Nucleon Tomography of Transverse Momentum Distributions

Feng Yuan Lawrence Berkeley National Laboratory





Semi-Inclusive Deep Inelastic Scattering

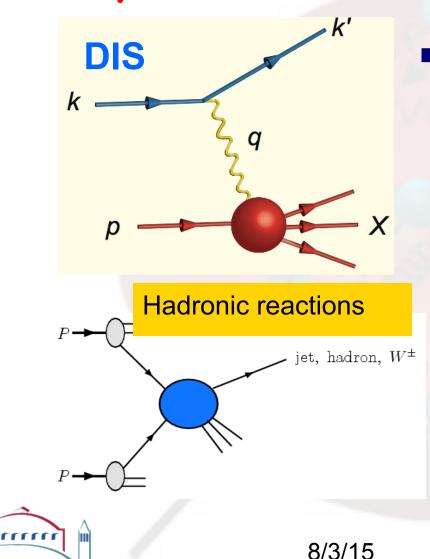
Feng Yuan RBRC, Brookhaven National Laboratory

Outline

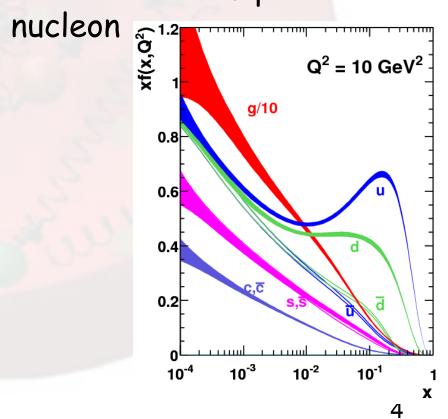
Basics Definitions, properties, ... QCD dynamics involved in these processes Universality, factorization, evolutions,... Applications Nucleon tensor charge TMDs at small-x



Feynman Parton: one-dimension



 Inclusive cross sections probe the momentum (longitudinal) distributions of partons inside

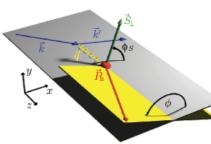


Extension to transverse direction...

- Semi-inclusive measurements (in DIS or Drell-Yan processes)
 - Transverse momentum dependent (TMD) parton distributions
- Deeply Virtual Compton Scattering and Exclusive processes
 - Generalized parton distributions (GPD)

Xiangdong's talk, this afternoon





TMD Parton Distribution:

The definition contains explicitly the gauge links

Collins-Soper 1981, Collins 2002, Belitsky-Ji-Yuan 2002

$$f(x,k_{\perp}) = \frac{1}{2} \int \frac{d\xi^{-} d^{2} \xi_{\perp}}{(2\pi)^{3}} e^{-i(\xi^{-} k^{+} - \vec{\xi}_{\perp} \cdot \vec{k}_{\perp})} \\ \times \langle PS | \overline{\psi}(\xi^{-},\xi_{\perp}) L_{\xi_{\perp}}^{\dagger}(\xi^{-}) \gamma^{+} L_{0}(0) \psi(0) | PS \rangle$$

QCD factorization has been proved for the hard processes in terms of TMDs Collins-Soper 1981, Ji-Ma-Yuan 2004, Collins 2011



Transverse momentum dependent parton distribution

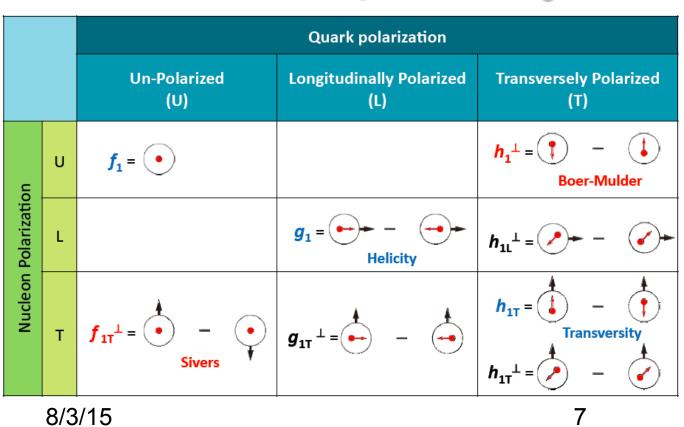
- Straightforward extension
 - Spin average, helicity, and transversity distributions
- P_{T} -spin correlations
 - Nontrivial distributions, S_TXP_T
- In quark model, depends on S- and P-wave

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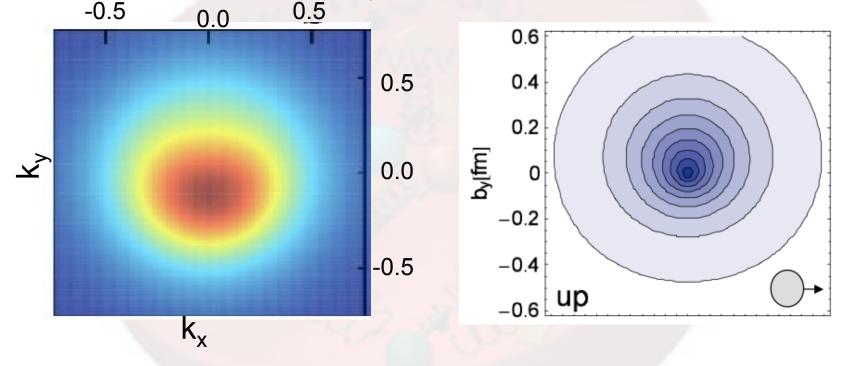
Leading Twist TMDs

