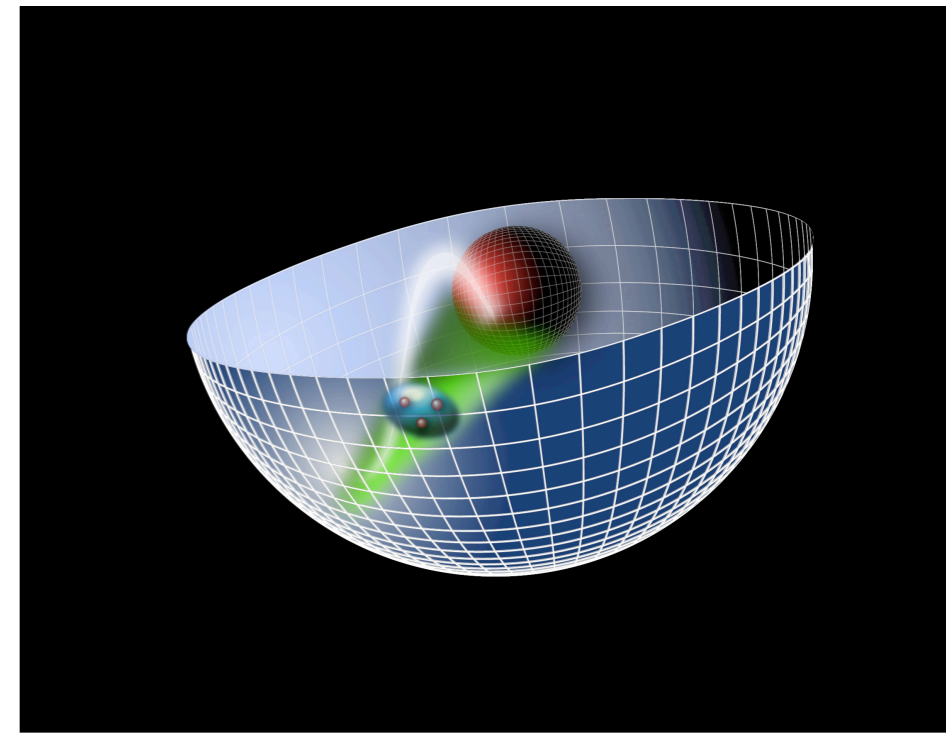
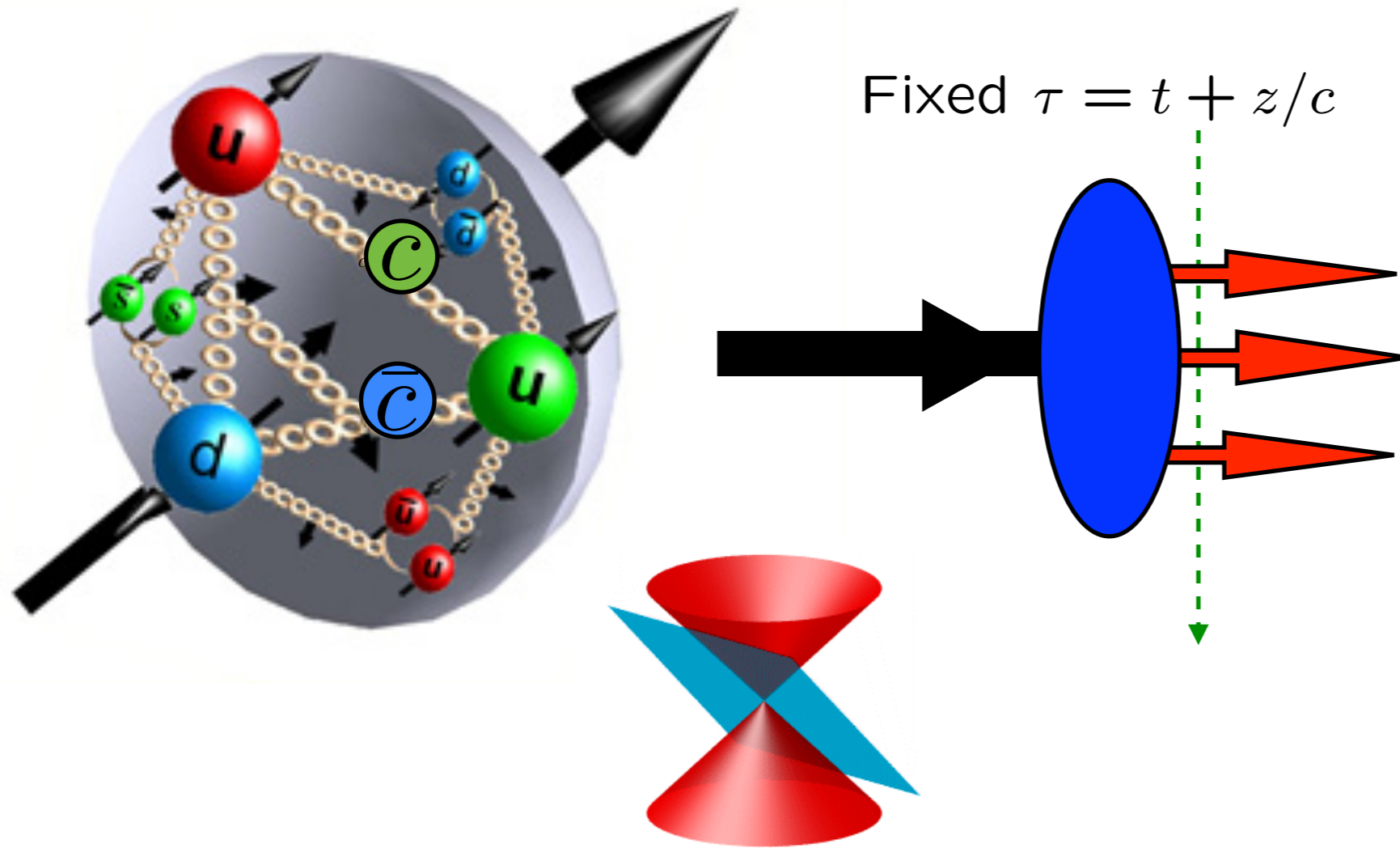


New Perspectives for Hadron Physics



*Kunshan, China
August 3, 2015*

Stan Brodsky



Goal: An analytic first approximation to QCD

- **As Simple as Schrödinger Theory in Atomic Physics**
- **Relativistic, Frame-Independent, Color-Confining**
- **Confinement in QCD -- What is the analytic form of the confining interaction?**
- **What sets the QCD mass scale?**
- **QCD Running Coupling at all scales**
- **Hadron Spectroscopy-Regge Trajectories**
- **Light-Front Wavefunctions**
- **Form Factors, Structure Functions, Hadronic Observables**
- **Constituent Counting Rules**
- **Hadronization at the Amplitude Level**
- **Insights into QCD Condensates**
- **Chiral Symmetry**
- **Systematically improvable**

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August 3, 2015



New Perspectives for Hadron Physics

Stan Brodsky

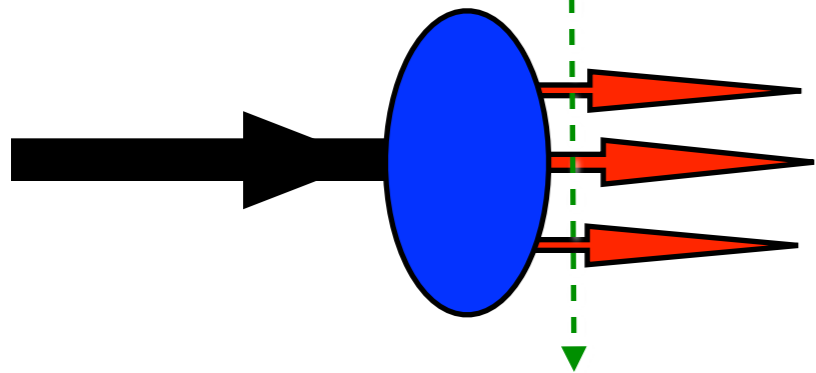
SLAC
NATIONAL ACCELERATOR LABORATORY

Bound States in Relativistic Quantum Field Theory:

