

# Physics with an Electron-Ion Collider

## *III: 1-d Structure: Parton Distributions*

Jian-ping Chen (陈剑平), Jefferson Lab, Virginia, USA

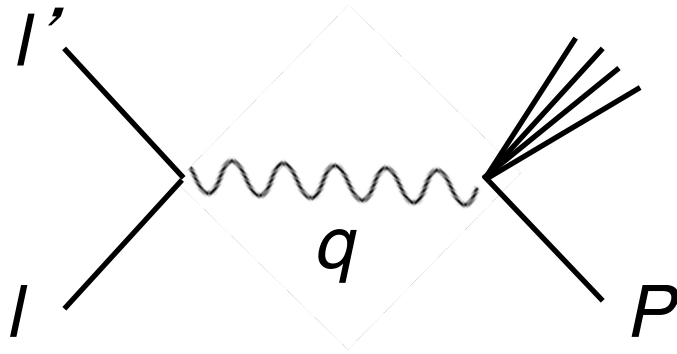
Pre-Workshop Lectures, Hadron-China2018, July 25, 2018

- Deep-Inelastic Scattering:
  - Unpolarized Structure Functions → Parton Distribution Functions
  - High-x physics
  - Sea asymmetry

# Deep-Inelastic Scattering Unpolarized Nucleon Structure

Parton Distributions:  
Flavor Structure: Valence and Sea  
Gluons

# Inelastic Scattering



Considerably more complex, indeed!

Simplify - consider *inclusive* inelastic scattering,

$$d\sigma \propto \langle |\mathcal{M}|^2 \rangle = \frac{g_e^4}{q^4} L_{\text{lepton}}^{\mu\nu} W_{\mu\nu \text{ nucleon}}, \quad W_{\mu\nu \text{ nucleon}}(p, q)$$

Again, two (parity-conserving, spin-independent) structure functions:

$$W_1, W_2 \text{ or, alternatively expressed, } F_1, F_2$$

which may depend on two invariants,

$$Q^2 = -q^2, \quad x = -\frac{q^2}{2q \cdot p}, \quad 0 < x < 1$$

So much for the structure, the physics is in the structure functions.

# Elastic scattering off Dirac Protons

Compare:

$$L_{\text{lepton}}^{\mu\nu} = 2 (k^\mu k'^\nu + k^\nu k'^\mu + g^{\mu\nu} (m^2 - k \cdot k'))$$

with:

$$K_{\mu\nu \text{ nucleon}} = K_1 \left( -g_{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) + \frac{K_2}{M^2} \left( p^\mu + \frac{1}{2} q^\mu \right) \left( p^\nu + \frac{1}{2} q^\nu \right)$$

which uses the relations between  $K_{1,2}$  and  $K_{4,5}$

Then, e.g. by substitution of  $k' = k - q$  in  $L$ :

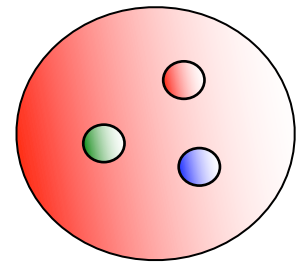
$$K_1 = -q^2, \quad K_2 = 4M^2$$

Note, furthermore, that inelastic cross section reduces to the elastic one for:

$$W_{1,2}(q^2, x) = -\frac{K_{1,2}(q^2)}{2Mq^2} \delta(x - 1)$$



# Elastic scattering off Dirac Partons



Imagine *incoherent* scattering off *Dirac* Partons (quarks)  $q$  :

$$W_1^q = \frac{e_q^2}{2m_q} \delta(x_q - 1), \quad W_2^q = -\frac{2m_q e_q^2}{q^2} \delta(x_q - 1) \quad \text{and} \quad x_q = -\frac{q^2}{2q \cdot p_q}$$

and, furthermore, suppose that the quarks carry a fraction,  $z$ , of the proton momentum

$$p_q = z_q p, \quad \text{so that} \quad x_q = \frac{x}{z_q} \quad (\text{also note } m_q = z_q M !)$$

which uses the relations between  $K_{1,2}$  and  $K_{4,5}$

Now,

$$MW_1 = M \sum_q \int_0^1 \frac{e_q^2}{2M} \delta(x - z_q) f_q(z_q) dz_q = \frac{1}{2} \sum_q e_q^2 f_q(x) \equiv F_1(x)$$
$$-\frac{q^2}{2Mx} W_2 = \sum_q \int_0^1 x e_q^2 \delta(x - z_q) f_q(z_q) dz_q = x \sum_q e_q^2 f_q(x) \equiv F_2(x)$$

Two important *observable* consequences,

Bjorken scaling:  $F_{1,2}(x)$ , not  $F_{1,2}(x, Q^2)$

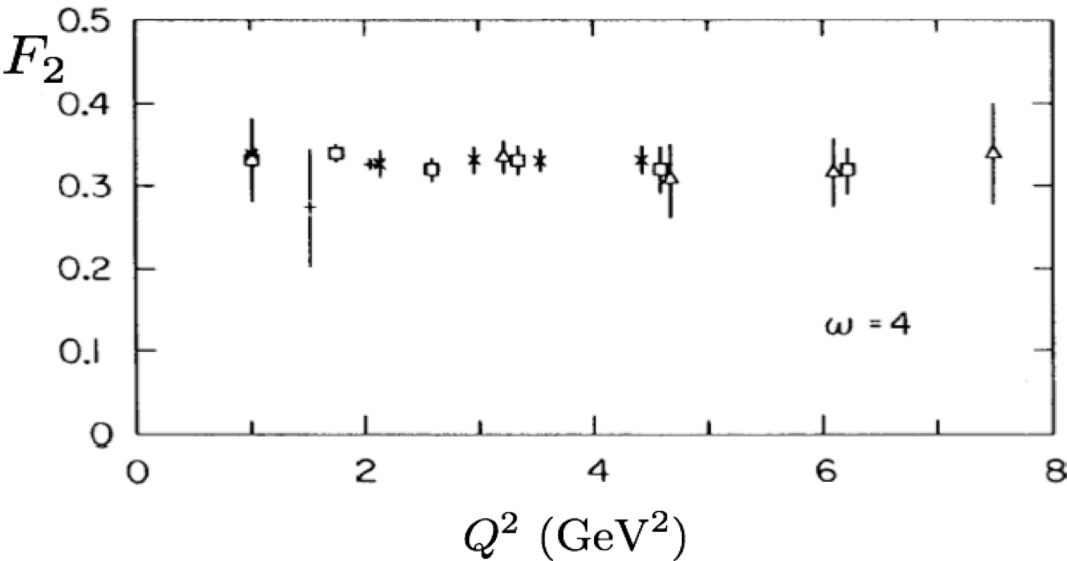
Callan-Gross relation:  $F_2 = 2xF_1(x)$

# Deep-Inelastic Electron Scattering

## Discovery of Quarks (Partons)

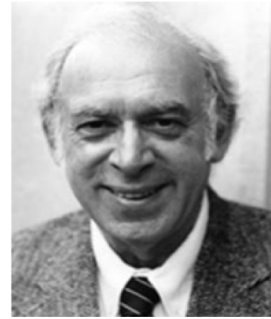
**Bjorken scaling:**

+ 6°    □ 18°  
 × 10°    △ 26°



*Point particles cannot be further resolved; their measurement does not depend on wavelength, hence  $Q^2$ ,*

*Spin-1/2 quarks cannot absorb longitudinally polarized vector bosons and, conversely, spin-0 (scalar) quarks cannot absorb transversely polarized photons.*



J.T. Friedman



R. Taylor

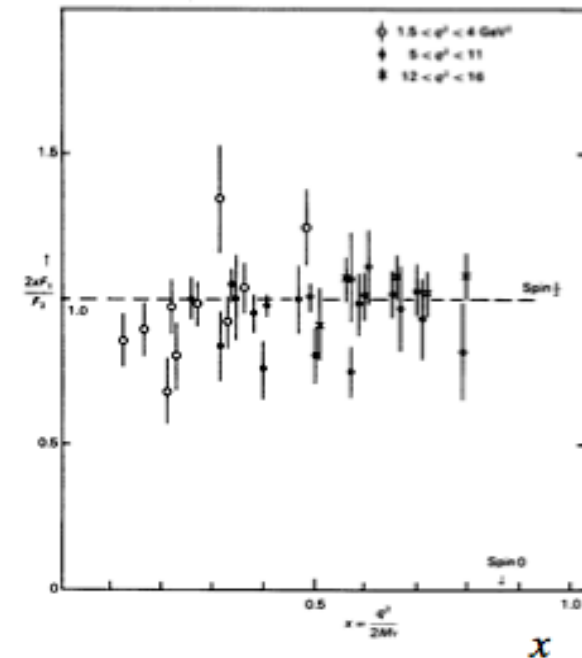


H.W. Kendall

**Nobel Prize 1990**

**Callan-Gross relation:**

$$\frac{2xF_1}{F_2}$$



spin 1/2

spin 0

# Deep-Inelastic Neutrino Scattering

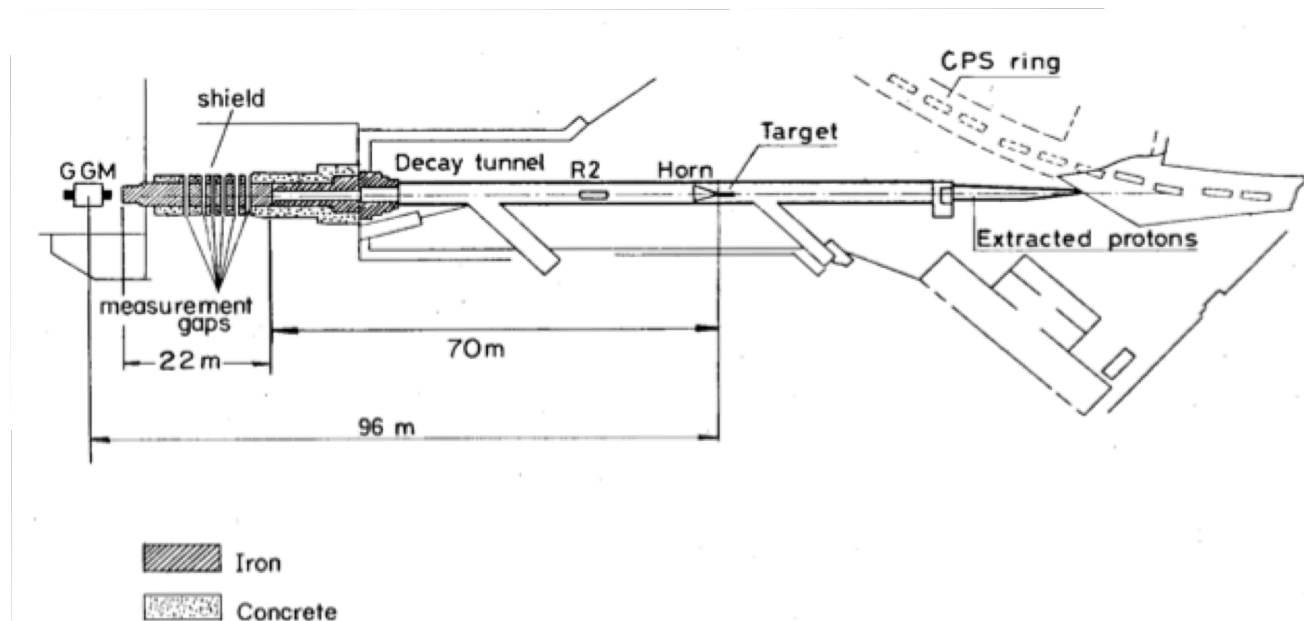


*Picture from CERN...*

*Gargamelle bubble chamber,  
observation of weak neutral  
current (1973).*

## Charged-current DIS!

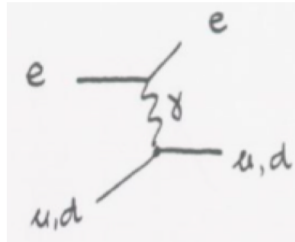
- Nucl.Phys. **B73** (1974) 1
- Nucl.Phys. **B85** (1975) 269
- Nucl.Phys. **B118** (1977) 218
- Phys.Lett. **B74** (1978) 134



# Deep-Inelastic Scattering - Fractional Electric Charges

**Neutral-current (photon)**

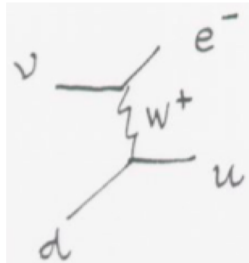
**DIS:**



$$F_2 = x \sum e_q^2 (q + \bar{q}), \quad p : uud, \quad n : ddu$$

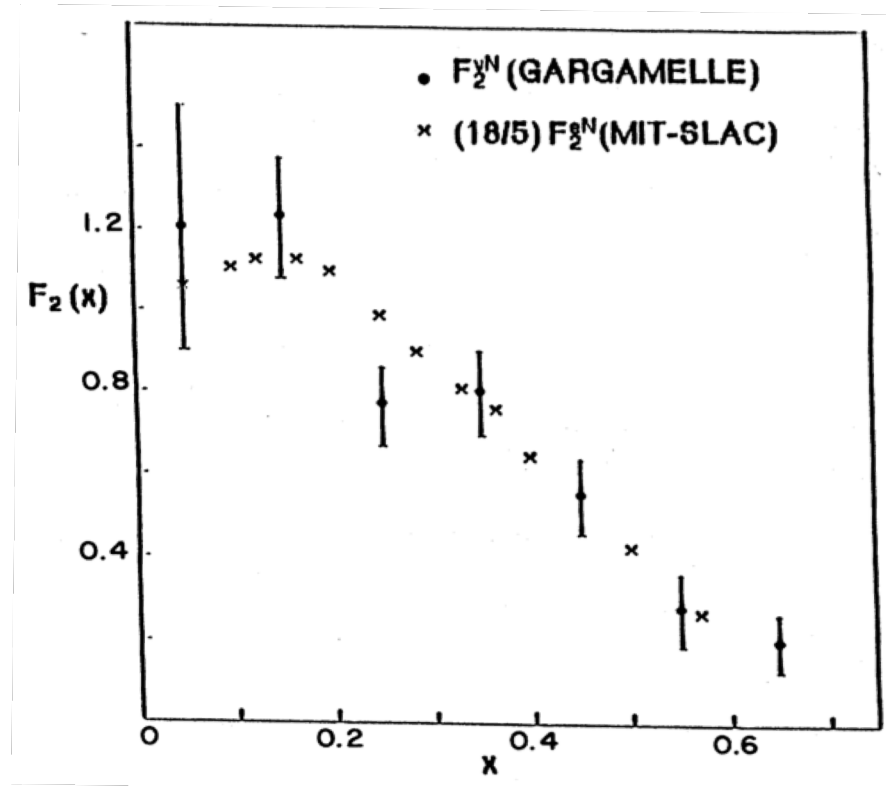
$$F_2^N = x \frac{e_u^2 + e_d^2}{2} (u + \bar{u} + d + \bar{d})$$

**Charged-current DIS:**



$$F_2^{\nu p} = 2x(d + \bar{u}), \quad F_2^{\nu n} = 2x(u + \bar{d})$$

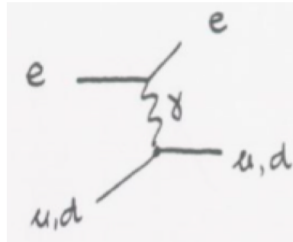
$$F_2^{\nu N} = x(u + \bar{u} + d + \bar{d})$$



# Deep-Inelastic Scattering - Fractional Electric Charges

**Neutral-current (photon)**

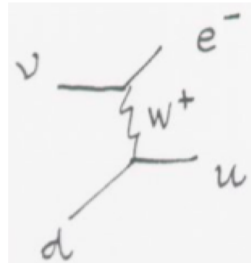
**DIS:**



$$F_2 = x \sum e_q^2 (q + \bar{q}), \quad p : uud, \quad n : ddu$$

$$F_2^N = x \frac{e_u^2 + e_d^2}{2} (u + \bar{u} + d + \bar{d})$$

**Charged-current DIS:**

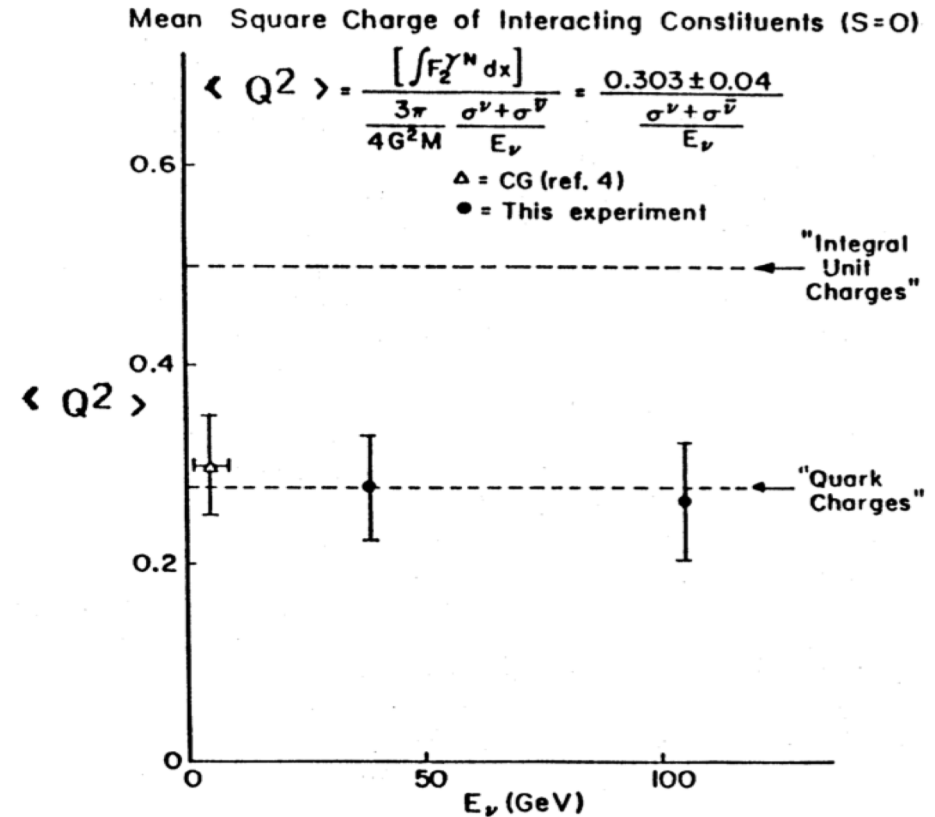


$$F_2^{\nu p} = 2x(d + \bar{u}), \quad F_2^{\nu n} = 2x(u + \bar{d})$$

$$F_2^{\nu N} = x(u + \bar{u} + d + \bar{d})$$

**Ratio:**

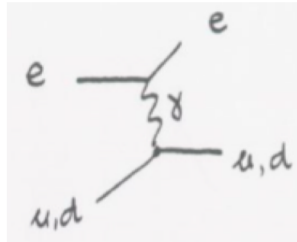
$$\frac{F_2^N}{F_2^{\nu N}} = \frac{1}{2} (e_u^2 + e_d^2) = \frac{5}{18} \simeq 0.28$$