- : Nucleon Spin





Deformation when nucleon is transversely polarized



Final State

nteractions

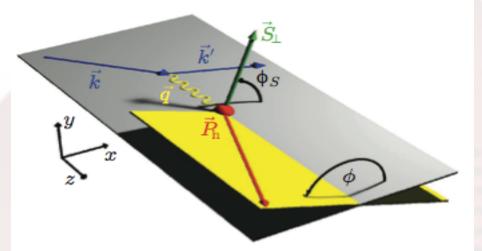
Burkardt's picture

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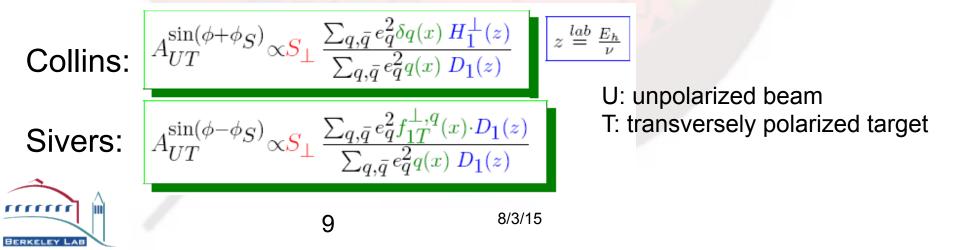
Quark Sivers function fit to the SIDIS Data, Anselmino, et al. 20009

Lattice Calculation of the tran. density of Up quark, QCDSF/UKQCD Coll., 2006

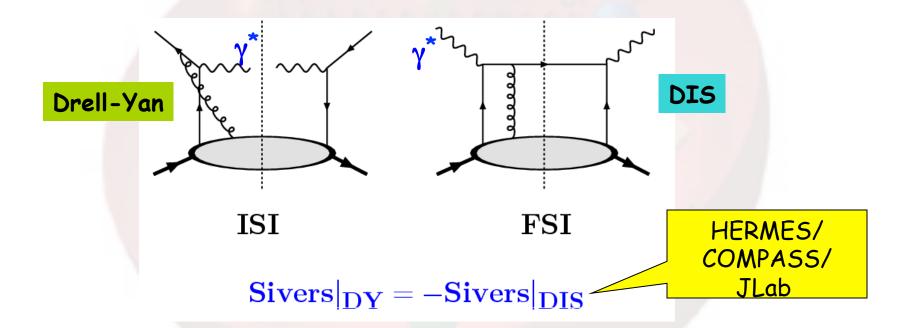
Collins and Sivers effects: Semi-inclusive DIS



Novel Single Spin Asymmetries



Sivers Asymmetries in DIS and Drell-Yan Initial state vs. final state interactions



"Universality": QCD prediction



TMD predictions rely on

- Non-perturbative TMDs constrained from experiments
- QCD evolutions, respect to the hard momentum scale Q
 - Strong theory/phenomenological efforts in the last few years
 - Iattice calculations may help as well



Collins-Soper-Sterman Resummation

Large Logs are resummed by solving the energy evolution equation of the TMDs

$$\frac{\partial}{\partial \ln Q} f(k_{\perp}, Q) = (K(q_{\perp}, \mu) + G(Q, \mu)) \otimes f(k_{\perp}, Q)$$

K and G obey the renormalization group eq.

$$\frac{\partial}{\partial \ln \mu} K = -\gamma_K = \frac{\partial}{\partial \ln \mu} G$$

