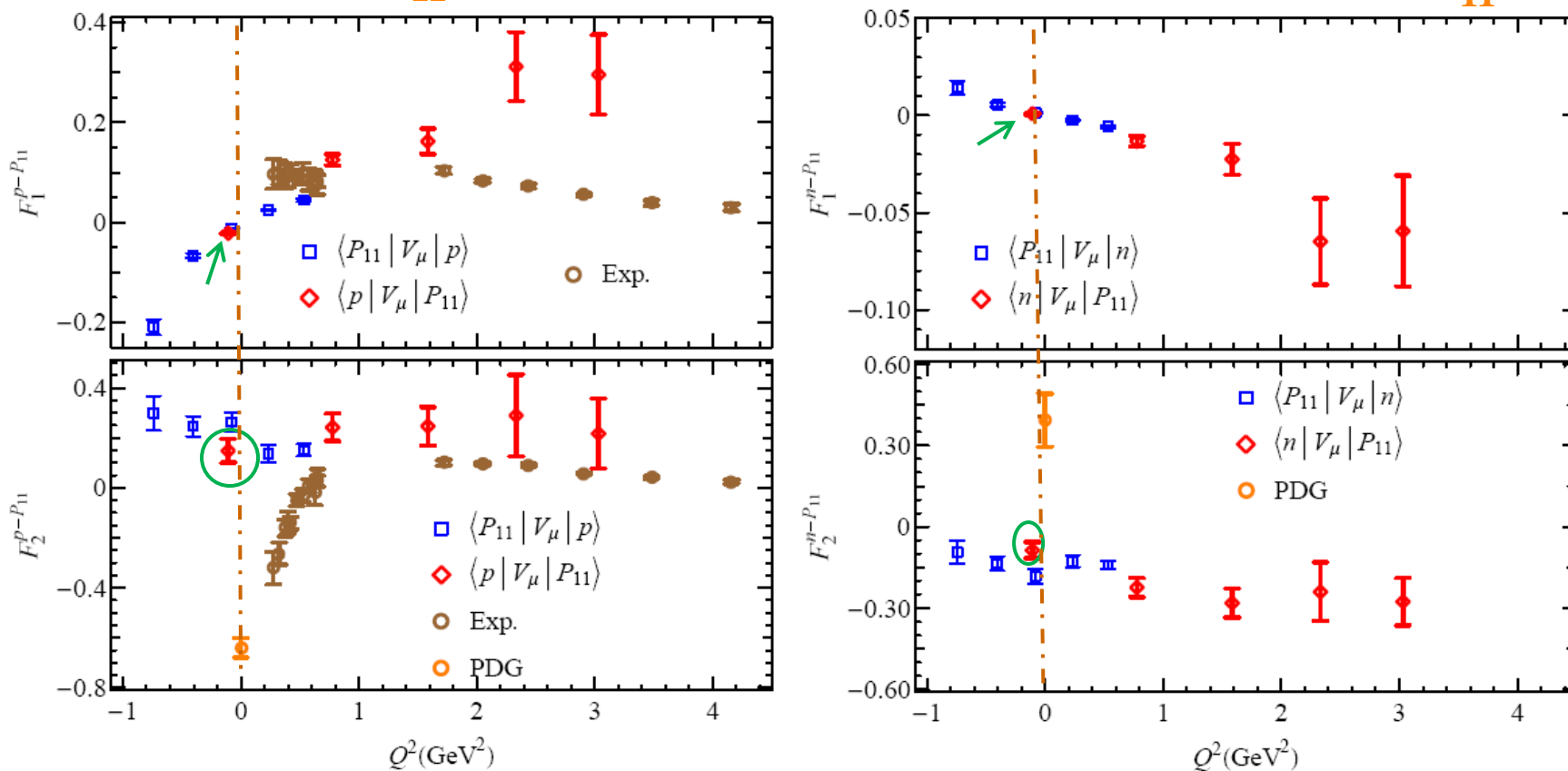
Nucleon-Roper Form Factors

- Completed exploratory study on quenched lattices [arXiv:0803.3020](https://arxiv.org/abs/0803.3020)

Proton- P_{11}

720 MeV Pion

Neutron- P_{11}



- Possible decaying state (circled above)
- 200 configurations give us reasonable signal
- Lower pion mass will shift the time-like region to space-like region

