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for



Pseudoscalar decay constants  
( $f_K, f_D, f_{D_s}$ )  
from  $N_f = 2$  Twisted Mass QCD

in collaboration with  
Benoit Blossier and Vittorio Lubicz

# Simulation Details

Gauge Action: **tree-level improved Symanzik**

- $N_f = 2$  dynamical fermions, degenerate in mass
- three values of  $\beta$  : 3.8 ( $a \approx 0.10\text{fm}$ ), 3.9 ( $a \approx 0.09\text{fm}$ ), 4.05 ( $a \approx 0.07\text{fm}$ )

Fermionic Action: **twisted mass (tm) at maximal twist**

## Advantages:

• Pseudoscalar meson masses and decay constants are **automatically  $O(a)$ -improved**

• Pseudoscalar **decay constants do not require** renormalization constant:  $f_p = 2\mu \frac{\langle 0|P^1(0)|P \rangle}{M_p \cdot \text{Sinh}(M_p)}$   
bare twisted mass

$\beta$	3.8	3.9	4.05
$v$	$24^3 \times 48$	$24^3 \times 48$	$32^3 \times 64$
$a$	0.10 fm	0.09 fm	0.07 fm
#confs	{240, 240}	{480, 240, 240, 240}	{128, 128, 128}
$a \cdot \mu_l$	{0.08, 0.011}	{0.004, 0.0064, 0.0085, 0.01}	{0.0030, 0.0060, 0.0080}
$a \cdot \mu_s$	{0.02, 0.025, 0.03, 0.036}	{0.022, 0.027, 0.032}	{0.015, 0.018, 0.022, 0.026}
$a \cdot \mu_c$	{0.27, 0.31, 0.355, 0.435, 0.52}	{0.25, 0.32, 0.39, 0.46}	{0.20, 0.23, 0.26, 0.315}