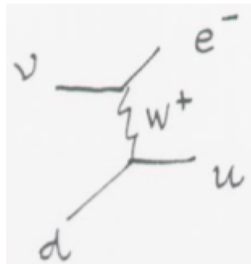
# Deep-Inelastic Scattering - Momentum Conservation

Neutral-current (photon) DIS:



$$F_2^N = x \frac{e_u^2 + e_d^2}{2} (u + \bar{u} + d + \bar{d})$$

Charged-current DIS:



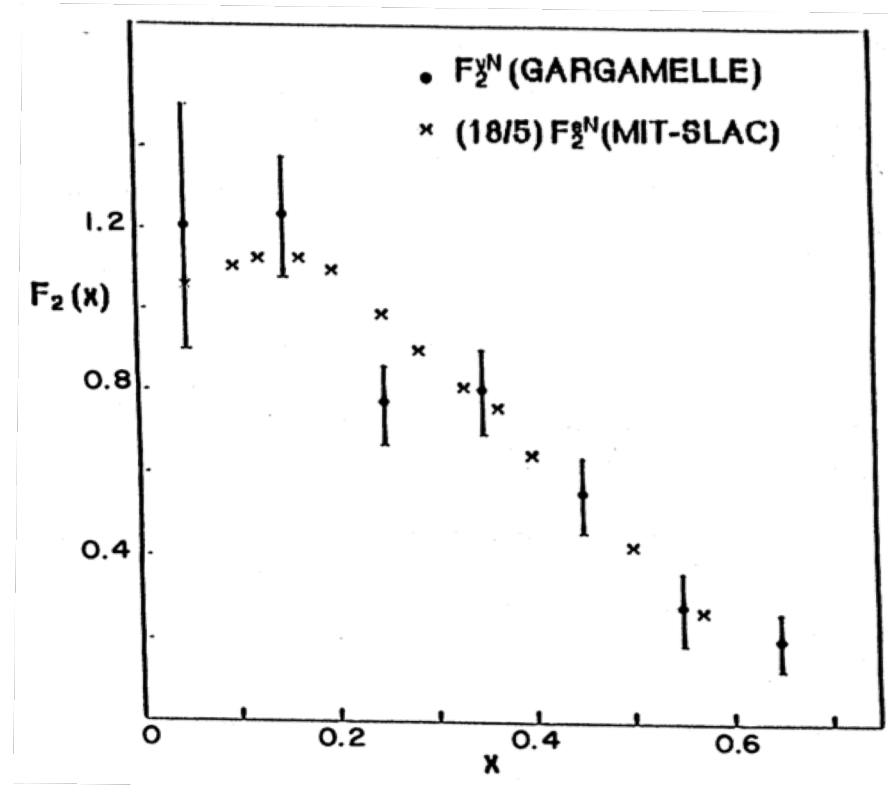
$$F_2^{\nu N} = x(u + \bar{u} + d + \bar{d})$$

Momentum fraction:

$$\int_0^1 F_2^N dx = \frac{e_u^2 + e_d^2}{2} \int_0^1 x(u + \bar{u} + d + \bar{d})$$

Gargamelle: 0.49 +/- 0.07

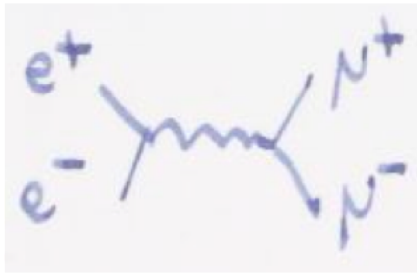
SLAC: 0.14 +/- 0.05



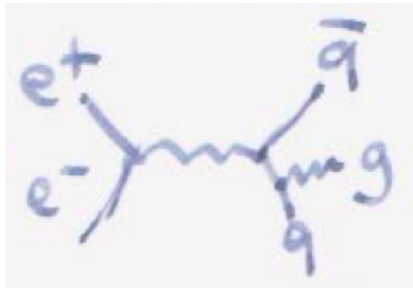
*Quarks carry half of the nucleon momentum!*

# 3-jet events at PETRA

Recall the intro on  
colour:

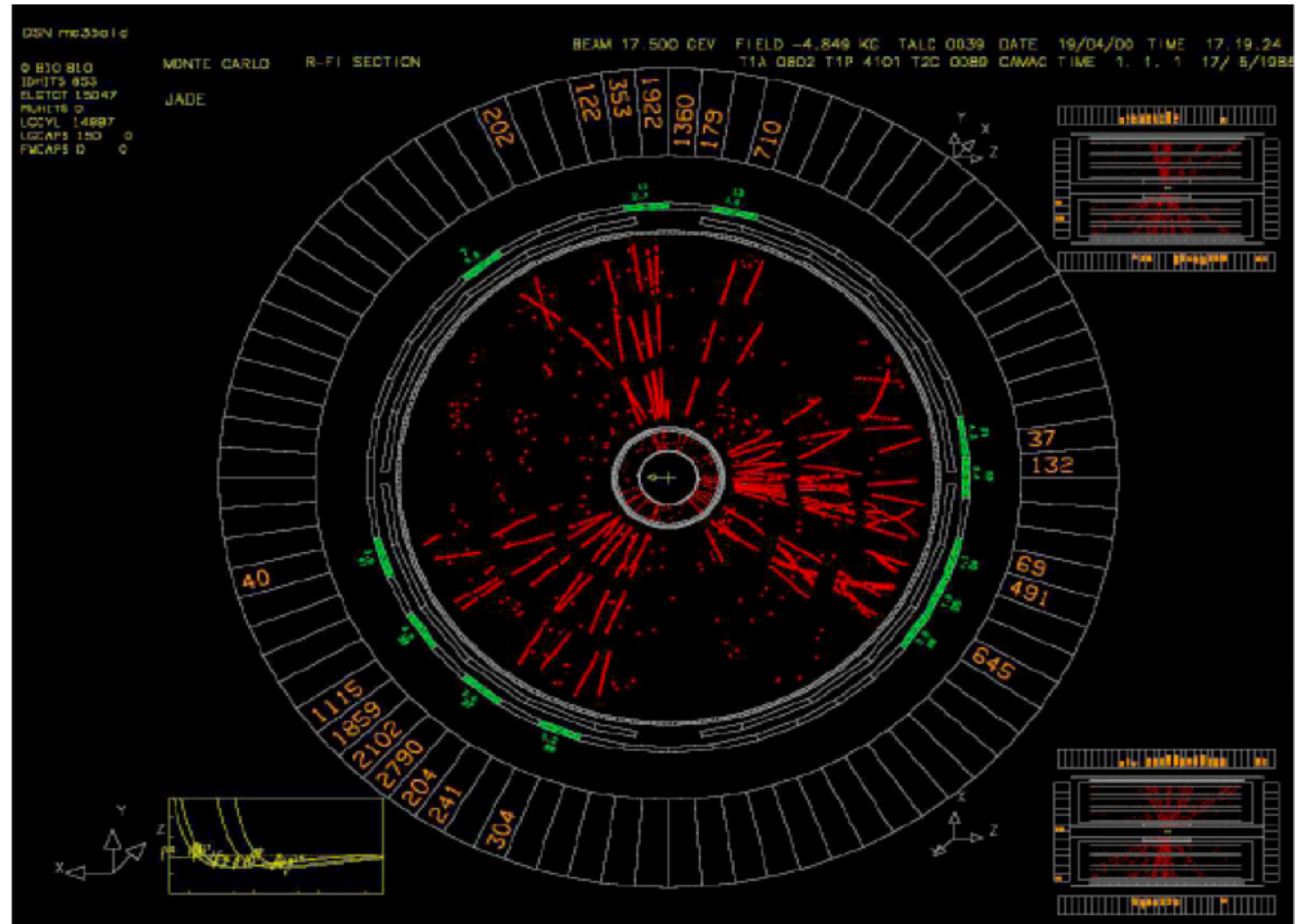


Observation of  
its higher order  
process,

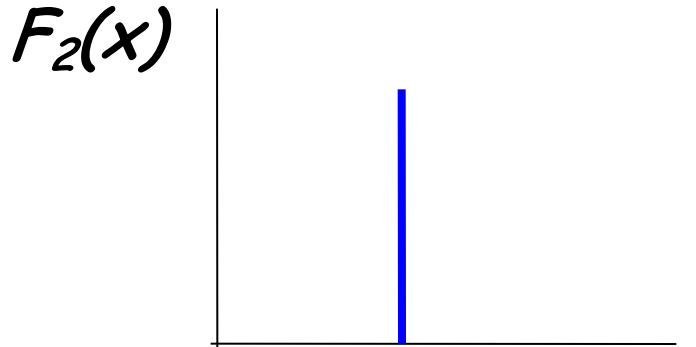


marks the discovery  
of the gluon.

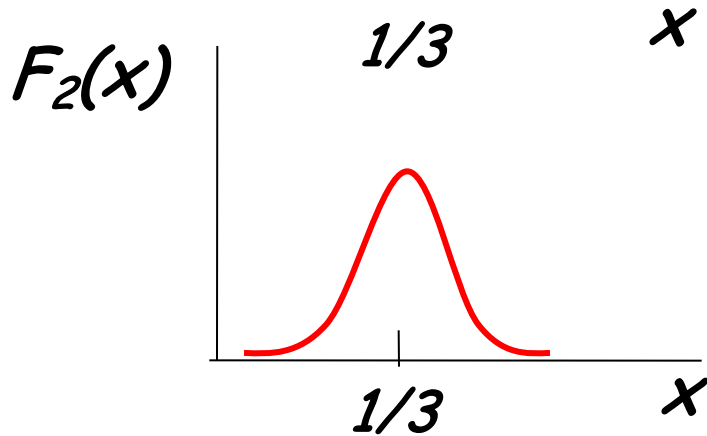
Mom. Conservation: *Gluons carry the other half of the nucleon momentum.*



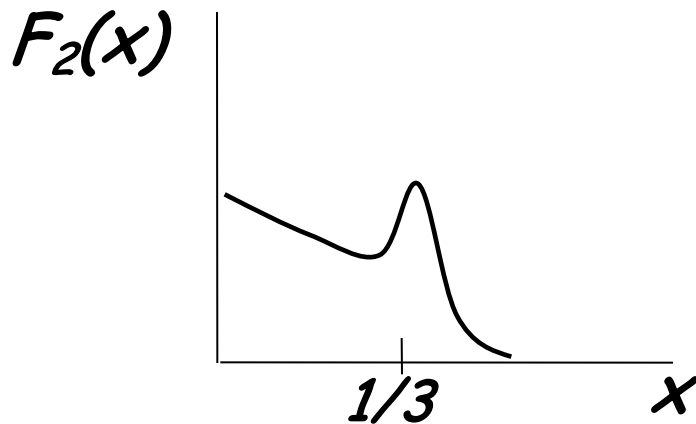
# ***Nucleon Structure***



***Three quarks with 1/3 of total proton momentum each.***



***Three quarks with some momentum smearing.***



***The three quarks radiate gluons to lower momentum fractions  $x$ .***

# Quark-Parton Model

$$F_1(x) = \frac{1}{2} \sum_i e_i^2 f_i(x) \quad g_1(x) = \frac{1}{2} \sum_i e_i^2 \Delta q_i(x)$$

$$f_i(x) = q_i^\uparrow(x) + q_i^\downarrow(x)$$

$$\Delta q_i(x) = q_i^\uparrow(x) - q_i^\downarrow(x)$$

$q_i(x)$  quark momentum distributions of flavor  $i$

$\uparrow(\downarrow)$  parallel (antiparallel) to the nucleon spin

$$F_2 = 2xF_1 \quad g_2 = 0$$

---

$$A_1(x) = \frac{g_1(x)}{F_1(x)} = \frac{\sum \Delta q_i(x)}{\sum f_i(x)}$$

## QCD Radiation

Schematically, DGLAP equations:

$$\frac{dq_f(x, Q^2)}{d \ln Q^2} = \alpha_s \left[ q_f \otimes P_{qq} + g \otimes P_{gq} \right]$$

convolution

strong coupling constant

That is, the change of quark distribution  $q$  with  $Q^2$  is given by the probability that  $q$  and  $g$  radiate  $q$ .

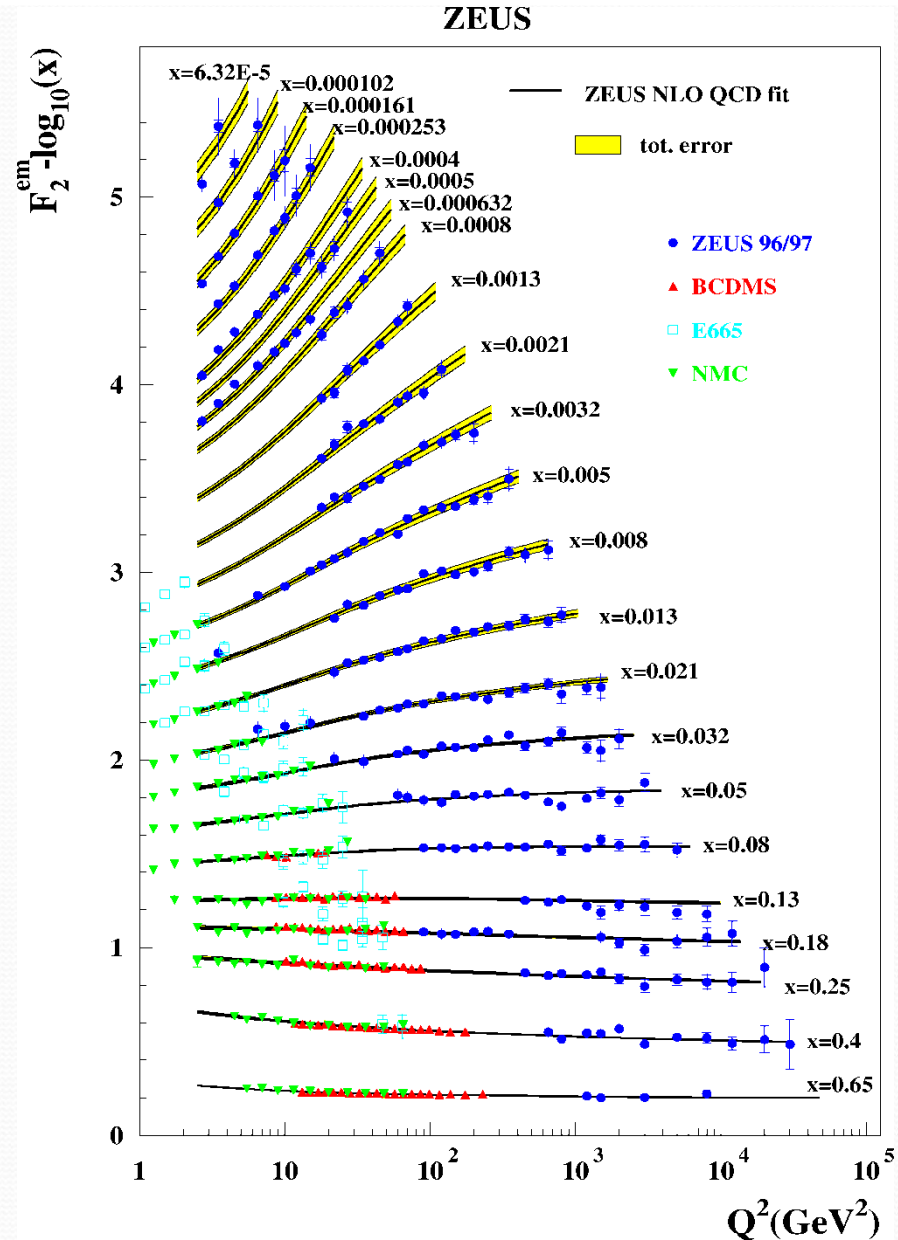
Similarly, for gluons:

$$\frac{dg(x, Q^2)}{d \ln Q^2} = \alpha_s \left[ \sum q_f \otimes P_{qg} + g \otimes P_{gg} \right]$$

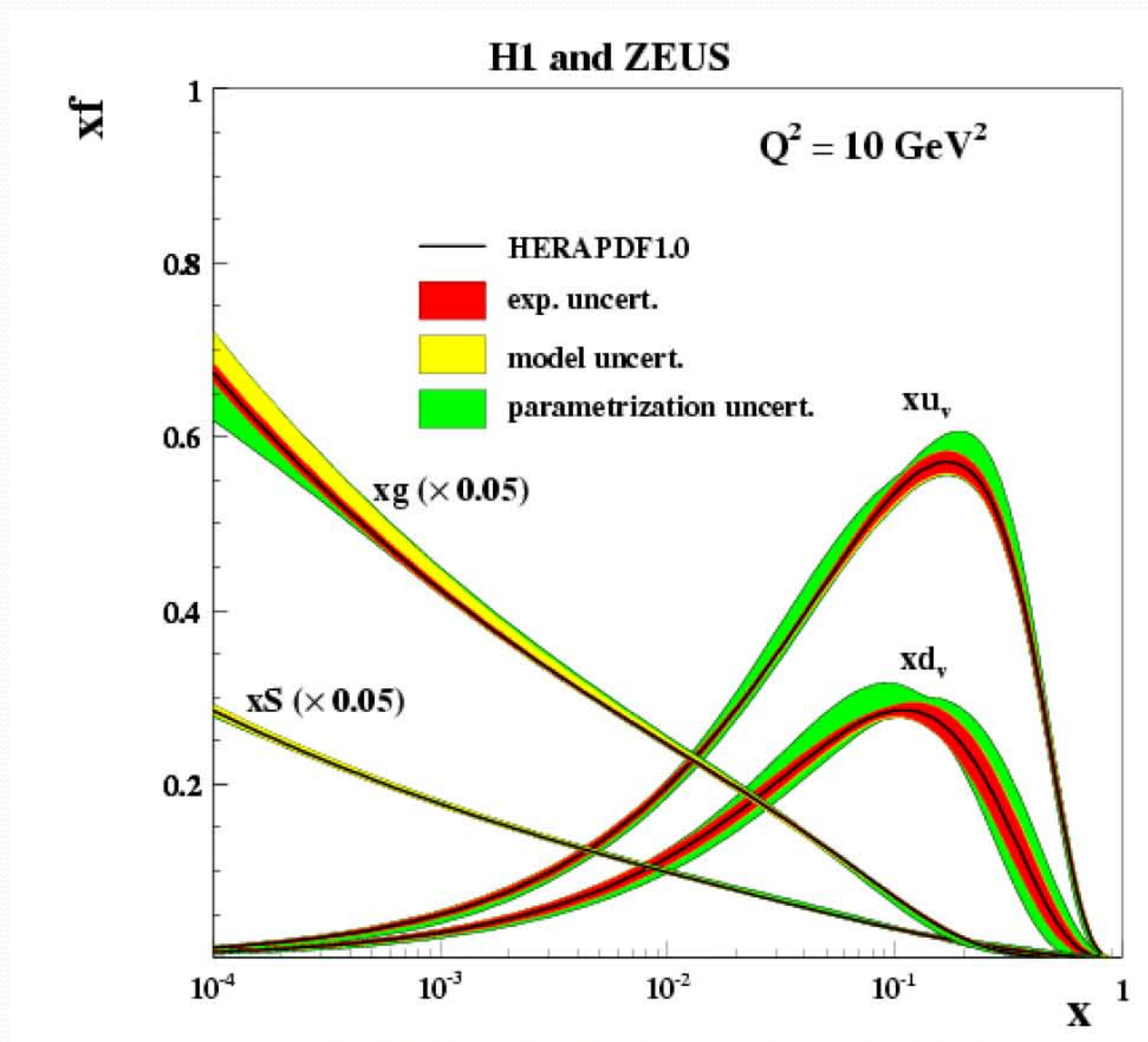


# Unpolarized Structure Function $F_2$

- Bjorken Scaling
- Scaling Violation
- Gluon radiation –
- QCD evolution
- NLO: Next-to-Leading-Order
- One of the best experimental tests of QCD



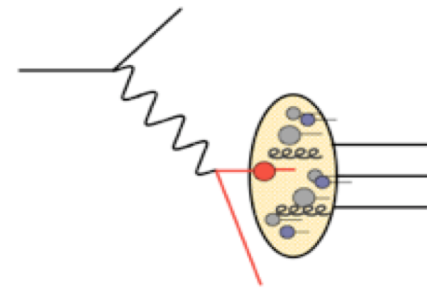
# Parton Distribution Functions (CTEQ6)



JHEP 1001: 109 (2010)

## **Brief Summary:**

DIS



DIS probe nucleon or nuclear structure,  
described in terms of quarks and gluons,

Feynman's parton model - point like partons, which

behave *incoherently* - combined with QCD radiation are remarkably successful in describing DIS cross sections.