Collins-Soper 81, Collins-Soper-Sterman 85



Solving the evolution equations

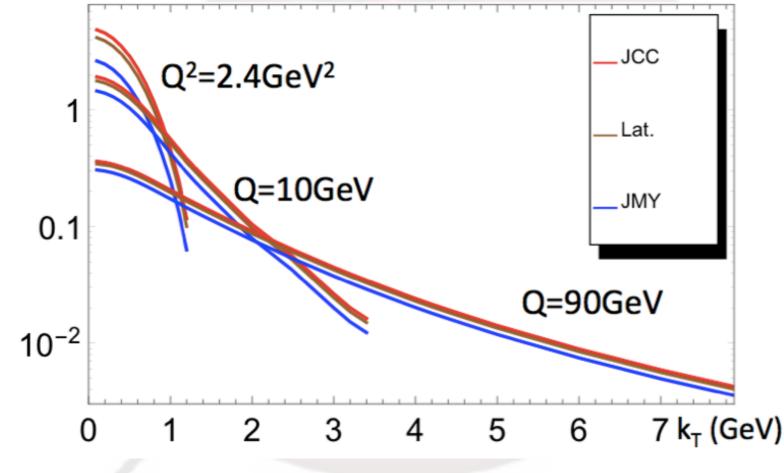
$$\begin{split} \widetilde{f}_{q}^{(sub.)}(x,b,\zeta^{2}=\rho Q^{2};\mu_{F}=Q) &= e^{-S_{pert}^{q}(Q,b_{*})-S_{NP}^{q}(Q,b)}\widetilde{\mathcal{F}}_{q}\left(\alpha_{s}(Q);\rho\right) \\ \text{Sudakov form factor (perturbative)} &\times \sum_{i} C_{q/i}(\mu_{b}/\mu) \otimes f_{i}(x,\mu) , \\ \text{Non-perturbative input} \\ \bullet \text{ Universal C-function} \\ C_{q/q'}(x) &= \delta_{qq'}\left[\delta(1-x) + \frac{\alpha_{s}}{2\pi}C_{F}(1-x)\right] \\ \bullet \text{ Scheme-dept.} & \widetilde{\mathcal{F}}_{q}^{\text{JCC}}\left(\alpha_{s}(Q)\right) = 1 + \mathcal{O}(\alpha_{s}^{2}) \\ \widetilde{\mathcal{F}}_{q}^{\text{JMY}}\left(\alpha_{s}(Q);\rho\right) = 1 + \frac{\alpha_{s}}{2\pi}C_{F}\left(\ln\rho - \frac{\ln^{2}\rho}{2} - \frac{\pi^{2}}{2} - 2\right) \\ \widetilde{\mathcal{F}}_{q}^{\text{Lat.}}\left(\alpha_{s}(Q)\right) = 1 + \frac{\alpha_{s}}{2\pi}C_{F}\left(-2\right) \end{split}$$



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Catani et al, 2000; Prokudin et al, 2015