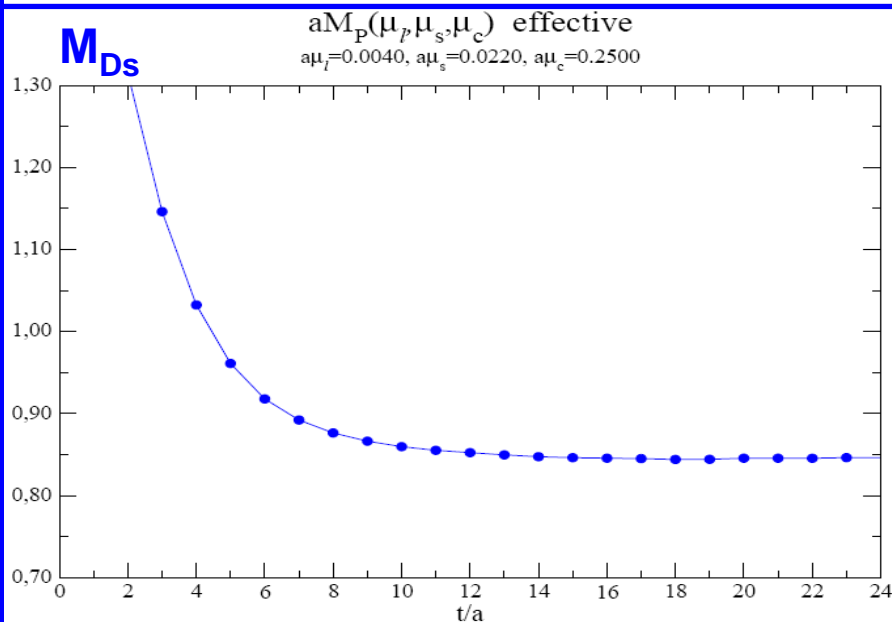
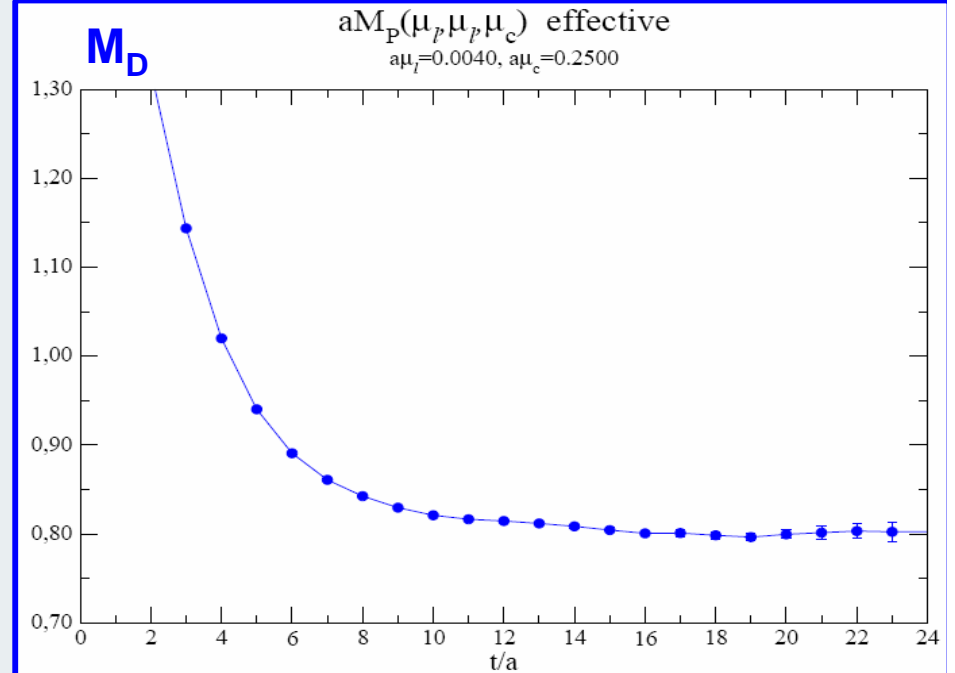
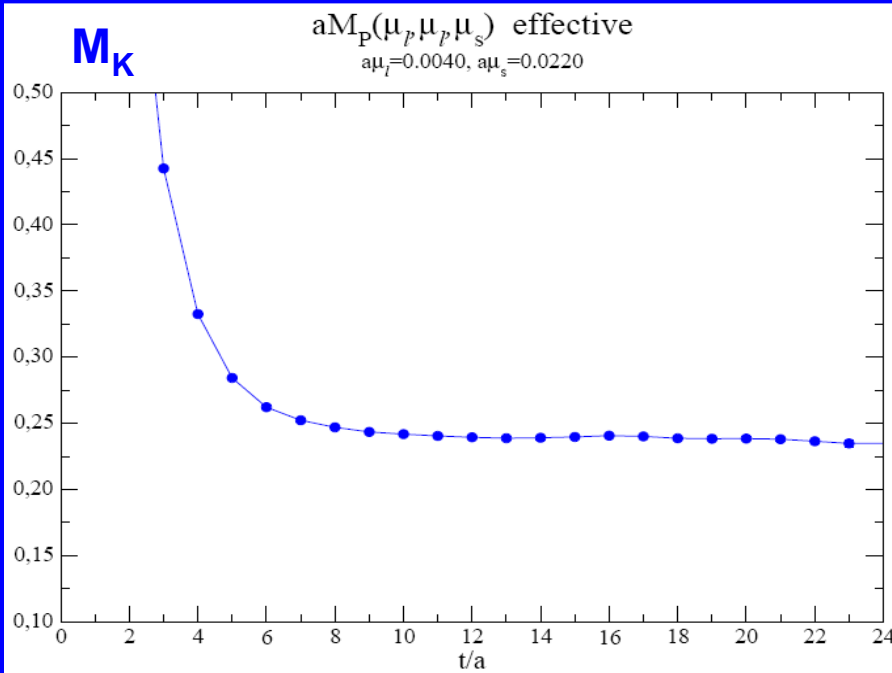
$$\begin{aligned}
 0.2 \cdot m_s^{\text{phys}} &\leq \mu_l \leq 0.5 \cdot m_s^{\text{phys}} \\
 0.9 \cdot m_s^{\text{phys}} &\leq \mu_s \leq 1.5 \cdot m_s^{\text{phys}} \\
 0.8 \cdot m_c^{\text{phys}} &\leq \mu_c \leq 1.5 \cdot m_c^{\text{phys}}
 \end{aligned}$$

$$\begin{aligned}
 &M(\mu_{\text{sea}}, \mu_{\text{val1}}, \mu_{\text{val2}}) \\
 &F(\mu_{\text{sea}}, \mu_{\text{val1}}, \mu_{\text{val2}})
 \end{aligned}$$

$$\begin{aligned}
 f_K &\leftrightarrow F(\overset{\text{sea}}{\mu_l}, \overset{\text{valence}}{\mu_l}, \mu_s) \\
 f_D &\leftrightarrow F(\mu_l, \mu_l, \mu_c) \\
 f_{D_s} &\leftrightarrow F(\mu_l, \mu_s, \mu_c)
 \end{aligned}$$

# Plateaux of effective pseudoscalar masses

e.g.  
 $\beta=3.9$



The **plateau quality** of  
2-point pseudoscalar corr. functions  
is **very good** !