QCD evolution allows one to relate quantitatively processes at different scales  $Q^2$ ,

Parton distributions  $f(x)$  are intrinsic properties of the nucleon and (thus) process independent.

*This is great for RHIC, LHC, and many other areas.*

Gluons are a very significant part of the nucleon

# Structure Functions at High $x$

## Valence Quark Distributions

# *Why Are PDFs at High $x$ Important?*

- Valence quark dominance: simpler picture
  - direct comparison with nucleon structure models
    - SU(6) symmetry, broken SU(6), di-quark, DSE
- $x \rightarrow 1$  region amenable to pQCD analysis
  - hadron helicity conservation?
  - Quark orbital angular momentum
- Clean connection with QCD, via lattice moments
- Input for search for physics BSM at high energy collider
  - evolution: high  $x$  at low  $Q^2 \rightarrow$  low  $x$  at high  $Q^2$
  - small uncertainties amplified
  - example: HERA 'anomaly' (1998)
- Input to nuclear and high energy calculations



# Predictions for High $x$

Proton Wavefunction (Spin and Flavor Symmetric)

$$|p\rangle = \frac{1}{\sqrt{2}} |u (ud)_{S=0}\rangle + \frac{1}{\sqrt{18}} |u (ud)_{S=1}\rangle - \frac{1}{3} |u \downarrow (ud)_{S=1}\rangle - \frac{1}{3} |d (uu)_{S=1}\rangle - \frac{\sqrt{2}}{3} |d \downarrow (uu)_{S=1}\rangle$$

Nucleon Model	$F_2^n/F_2^p$	d/u	$\Delta u/u$	$\Delta d/d$	$A_1^n$	$A_1^p$
SU(6)	2/3	1/2	2/3	-1/3	0	5/9
Scalar diquark	1/4	0	1	-1/3	1	1
pQCD	3/7	1/5	1	1	1	1

# Hadronic physics output 1: d/u ratio

→  $d/u$  ratio at high  $x$  of interest for nonperturbative models of nucleon

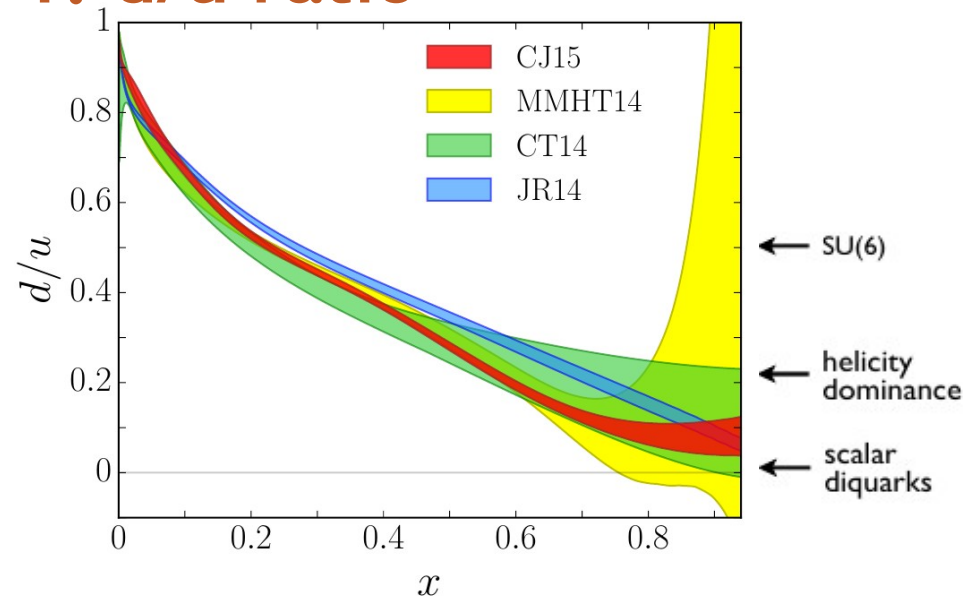
→ **CJ15:**

more flexible parametrization

$$d \rightarrow d + b x^c u$$

allows finite, nonzero  $x = 1$  limit

(standard PDF form gives 0 or  $\infty$  unless  $a_2^d = a_2^u$ )



**MMHT14:** fitted deuteron corrections  
standard  $d$  parametrization  
→ “UNDERCONSTRAINED”

**JR14 (and ABM12):**

Similar deuteron corrections  
standard  $d$ ; no lepton/W asym.  
→ “OVERCONSTRAINED”

**CT14:**  $\beta_u = \beta_d \implies d/u$  finite  
No nuclear corrections

Longstanding issue in proton structure

# Proton PVDIS: $d/u$ at high $x$

$$A_{PV} = \frac{G_F Q^2}{\sqrt{2}\pi\alpha} [a(x) + f(y)b(x)]$$

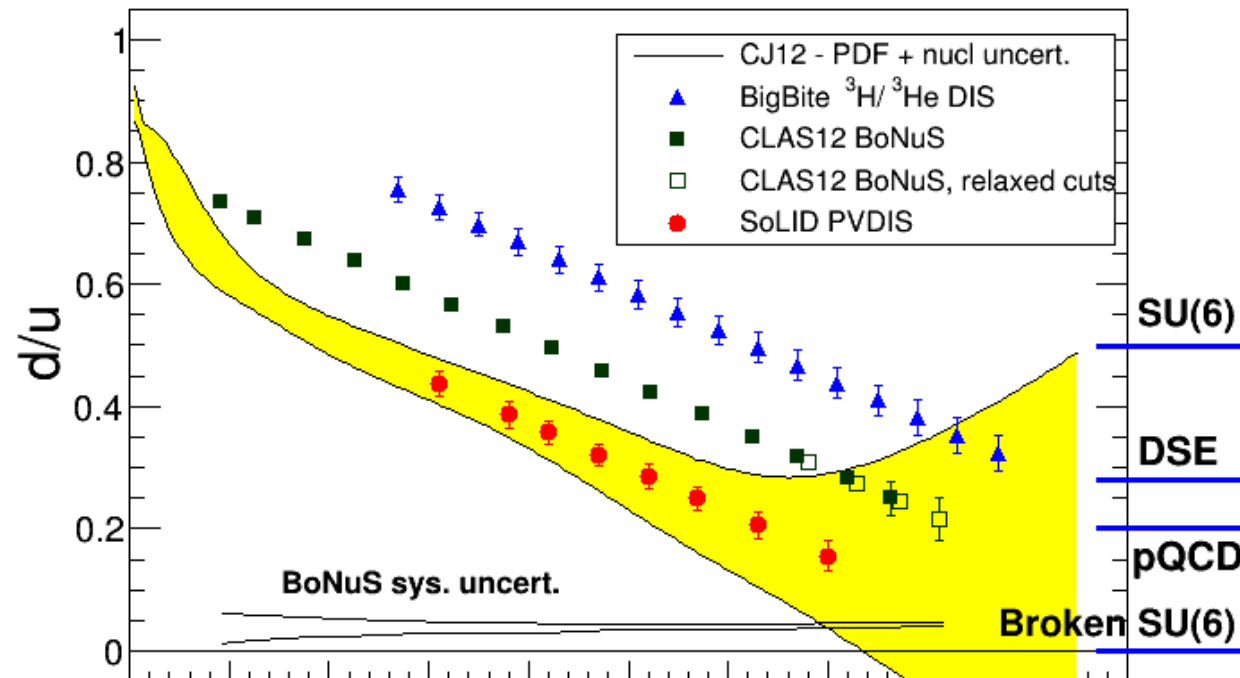
$$a^P(x) \approx \frac{u(x) + 0.91d(x)}{u(x) + 0.25d(x)}$$

SU(6):  $d/u \sim 1/2$

Broken SU(6):  $d/u \sim 0$

Perturbative QCD:  $d/u \sim 1/5$

Projected 12 GeV  $d/u$  extractions



- 3 JLab 12 GeV experiments:

- CLAS12 BoNuS spectator tagging
- BigBite – DIS  $^3\text{H}/^3\text{He}$  ratio
- **SoLID – PVDIS  $ep$**

- **The SoLID extraction of  $d/u$  is directly from  $ep$  DIS:**

- **No nuclear corrections**
- **No assumption of charge symmetry**

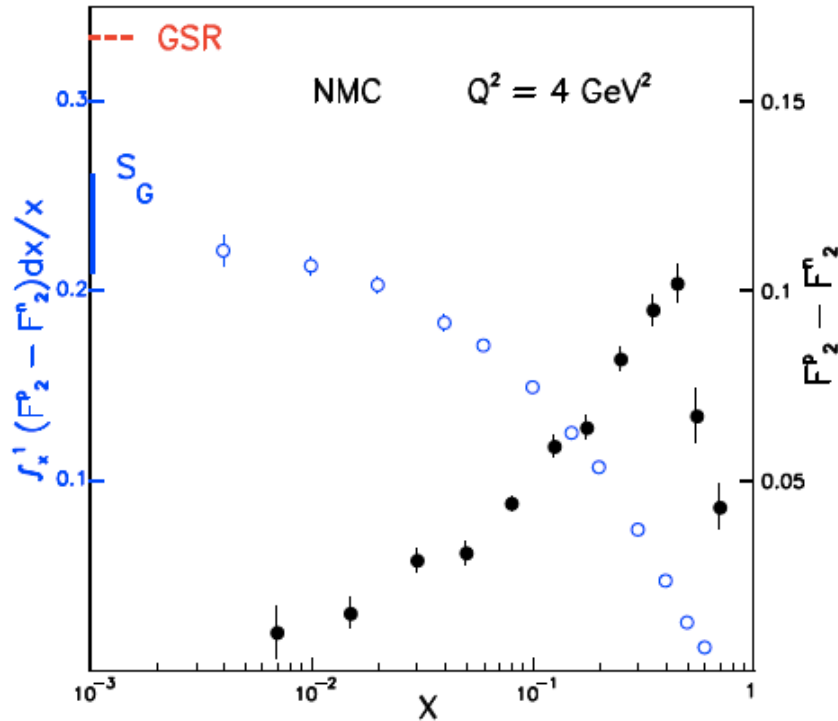
# Sea Quark Distributions

Sea Asymmetry ( $\bar{d}/\bar{u}$ )

# Flavor structure of the proton sea

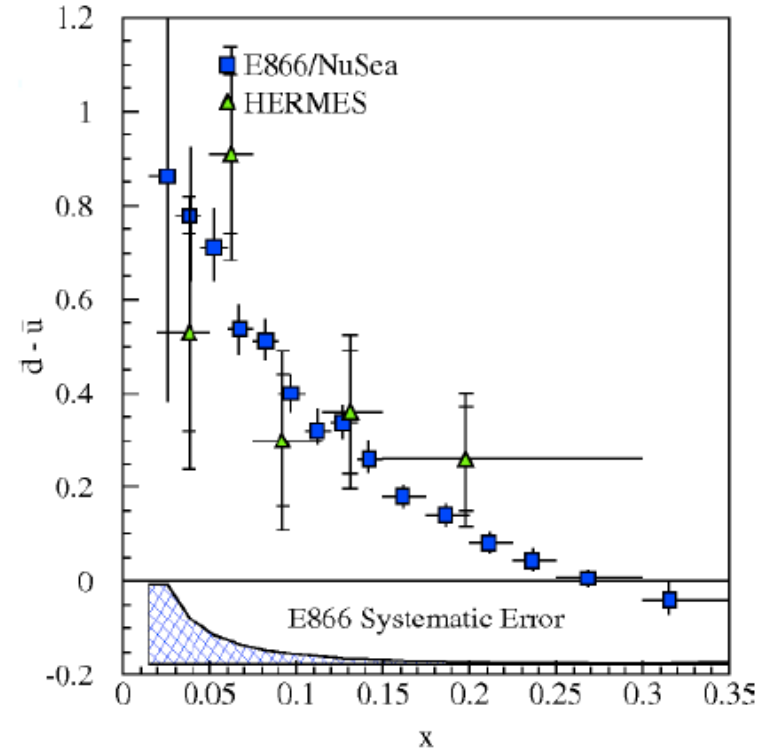
□ The proton sea is not SU(3) symmetric!

Violation of Gottfried sum rule



$$\begin{aligned}
 S_G &= \int_0^1 [(F_2^p(x) - F_2^n(x)) / x] dx \\
 &= \frac{1}{3} + \frac{2}{3} \int_0^1 (\bar{u}_p(x) - \bar{d}_p(x)) dx \\
 &= \frac{1}{3} \quad (\text{if } \bar{u}_p = \bar{d}_p) \quad \text{NMC: } S_G = 0.235 \pm 0.026
 \end{aligned}$$

Confirmed by Drell-Yan exp't



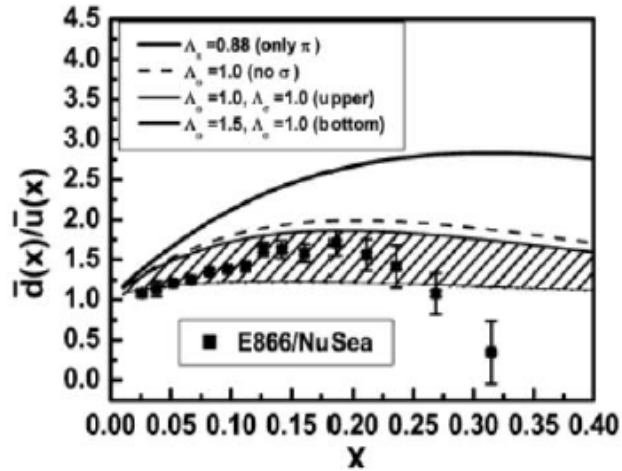
Why  $\bar{d}(x) \neq \bar{u}(x)$  ?

Why does  $\bar{d}(x) - \bar{u}(x)$  change sign?

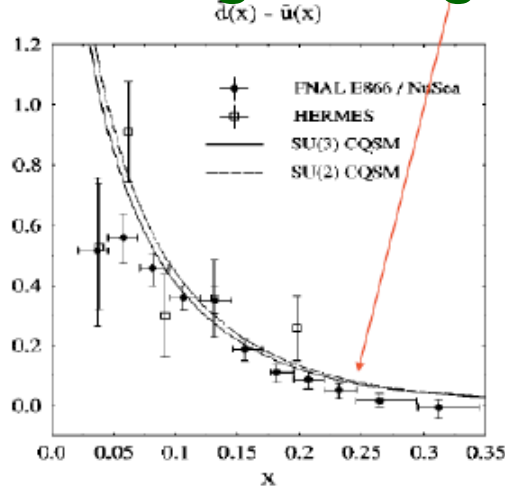


# Challenges for $\bar{d}(x) - \bar{u}(x)$

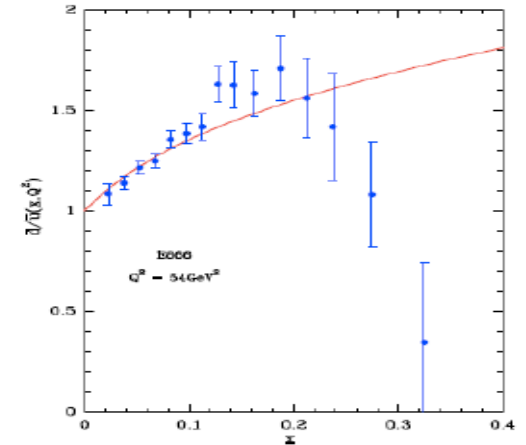
□ All known models predict no sign change!



Meson cloud

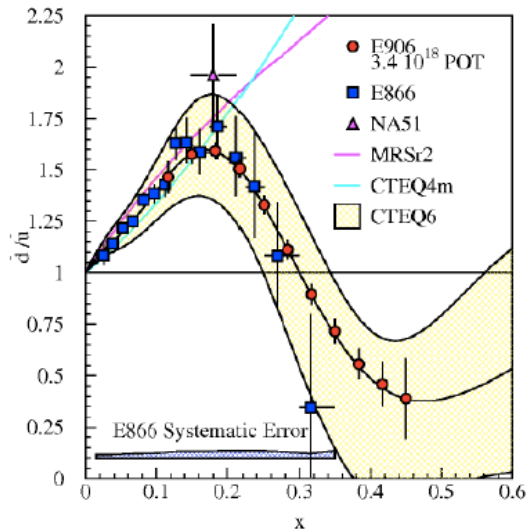


Chiral-quark soliton model

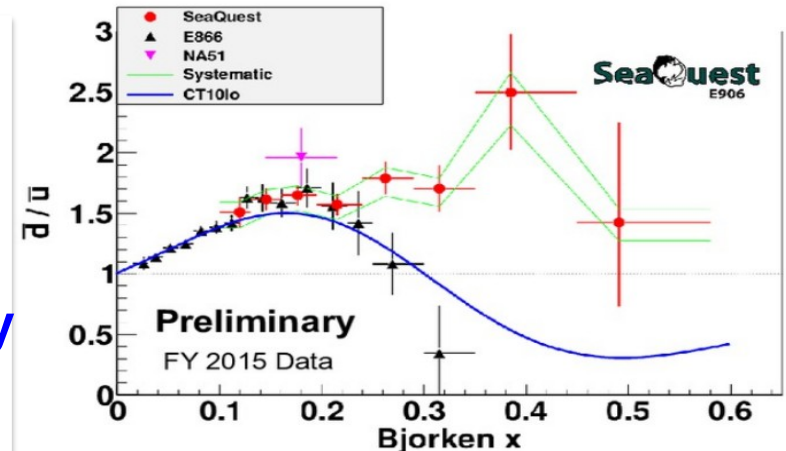


Statistic model

□ Future experiments – E906, ...



