

Lambda-nucleon force from lattice QCD

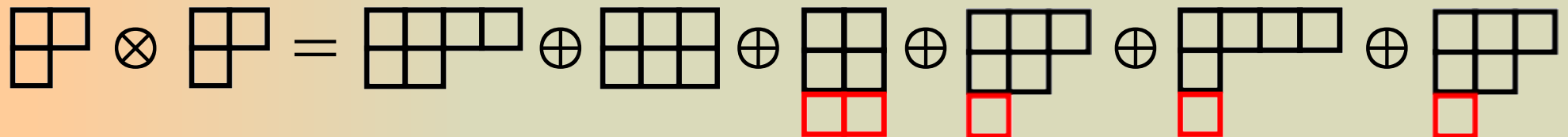
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for PACS-CS Collaboration

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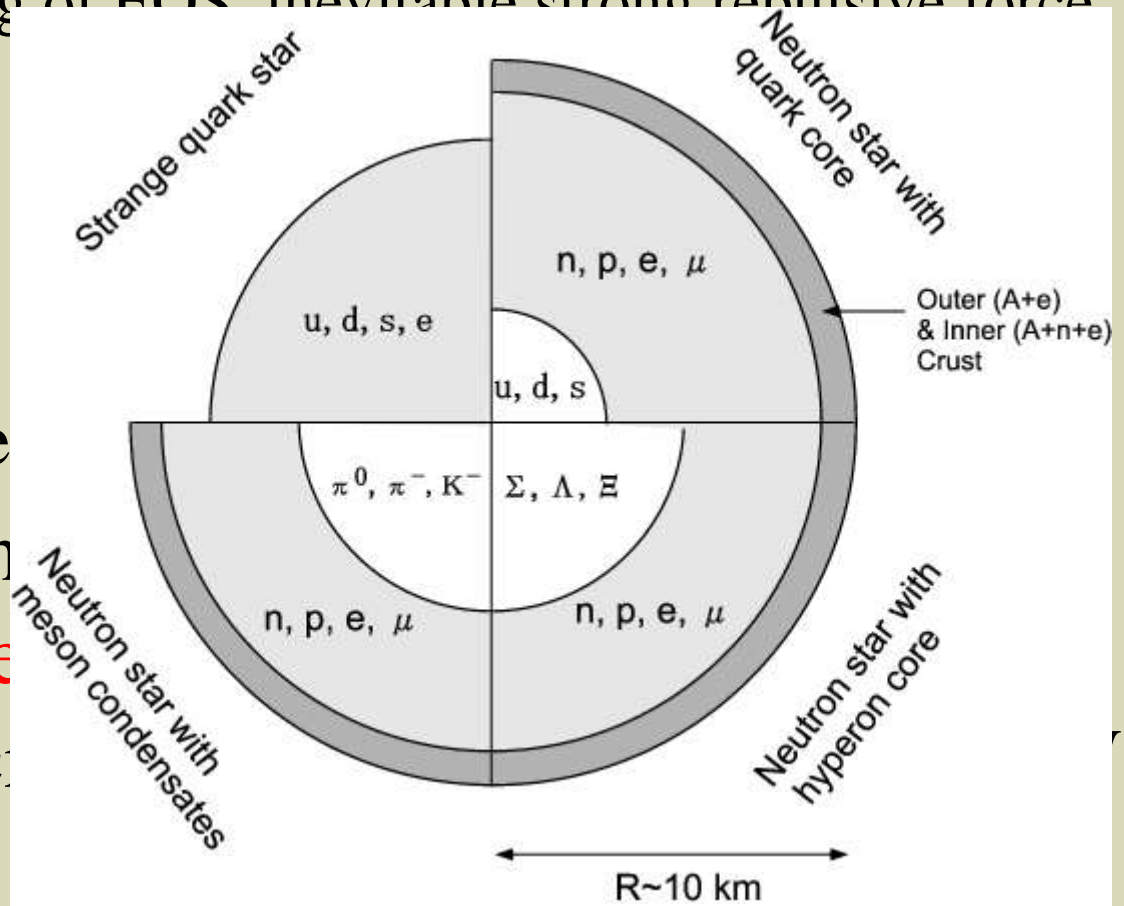


Introduction:

- ⊗ Study of **hyperon-nucleon (YN)** and **hyperon-hyperon (YY)** interactions is one of the important subjects in the nuclear physics.
 - ⊗ Structure of the neutron-star core,
 - ⊗ Hyperon mixing, softning of EOS, inevitable strong repulsive force,
 - ⊗ H-dibaryon problem,
 - ⊗ To be, or not to be,
- ⊗ The project at J-PARC:
 - ⊗ Explore the multistrange world,
- ⊗ However, the phenomenological description of YN and YY interactions has **large uncertainties**, which is in sharp contrast to the nice description of phenomenological NN potential.

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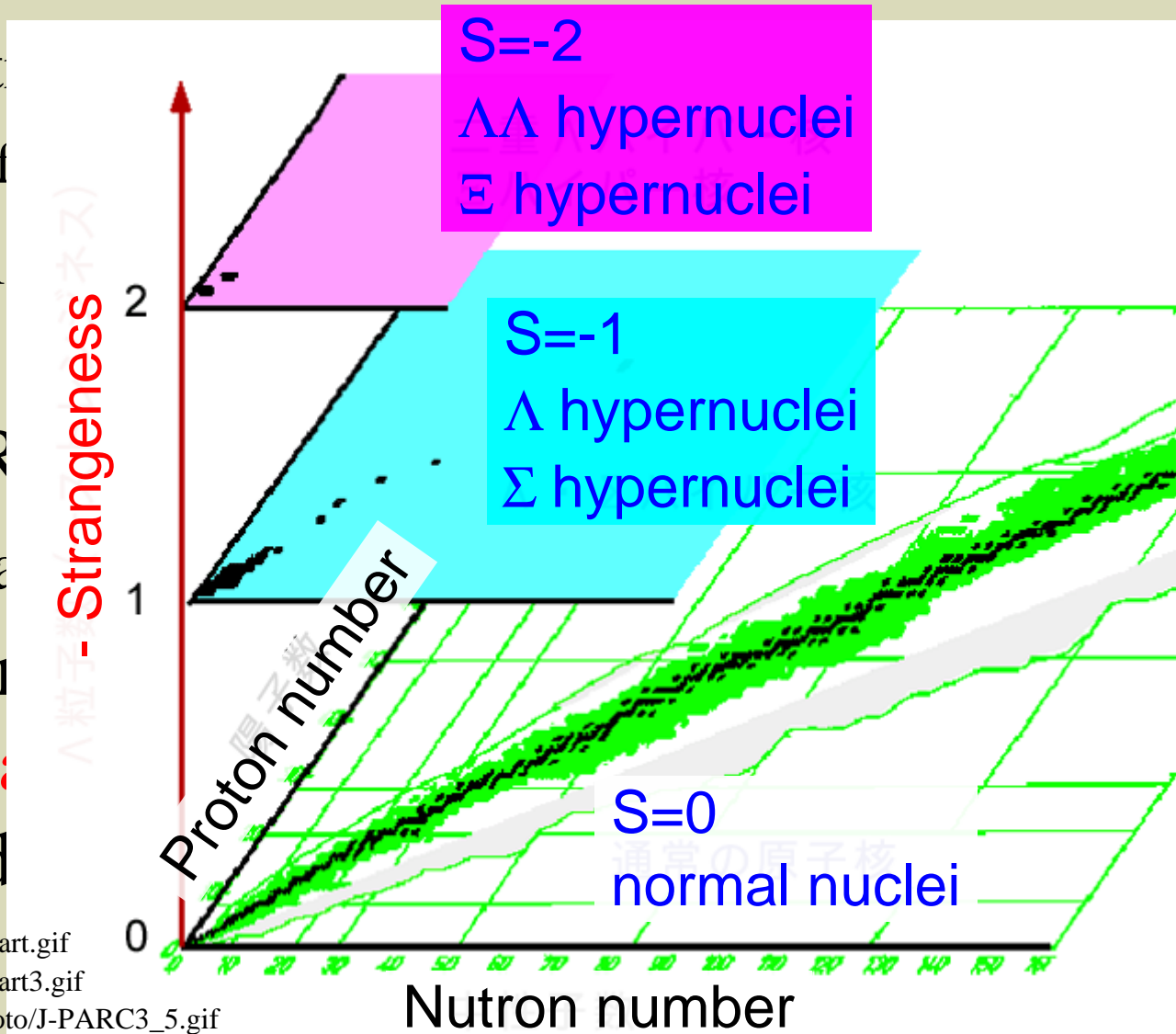
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- However, the phenomenon of **YY** interactions has long been a puzzle in contrast to the nice description of **YN** interactions.



http://www.rarf.riken.go.jp/RIBF/nuclear_chart.gif

<http://www.war.kph.uni-mainz.de/Hyp2006/chart3.gif>

http://www.kek.jp/newskek/2005/marapr/photo/J-PARC3_5.gif

Extension from NN to YN and YY:

- ⊗ If we take only non-strange sector, there are only 2 representations for isospin space.

$$\begin{array}{ccccccc}
 2 & & 2 & 3 & & 1 & \\
 \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \otimes & \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & = & \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \end{array} & \oplus & \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} \\
 I=\frac{1}{2} & & I=\frac{1}{2} & & I=1 & & I=0
 \end{array}$$