Unpolarized quark distribution

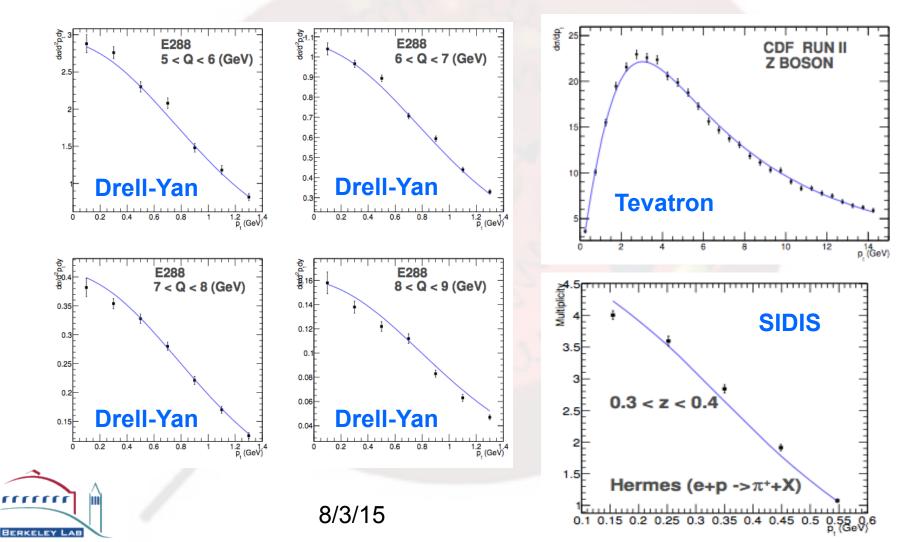


Prokudin-Sun-Yuan, 1505.05588



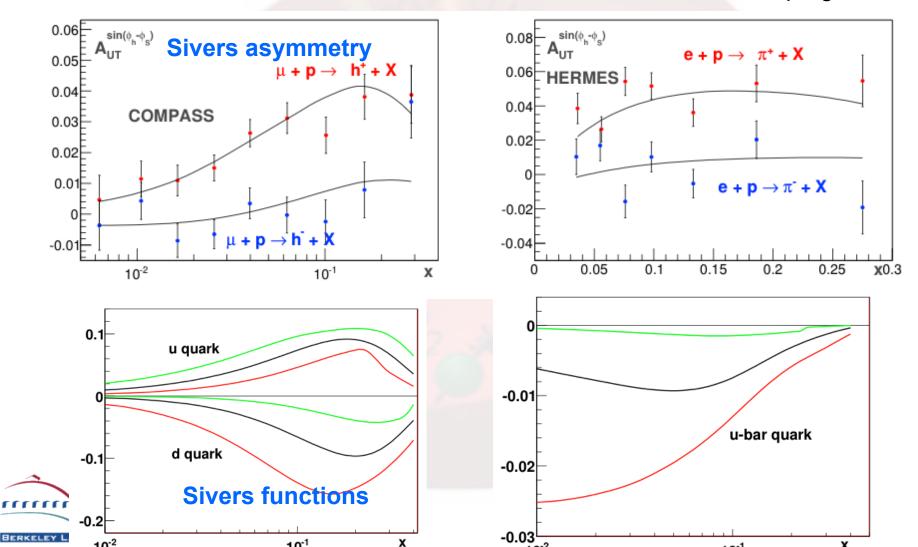
Describe well the exp. data

Sun-Issacson-Yuan-Yuan, 2014



Sivers asymetries in SIDIS with **Evolution** Sun, Yuan, PRD 2013

Prokudin-Sun-Yuan, in progress



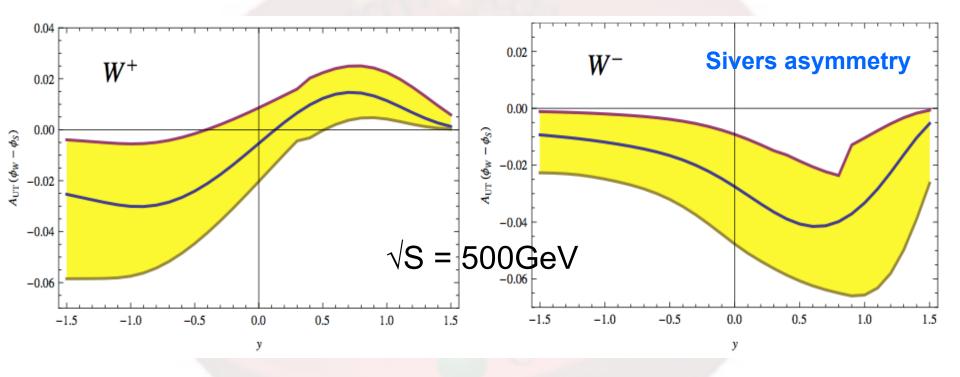
Complementary proposals for Drell-Yan experiments

- COMPASS@CERN
 - □ Kinematics (x_B) in the same range as existing SIDIS measurements
- SeaQuest@Fermilab
 - Polarized beam (valence) and target (sea)
- RHIC@BNL
 - □ W-boson, unique flavor dependence
 - □ Unique on the Q-evolution



Predictions at RHIC

Sun, Yuan, PRD 2013



Additional theory uncertainties: x-dependence of the TMDs comes from a fit to fixed target Drell-Yan and w/z production at Tevatron by Nadolsky et al.



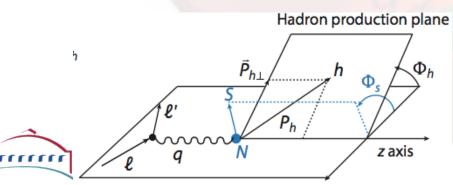
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Transversity and Collins FF

Kang-Prokudin-Sun-Yuan 2014

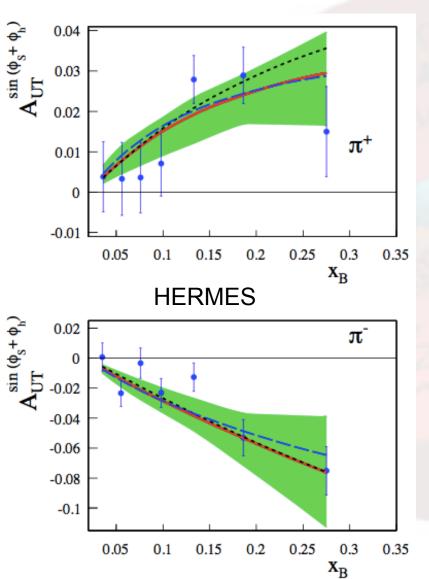
Lepton plane

- First comprehensive analysis with TMD evolution at next-to-leading logarithm
- Consistent treatment of both unpolarized and polarized cross sections
- Nucleon tensor charge determined with much theoretical improvement

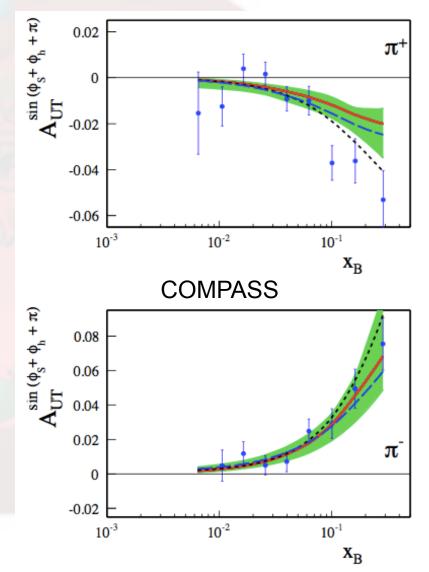


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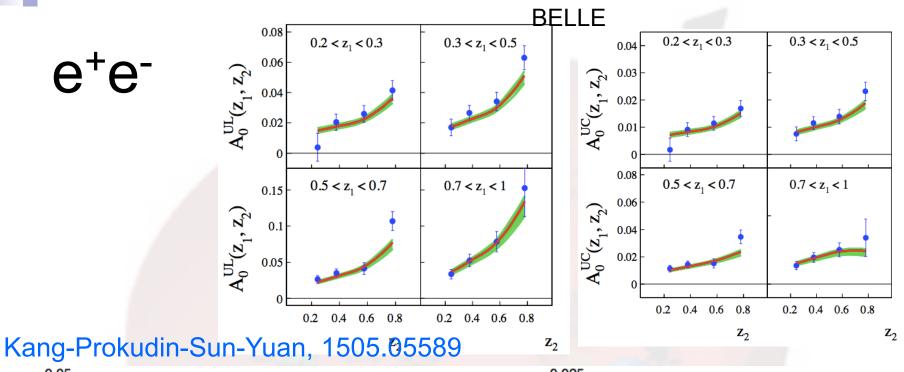
Fit to SIDIS