Preliminary  
From  
Seaquest



B.Kerns, DNP April 2016

# Summary

- Electron Scattering to study Nucleon Structure
- Deep-Inelastic Scattering
  - Unpolarized Structure functions → Parton Distributions
  - High- $x$ , valance quark distributions,  $d/u$
  - Sea distributions:  $d_{\text{bar}}/u_{\text{bar}}$  asymmetry

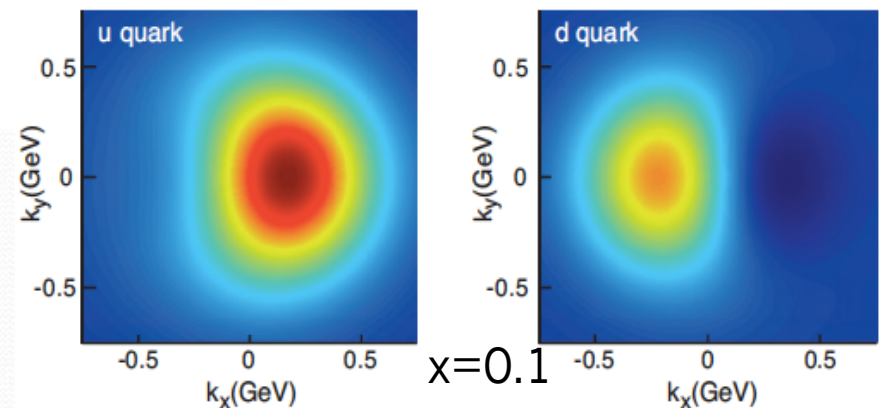
# Physics with an Electron-Ion Collider

## *II: Nucleon 3-d Structure*

Jian-ping Chen (陈剑平), Jefferson Lab, Virginia, USA

Pre-Workshop Lectures, Hadron-China2018, July 25, 2018

- Unified Picture of Nucleon Structure: Wigner Distribution
- GPDs: 3-d (2-d spatial+1-d momentum) distributions
- TMDs: 3-d momentum distributions
- Transversity and tensor charge
- SoLID program
- EIC program





# Nucleon Spin Structure Study

- 1980s: EMC (CERN) + early SLAC

quark contribution to proton spin is very small

$$\Delta\Sigma = (12 + -9 + -14)\% ! \quad \text{'spin crisis'}$$

- 1990s: SLAC, SMC (CERN), HERMES (DESY)

$\Delta\Sigma = 20-30\%$ , the rest: gluon and quark orbital angular momentum

$$\frac{1}{2} = \frac{1}{2} \sum_f (q_f^+ - q_f^-) + L_q + \Delta G + L_g$$

gauge invariant

$$(\frac{1}{2})\Delta\Sigma + \mathcal{L}q + J_G = 1/2$$

**Bjorken Sum Rule** verified to <10% level

- 2000s: COMPASS (CERN), HERMES, RHIC-Spin, JLab, ... :

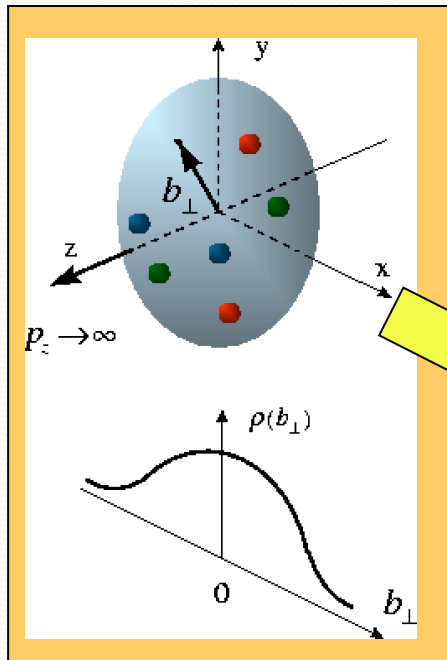
$\Delta\Sigma \sim 30\%$ ;  $\Delta G$  contributes, **orbital angular momentum significant**

Large-x (valence quark) behavior; Moments and sum rules

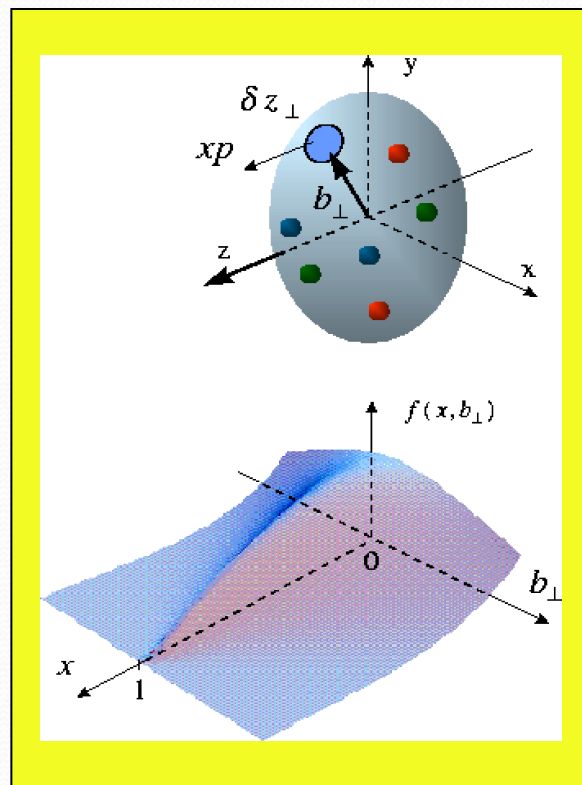
**Needs 3-d structure information to complete the proton spin puzzle**

# Generalized Parton Distributions (GPDs)

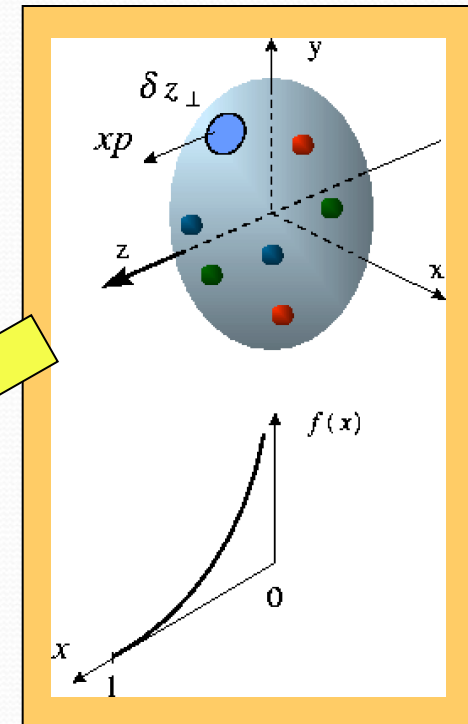
X. Ji, D. Mueller, A. Radyushkin (1994-1997)



Proton form factors,  
**transverse** charge &  
current densities



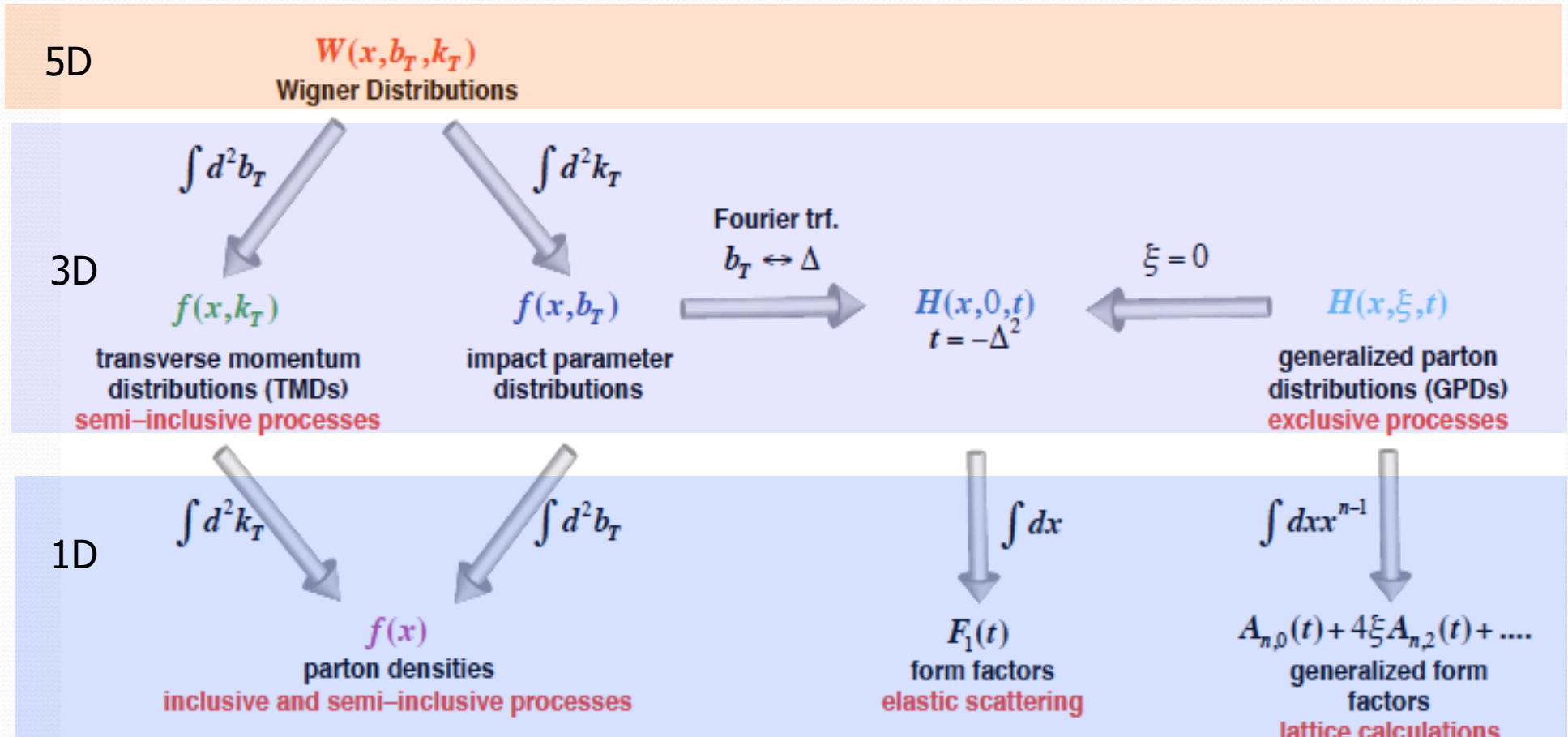
**Correlated** quark momentum  
and helicity distributions in  
transverse space - **GPDs**



Structure functions,  
quark **longitudinal**  
momentum & helicity  
distributions

# Unified View of Nucleon Structure

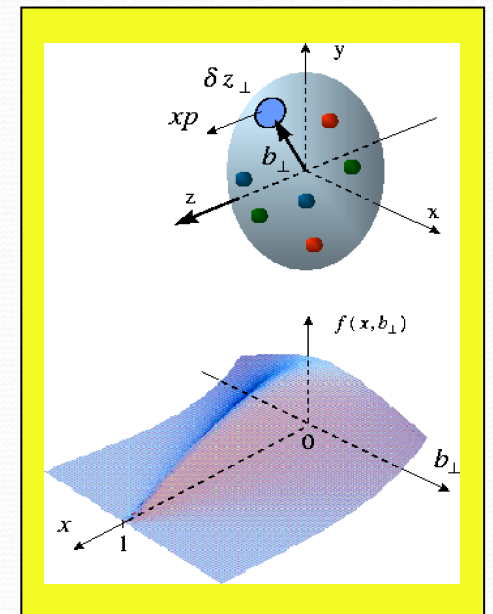
## □ Wigner distributions





# 3-D Structure I

## Generalized Parton Distributions

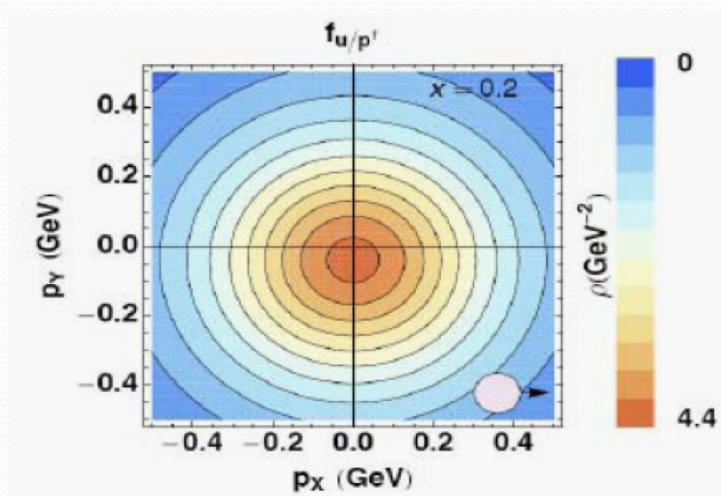




# 3-D Imaging - Two Approaches

## TMDs

2+1 D picture in **momentum space**

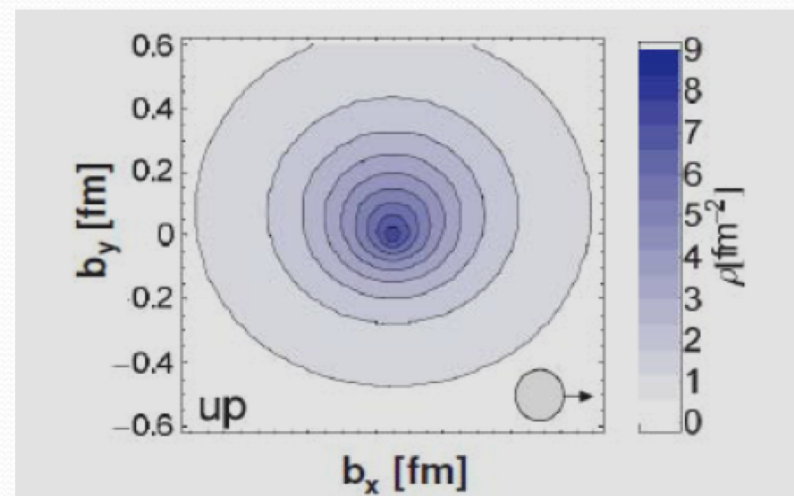


Bacchetta, Conti, Radici

- intrinsic transverse motion
- spin-orbit correlations- relate to OAM
- non-trivial factorization
- accessible in SIDIS (and Drell-Yan)

## GPDs

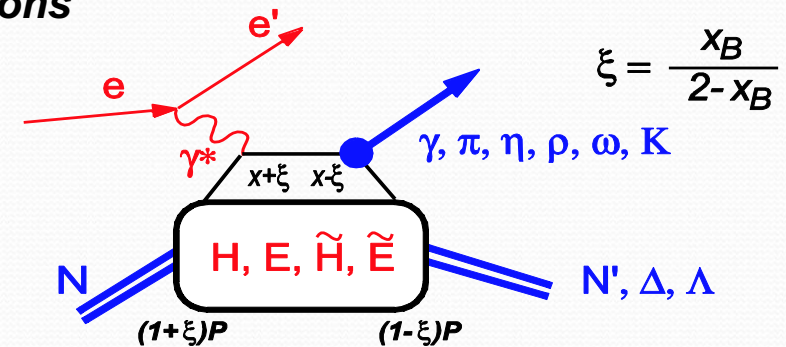
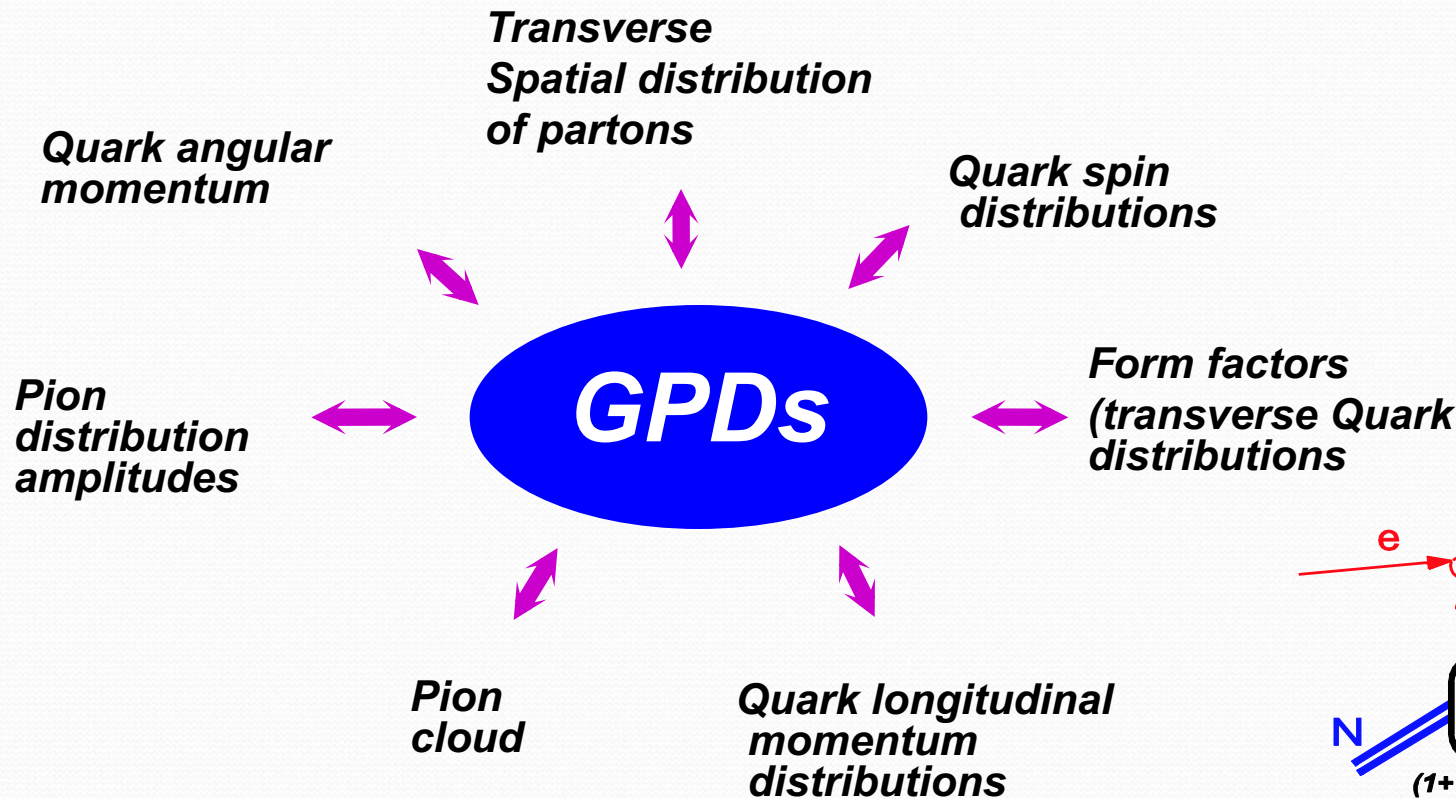
2+1 D picture in **impact-parameter space**