Light-Front Wavefunctions

Dirac's Front Form: Fixed $\tau = t + z/c$

Fixed $\tau = t + z/c$



$$\psi(x_i, \vec{k}_{\perp i}, \lambda_i)$$

$$x = \frac{k^+}{P^+} = \frac{k^0 + k^3}{P^0 + P^3}$$

Invariant under boosts. Independent of P^μ

$$H_{LF}^{QCD} |\psi\rangle = M^2 |\psi\rangle$$

Direct connection to QCD Lagrangian

Off-shell in invariant mass

Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space

Each element of
flash photograph
illuminated
at same LF time

$$\tau = t + z/c$$

Causal, frame-independent

Evolve in LF time

$$P^- = i \frac{d}{d\tau}$$

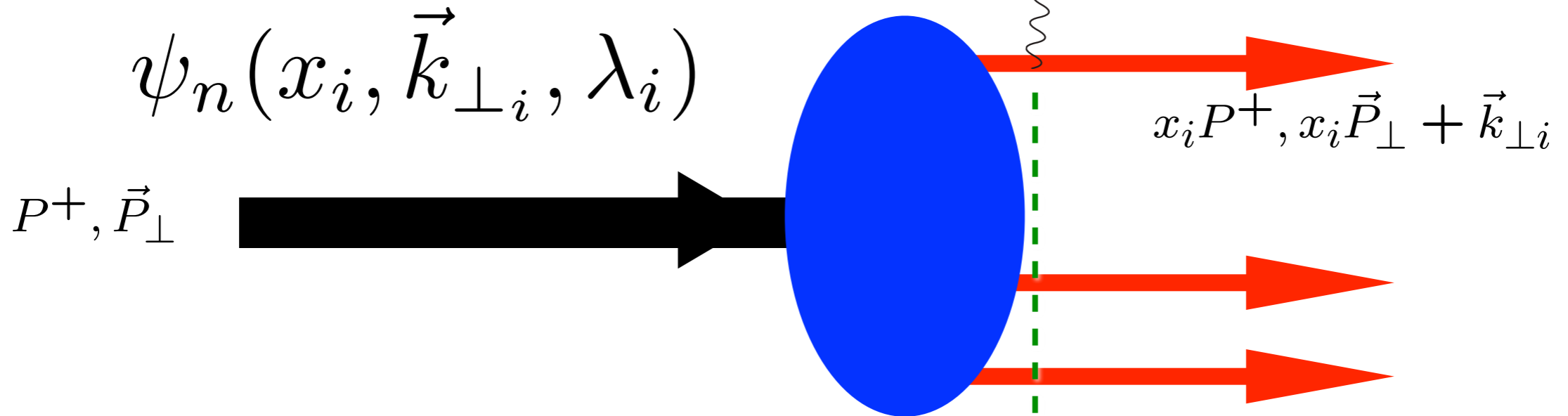
Eigenstate -- independent of τ

$$H_{LF} = P^+ P^- - \vec{P}_\perp^2$$

$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$



$$x = \frac{k^+}{P^+} = \frac{k^0 + k^3}{P^0 + P^3}$$



Measurements of hadron LF wavefunction are at fixed LF time

Like a flash photograph

$$x_{bj} = x = \frac{k^+}{P^+}$$

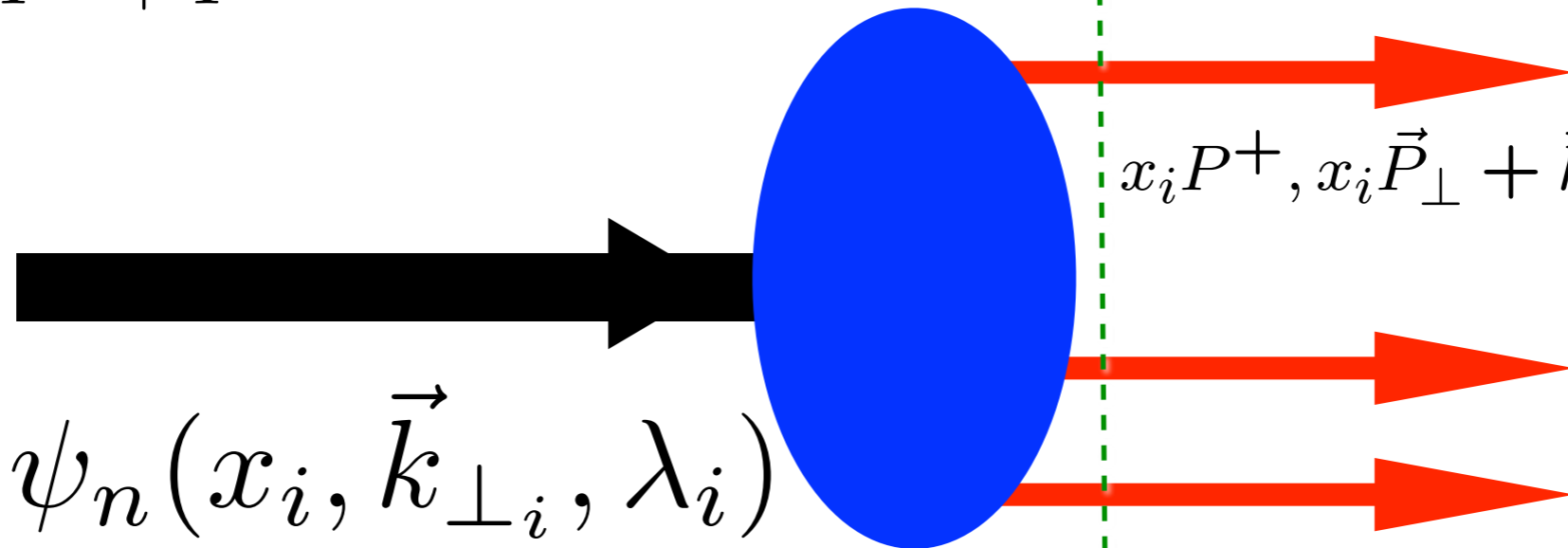
Light-Front Wavefunctions: **rigorous** representation of composite systems in quantum field theory

Eigenstate of LF Hamiltonian

$$x = \frac{k^+}{P^+} = \frac{k^0 + k^3}{P^0 + P^3}$$

Fixed $\tau = t + z/c$

P^+, \vec{P}_\perp



$x_i P^+, x_i \vec{P}_\perp + \vec{k}_{\perp i}$

$$\sum_i^n x_i = 1$$

$$\sum_i^n \vec{k}_{\perp i} = \vec{0}_\perp$$

$$\psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

$$\int \psi_{BS}(p, k) dk^- \rightarrow \psi_{LF}$$

$$|p, J_z \rangle = \sum_{n=3} \psi_n(x_i, \vec{k}_{\perp i}, \lambda_i) |n; x_i, \vec{k}_{\perp i}, \lambda_i \rangle$$

Invariant under boosts! Independent of P^μ

Causal, Frame-independent. Creation Operators on Simple Vacuum, Current Matrix Elements are Overlaps of LFWFS