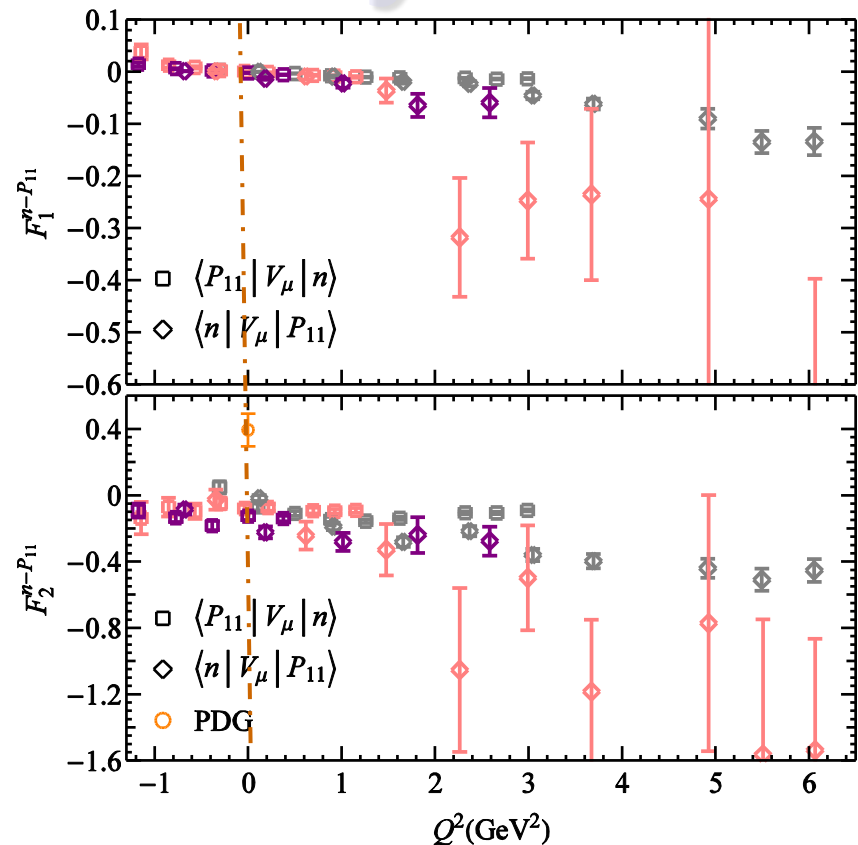
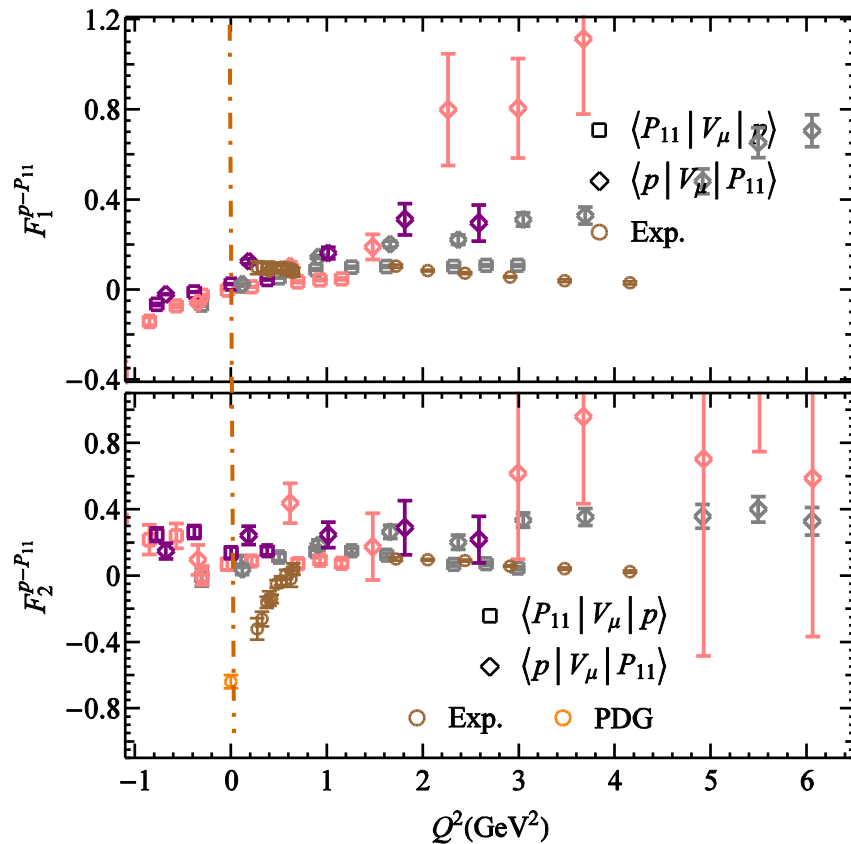
Nucleon-Roper Form Factors

- ◆ Add two more mass points at $m_\pi \sim 480$ and 1100 MeV

Proton- P_{11}

Neutron- P_{11}

Preliminary



- ◆ Need to remove the points with potential decay kinematics
- ◆ May want to calculate on a second volume

Summary and Outlook

Solving the form-factor challenges in Lattice QCD calculation...

- ◆ Large- Q^2 momentum N - N form factors; no sign of negative G_E for 5.5 GeV²
- ◆ We demonstrate a method to determine N - N^* form factors
- ◆ Test case is in a small “quenched” box with large pion mass

Further along our roadmap...

- ◆ Starting full-QCD anisotropic lattice calculations this summer
- ◆ Search over low and larger- Q^2 regions
- ◆ Implement group theory operators for baryons
- ◆ Other N - N^* form factors. The methodology developed can be applied to many other excited-nucleon form factors.