(Statistical accuracy improved using  
a stochastic method to include all spatial sources)  
C.McNeile and C.Michael [UKQCD], hep-lat/0603007



# **KAON SECTOR**

# Progress w.r.t. ETMC'07: 0709.4574

2007	2008
One value of $\beta$ : 3.9	Three values of $\beta$ : 3.8, 3.9, 4.05 → continuum limit
Chiral fit based on SU(3)-ChPT	Chiral fit based on SU(2)-ChPT (SU(3)-ChPT for comparison)
Statistics: 240 confs for each $\mu_l$ , at $\beta=3.9$	Statistics: 480 confs at the lightest $\mu_l$ , at $\beta=3.9$
Study of the dependence of $f_K$ on quark masses	Study of the dependence of $f_K$ on meson masses (with quark masses it is fundamental to have $O(a^2)$ under control in quark mass Renormalization Constants → study in progress)

# Combined analysis

## Fit of data at various $\beta$

At the lowest  $\beta$  (3.8) the tuning of the critical mass is found to be problematic

→  $O(a)$  improvement is not guaranteed at the lightest quark mass

→ **we do not include data at  $\beta=3.8$**  in the analysis

(we consider them only for comparison, finding **perfect compatibility**)

## The fit formulae for $f_K$ include both:

• the meson mass dependence → chiral limit

• an  $O(a^2)$  term → continuum limit



(small but visible)

The dependence of  $f_K$  on meson masses is well described by (NLO) SU(2)-ChPT

**IDEA:** in kaons the **strange quark is not light enough** to apply SU(3)-ChPT  
chiral symmetry properties are satisfied only by the light quark