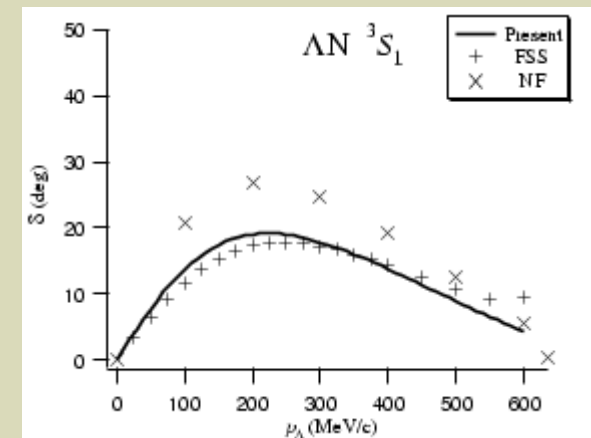
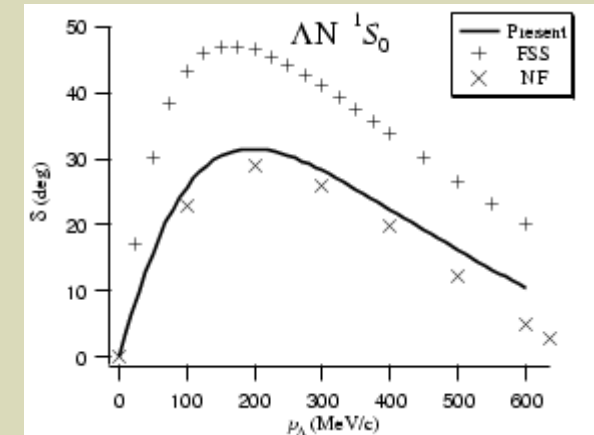
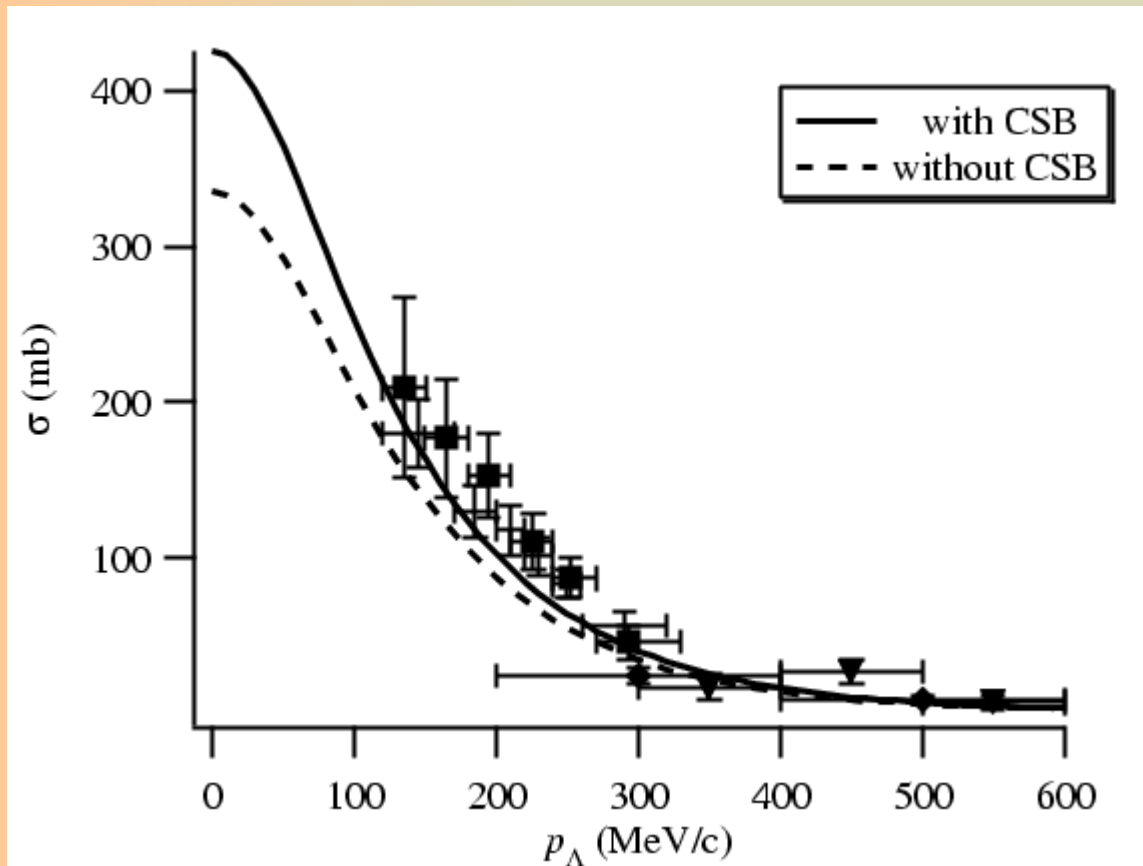
- ⊗ On the other hand, if we take account of strange degree of freedom, other representations should be included.

$$\begin{array}{cccccccccccc}
 8 & & 8 & 27 & 10^* & 1 & 8 & & 10 & 8 & \\
 \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \otimes & \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & = & \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \end{array} & \oplus & \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \oplus & \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \oplus & \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \oplus & \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \square & & \square \\ \hline \end{array} & \oplus & \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \square & & \square \\ \hline \end{array}
 \end{array}$$

- ⊗ This means that the YN and YY interactions cannot be determined from the precise NN experimental data even if we assume the flavor SU(3) symmetry.
- ⊗ **Lattice QCD** is desirable for the study of the YN and YY interaction, because this is *ab initio* numerical simulation.

Experimental data for ΛN interaction:

- ⊗ Only total cross section.
- ⊗ No phase shift analysis is available.
- ⊗ Spin-dependence is unclear



Recent impressive works of lattice QCD:

⊗ S. Aoki, *et al.*, PRD71, 094504 (2005);

π - π scattering length from the wave function.

⊗ N. Ishii, *et al.*, PRL99, 022001 (2007); nucl-th/0611096;

NN potential from the wave function.

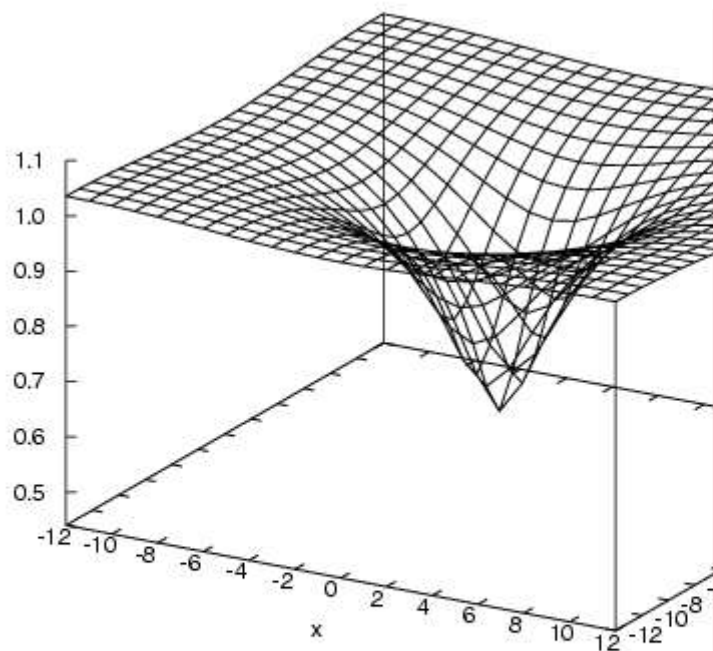
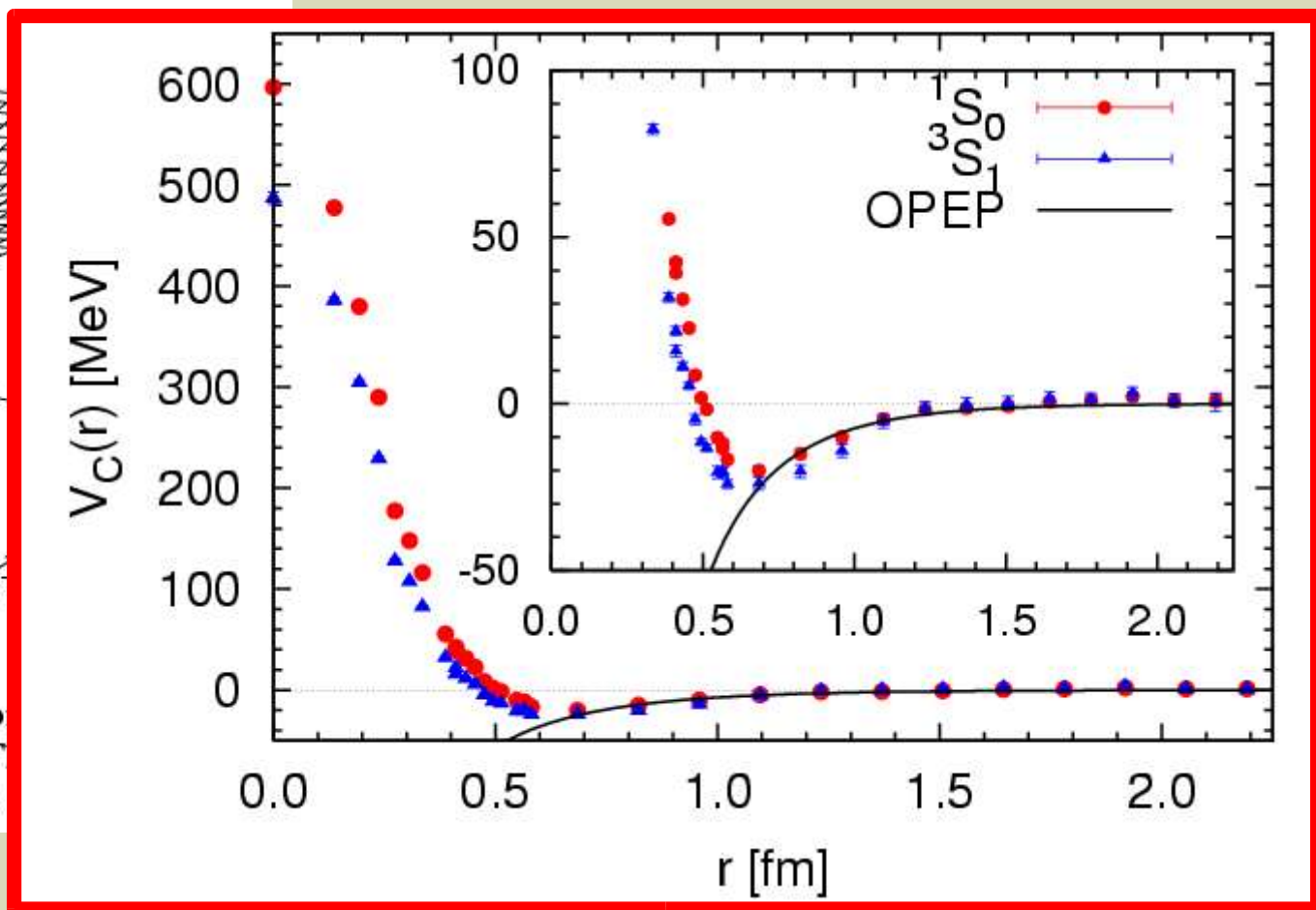