$$\langle p + q | j^+(0) | p \rangle = 2p^+ F(q^2)$$

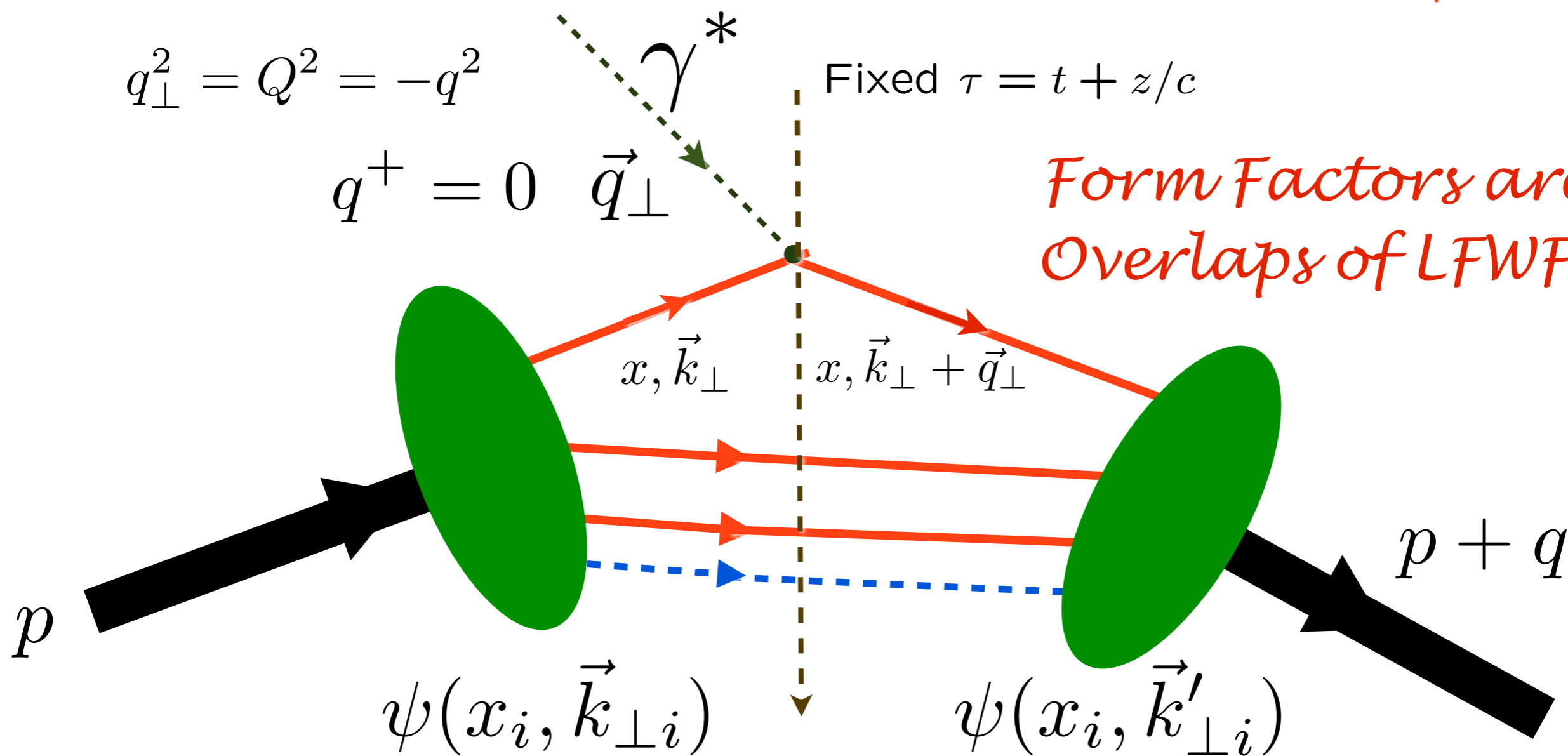
Interaction picture

$$q_{\perp}^2 = Q^2 = -q^2$$

$$q^+ = 0 \quad \vec{q}_{\perp}$$

Fixed $\tau = t + z/c$

Form Factors are Overlaps of LFWFs



$$\psi(x_i, \vec{k}_{\perp i})$$

$$\psi(x_i, \vec{k}'_{\perp i})$$

struck $\vec{k}'_{\perp i} = \vec{k}_{\perp i} + (1 - x_i)\vec{q}_{\perp}$

spectators $\vec{k}'_{\perp i} = \vec{k}_{\perp i} - x_i\vec{q}_{\perp}$

**Drell & Yan, West
Exact LF formula!**

No comparable formula in instant form

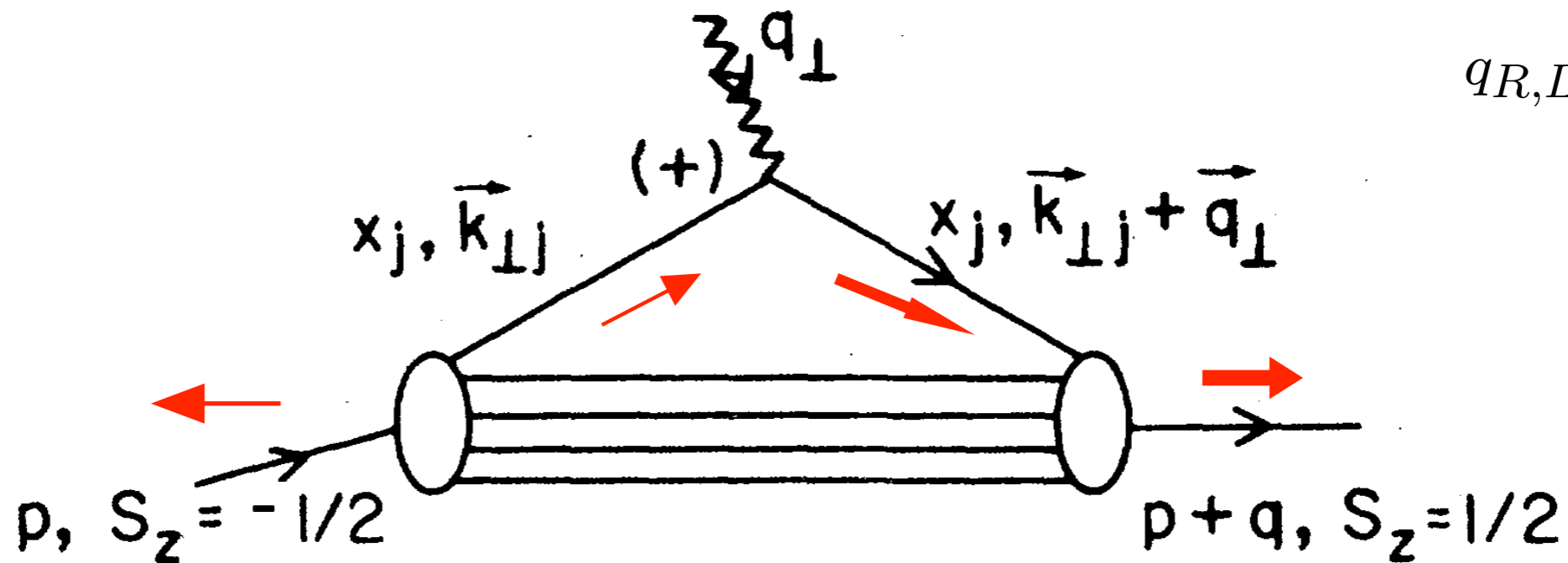
Exact LF Formula for Pauli Form Factor

$$\frac{F_2(q^2)}{2M} = \sum_a \int [dx][d^2\mathbf{k}_\perp] \sum_j e_j \frac{1}{2} \times$$

Drell, sjb

$$\left[-\frac{1}{q^L} \psi_a^{\uparrow*}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \psi_a^\downarrow(x_i, \mathbf{k}_{\perp i}, \lambda_i) + \frac{1}{q^R} \psi_a^{\downarrow*}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \psi_a^\uparrow(x_i, \mathbf{k}_{\perp i}, \lambda_i) \right]$$

$$\mathbf{k}'_{\perp i} = \mathbf{k}_{\perp i} - x_i \mathbf{q}_\perp \qquad \mathbf{k}'_{\perp j} = \mathbf{k}_{\perp j} + (1 - x_j) \mathbf{q}_\perp$$



$$q_{R,L} = q^x \pm iq^y$$

Must have $\Delta l_z = \pm 1$ to have nonzero $F_2(q^2)$

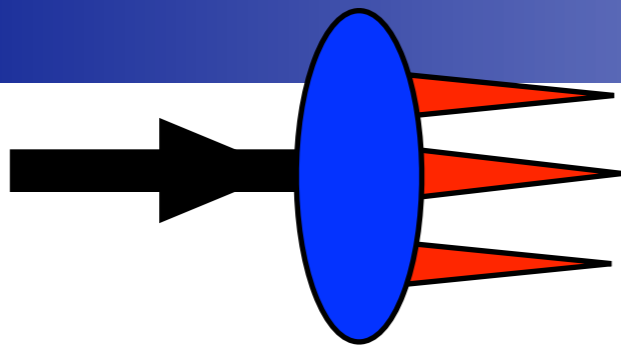
*Nonzero Proton Anomalous Moment -->
Nonzero orbital quark angular momentum*

Advantages of the Dirac's Front Form for Hadron Physics

- **Measurements are made at fixed τ**
- **Causality is automatic**
- **Structure Functions are squares of LFWFs**
- **Form Factors are overlap of LFWFs**
- **LFWFs are frame-independent -- no boosts!**
- **No dependence on observer's frame**
- **LF Holography: Dual to AdS space**
- **LF Vacuum trivial -- no condensates!**
- **Profound implications for Cosmological Constant**



• *Light Front Wavefunctions:*



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

GTMDs

$$x, \vec{k}_{\perp}, \vec{b}_{\perp}$$

Momentum space $\vec{k}_{\perp} \leftrightarrow \vec{z}_{\perp}$ Position space
 $\vec{\Delta}_{\perp} \leftrightarrow \vec{b}_{\perp}$