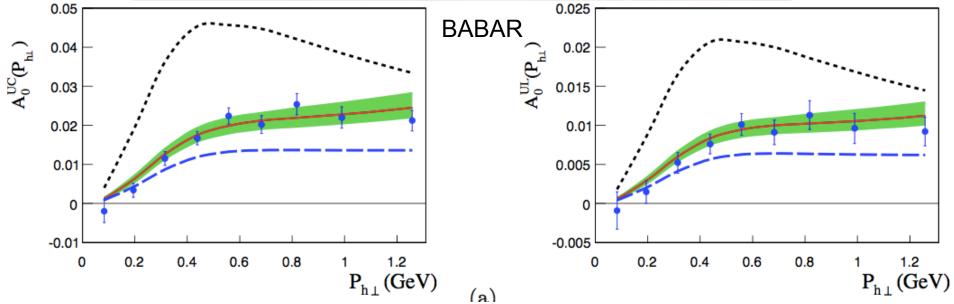


Kang-Prokudin-Sun-Yuan, 1505.05589



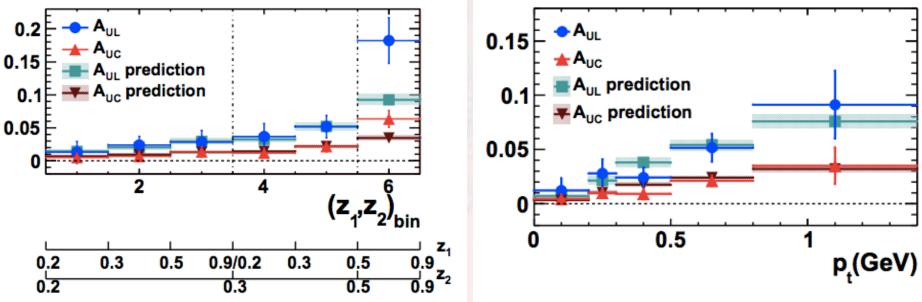
e⁺e⁻





BESIII @ BEPCII (Q=3.65GeV)

1507.06824

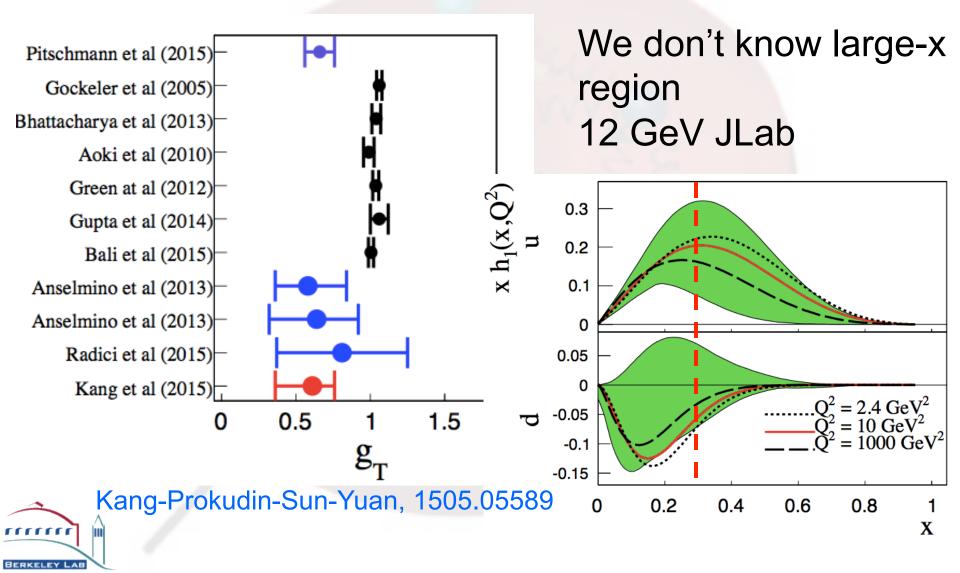


*Data not in the fit

Energy is between DIS and previous e+e- annihilation exp. Explore the QCD evolution!!