FIG. 1. Two-pion wave function $\phi(\vec{x}; k)$ on $(t, z) = (52, 0)$ plane for $m_\pi^2 = 0.273 \text{ GeV}^2$. Torus is set at $\vec{x}_0 = (7, 5, 2)$ ($x_0 = |\vec{x}_0| = 8.832$).

⊗ This work;

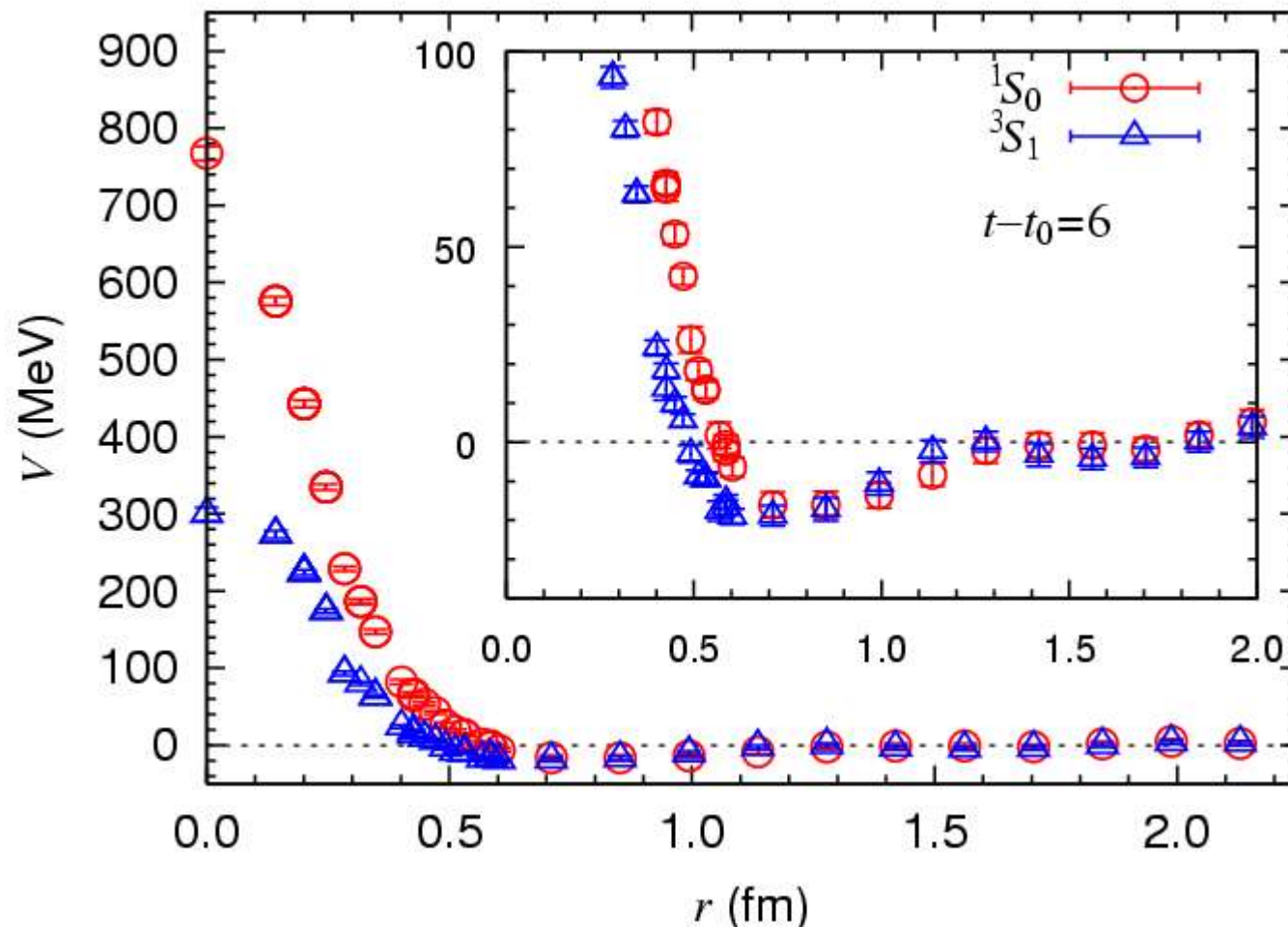
YN and YY potentials by applying these techniques.



The first $N\Xi$ calculation from lattice QCD

arXiv:0806.1094 [nucl-th]

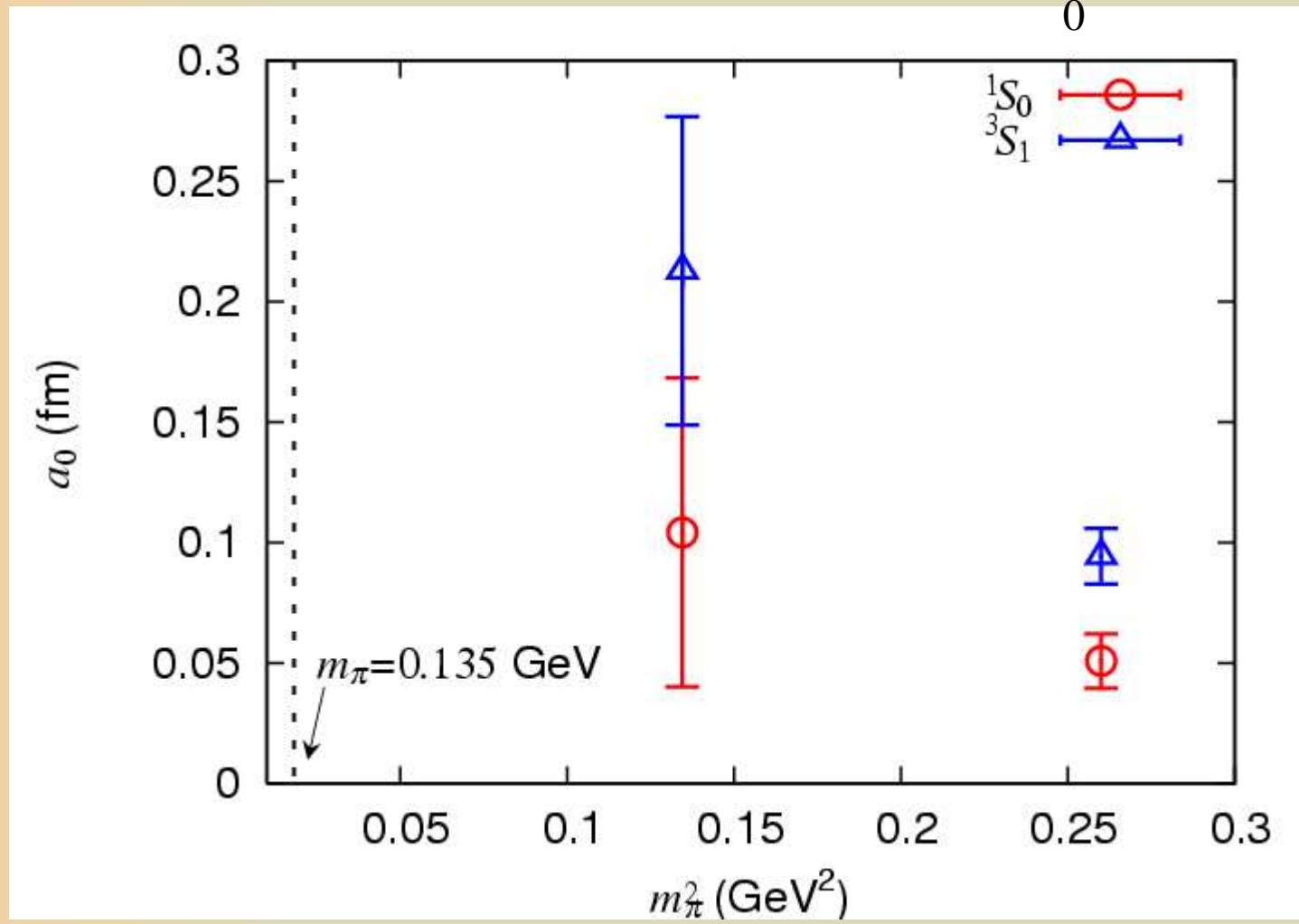
- ⊗ Quenched QCD calculation
- ⊗ Volume: $32^3 \times 32$
- ⊗ Lattice scale: $a=0.14\text{fm}$ ($L \sim 4.5\text{ fm}$)
- ⊗ Isospin $I=1$ channel: attractive in both 1S_0 and 3S_1 .