Transverse density in momentum space

Transverse density in position space

TMDs

$$x, \vec{k}_{\perp}$$

TMFFs

$$\vec{k}_{\perp}, \vec{b}_{\perp}$$

GPDs

$$x, \vec{b}_{\perp}$$

TMSDs

$$\vec{k}_{\perp}$$

PDFs

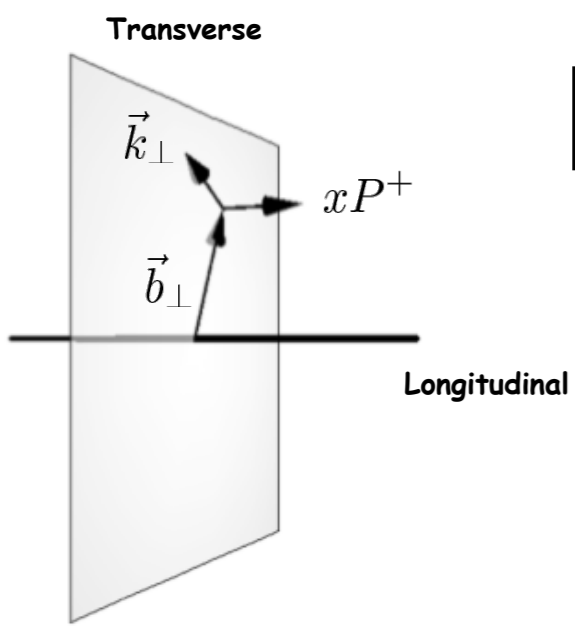
$$x,$$

FFs

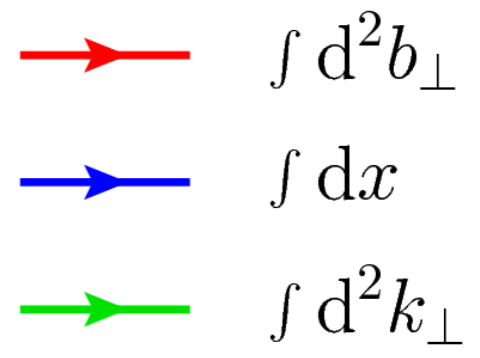
$$\vec{b}_{\perp}$$

Charges

Lorce, Pasquini



Sivers, T-odd from lensing



Single-spin asymmetries

Leading Twist Sivers Effect

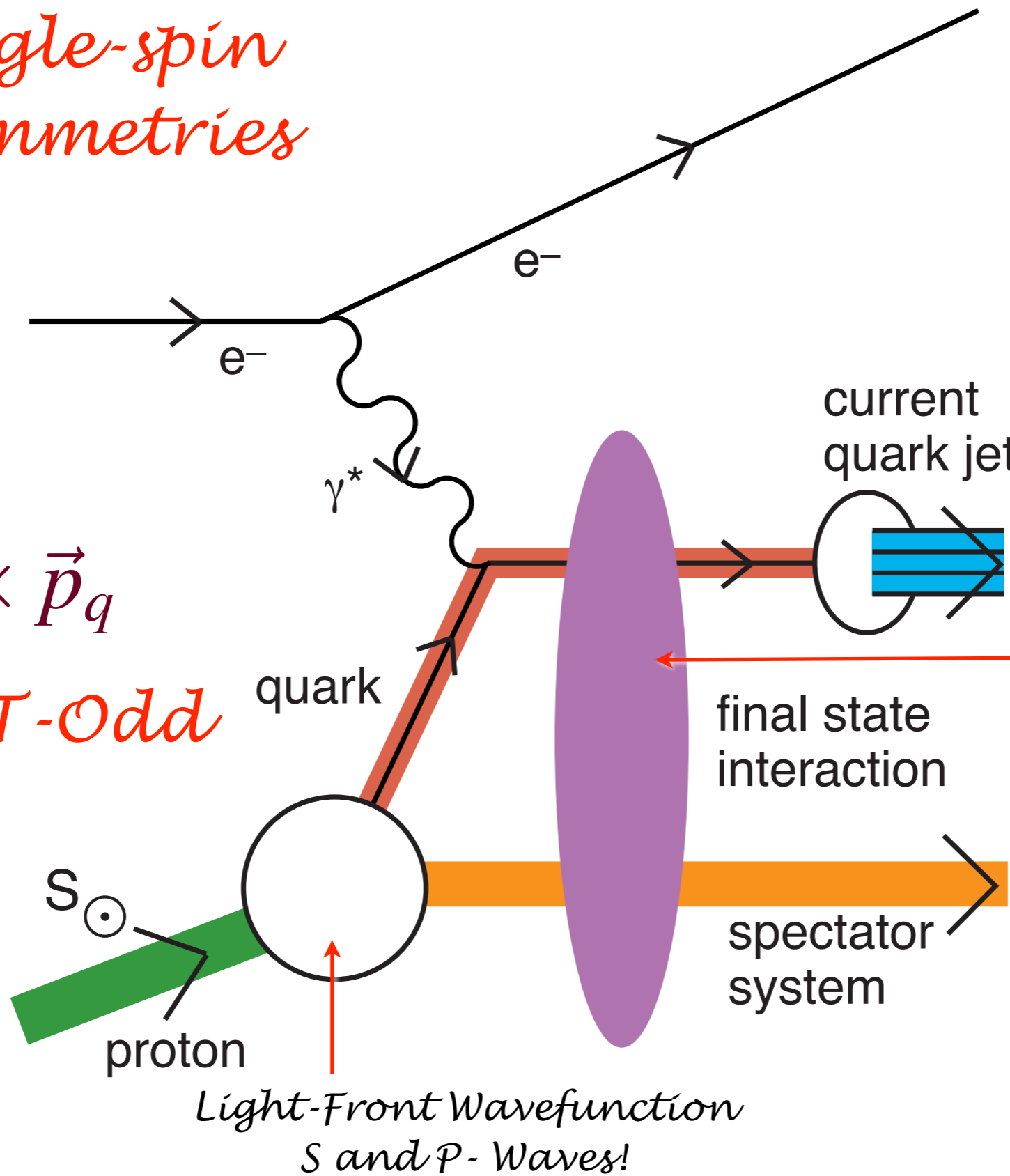
Hwang, Schmidt, sjb

Collins, Burkardt, Ji, Yuan. Pasquini, ...

QCD S- and P-Coulomb Phases --Wilson Line

“Lensing Effect”

Leading-Twist Rescattering Violates pQCD Factorization!



$$i \vec{S}_p \cdot \vec{q} \times \vec{p}_q$$

Pseudo-T-Odd

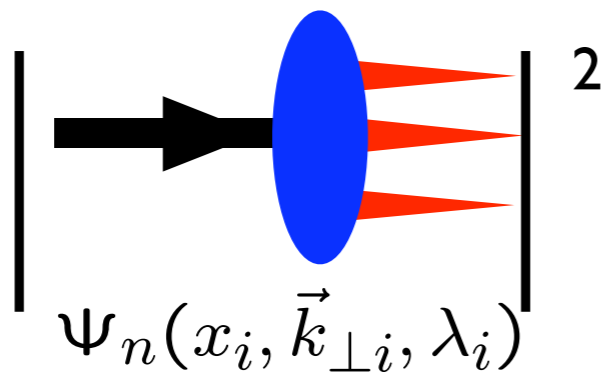
“Lensing” involves soft scales

Light-Front Wavefunction S and P-Waves!

Sign reversal in DY!

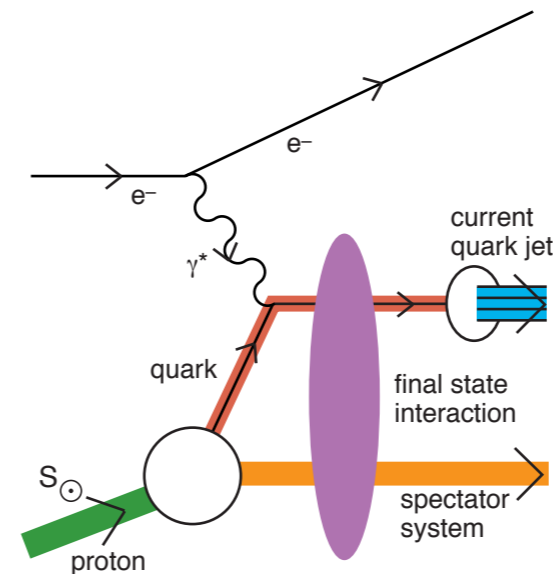
Static

- Square of Target LFWFs
- No Wilson Line
- Probability Distributions
- Process-Independent
- T-even Observables
- No Shadowing, Anti-Shadowing
- Sum Rules: Momentum and J^z
- DGLAP Evolution; mod. at large x
- No Diffractive DIS