QCDSF collaboration

- collinear but long. momentum transfer
- indicator of OAM; access to Ji's total  $J_{q,g}$
- existing factorization proofs
- DVCS, exclusive vector-meson production

# Description of Hadron Structure via Generalized Parton Distributions



H, E - unpolarized,  $\tilde{H}, \tilde{E}$  - polarized GPD  
 The GPDs Define Nucleon Structure



# known information on **GPDs**

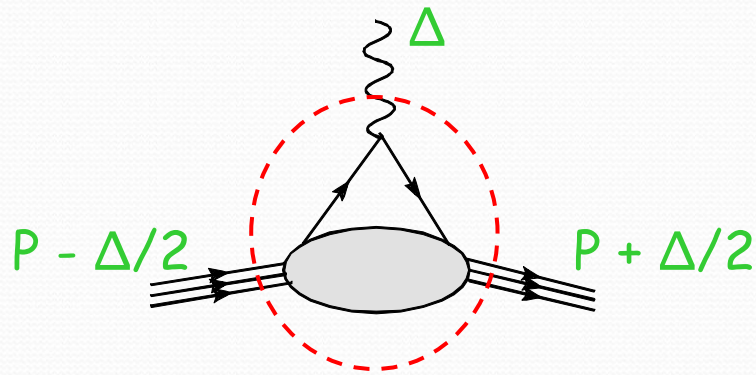
➔ forward limit : ordinary **parton distributions**

$$H^q(x, \xi = 0, t = 0) = q(x) \quad \text{unpolarized quark distribution}$$

$$\tilde{H}^q(x, \xi = 0, t = 0) = \Delta q(x) \quad \text{polarized quark distribution}$$

$E^q, \tilde{E}^q$  : do NOT appear in DIS ➔ additional information

➔ first moments : nucleon **electroweak form factors**



$\xi$  independence : Lorentz invariance

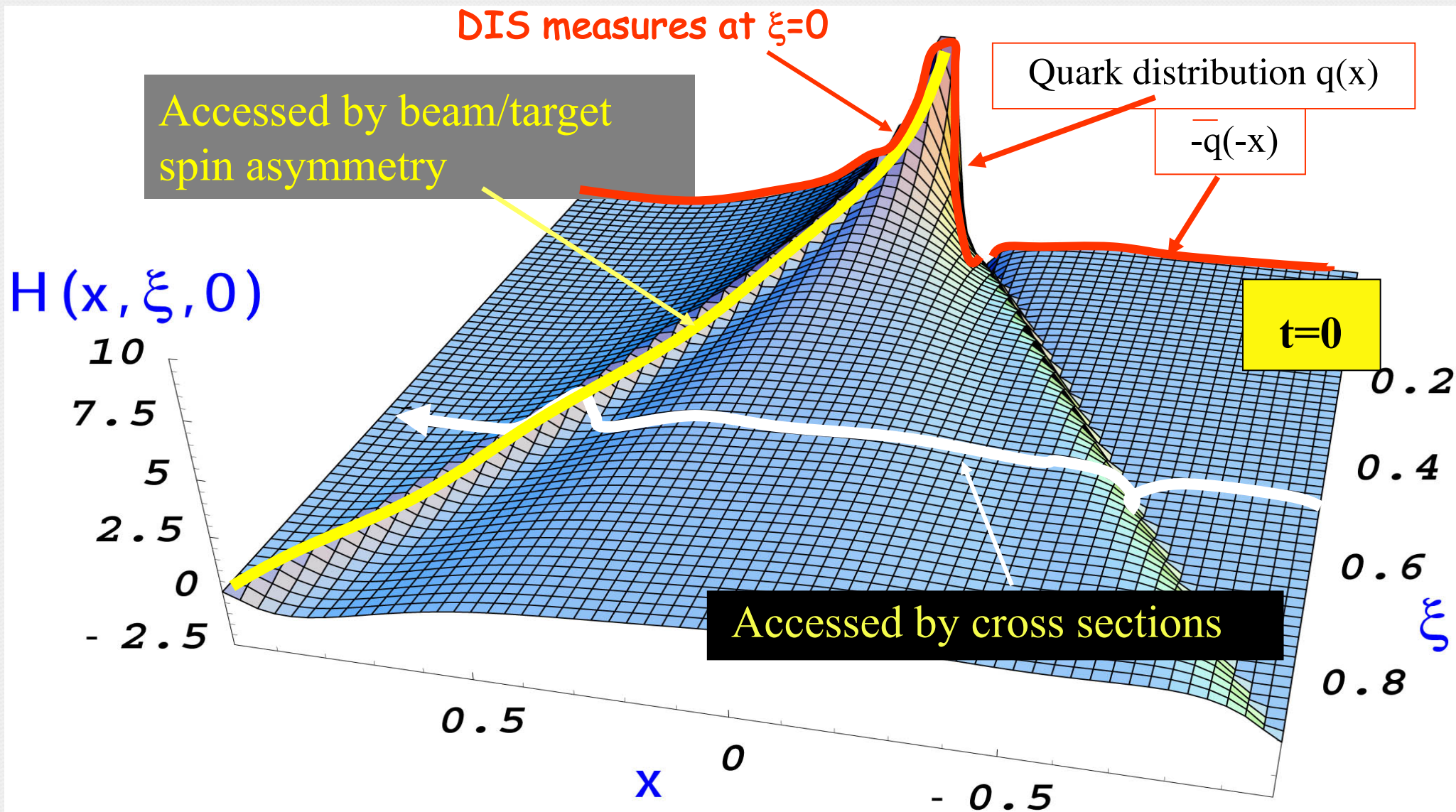
$$\int_{-1}^1 dx H^q(x, \xi, t) = F_1^q(t) \quad \text{Dirac}$$

$$\int_{-1}^1 dx E^q(x, \xi, t) = F_2^q(t) \quad \text{Pauli}$$

$$\int_{-1}^1 dx \tilde{H}^q(x, \xi, t) = G_A^q(t) \quad \text{axial}$$

$$\int_{-1}^1 dx \tilde{E}^q(x, \xi, t) = G_P^q(t) \quad \text{pseudo-scalar}$$

# Access GPDs through DVCS x-section & asymmetries

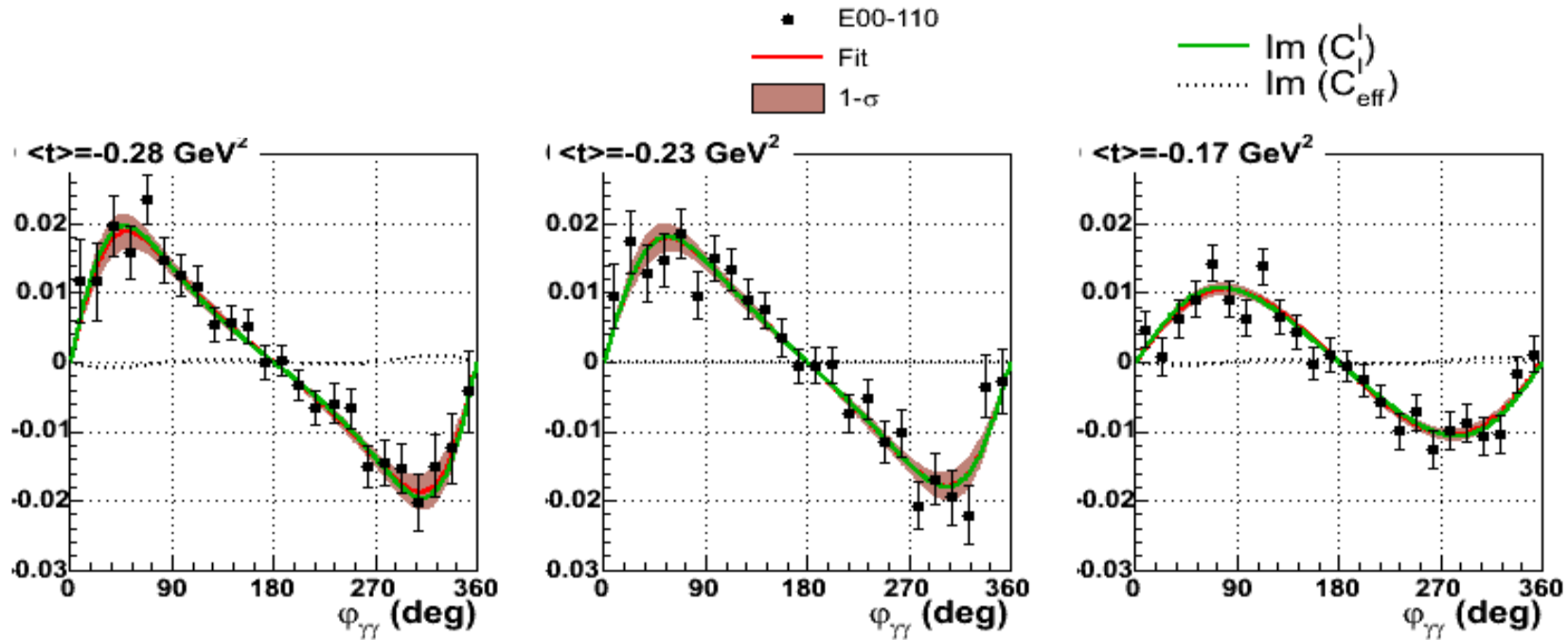




# Hall A DVCS Experiment

## Handbag Dominance at Modest $Q^2$

$$\frac{d^4\sigma^+}{dx_B dQ^2 d\varphi dt} - \frac{d^4\sigma^-}{dx_B dQ^2 d\varphi dt} \quad [\text{nb/GeV}^4]$$



— Twist 2 contribution  
 - - - - - Twist 3 contribution **strongly suppressed**

**The Twist-2 term can be extracted accurately from the cross-section difference**  
**Dominance of twist-2  $\Rightarrow$  handbag dominance  $\Rightarrow$  DVCS interpretation**

# 3D Images of the Proton's Quark Content

M. Burkardt PRD 66, 114005 (2002)

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\Delta_\perp \cdot \mathbf{b}_\perp} H(x, 0, -\Delta_\perp^2).$$

$\mathbf{b}_\perp$  - Impact parameter

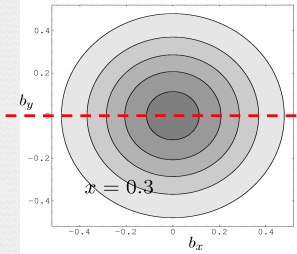
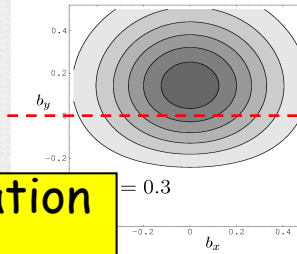
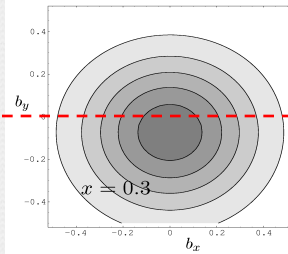
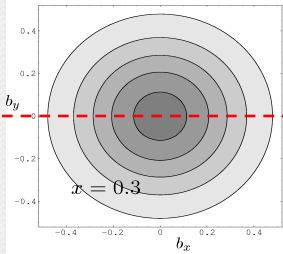
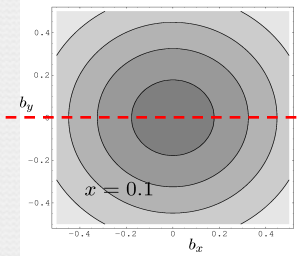
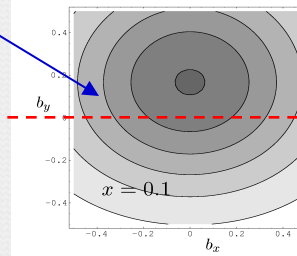
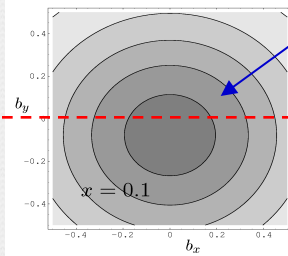
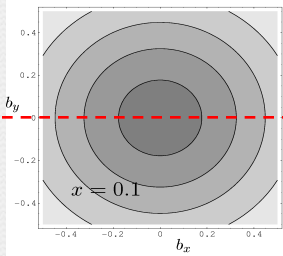
transverse polarized target

$u(x, \mathbf{b}_\perp)$

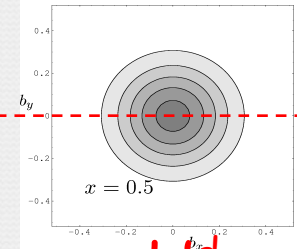
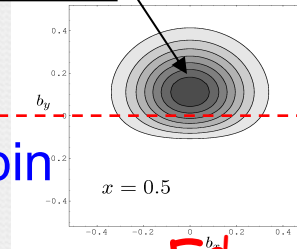
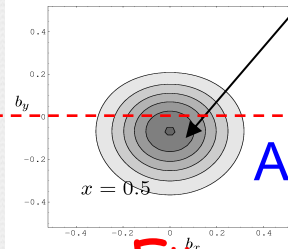
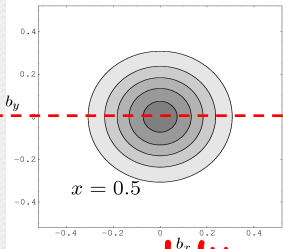
$u_x(x, \mathbf{b}_\perp)$

$d_x(x, \mathbf{b}_\perp)$

$d(x, \mathbf{b}_\perp)$



quark flavor polarization



Accessed in Single Spin Asymmetries.

Needs:

$H^u$

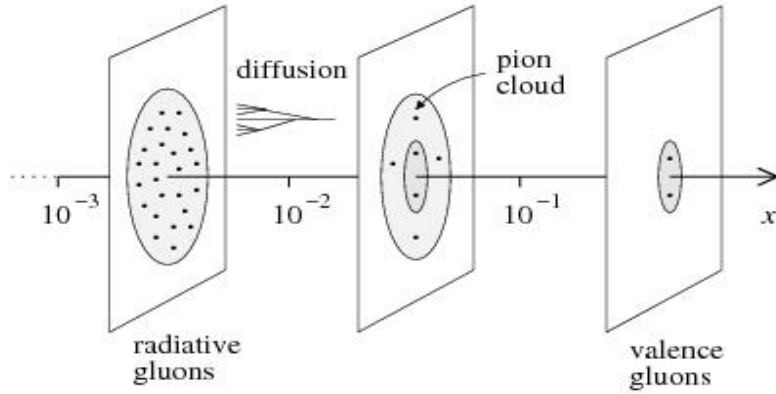
$E^u$

$E^d$

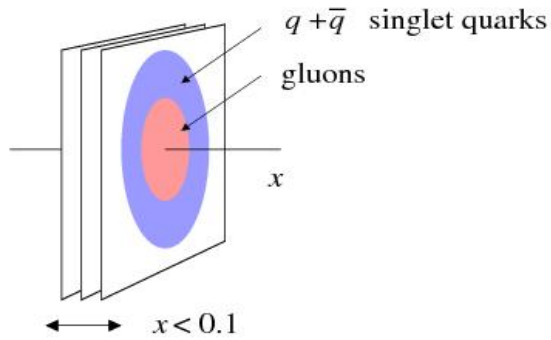
$H^d$



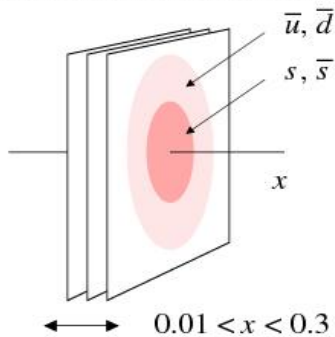
# Detailed differential images from nucleon's partonic structure



Weiss, Hyde, Horn



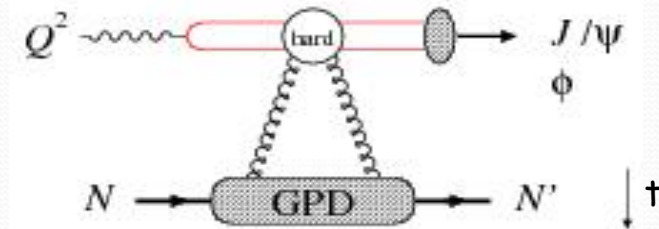
Fazio



Horn

- $Q^2 > 10 \text{ GeV}^2$  for factorization
- Statistics hungry at high  $Q^2$ !

EIC: Gluon size from  $J/\Psi$  and  $\phi$  electroproduction ( $Q^2 > 10 \text{ GeV}^2$ )



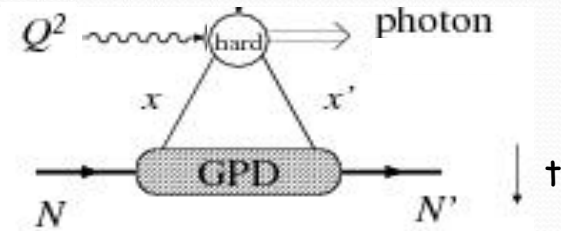
[Transverse distribution derived directly from  $t$  dependence]

Hints from HERA:

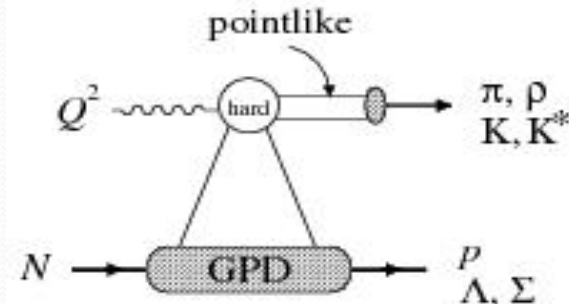
Area ( $q + \bar{q}$ )  $>$  Area (g)

Dynamical models predict difference: pion cloud, constituent quark picture

EIC: singlet quark size from deeply virtual compton scattering



EIC: strange and non-strange (sea) quark size from  $\pi$  and K production



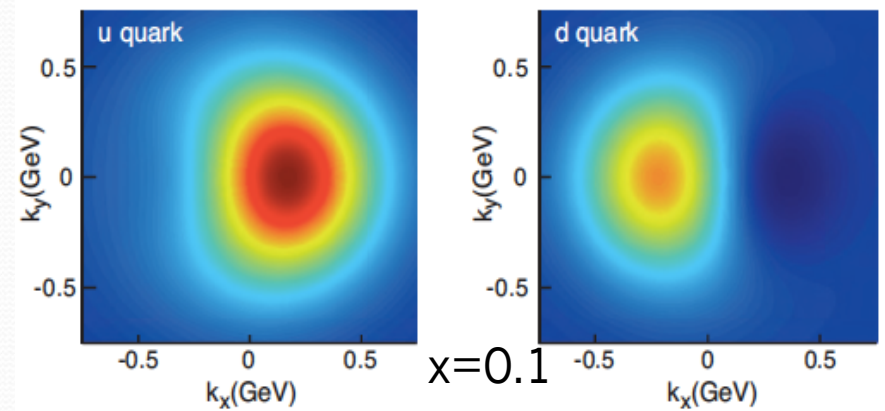


# ***GPD Study at EIC@HIAF***

- Unique opportunity for DVMP (pion/Kaon)  
flavor decomposition needs DVMP  
energy reach  $Q^2 > 5-10 \text{ GeV}^2$ , scaling region for exclusive light meson production  
(JLab12 energy not high enough to have clean light meson deep exclusive process)
- Significant increase in range for DVCS  
combination of energy and luminosity
- Other opportunities: vector meson, heavy flavors?