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The purpose of this work

- ⊗ $N \Lambda$ force from lattice QCD
- ⊗ Spin dependence (Scattering lengths)
- ⊗ Potential (explore the flavor dependence of baryon potentials)
- ⊗ Numerical calculation is twofold:
 - ⊗ Full lattice QCD by using $N_F=2+1$ PACS-CS full QCD gauge configurations with the spatial lattice volume $(2.86 \text{ fm})^3$
 - ⊗ Quenched lattice QCD with larger spatial lattice volume $(4.5 \text{ fm})^3$

A recipe for $N\Lambda$ potential:

⊗ More accurate explanation, see, e.g., [arXiv:0805.2462\[hep-ph\]](https://arxiv.org/abs/0805.2462).

⊗ Start from an effective Schroedinger eq for the equal-time Bethe-Salpeter wave function:

$$-\frac{1}{2\mu} \nabla^2 \phi(\vec{r}) + \int d^3 r' U(\vec{r}, \vec{r}') = E \phi(\vec{r})$$

$$U(\vec{r}, \vec{r}') = V_{N\Lambda}(\vec{r}, \nabla) \delta(\vec{r} - \vec{r}')$$

⊗ A general expression of the potential:

$$\begin{aligned} V_{N\Lambda} = & V_0(r) + V_\sigma(r) (\vec{\sigma}_N \cdot \vec{\sigma}_\Lambda) \\ & + V_T(r) S_{12} + V_{LS}(r) (\vec{L} \cdot \vec{S}_+) \\ & + V_{ALS}(r) (\vec{L} \cdot \vec{S}_-) + O(\nabla^2) \end{aligned}$$

A recipe for $N\Lambda$ potential:

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⊗ The equal time BS wave function in the S -wave on the lattice,

$$\phi(\vec{r}) = \frac{1}{24} \sum_{R \in O} \frac{1}{L^3} \sum_{\vec{x}} P_{\alpha\beta}^{\sigma} \langle 0 | p_{\alpha}(R[\vec{r}] + \vec{x}) \Lambda_{\beta}(\vec{x}) | p \Lambda; k \rangle$$

$$p_{\alpha}(x) = \varepsilon_{abc} (u_a(x) C \gamma_5 d_b(x)) u_{c\alpha}(x),$$

$$\Lambda_{\alpha}(x) = \varepsilon_{abc} \left\{ (d_a C \gamma_5 s_b) u_{c\alpha} + (s_a C \gamma_5 u_b) d_{c\alpha} - 2(u_a C \gamma_5 d_b) s_{c\alpha} \right\}$$

⊗ The **4-point $N\Lambda$ correlator** on the lattice,

$$\begin{aligned} F_{p\Lambda}(\vec{x}, \vec{y}, t; t_0) &= \langle 0 | p_{\alpha}(\vec{x}, t) \Lambda_{\beta}(\vec{y}, t) \overline{J}_{p\Lambda}(t_0) | 0 \rangle \\ &= \sum_n A_n \langle 0 | p_{\alpha}(\vec{x}) \Lambda_{\beta}(\vec{y}) | n \rangle e^{-E_n(t-t_0)} \end{aligned}$$

$$\overline{J}_{p\Lambda}(t_0)$$

wall source at $t=t_0$

A recipe for $N\Lambda$ potential:

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⊗ Calculate the **4-point $N\Lambda$ correlator** on the lattice,

$$\phi_{N\Lambda}(\mathbf{x}-\mathbf{y}) e^{-E(t-t_0)} \propto \langle p_\alpha(\mathbf{x}, t) \Lambda_\beta(\mathbf{y}, t) \overline{\Lambda}_{\beta'}(\mathbf{0}, t_0) \overline{p}_{\alpha'}(\mathbf{0}, t_0) \rangle$$

⊗ Which has the physical meanings of,

⊗ Create a $N\Lambda$ state and making imaginary time evolution, in order to have the lowest state of the $N\Lambda$ system.

⊗ Take the **amplitude $\phi(\mathbf{x}-\mathbf{y})$** , which can be understood as a wave function of the non-relativistic quantum mechanics.

⊗ Obtain the **effective central potential** from the **effective Schroedinger equation**.

$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right) \phi(r) = E \phi(r)$$



$$V(r) = E + \frac{\hbar^2}{2\mu} \frac{\nabla^2 \phi(r)}{\phi(r)}$$

Full QCD calculations by using $N_F=2+1$

PACS-CS gauge configurations:

- ⊗ S. Aoki, et al., (PACS-CS Collaboration),
arXiv:0807.1661 [hep-lat].
- ⊗ Iwasaki gauge action at $\beta=1.90$ on $32^3 \times 64$ lattice
- ⊗ $O(a)$ improved Wilson quark action
- ⊗ $1/a = 2.17 \text{ GeV}$ ($a = 0.0907 \text{ fm}$)
- ⊗ $\kappa_{ud} = 0.13770$, $\kappa_s = 0.13640$
- ⊗ $m_\pi = 0.30 \text{ GeV}$, $m_K = 0.59 \text{ GeV}$,
 $m_\rho = 0.84 \text{ GeV}$, $m_{K^*} = 0.97 \text{ GeV}$
 $m_p = 1.1 \text{ GeV}$, $m_\Lambda = 1.2 \text{ GeV}$, $m_\Sigma = 1.3 \text{ GeV}$



Quenched calculation with larger spatial volume:

- ⊗ Plaquette gauge action and Wilson fermion action
- ⊗ Gauge coupling $\beta=5.7$
- ⊗ Volume: $32^3 \times 48$ ($L \sim 4.5$ fm).
- ⊗ Lattice spacing: $a \sim 0.14$ fm. ($1/a \sim 1.4$ GeV.)
- ⊗ The lattice calculations were performed by using KEK Blue Gene/L supercomputer.
- ⊗ The main results are obtained with

⊗ $\kappa_{ud} = 0.1665$ for the u and d quarks, and

⊗ $\kappa_s = 0.1643$ for s quark.



Meson masses:

$$m_{\pi} \sim 0.511.2(6) \text{ GeV}$$

$$m_{\rho} \sim 0.861(2) \text{ GeV}$$

$$m_K \sim 0.605.3(5) \text{ GeV}$$

$$m_{K^*} \sim 0.904(2) \text{ GeV}$$



Results

$N_f = 2 + 1$

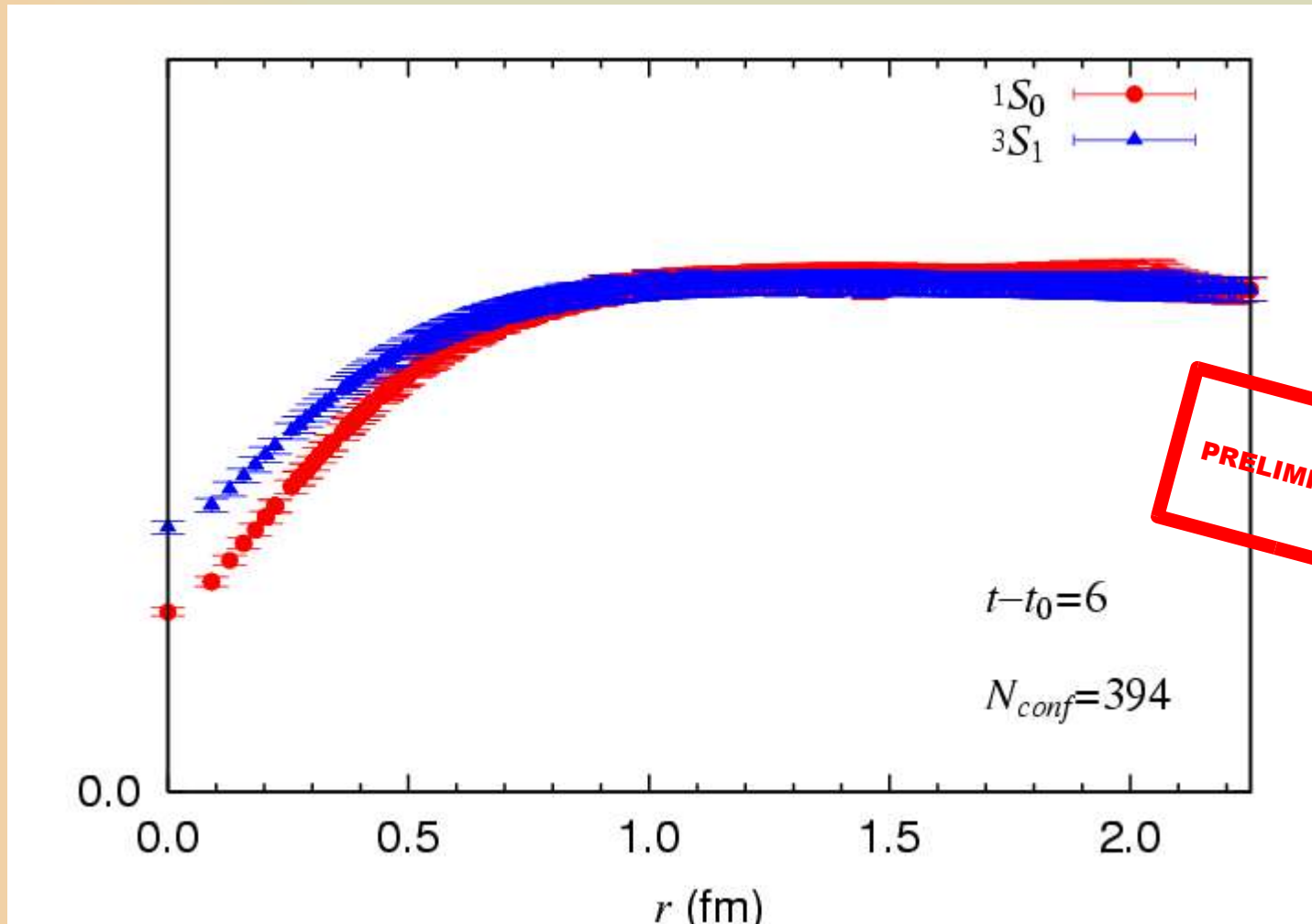
Full QCD

by using

PACS-CS

Results — wave function

- ⊗ Suggests the **repulsive core** in short range for both spin $S=0$ and 1.

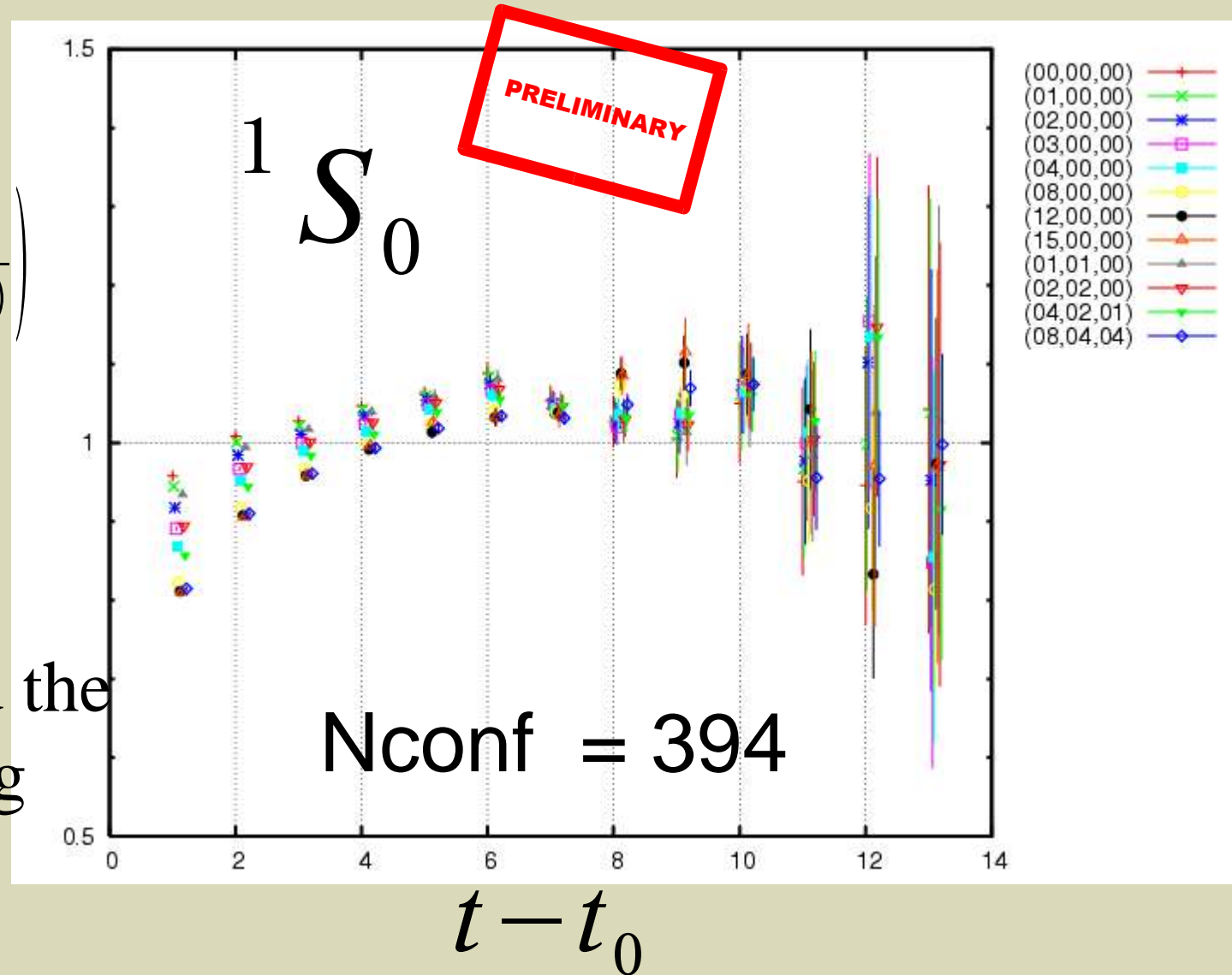


Results — “effective mass”

- ⊗ Time dependence of 4-point correlator to find the ground state (plateaux in the effective mass)

$$m_{eff}(t-t_0, \vec{r} = \vec{x} - \vec{y}) \equiv \log \left(\frac{F_{p\Lambda}(\vec{x}, \vec{y}, t; t_0)}{F_{p\Lambda}(\vec{x}, \vec{y}, t+1; t_0)} \right)$$

- ⊗ Still hardly find the plateaux starting in the 1S_0 channel.