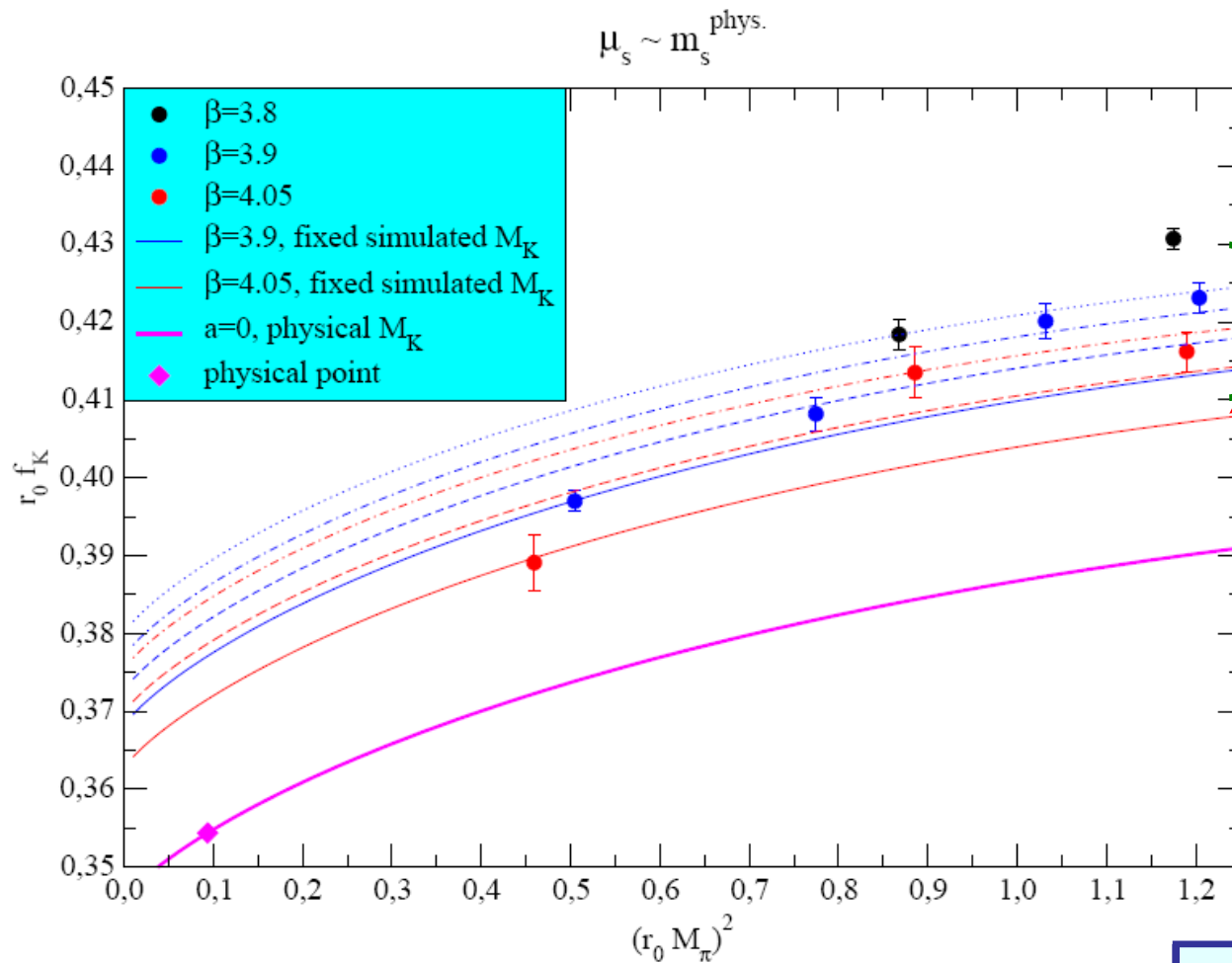
RBC/UKQCD Collaborations, 0804.0473

The SU(2)-ChPT formulae can be obtained, for the decay constant, from SU(3)-ChPT by expanding in  $\mu_l/\mu_s$

A fit based on NLO SU(3)-ChPT + quadratic terms is considered for comparison, finding **perfect compatibility**



# Chiral Behaviour: dependence of $f_K$ on $m_\pi^2$



**Finite size effects**  
(from 1-loop ChPT)  
[D. Becirevic, G. Villadoro,  
hep-lat/0506004]  
turn out to be **negligible**  
in the **kaon region**

- Data with  $\mu_s \sim m_s^{\text{phys.}}$
- Curves at simulated  $\beta$  and  $M_K$
- Curve with  $a=0$ , physical  $M_K$

The interpolation to the physical  $M_K$  is well described by a **linear** dependence in  $M_K^2$

# Results

still preliminary systematic uncertainty!

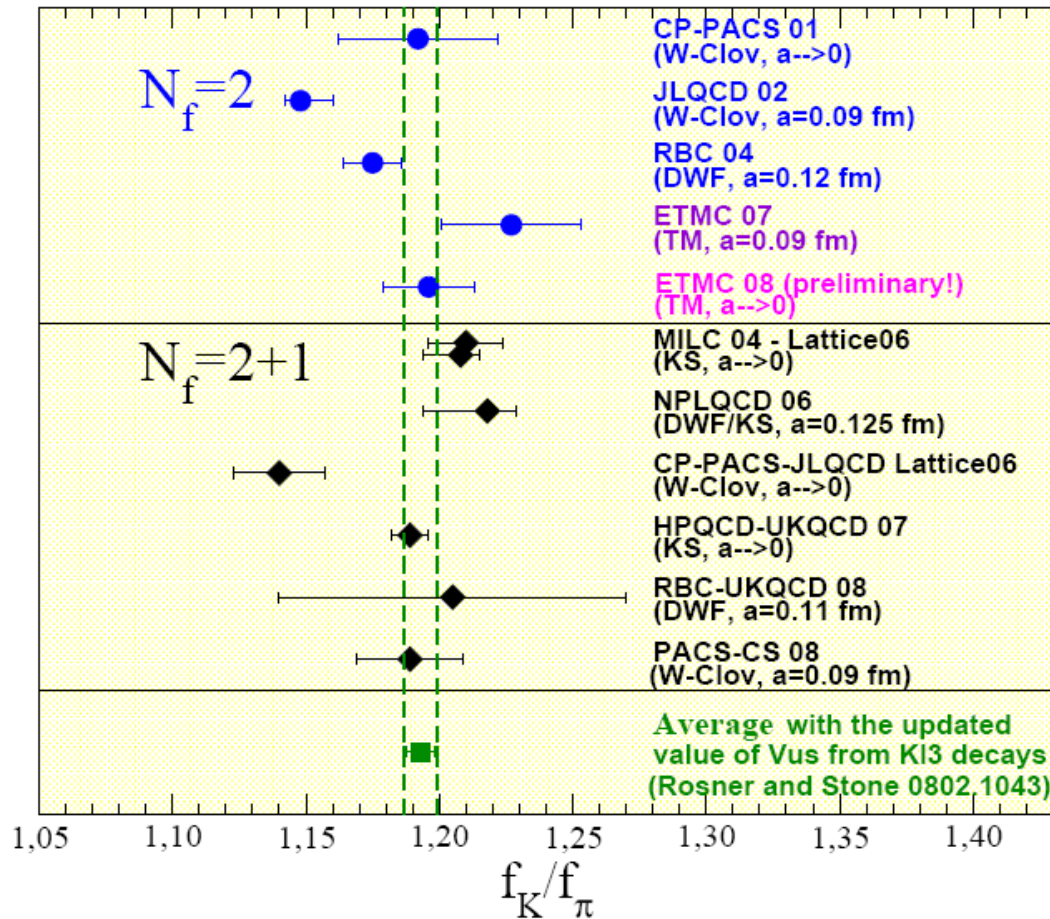
$$f_K = 156.3 (1.7)(0.9)(1.0)(?) \text{ MeV}$$