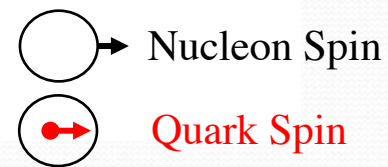
## 3-D Structure II

### Transverse Momentum-Dependent Distributions



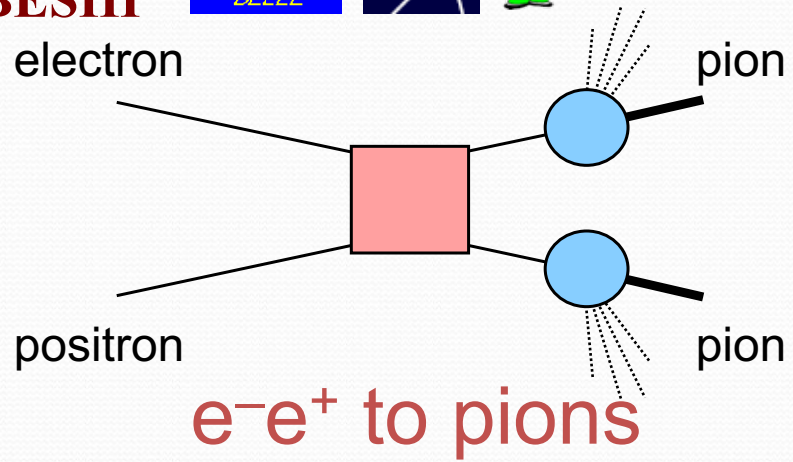
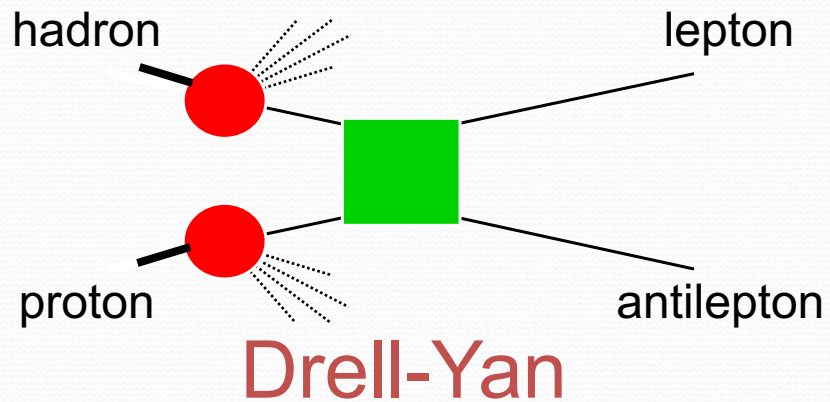
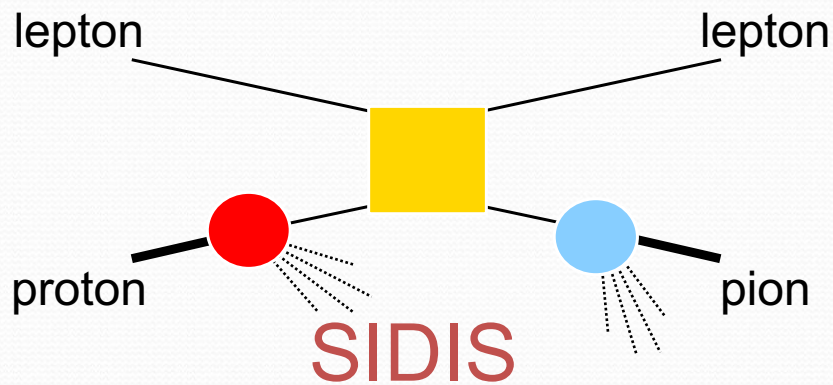


# Leading-Twist TMD PDFs



		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1$		$h_1^\perp$ Boer-Mulders
	L		$g_1$ Helicity	$h_{1L}^\perp$ Long-Transversity
	T	$f_{1T}^\perp$ Sivers	$g_{1T}$ Trans-Helicity	$h_1$ Transversity $h_{1T}^\perp$ Pretzelosity

# Access TMDs through Hard Processes

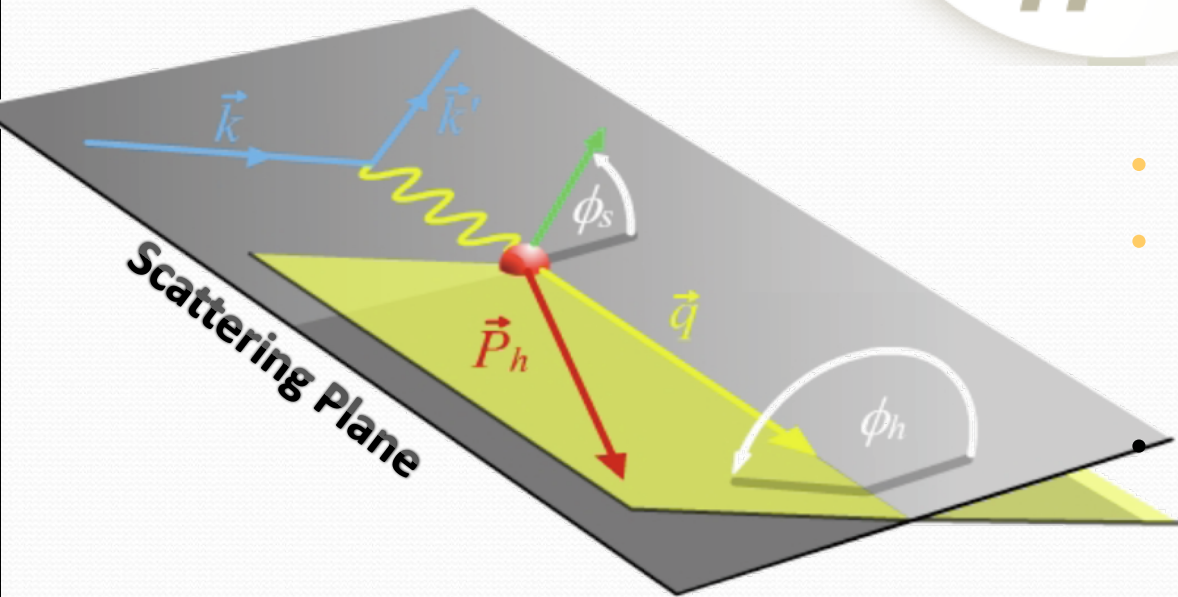
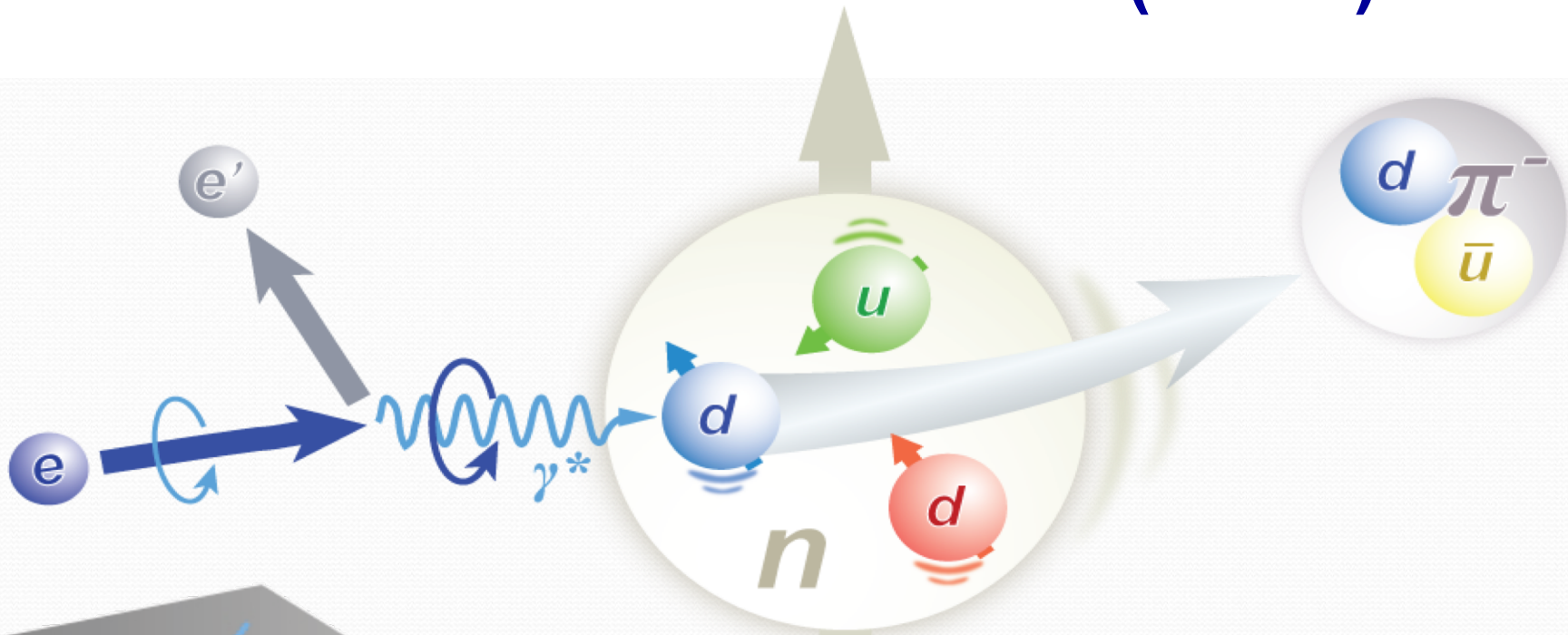


- Partonic scattering amplitude
- Fragmentation amplitude
- Distribution amplitude

$$f_{1T}^{\perp q}(\text{SIDIS}) = -f_{1T}^{\perp q}(\text{DY})$$

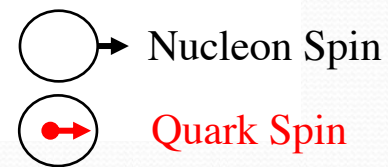










# Tool: Semi-inclusive DIS (SIDIS)



- Gold mine for TMDs
- Access all eight leading-twist TMDs through spin-comb. & azimuthal-modulations
- Tagging quark flavor/kinematics

# Leading-Twist TMD PDFs

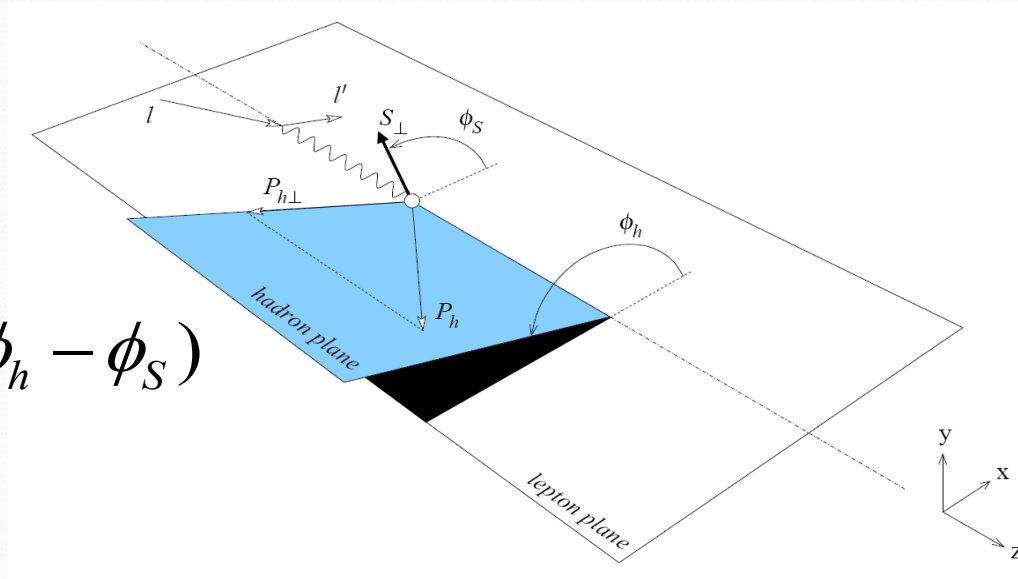


		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1$ 		$h_1^\perp$  <b>Boer-Mulders</b>
	L		$g_1$  <b>Helicity</b>	$h_{1L}^\perp$  <b>Long-Transversity</b>
	T	$f_{1T}^\perp$  <b>Sivers</b>	$g_{1T}$  <b>Trans-Helicity</b>	$h_1$  <b>Transversity</b> $h_{1T}^\perp$  <b>Pretzelosity</b>



# Separation of Collins, Sivers and pretzelosity effects through angular dependence

$$\begin{aligned}
 A_{UT}(\varphi_h^l, \varphi_S^l) &= \frac{1}{P} \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} \\
 &= A_{UT}^{\text{Collins}} \sin(\phi_h + \phi_S) + A_{UT}^{\text{Sivers}} \sin(\phi_h - \phi_S) \\
 &+ A_{UT}^{\text{Pretzelosity}} \sin(3\phi_h - \phi_S)
 \end{aligned}$$



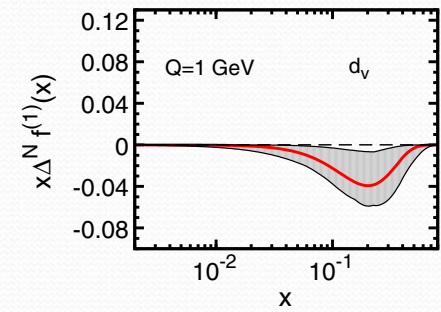
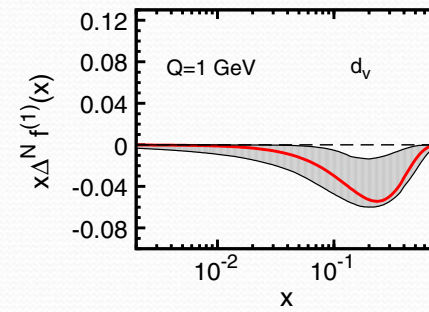
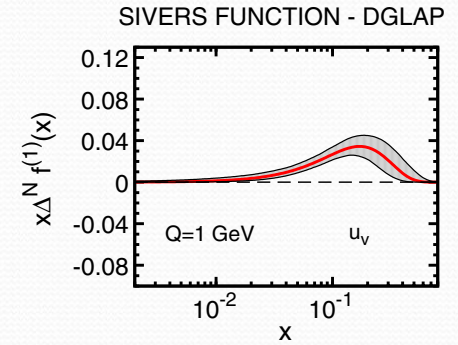
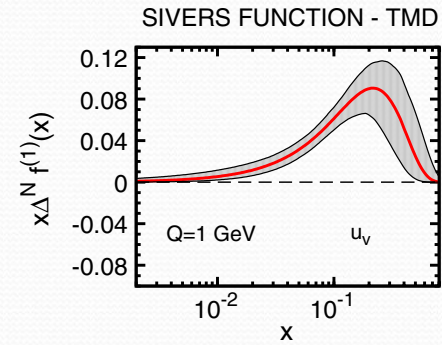
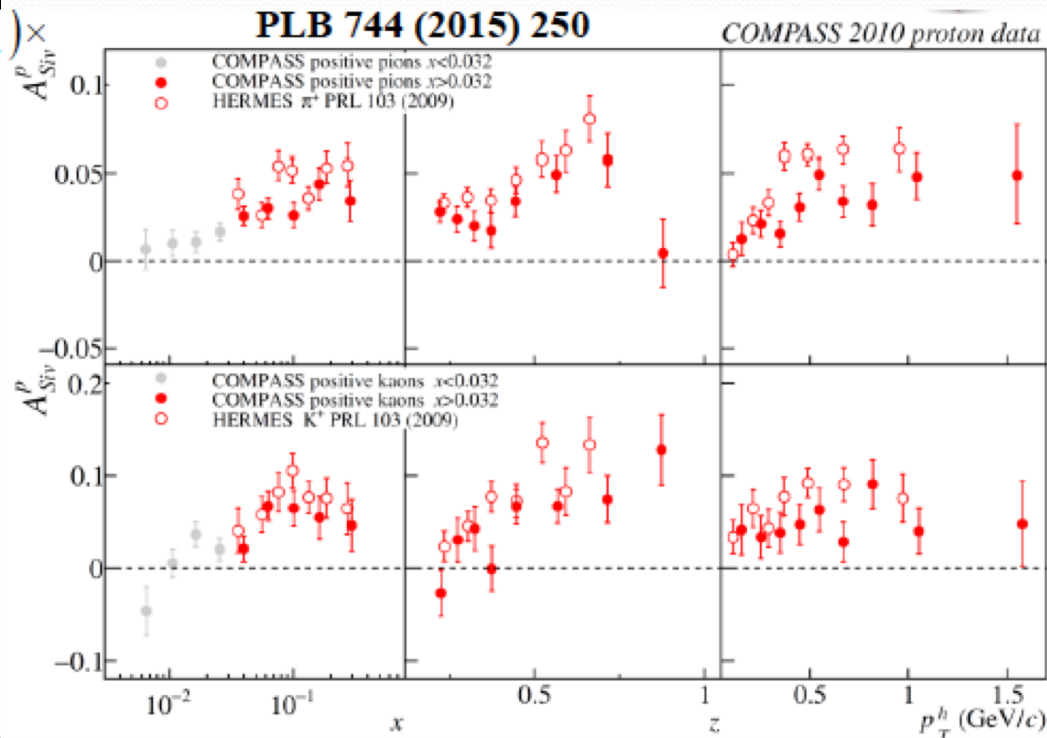
$$A_{UT}^{\text{Collins}} \propto \langle \sin(\phi_h + \phi_S) \rangle_{UT} \propto h_1 \otimes H_1^\perp$$

$$A_{UT}^{\text{Sivers}} \propto \langle \sin(\phi_h - \phi_S) \rangle_{UT} \propto f_{1T}^\perp \otimes D_1$$

$$A_{UT}^{\text{Pretzelosity}} \propto \langle \sin(3\phi_h - \phi_S) \rangle_{UT} \propto h_{1T}^\perp \otimes H_1^\perp$$



# COMPASS/HERMES: Sivers Asymmetries and Extraction of Sivers Function



Also other TMDs.