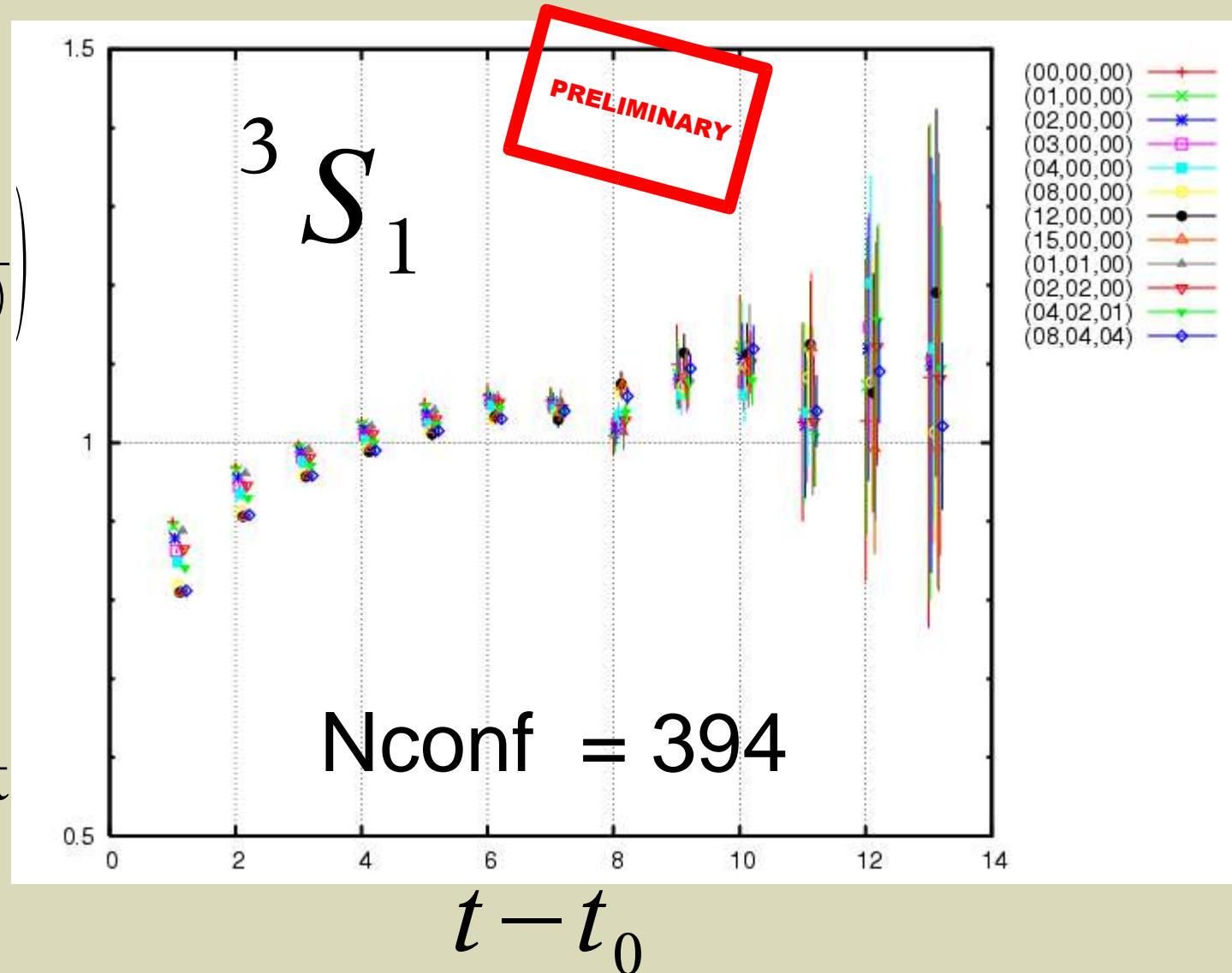


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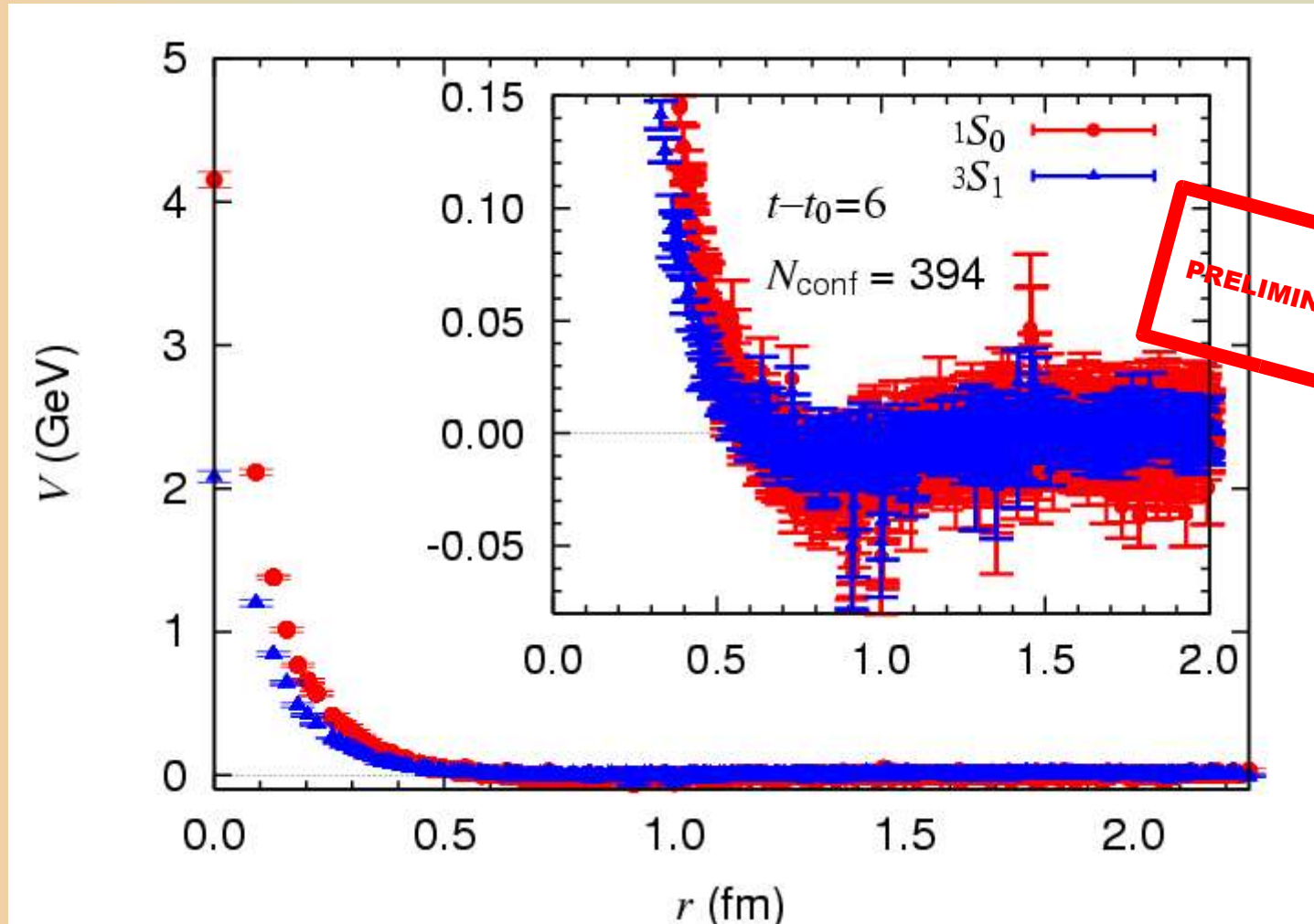
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- ⊗ The plateaux starting would appear at $t - t_0 = 6$.



Results — potential

⊗ $N\Lambda$ potential, from lattice QCD for the first time.



⊗ Strong repulsive core in spin $S=0$ channel.

⊗ Spin dependence.

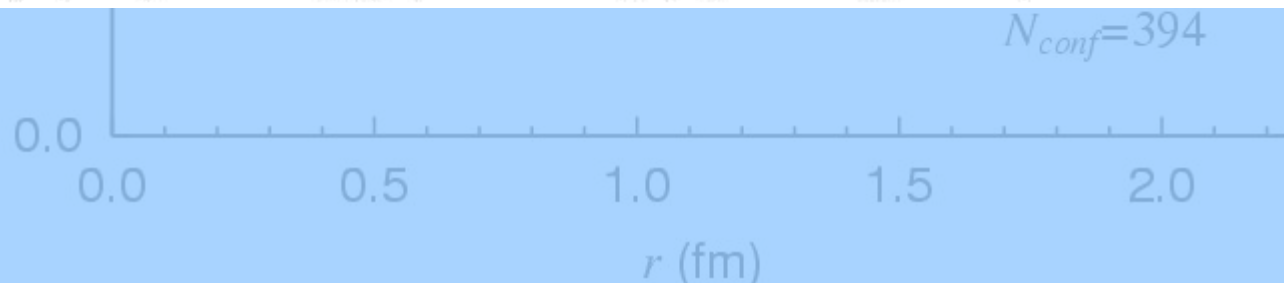
Results — wave function

The non-relativistic energy $E=k^2/(2\mu)$ can be accurately determined by fitting the wave function in the asymptotic region in terms of the lattice Green's function:

$$G(\vec{r}, k^2) = \frac{1}{L^3} \sum_{\vec{p} \in \Gamma} \frac{1}{p^2 - k^2} e^{i\vec{p} \cdot \vec{r}}, \quad \Gamma = \left\{ \vec{p}; \vec{p} = \vec{n} \frac{2\pi}{L}, \vec{n} \in \mathbf{Z}^3 \right\},$$

which is a solution of

$(\Delta + k^2)G(\vec{r}, k^2) = -\delta_L(\vec{r})$ with $\delta_L(\vec{r})$ being the periodic delta function.



Results — wave function

The energies are almost zero. (although slightly negative (**attractive**)).
 Interaction in the 1S_0 seems to be more attractive than that in 3S_1 .

A very preliminary results for fitting the k^2 (in lattice unit)

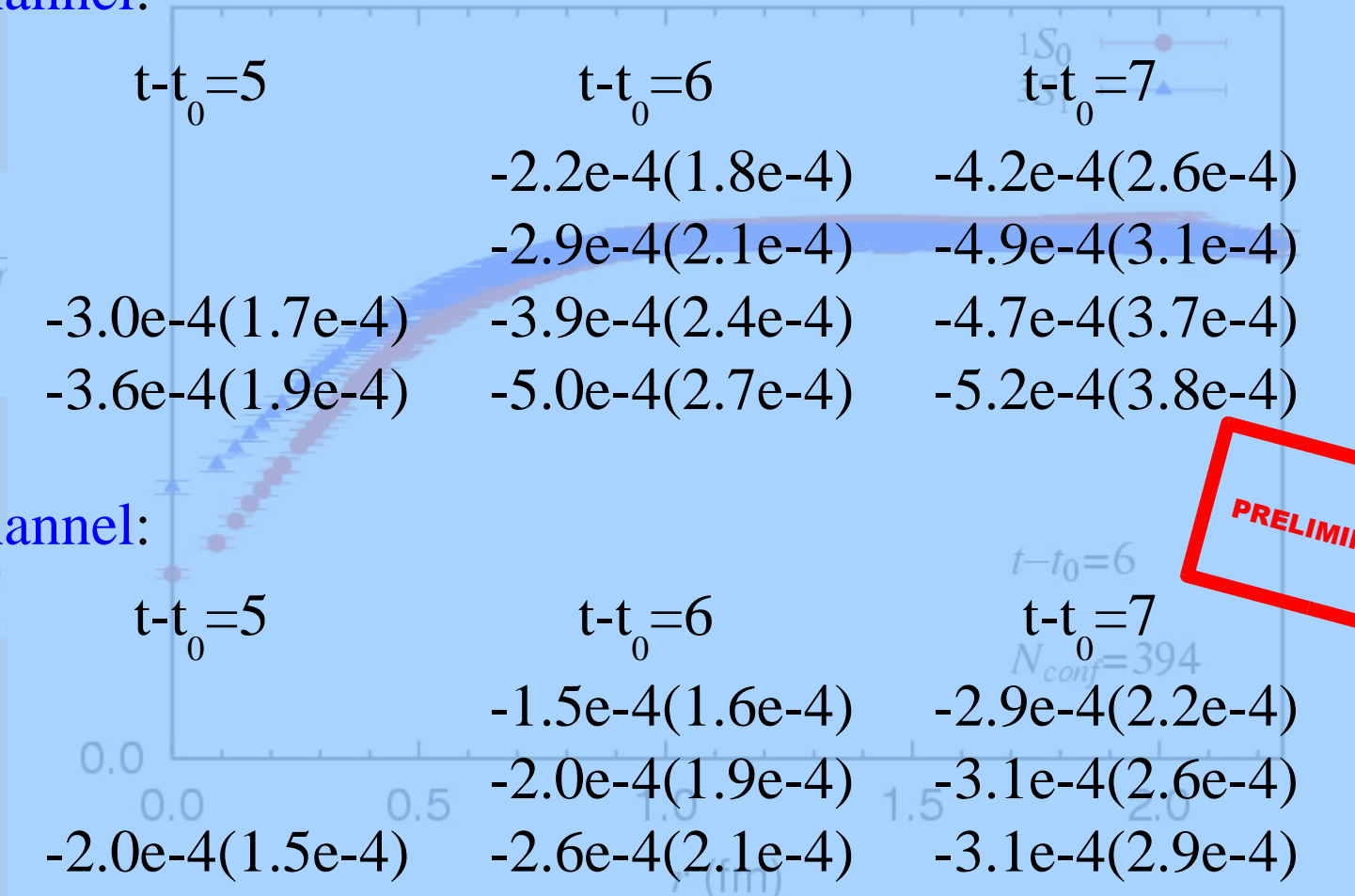
1S_0 channel:

	$t-t_0=5$	$t-t_0=6$	$t-t_0=7$
$r=13$		$-2.2e-4(1.8e-4)$	$-4.2e-4(2.6e-4)$
$r=14$		$-2.9e-4(2.1e-4)$	$-4.9e-4(3.1e-4)$
$r=15$	$-3.0e-4(1.7e-4)$	$-3.9e-4(2.4e-4)$	$-4.7e-4(3.7e-4)$
$r=16$	$-3.6e-4(1.9e-4)$	$-5.0e-4(2.7e-4)$	$-5.2e-4(3.8e-4)$

3S_1 channel:

	$t-t_0=5$	$t-t_0=6$	$t-t_0=7$
$r=13$		$-1.5e-4(1.6e-4)$	$-2.9e-4(2.2e-4)$
$r=14$		$-2.0e-4(1.9e-4)$	$-3.1e-4(2.6e-4)$
$r=15$	$-2.0e-4(1.5e-4)$	$-2.6e-4(2.1e-4)$	$-3.1e-4(2.9e-4)$
$r=16$	$-2.5e-4(1.7e-4)$	$-2.9e-4(2.4e-4)$	$-3.2e-4(3.2e-4)$

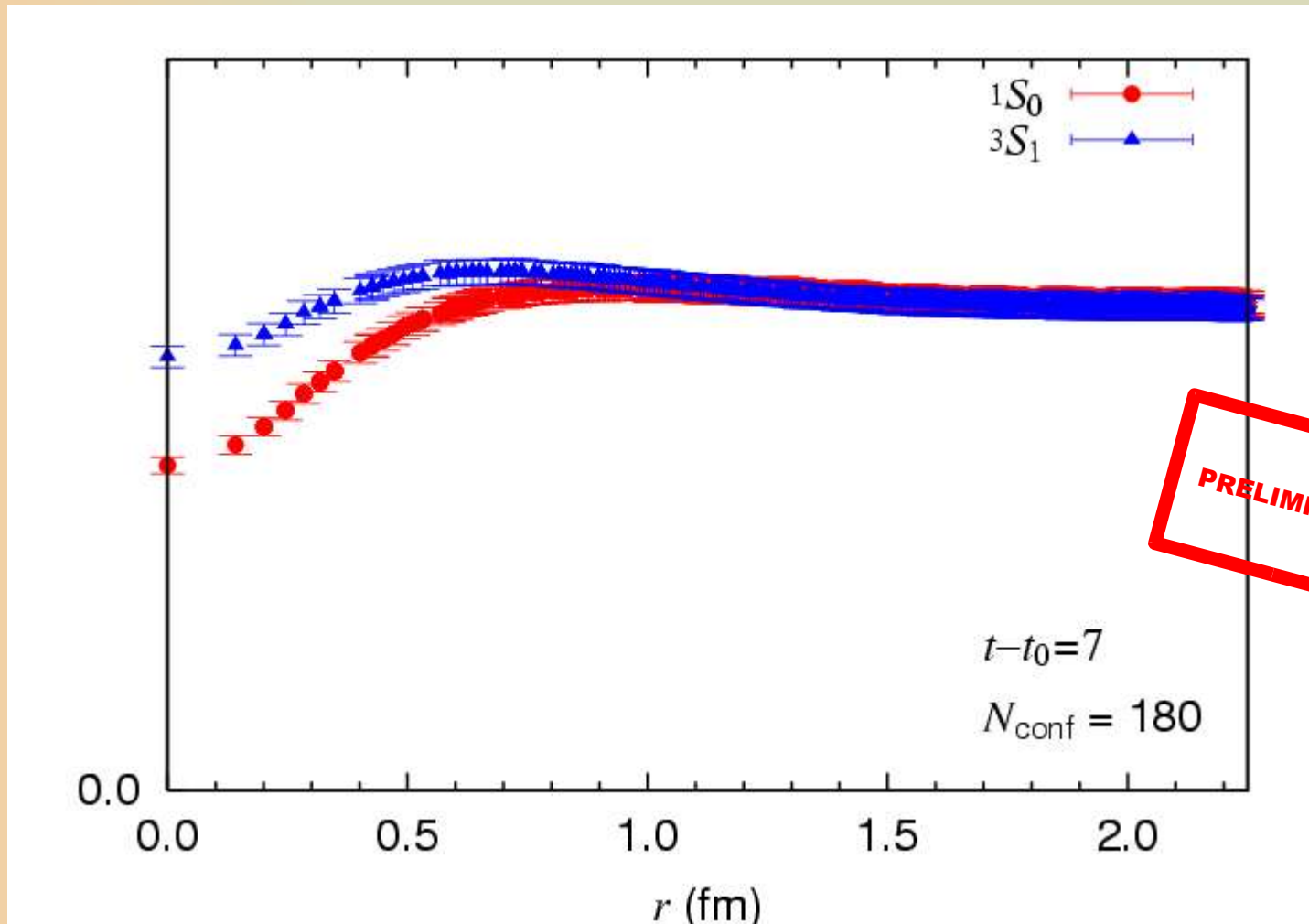
PRELIMINARY



**Quenched
QCD
with larger
spatial volume**

Results — wave function

- ⊗ Suggests the **repulsive core** in short range for both spin $S=0$ and 1.

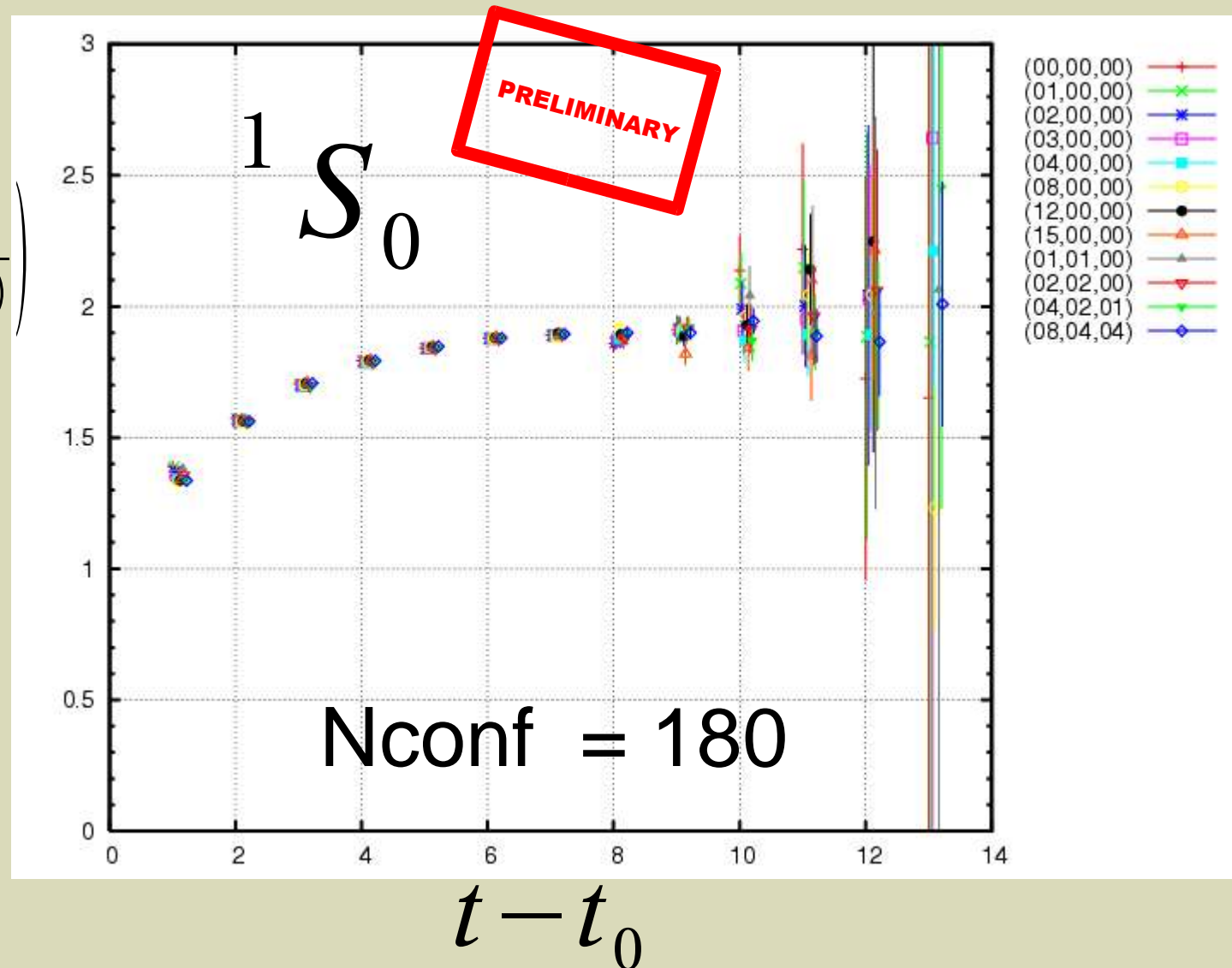


Results — “effective mass”

- ⊗ Time dependence of 4-point correlator to find the ground state (plateaux in the effective mass)

$$m_{eff}(t-t_0, \vec{r} = \vec{x} - \vec{y}) \equiv \log \left(\frac{F_{p\Lambda}(\vec{x}, \vec{y}, t; t_0)}{F_{p\Lambda}(\vec{x}, \vec{y}, t+1; t_0)} \right)$$

- ⊗ The plateaux starting appear at $t - t_0 = 7$.