$$f_K/f_\pi = 1.196 (13) (7) (8) (?)$$

statistics

including  $\beta=3.8$

SU(3)-ChPT



## ETMC

2007 → 2008

$$f_K: 161.7(1.2)(3.1) \rightarrow 156.3(1.7)(?) \text{ MeV}$$

$$f_K/f_\pi: 1.227(9)(24) \rightarrow 1.196(13)(?)$$

Small decrease due to the continuum limit



**D-SECTOR**

# Method for $f_D, f_{D_s}$ : **combined analysis** (similarly to the analysis of $f_K$ )

**The fit formulae for  $f_D, f_{D_s}$  include both:**

- the meson mass dependence → **chiral limit**
- an  $O(a^2)$  and an  $O(a^2 \cdot M_{D_s})$  term → **continuum limit**  
(expected and found to be large in this case)  
(data at  $\beta=3.8$  considered only for comparison  
→ **good compatibility**)

S. Sharpe and Y. Zhang, hep-lat/9510037

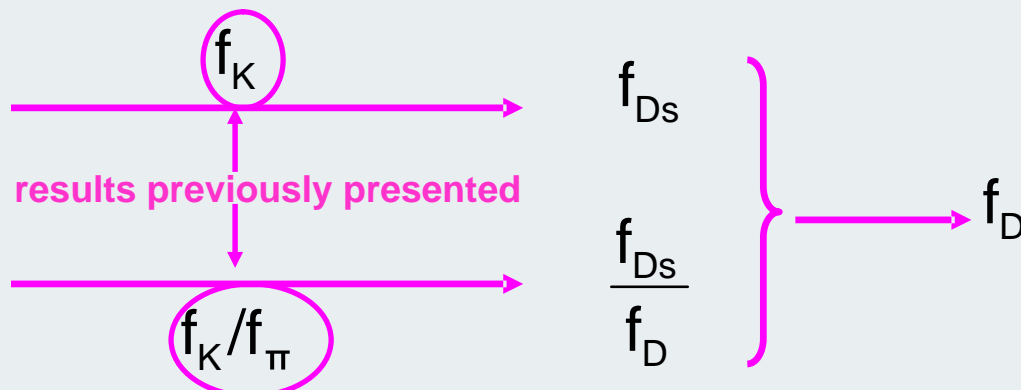
**Chiral Behaviour based on SU(2)-HMChPT**  
(i.e. HMChPT formulae expanded in  $\mu_l/\mu_s$ )

A fit based on HMChPT  
is considered for comparison,  
finding **perfect compatibility**