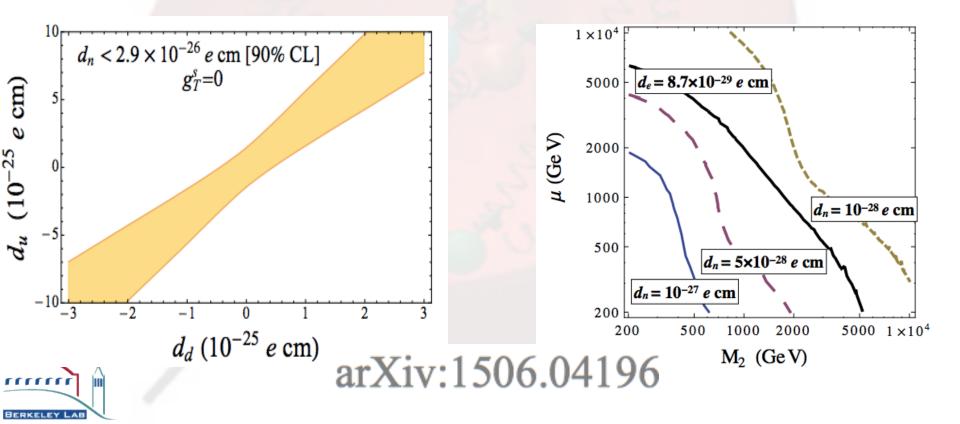
Nucleon tensor charge



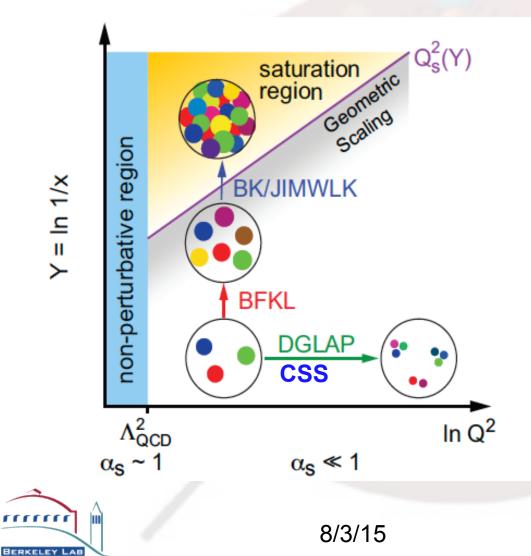
Implication for neutron EDM and new physics search

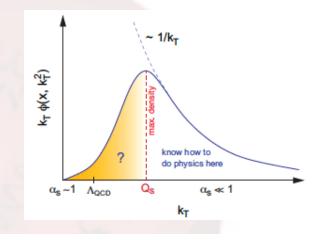
Neutron Electric Dipole Moment and Tensor Charges from Lattice QCD

Tanmoy Bhattacharya,^{1, *} Vincenzo Cirigliano,^{1, †} Rajan Gupta,^{1, ‡} Huey-Wen Lin,^{2, §} and Boram Yoon^{1, ¶}



TMDs at small-x





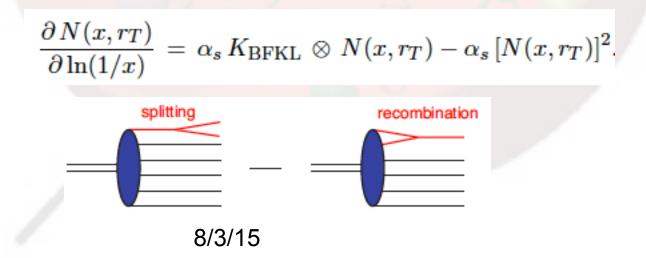
- Hard processes probe the kt-dependent gluon distributions directly
- Saturation phenomena manifest in the observables
 - □ Marquet, 2007
 - Dominguiz-Marquet-Xiao-Yuan, 2010

Non-linear term at high density

Balitsky-Fadin-Lipatov-Kuraev, 1977-78

 $\frac{\partial N(x, r_T)}{\partial \ln(1/x)} = \alpha_s K_{\text{BFKL}} \otimes N(x, r_T)$

Balitsky-Kovchegov: Non-linear term, 98



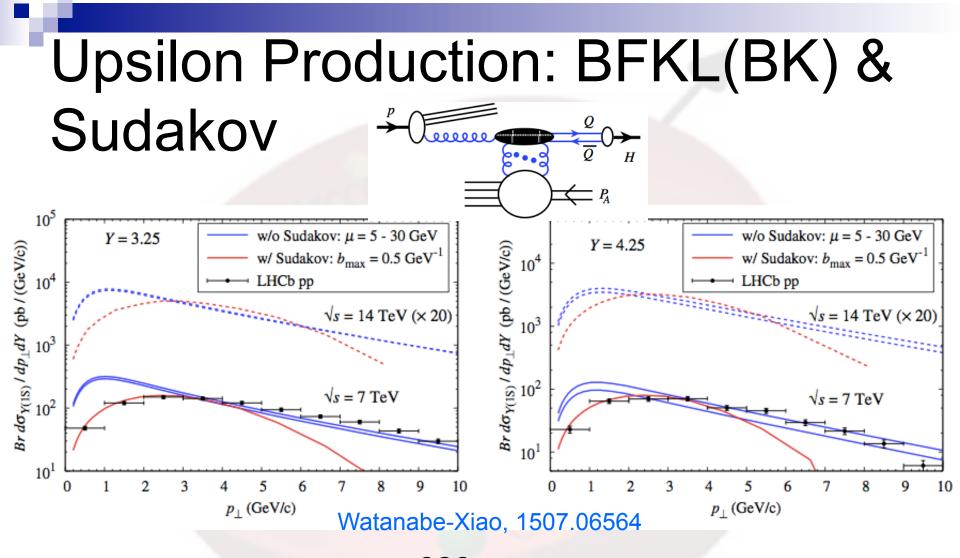
Therefore

- x-dependence of the TMDs at small-x, in principle, can be calculated from the QCD evolution (BK-JIMWLK)
- How about Q²