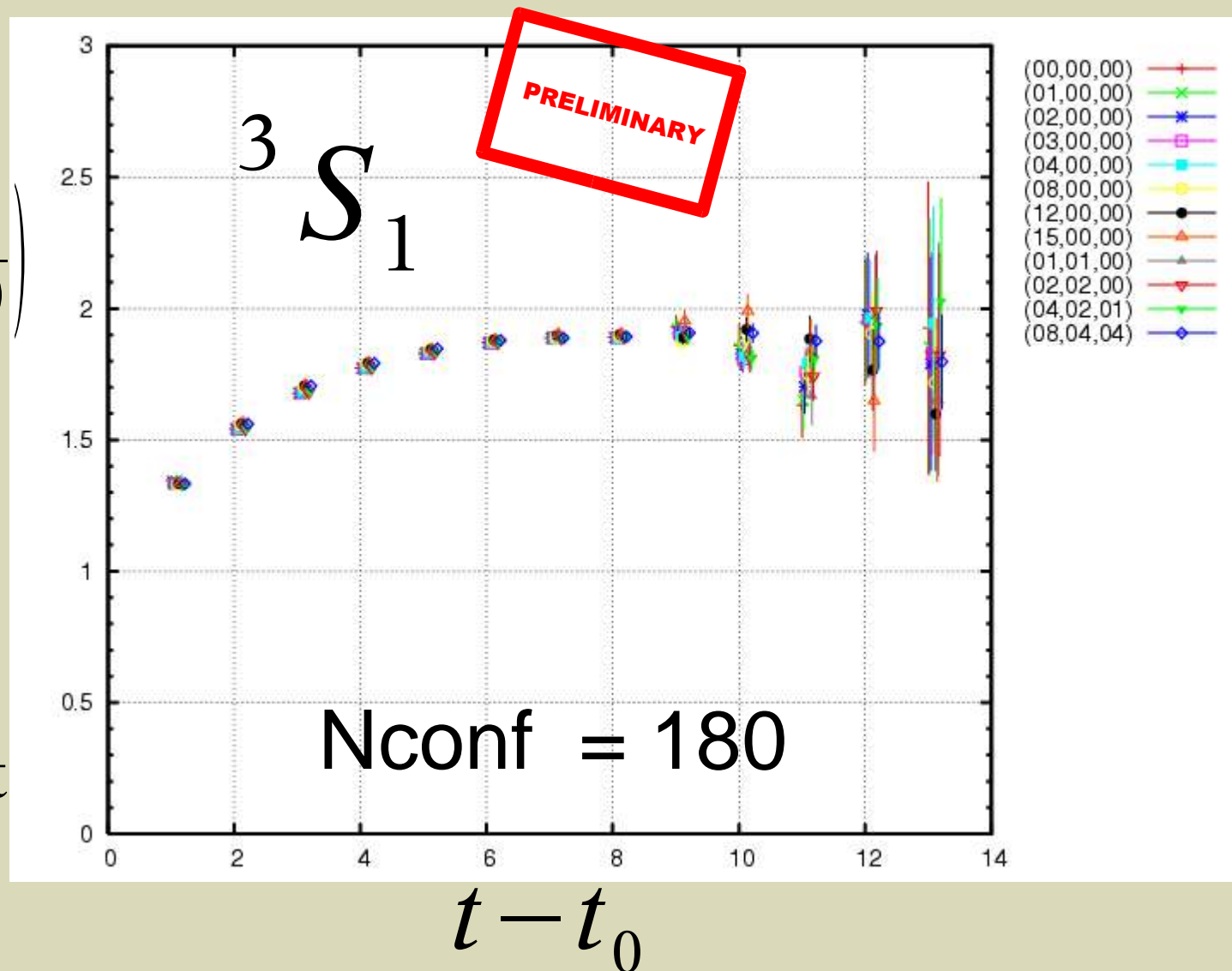


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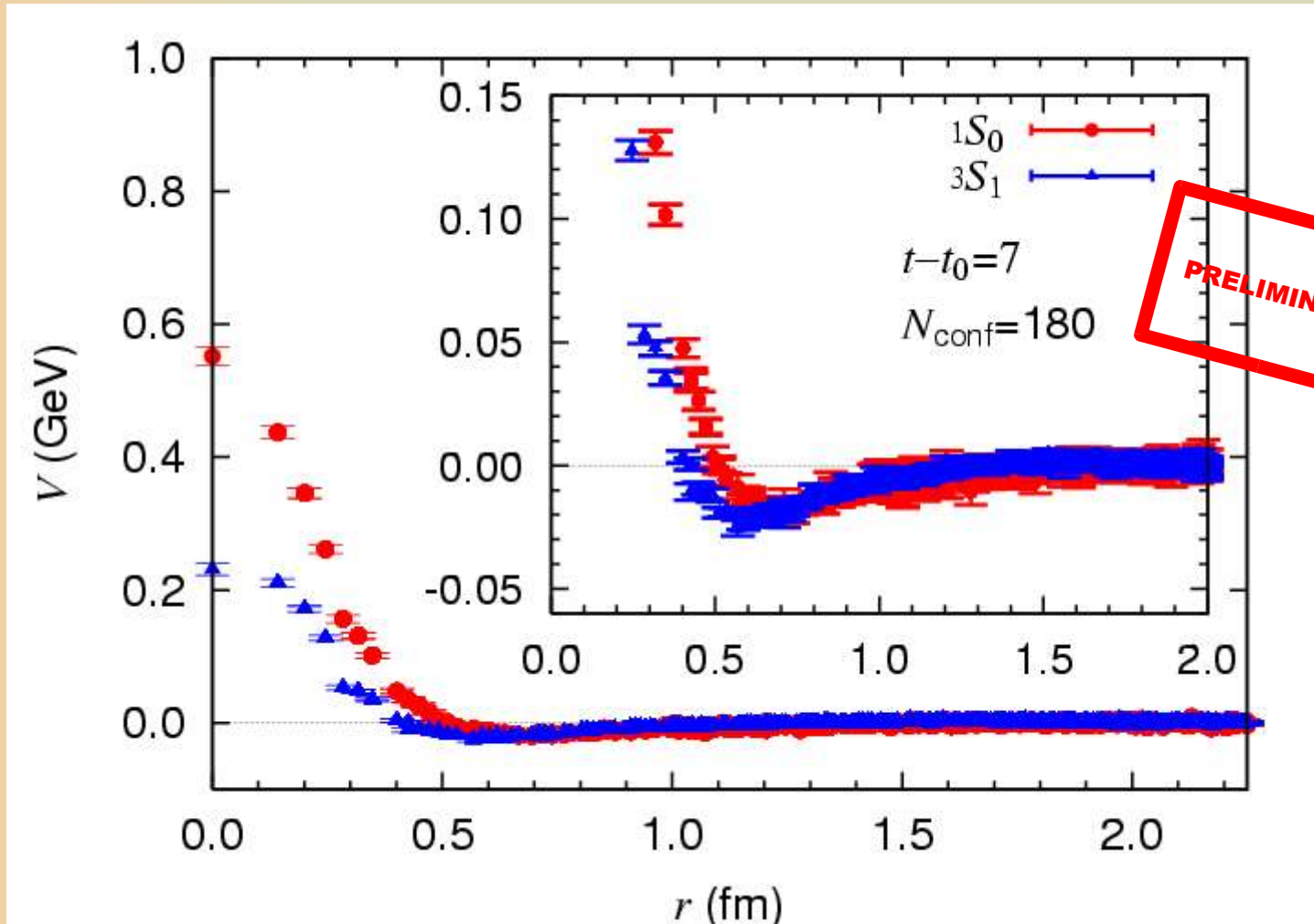
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- ⊗ The plateaux starting would appear at $t - t_0 = 7$.



Results — potential

⊗ $N\Lambda$ potential, from lattice QCD for the first time.



- ⊗ Strong repulsive core in spin $S=0$ channel.
(but relatively weaker than that from the full QCD)
- ⊗ Spin dependence.

Summary:

- ⊗ Study the $N \Lambda$ force by using lattice QCD.
- ⊗ $N_f=2+1$ Full QCD with PACS-CS:
 - ⊗ Strong repulsive core in 1S_0 .
 - ⊗ Spin dependence in short distance region.
 - ⊗ Scattering lengths will be small and attractive ($\sim 0.1\text{fm}$),
 - ⊗ We need to more statistics and to check volume dependence to see the spin-dependence.
- ⊗ Quenched QCD with larger spatial volume:
 - ⊗ Results are qualitatively similar to those from full QCD.
- ⊗ We will study further with:
 - ⊗ Energy dependence (coupled channel with ΣN)
 - ⊗ Tensor force
 - ⊗ Spin-orbit (LS) force, antisymmetric LS force.