**We introduce and fit two suitable ratios,**  
**where the light meson dependence is softened**

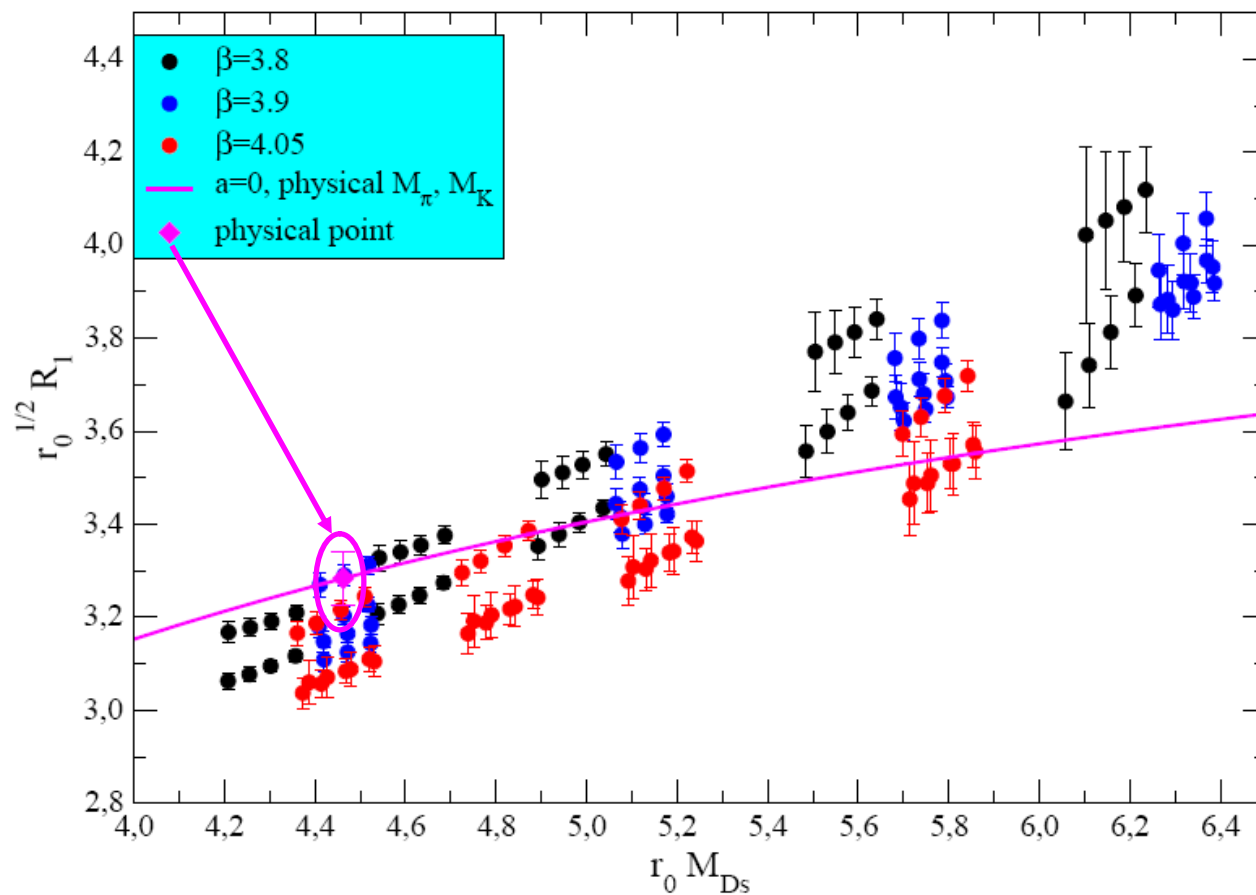
$$R_1 \equiv \frac{f_{D_s} \sqrt{m_{D_s}}}{f_K}$$

$$R_2 \equiv \frac{f_{D_s} \sqrt{m_{D_s}}}{f_K} \cdot \frac{f_\pi}{f_D \sqrt{m_D}}$$