It has been shown that the Sudakov double log resummation (which controls Q-evolution) can be performed consistently in the small-x formalism

Mueller, Xiao, Yuan, PRL110,082301 (2013); Phys.Rev. D88 (2013) 114010 Kovchegov-Sievert, 1505.01176 Balitsky-Tarasov,,1505.02151





collinear + Sudakov: Berger-Qiu-Wang, hep-ph/0404158 Sun-Yuan-Yuan, 1210.3432

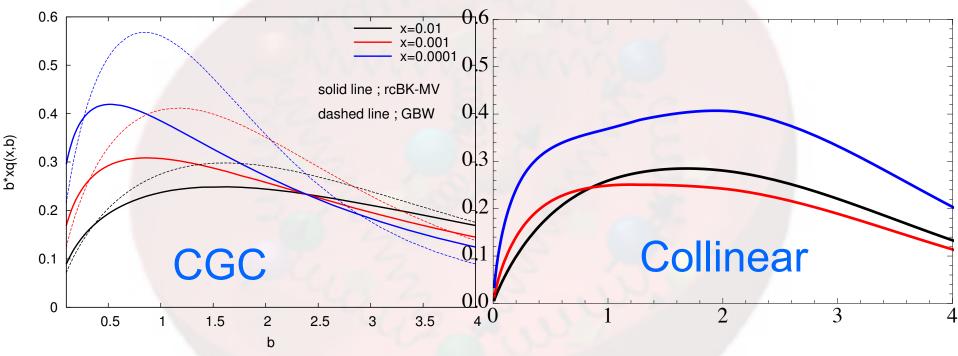
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CGC: Fujii-Watanabe, 1304.2221; Qiu-Sun-Xiao-Yuan, 1310.2230; Kang-Ma-Venugopalan, 1309.7337; Ma-Rajugopalan,1408.4075; Duclouse-Lappi-Mantysaari,1503.02789 28

Merging Framework: Spin and Small-x communities



 Realistic comparison will shed light on the TMD quarks at small-x (work in progress)

rrrr

Summary: What we have learned

- Unpolarized TMDs from, mainly, Drell-Yan, W/Z boson productions, partially from SIDIS processes
- Indications of polarized quark distributions (Sivers and cousins), from low energy SIDIS (HERMES, COMPASS, Jlab)



What we are missing

- Precision and detailed mapping of polarized quark/gluon distribution
 Universality/evolution more evident
- Spin correlation in momentum and coordinate space/tomography
- Systematics at small-x
- Jlab 12 and EIC are essential to answer these questions

All the talks yesterday!!

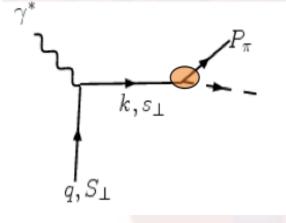


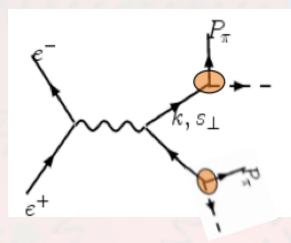
BACK-UP

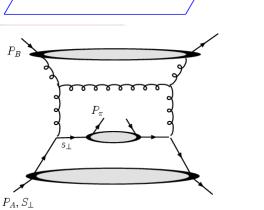




Universality of the Collins Fragmentation







ep--> e Pi X

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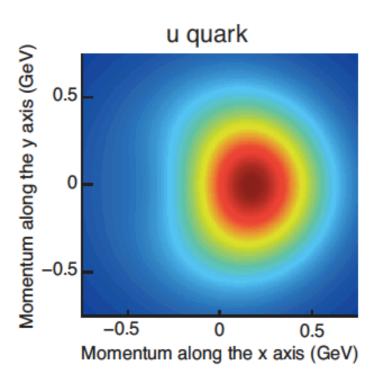
Metz 02, Collins-Metz 02, Yuan 07, Gamberg-Mukherjee-Mulders 08,10 Meissner-Metz 0812.3783 Yuan-Zhou, 0903.4680