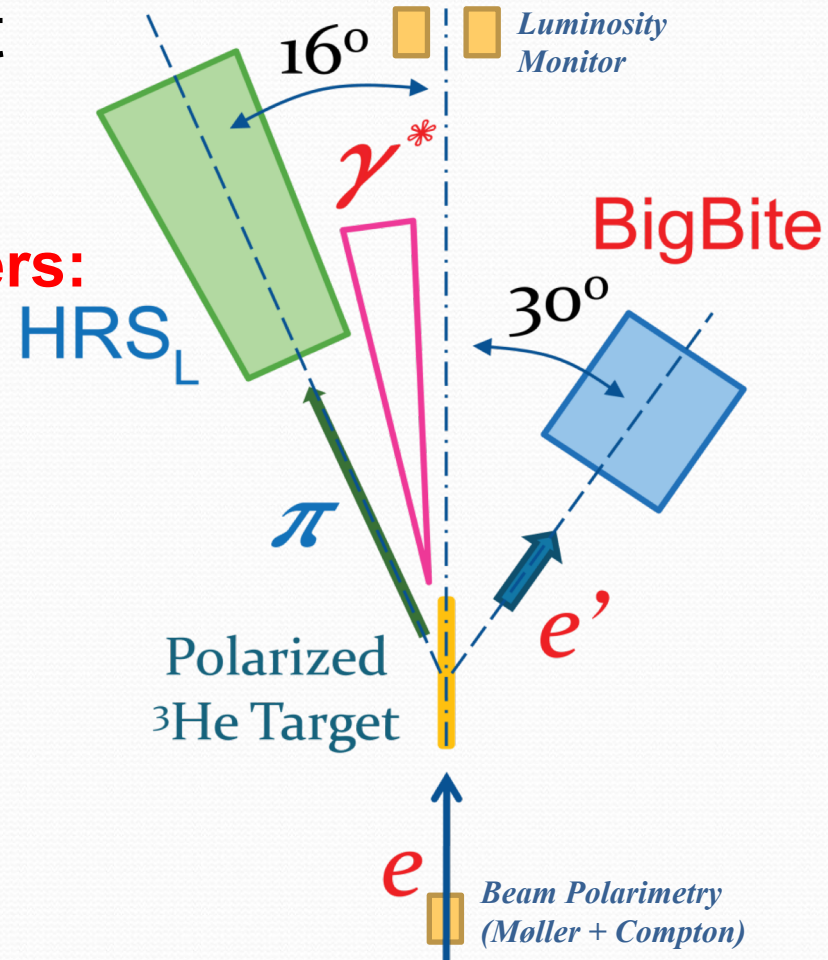
$N \backslash q$	U	L	T
U	$f_1$		$h_1^\perp$
L		$g_1$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1 \quad h_{1T}^\perp$

*M. Anselmino, M. Boglione, and S. Melis  
 Phys. Rev. D 86, 014028 (2012)*

# JLab 6 GeV Experiment E06-010

- **First measurement on n ( $^3\text{He}$ )**
- Transversely Polarized  $^3\text{He}$  Target
- Polarized Electron Beam, 5.9 GeV
- **Results published in 8 PRL/PRC papers:**
  - ✓  $\pi^{\pm}$  Collins/Sivers asymmetries: PRL 107:072003(2011)
  - ✓  $\pi^{\pm}$  worm-gear asymmetries: PRL 108, 052001 (2012)
  - ✓  $\pi^{\pm}$  pretzelosity asymmetries: PRC 90 5, 055209(2014)
  - ✓  $K^{\pm}$  Collins/Sivers asymmetries: PRC 90 5, 05520 (2014)
  - ✓ Inclusive hadron SSA: PRC 89, 042201 (2014)
  - ✓ Inclusive electron SSA: PRL 113, 022502 (2014)
  - ✓ Inclusive hadron DSA: PRC 92, 015207 (2015)
  - ✓  $\pi^{\pm}$  SIDIS cross sections: PRC 95, 035209 (2017)

$${}^3\text{He}^{\uparrow} (\vec{e}, e' \pi^{\pm}) X$$
$${}^3\text{He}^{\uparrow} (\vec{e}, e' K^{\pm}) X$$

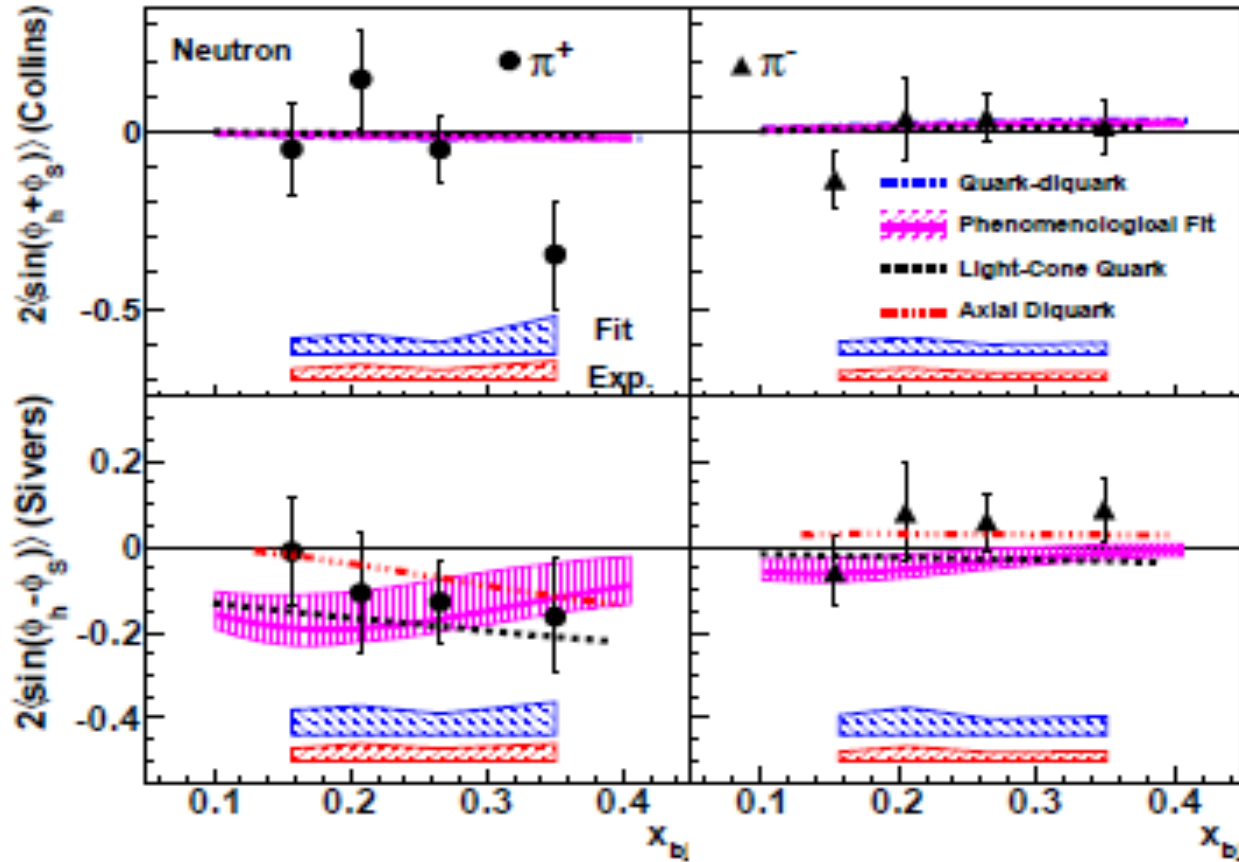




# $^3\text{He}$ (n) Target Single-Spin Asymmetry in SIDIS

E06-010 collaboration, X. Qian et al., PRL 107:072003(2011)

$$n^\uparrow(e, e'h), h = \pi^+, \pi^-$$



**Blue band:** model (fitting) uncertainties  
**Red band:** other systematic uncertainties

neutron Collins SSA small  
 Non-zero at highest x for  $\pi^+$

$Z \backslash q$	U	L	T
U	$f_1$		$h_1^\perp$
L		$g_1$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1^\perp, h_{1T}^\perp$

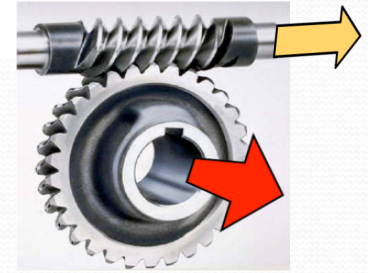
neutron Sivers SSA:  
 negative for  $\pi^+$ ,  
 Agree with Torino Fit

# Asymmetry $A_{LT}$ Result

J. Huang et al., PRL. 108, 052001 (2012).

To leading twist:

$$A_{LT}^{\cos(\phi_h - \phi_s)} \propto F_{LT}^{\cos(\phi_h - \phi_s)} \propto g_{1T}^q \otimes D_{1q}^h$$

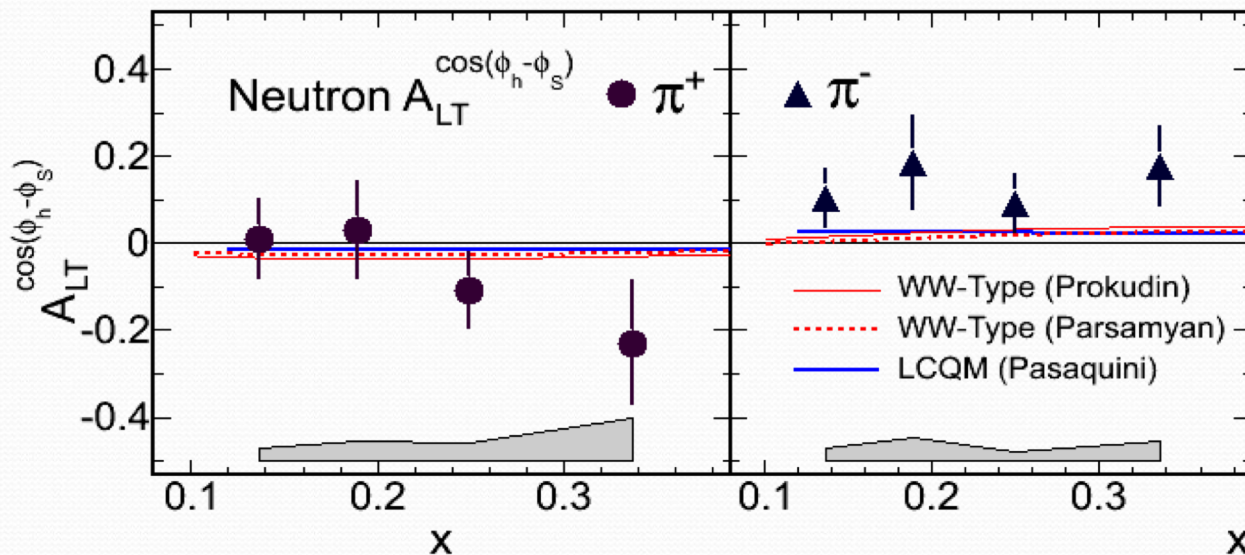


Worm-Gear  
Trans helicity

Dominated by  $L=0$  (S) and  $L=1$  (P) interference

- neutron  $A_{LT}$ : Positive for  $\pi^-$
- Consist w/ model in signs, suggest larger asymmetry

$Z \backslash q$	U	L	T
U	$f_1$		$h_1^\perp$
L		$g_1$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1, h_{1T}^\perp$



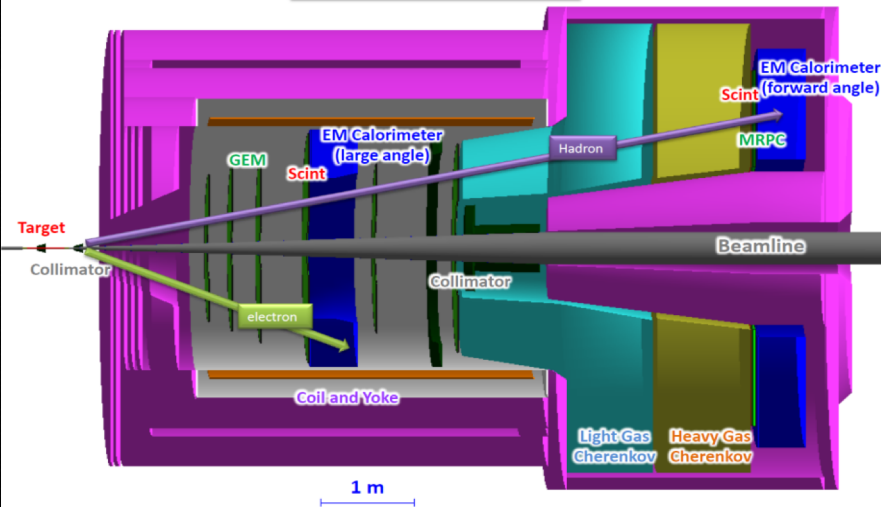


# ***Precision Study of TMDs: JLab 12 GeV, EIC***

- Explorations: HERMES, COMPASS, RHIC-spin, JLab6,...
- From exploration to **precision** study
  - JLab12: valence region; EIC: sea and gluons
- Transversity: fundamental *PDFs*, tensor charge
- *TMDs*: 3-d momentum structure of the nucleon
  - information on quark orbital angular momentum
  - information on QCD dynamics
- **Multi-dimensional** mapping of *TMDs*
- Precision → high statistics
  - **high luminosity and large acceptance**

# SoLID-Spin: SIDIS on $^3\text{He}$ /Proton @ 11 GeV

SoLID (SIDIS &  $J/\psi$ )

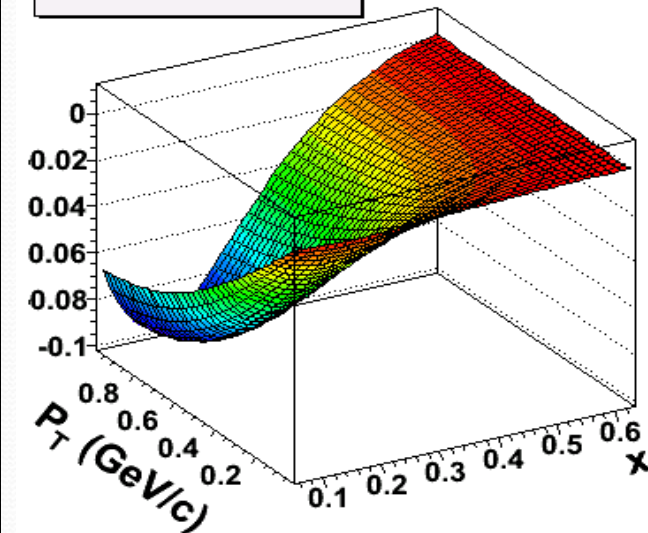


**E12-10-006:** Single Spin Asymmetry on Transverse  $^3\text{He}$ , **rating A**

**E12-11-007:** Single and Double Spin Asymmetries on  $^3\text{He}$ , **rating A**

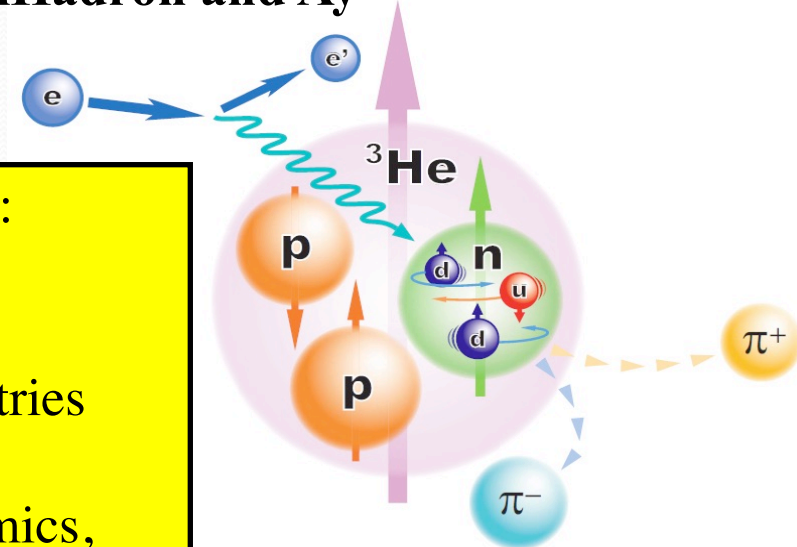
**E12-11-108:** Single and Double Spin Asymmetries on Transverse Proton, **rating A**

Sivers  $\pi^-$  @  $z = 0.55$



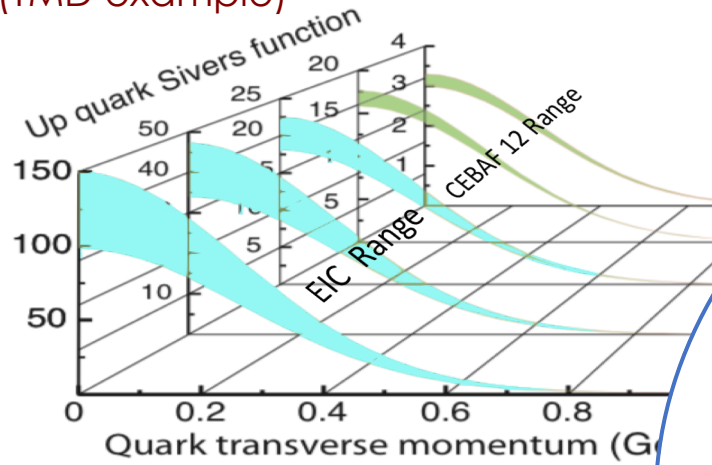
Two run group experiments DiHadron and  $A_y$

Key of SoLID-Spin program:  
 Large Acceptance  
 + High Luminosity  
 → 4-D mapping of asymmetries  
 → Tensor charge, TMDs ...  
 → Lattice QCD, QCD Dynamics, Models.