# Dependence on the heavy meson mass: $R_1$ vs $m_{D_s}$

$$R_1 \equiv \frac{f_{D_s} \sqrt{m_{D_s}}}{f_K}$$

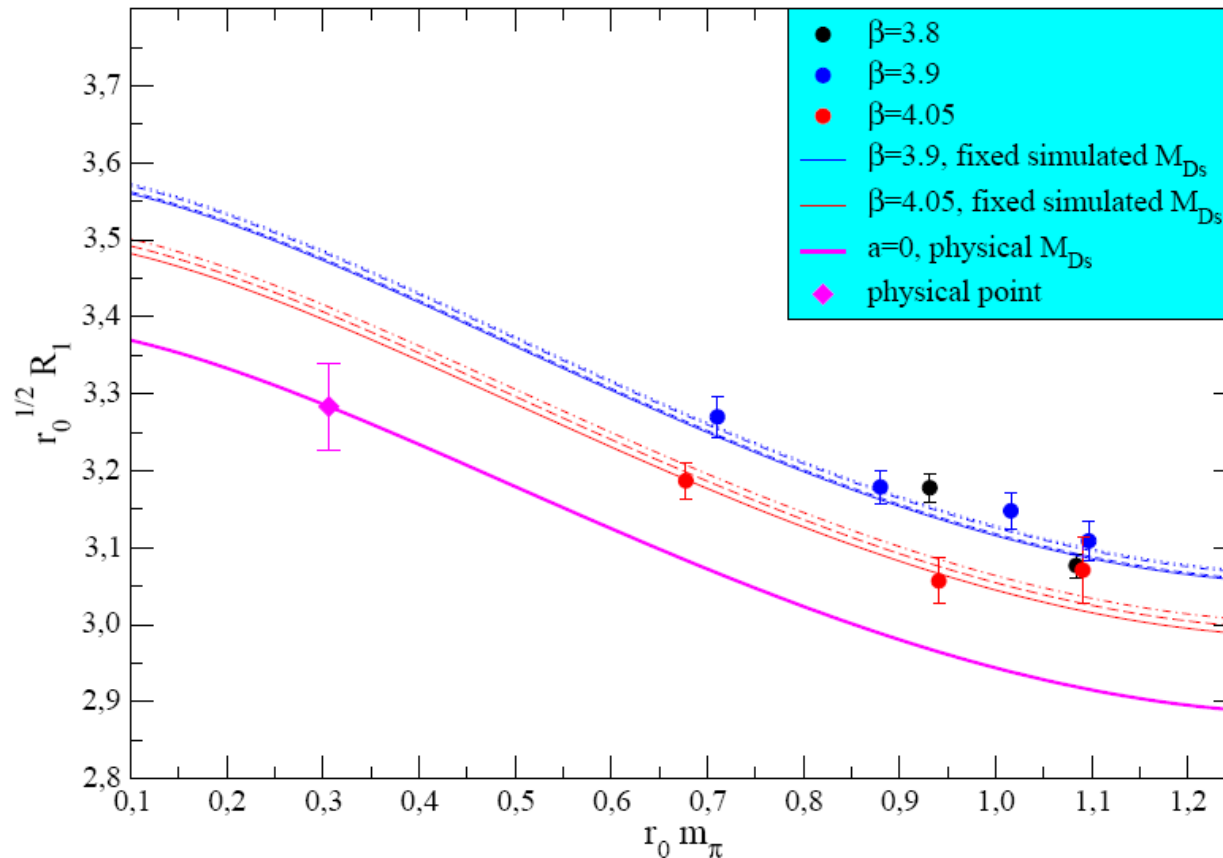


- All data
- Curve with  $a=0$  and physical  $M_\pi, M_K$

# Chiral Behaviour: dependence of $R_1$ on $m_\pi$

$$R_1 \equiv \frac{f_{D_s} \sqrt{m_{D_s}}}{f_K}$$

$$\mu_s \sim m_s^{\text{phys.}}, \quad \mu_c \sim m_c^{\text{phys.}}$$

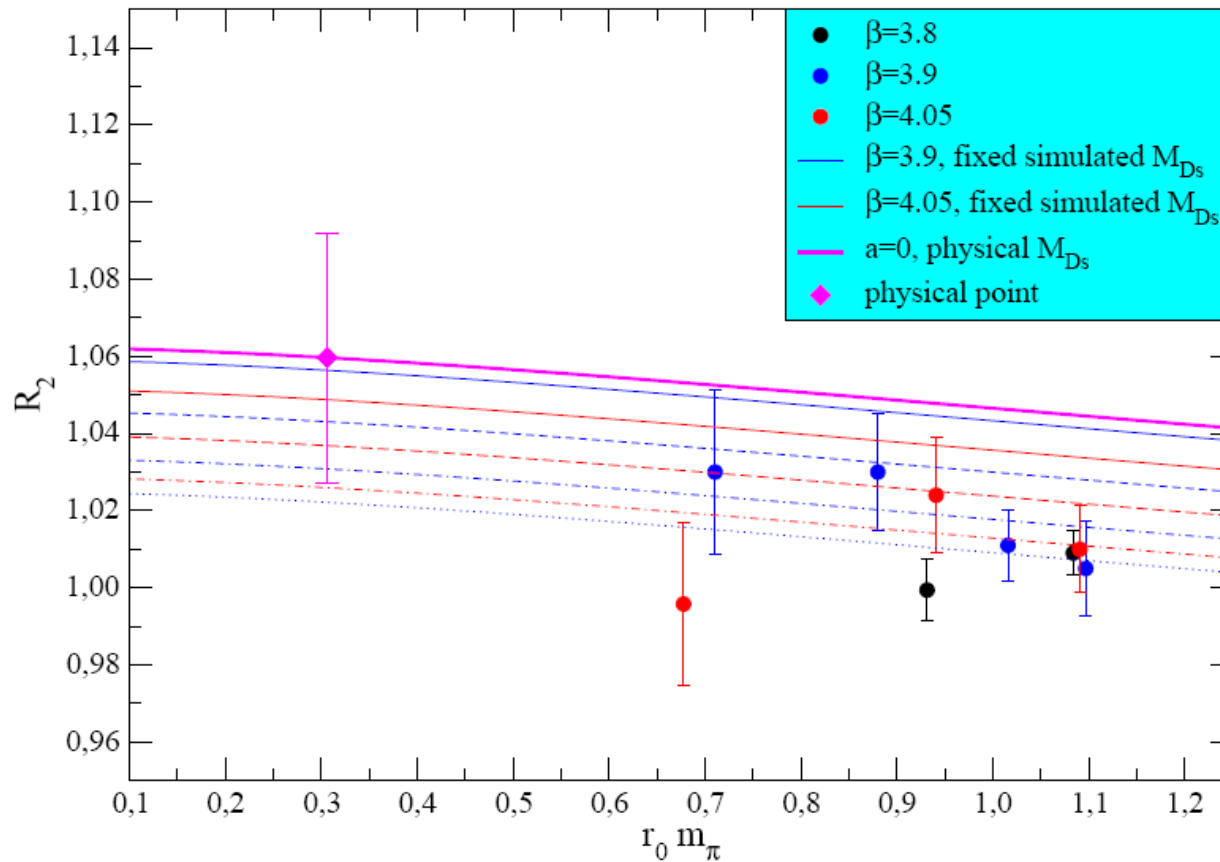


- Data with  $\mu_s \sim m_s^{\text{phys.}}$   
 $\mu_c \sim m_c^{\text{phys.}}$
- Curves at simulated  $\beta$  and  $M_K, M_{D_s}$
- Curve with  $a=0$  and physical  $M_K, M_{D_s}$

# Chiral Behaviour: dependence of $R_2$ on $m_\pi$

$$R_2 \equiv \frac{f_{D_s} \sqrt{m_{D_s}}}{f_K} \cdot \frac{f_\pi}{f_D \sqrt{m_D}}$$

$\mu_s \sim m_s^{\text{phys.}}$ ,  $\mu_c \sim m_c^{\text{phys.}}$