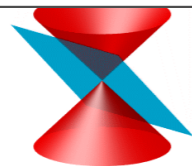


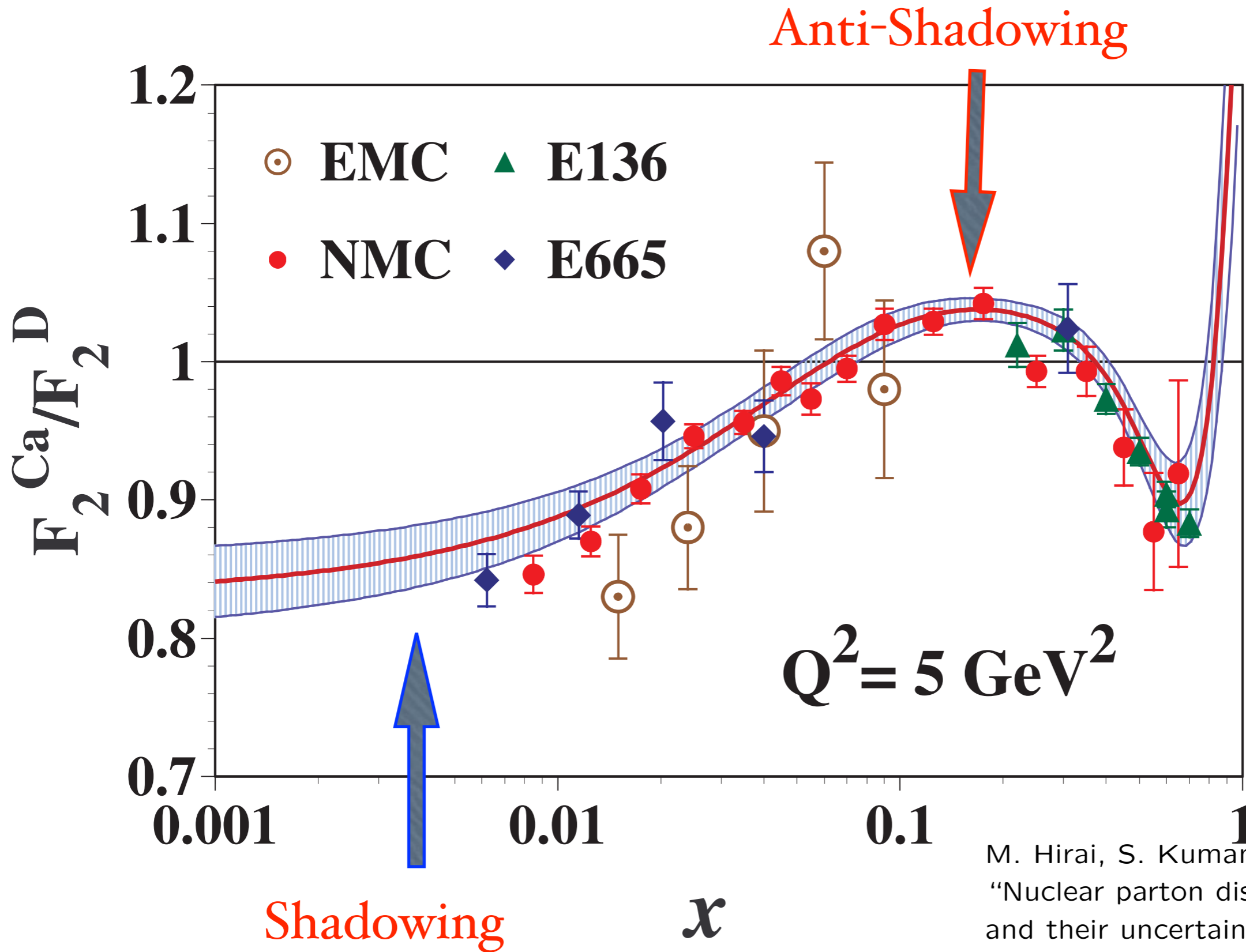
Dynamic

- Modified by Rescattering: ISI & FSI
- Contains Wilson Line, Phases
- No Probabilistic Interpretation
- Process-Dependent - From Collision
- T-Odd (Sivers, Boer-Mulders, etc.)
- Shadowing, Anti-Shadowing, Saturation
- Sum Rules Not Proven
- DGLAP Evolution
- Hard Pomeron and Odderon Diffractive DIS



**Hwang,
Schmidt, sjb,
Mulders, Boer
Qiu, Sterman
Collins, Qiu
Pasquini, Xiao,
Yuan, sjb**





M. Hirai, S. Kumano and T. H. Nagai,
 "Nuclear parton distribution functions
 and their uncertainties,"
 Phys. Rev. C **70**, 044905 (2004)
 [arXiv:hep-ph/0404093].

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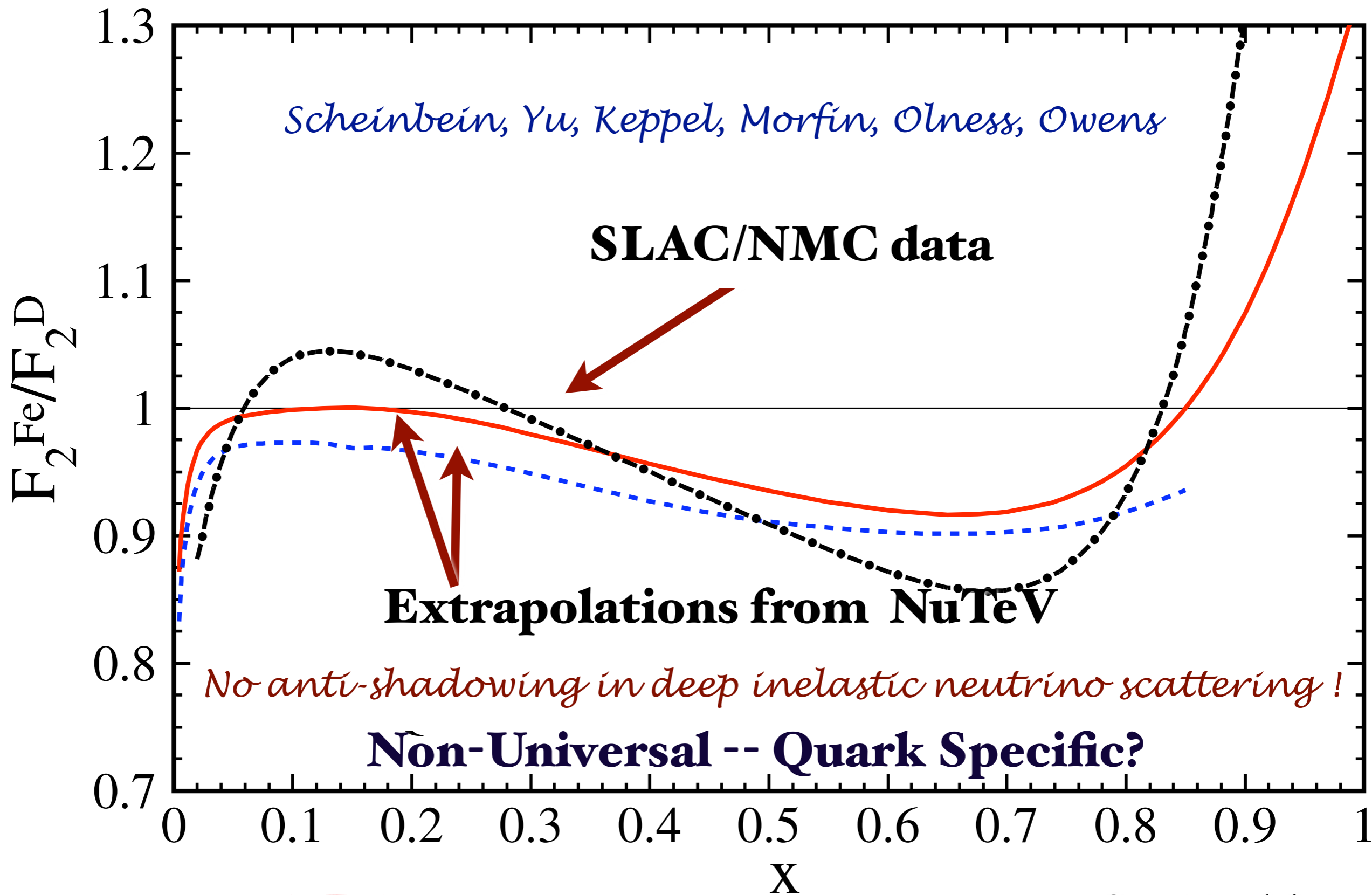