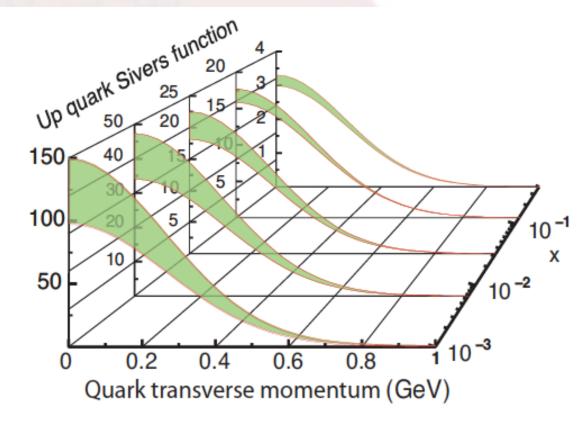
e⁺e⁻--> Pi Pi X

pp--> jet(->Pi) X

Exps: BELLE, HERMES, RHIC

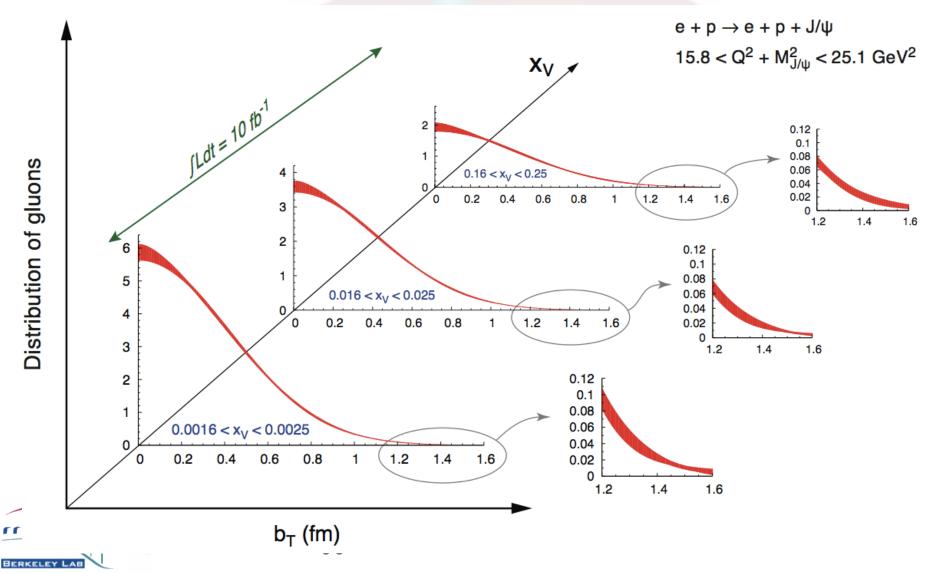
Kt-dependence



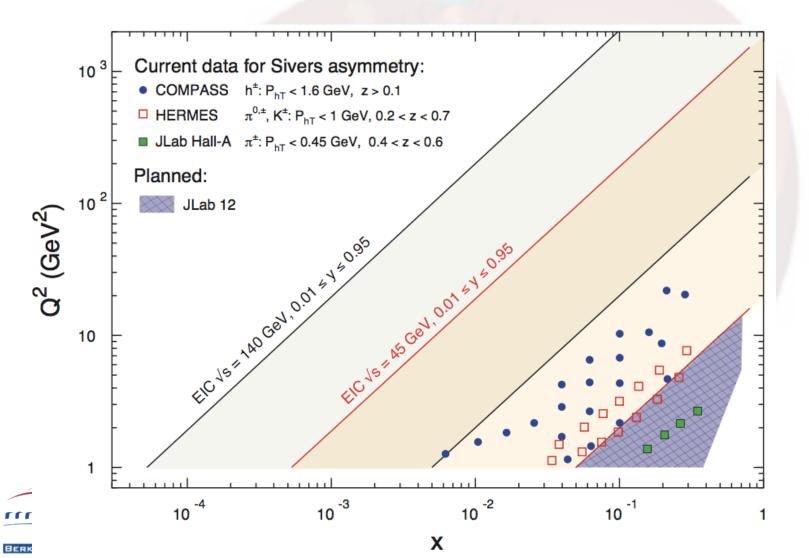




Tomography



EIC kinematics



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Computational access

Lattice calculations of the parton distributions (as functions of x) become possible by a recent proposal

$$egin{aligned} q(x,\mu^2,P^z) &= \int rac{dz}{4\pi} e^{izk^z} \langle P|\overline{\psi}(z)\gamma^z \ & imes \exp\left(-ig\int_0^z dz' A^z(z')
ight)\psi(0)|P
ight) \end{aligned}$$

Ji, PRL110 (2013) 262002



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TMDs on lattice

Ji-Sun-Xiong-Yuan, PRD 2015

$$egin{aligned} q(x_z,k_\perp;P_z) \ &= \ rac{1}{2} \int rac{d^3 z}{(2\pi)^3} e^{ik\cdot z} \langle PS | \overline{\psi}(0) \mathcal{L}^\dagger_{n_z(0,-\infty)} \gamma^z \mathcal{L}_{n_z(z,-\infty)} \psi(z) | PS
angle \ & imes \sqrt{rac{S^{n_x,n_y}(b_\perp)}{S^{n_x,n_z}(b_\perp) S^{n_z,n_y}(b_\perp)}} \ , \end{aligned}$$

Nontrivial soft factor subtraction is essential to achieve the factorization
 All unwanted divergences cancelled out
 Calculable on lattice

 shall be extended to small-x
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TMD factorization in quasi-pdfs

For example, for Drell-Yan lepton pair production

$$W(Q, b_{\perp}) = q^{sub.}(x_{z}, b_{\perp})\overline{q}^{sub.}(x'_{z}, b_{\perp})H(Q, \mu)$$

$$= q^{unsub.}(x_{z}, b_{\perp})\overline{q}^{unsub.}(x'_{z}, b_{\perp})H(Q, \mu)S(b_{\perp}, \mu)$$
Soft factor
Lattice calculation
Quasi-pdfs
$$\sqrt{\frac{S^{n_{x},n_{y}}(b_{\perp})}{S^{n_{x},n_{z}}(b_{\perp})S^{n_{z},n_{y}}(b_{\perp})}}$$

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$$39$$