- Data with  $\mu_s \sim m_s^{\text{phys.}}$   
 $\mu_c \sim m_c^{\text{phys.}}$
- Curves at simulated  $\beta$  and  $M_K, M_{D_s}$
- Curve with  $a=0$  and physical  $M_K, M_{D_s}$

# Results

still preliminary systematic uncertainty!

$$f_D = 197(7)(4)(0)(11) \text{ MeV}$$

$$f_{D_s} = 244(4)(3)(2)(10) \text{ MeV}$$

$$f_{D_s}/f_D = 1.24(4)(1)(1)(2)$$

In agreement with the **NEW** exper. measurement:

$$f_D^{\text{exp.}} = (205.8 \pm 8.5 \pm 2.5) \text{ MeV}$$

Cleo-c (0806.2112)

**Smaller than the exper. average!!!**

$$f_{D_s}^{\text{exp.}} = (277 \pm 9) \text{ MeV}$$

Cleo-c+BaBar+Belle

→ (B.A. Dobrescu and A. Kronfeld '08)

statistics

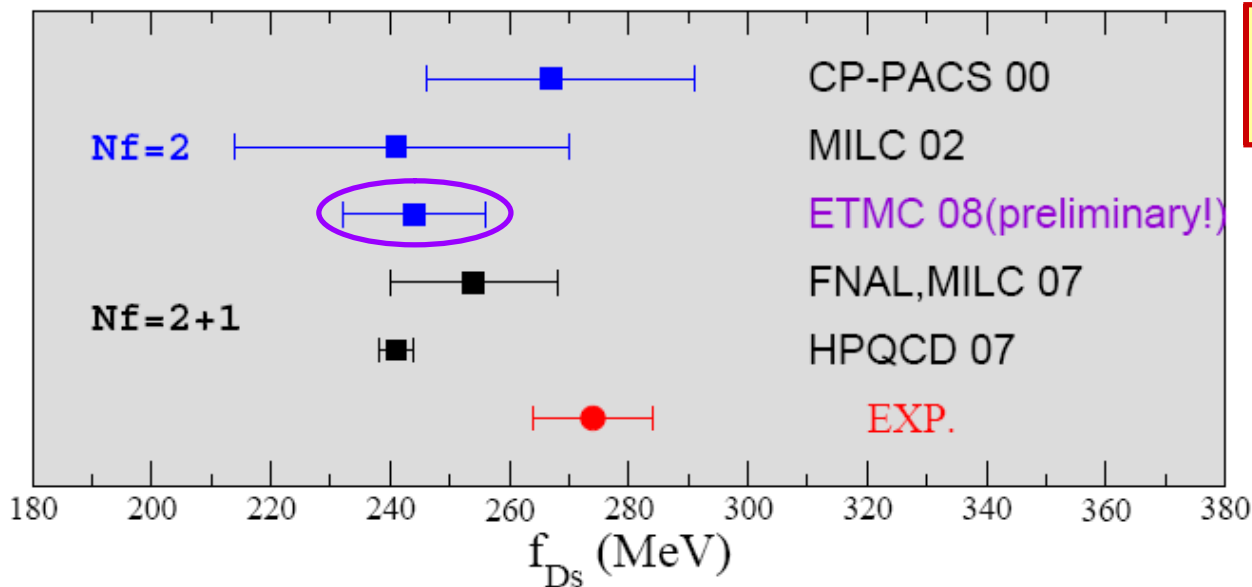
including  $\beta=3.8$

SU(3)-HMChPT

only  $\beta=4.05$



**Compatible with other recent unquenched results**





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