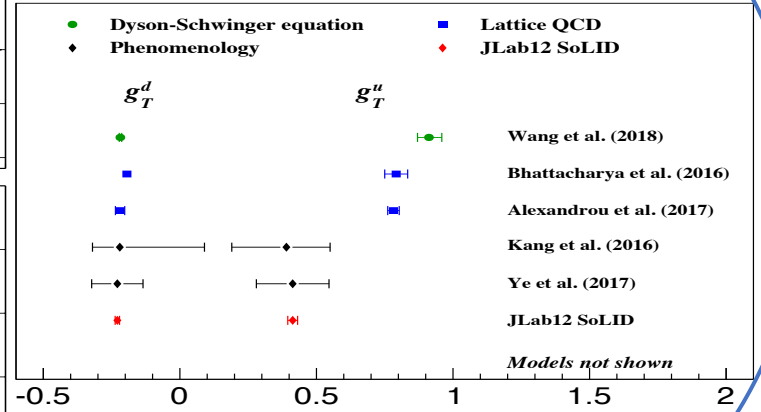
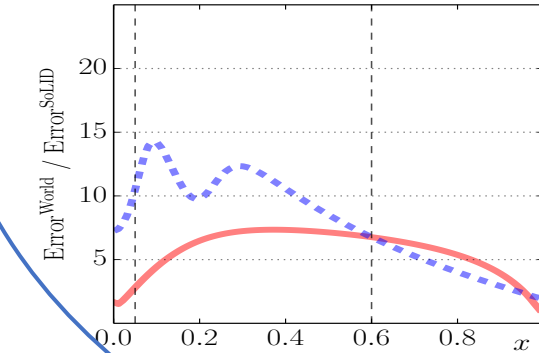
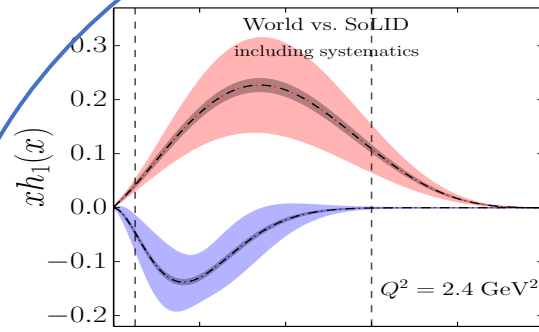
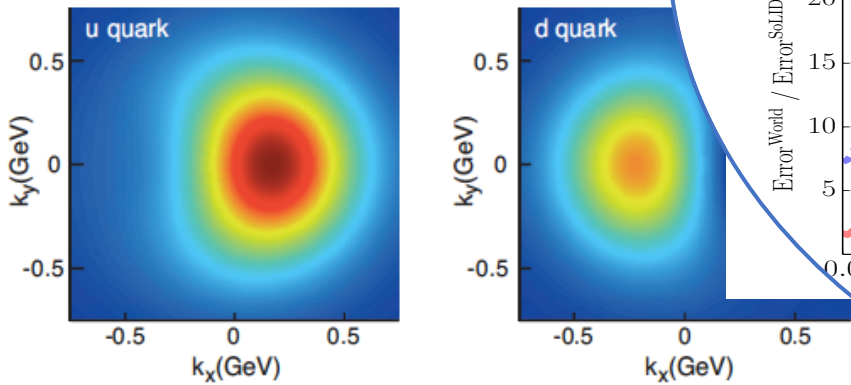


# SoLID and EIC: full imaging of nucleons and study QCD

(TMD example)



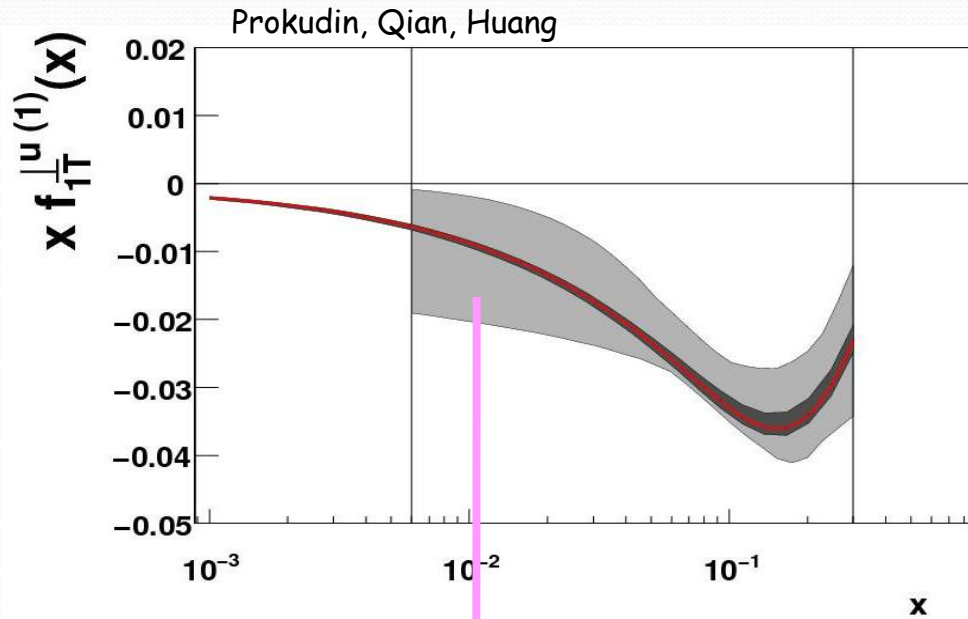
Polarized Quark 3D Momentum distributions



Transversity: valence quark effect

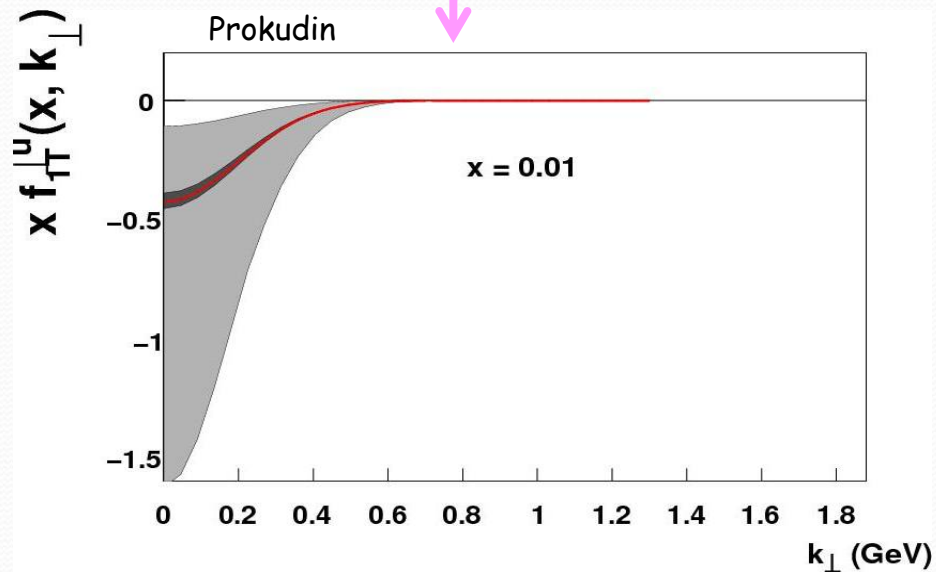


# Image the Transverse Momentum of the Quarks



Only a small subset of the  $(x, Q^2)$  landscape has been mapped here.

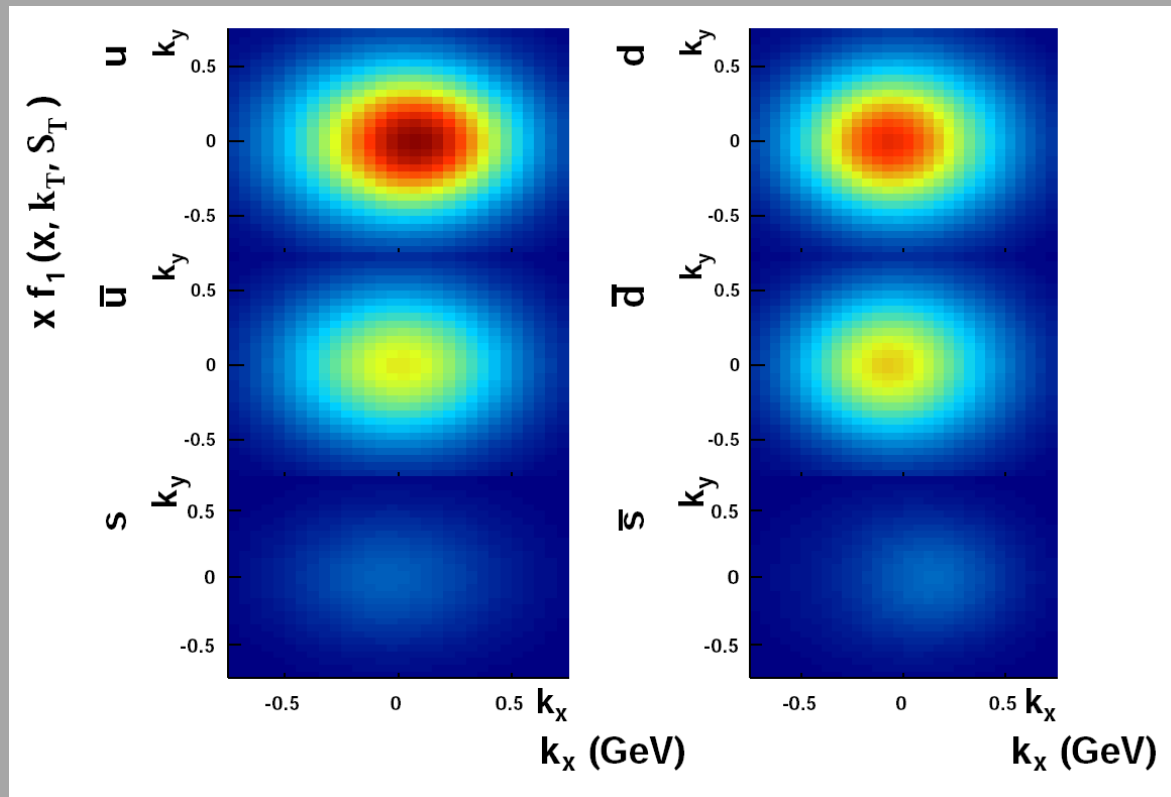
An EIC with good luminosity & high transverse polarization is the optimal tool to study this!



Exact  $k_{\perp}$  distribution presently essentially unknown!

# What do we learn from 3D distributions?

$$f(x, \mathbf{k}_T, \mathbf{S}_T) = f_1(x, \mathbf{k}_T^2) - f_{1T}^\perp(x, \mathbf{k}_T^2) \frac{\mathbf{k}_{T1}}{M}$$



The slice is at:

$$x = 0.1$$

Low- $x$  and high- $x$  region  
is uncertain

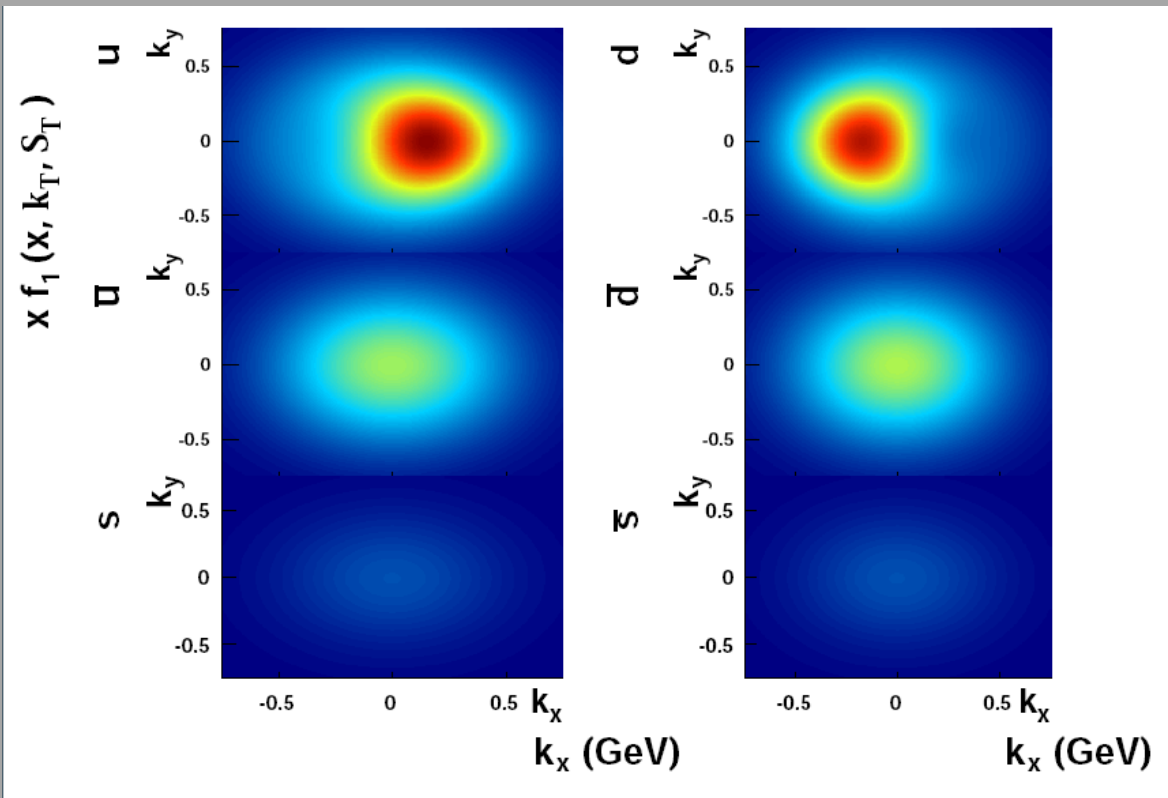
JLab 12 and EIC will  
contribute

No information on sea  
quarks

Picture is still quite  
uncertain

# What do we learn from 3D distributions?

$$f(x, \mathbf{k}_T, \mathbf{S}_T) = f_1(x, \mathbf{k}_T^2) - f_{1T}^\perp(x, \mathbf{k}_T^2) \frac{\mathbf{k}_{T1}}{M}$$



The slice is at:

$$x = 0.1$$

Low- $x$  and high- $x$  region  
is uncertain

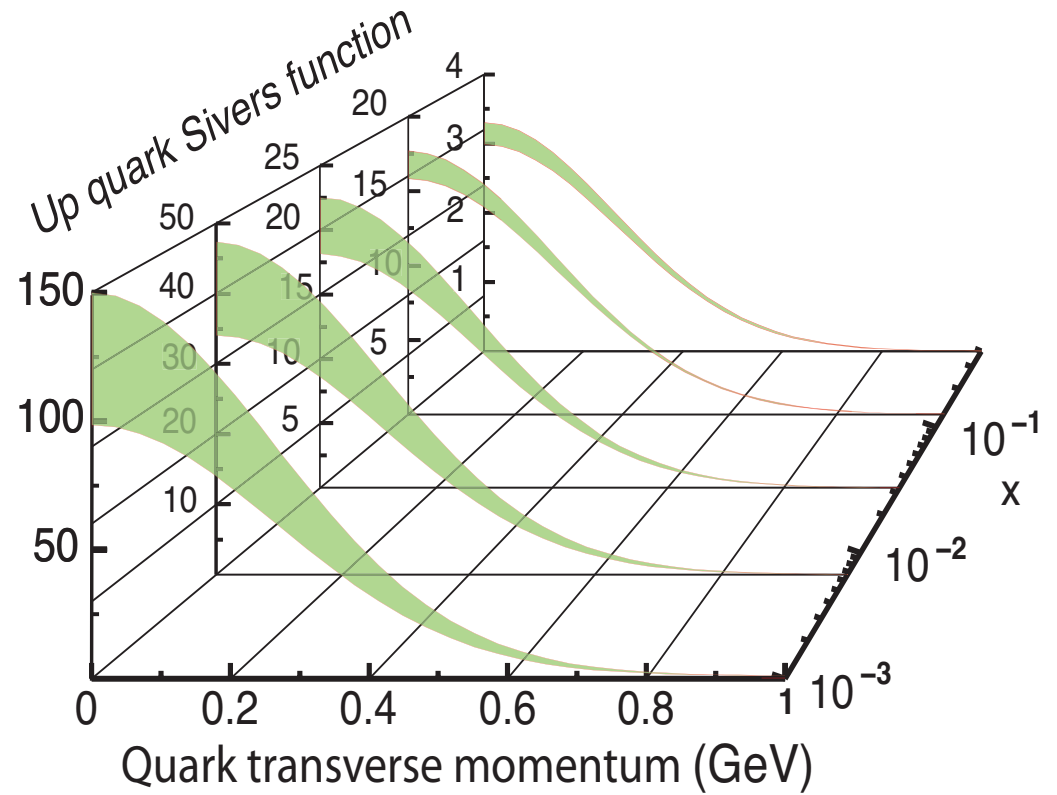
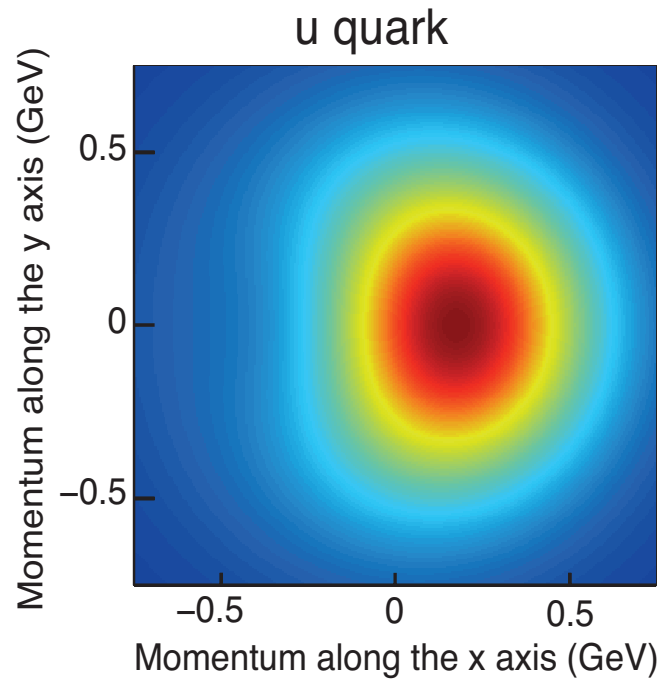
JLab 12 and EIC will  
contribute

No information on sea  
quarks

In future we will obtain  
much clearer picture



# EIC Imaging in 3-d momentum space



# Summary on TMD Program

- Exploratory results from 6 GeV neutron experiment
- **Unprecedented precision *multi-d* mapping of SSA in valence quark region with SoLID at 12 GeV JLab**
- Both polarized n ( $^3\text{He}$ ) and polarized proton
  - Three “A” rated experiments approved
  - + three run-group experiments
- Combining with the world data (fragmentation functions)
  - extract transversity for both  $u$  and  $d$  quarks
  - determine tensor charges  $\rightarrow$  LQCD
  - learn quark orbital motion and QCD dynamics
  - 3-d imaging
- Global efforts (experimentalists and theorists), global analysis
  - much better understanding of 3-d nucleon structure and QCD
- **Long-term future: EIC to map sea and gluon SSAs**

# Summary

- Nucleon Structure Study: Discoveries and Surprises
  - Understand strong interaction/nucleon structure: remains a challenge
- Highlights of JLab and EIC program
  - Precision EM form factors, proton radius
  - Nucleon spin-flavor structure (unpolarized and polarized, valence, sea)
  - 3-d Structure: GPDs
  - 3-d Structure: TMDs, SoLID program
- EIC opens up a new window to study/understand nucleon structure, especially the sea quarks and gluons

Exciting new opportunities → lead to breakthroughs?