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$$Q^2 = 5 \text{ GeV}^2$$



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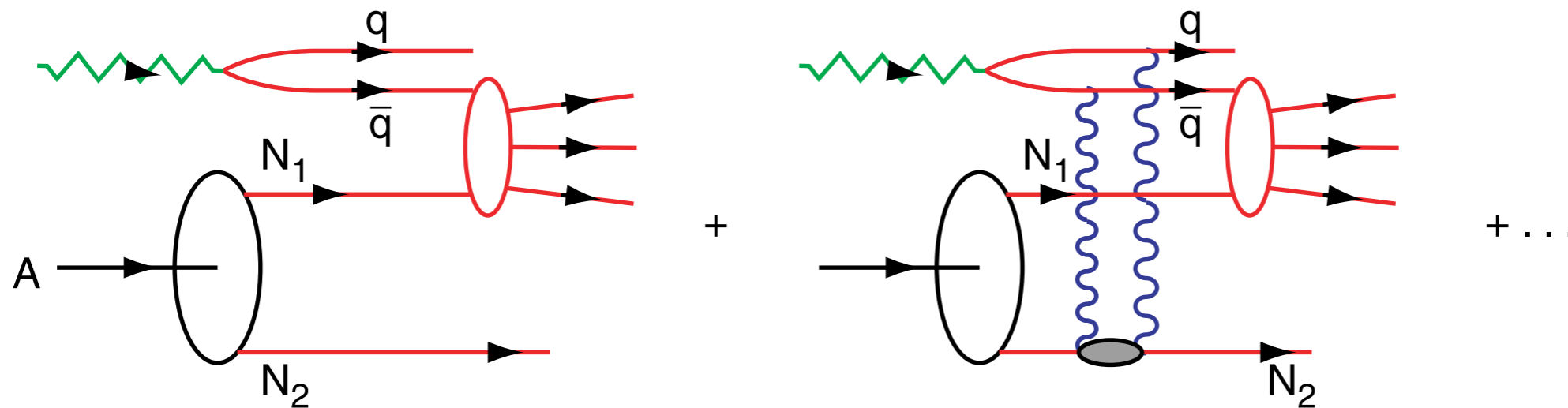


New Perspectives for Hadron Physics

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Nuclear Shadowing in QCD



Shadowing depends on understanding leading twist-diffraction in DIS

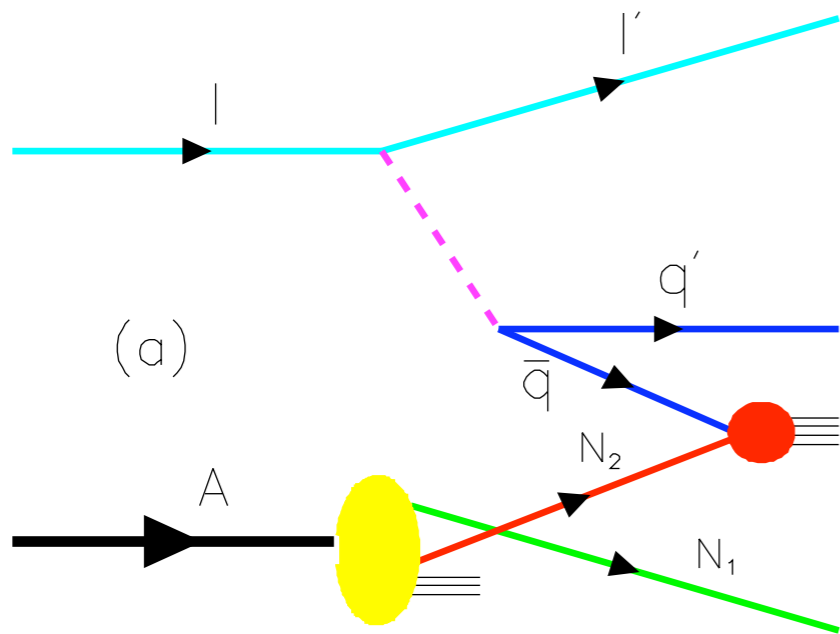
Nuclear Shadowing not included in nuclear LFWF !

Dynamical effect due to virtual photon interacting in nucleus

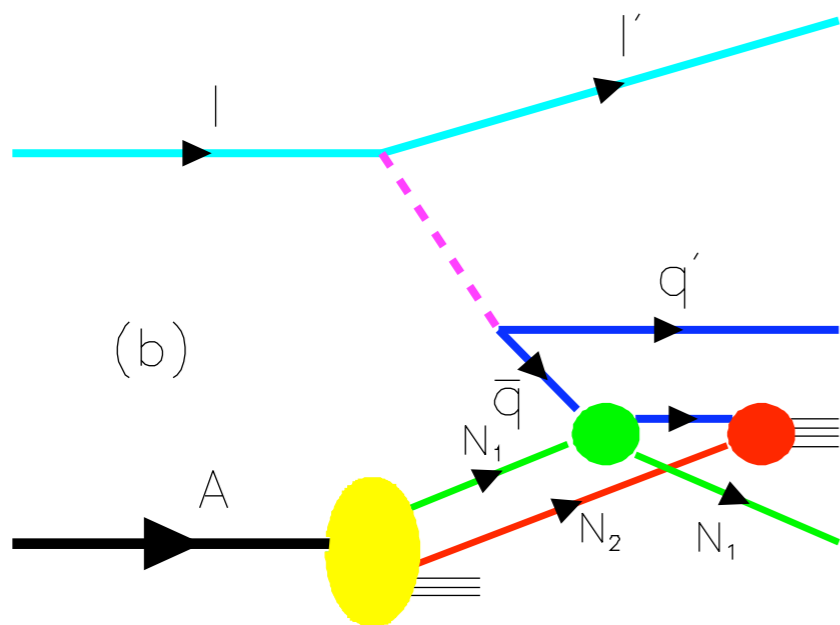
Diffraction via Reggeon gives constructive interference!

Anti-shadowing not universal





The one-step and two-step processes in DIS on a nucleus.



Coherence at small Bjorken x_B :
 $1/Mx_B = 2\nu/Q^2 \geq L_A$.

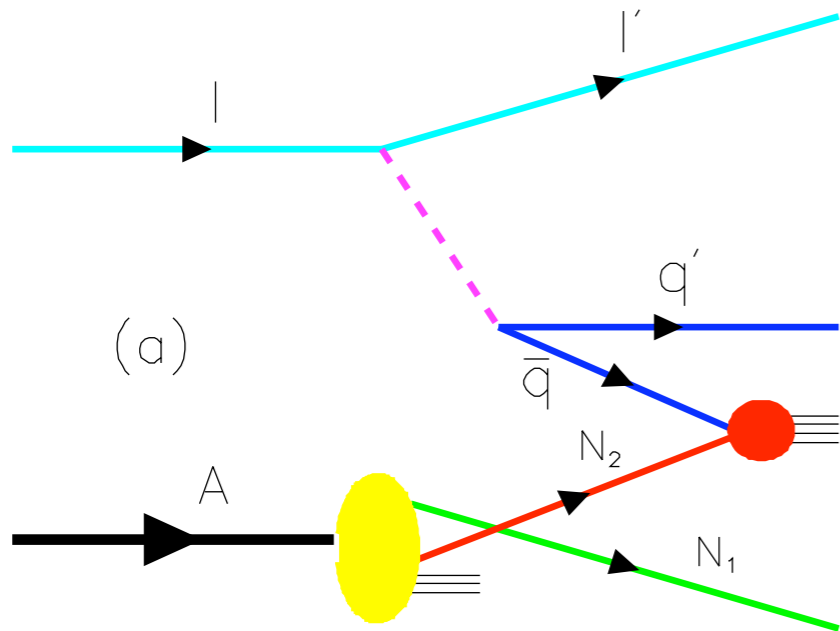
If the scattering on nucleon N_1 is via pomeron exchange, the one-step and two-step amplitudes are opposite in phase, thus diminishing the \bar{q} flux reaching N_2 .

→ Shadowing of the DIS nuclear structure functions.

Diffraction via Pomeron gives destructive interference!

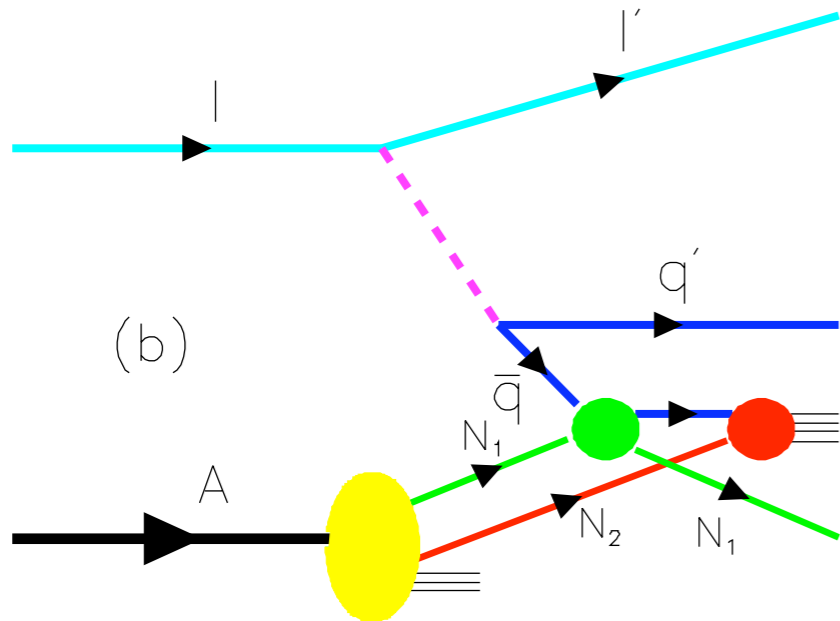
Shadowing





The one-step and two-step processes in DIS on a nucleus.

Coherence at small Bjorken x_B :
 $1/Mx_B = 2\nu/Q^2 \geq L_A$.



Regge

If the scattering on nucleon N_1 is via ~~pomeron~~ exchange, the one-step and two-step amplitudes are ~~opposite in phase~~, thus ~~diminishing the \bar{q} flux reaching N_2~~ .

constructive in phase
thus increasing the flux reaching N_2

Reggeon DDIS produces nuclear flavor-dependent anti-shadowing



Reggeon Exchange

Phase of two-step amplitude relative to one step:

$$\frac{1}{\sqrt{2}}(1 - i) \times i = \frac{1}{\sqrt{2}}(i + 1)$$

Constructive Interference

Depends on quark flavor!

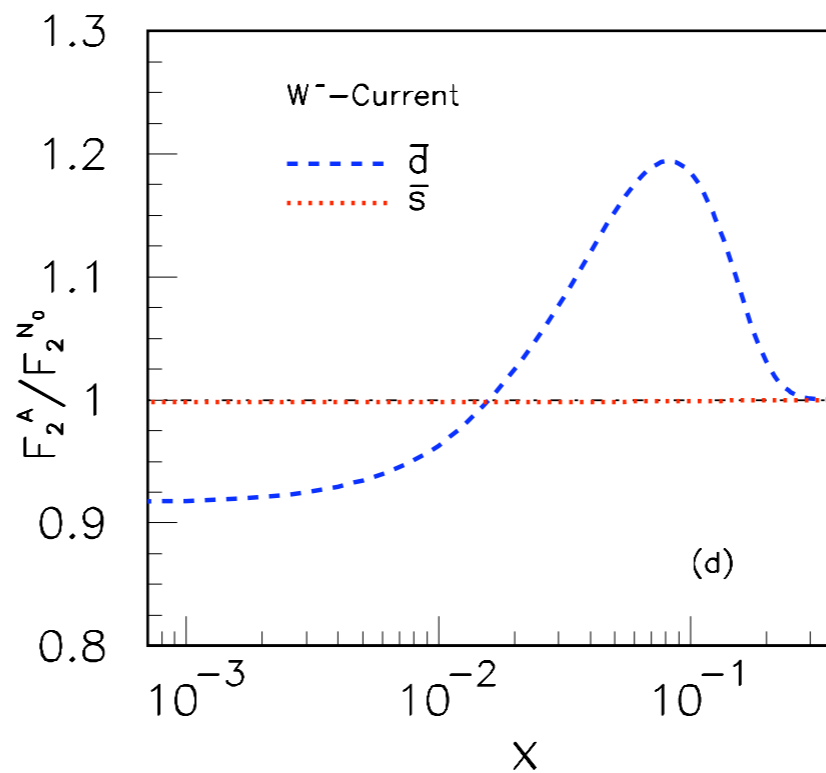
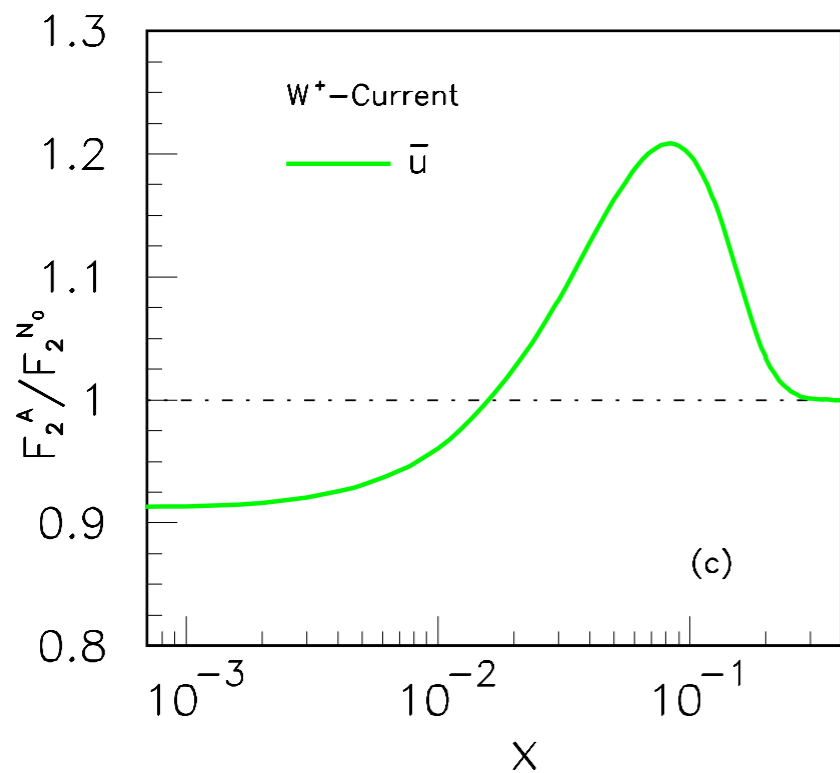
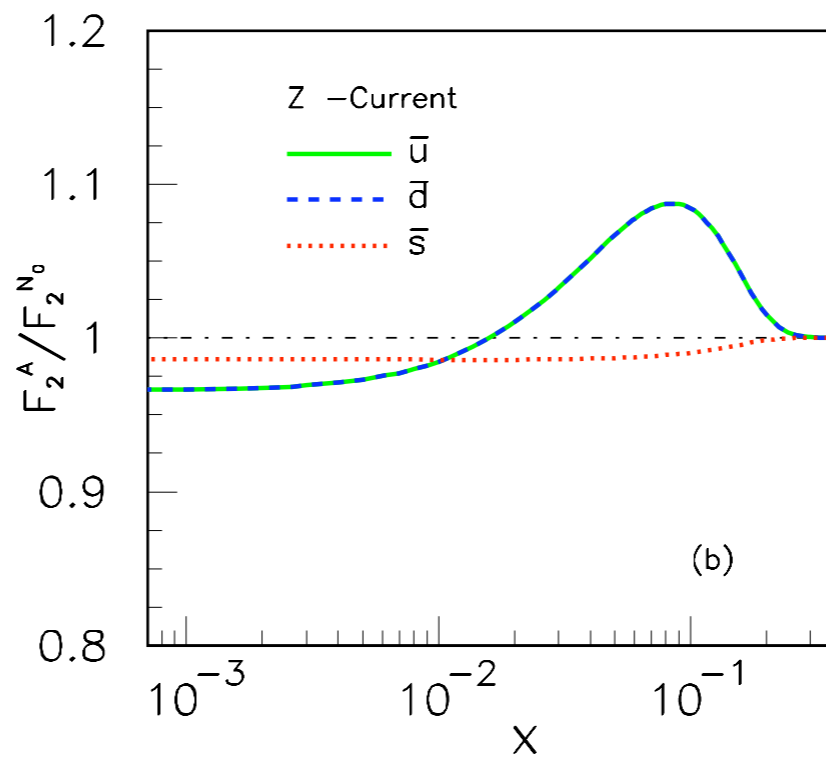
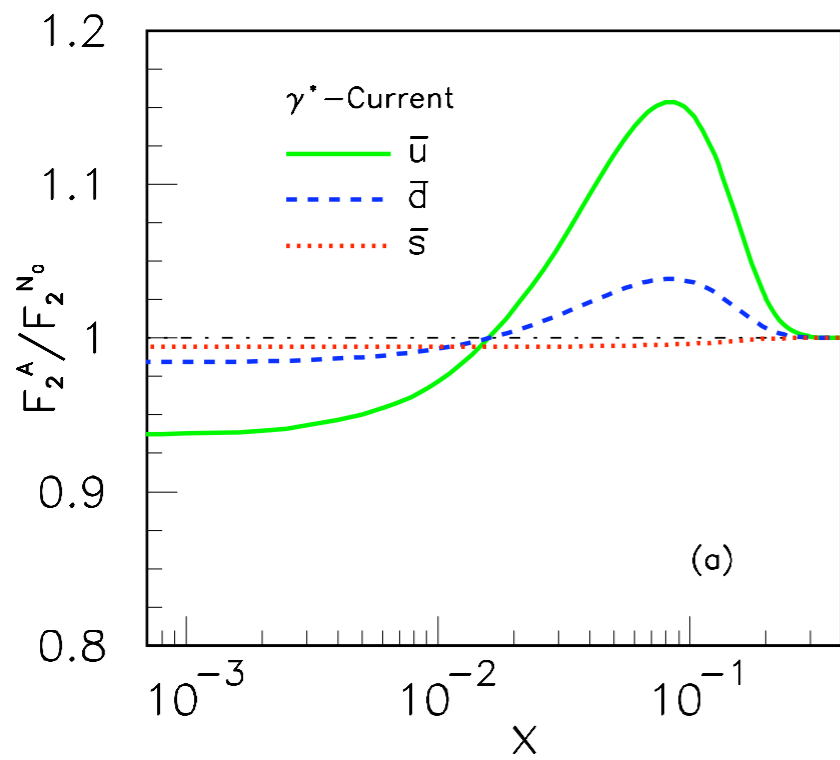
Thus antishadowing is not universal

Different for couplings of γ^* , Z^0 , W^\pm

Critical test: Tagged Drell-Yan



Schmidt, Yang; sjb



Modifies
NuTeV extraction of
 $\sin^2 \theta_W$

Test in flavor-tagged
DIS at the EIC

Nuclear Antishadowing not universal !

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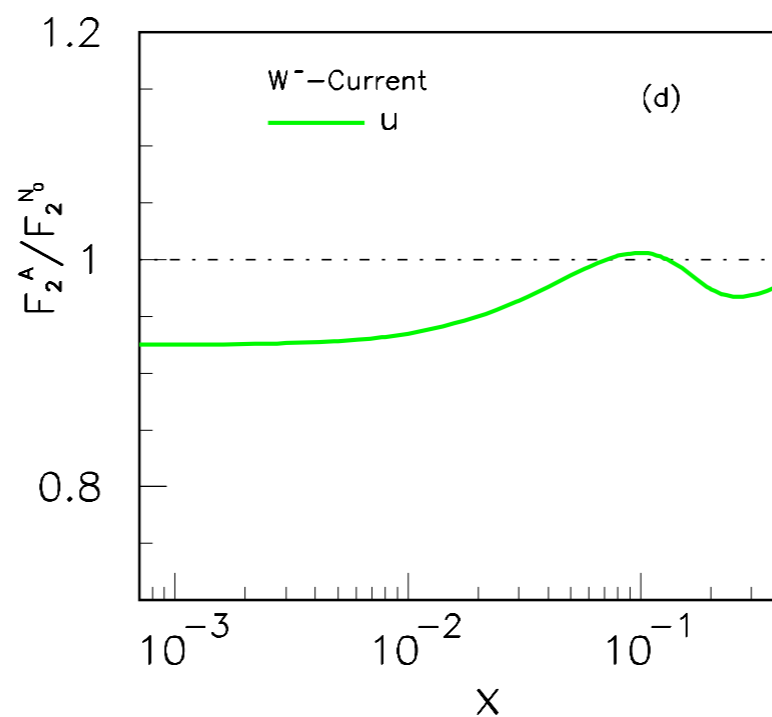
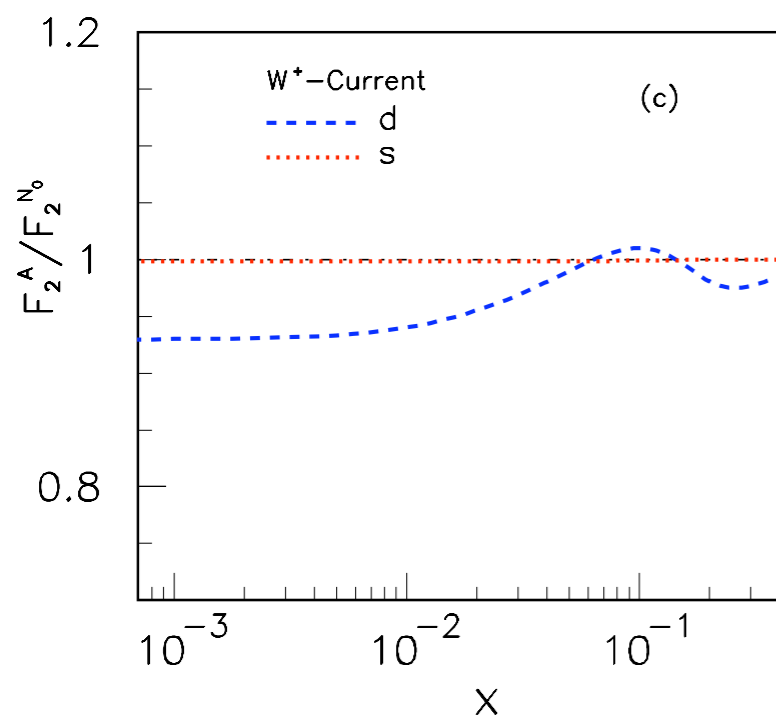
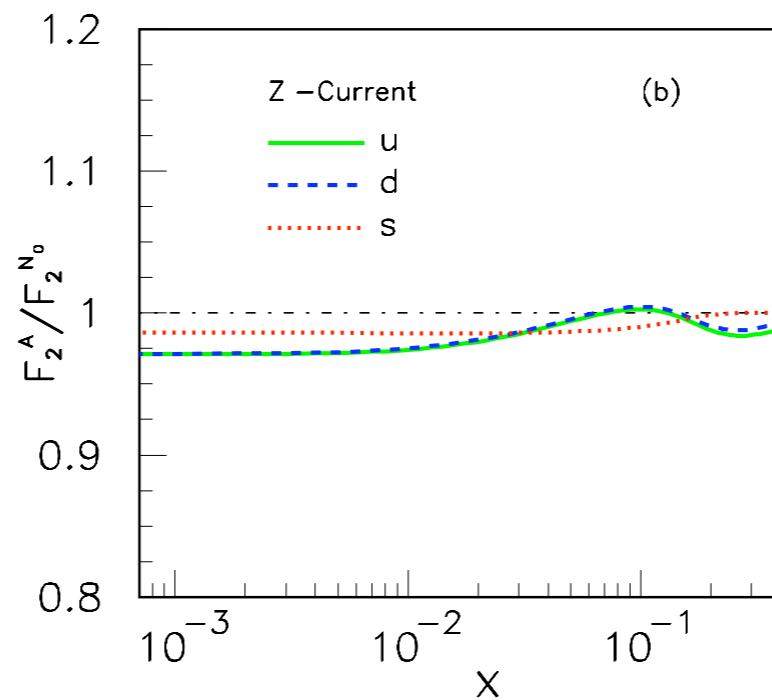
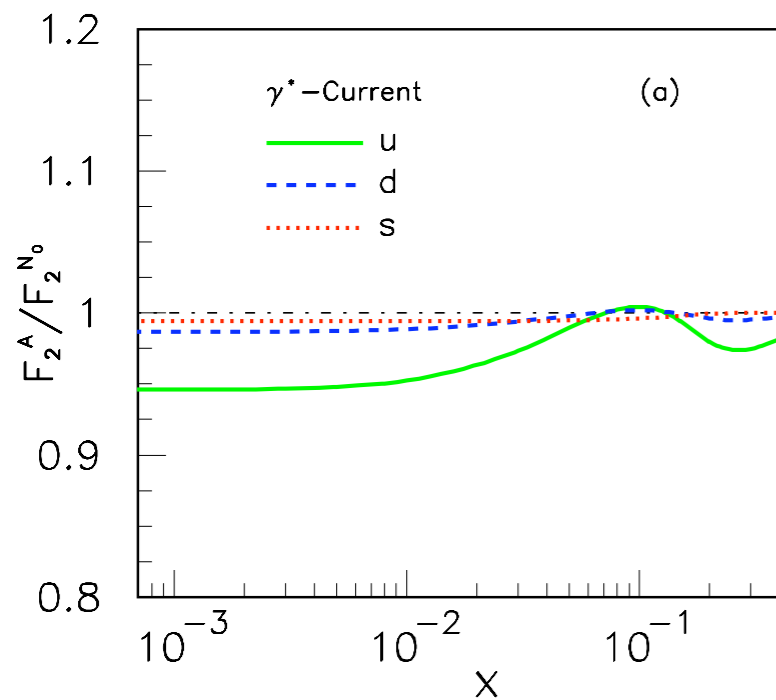


New Perspectives for Hadron Physics

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Shadowing and Antishadowing of DIS Structure Functions



S. J. Brodsky, I. Schmidt and J. J. Yang,
 “Nuclear Antishadowing in
 Neutrino Deep Inelastic Scattering,”
 Phys. Rev. D 70, 116003 (2004)
 [arXiv:hep-ph/0409279].

Modifies
NuTeV extraction of
 $\sin^2 \theta_W$

Test in flavor-tagged
lepton-nucleus collisions

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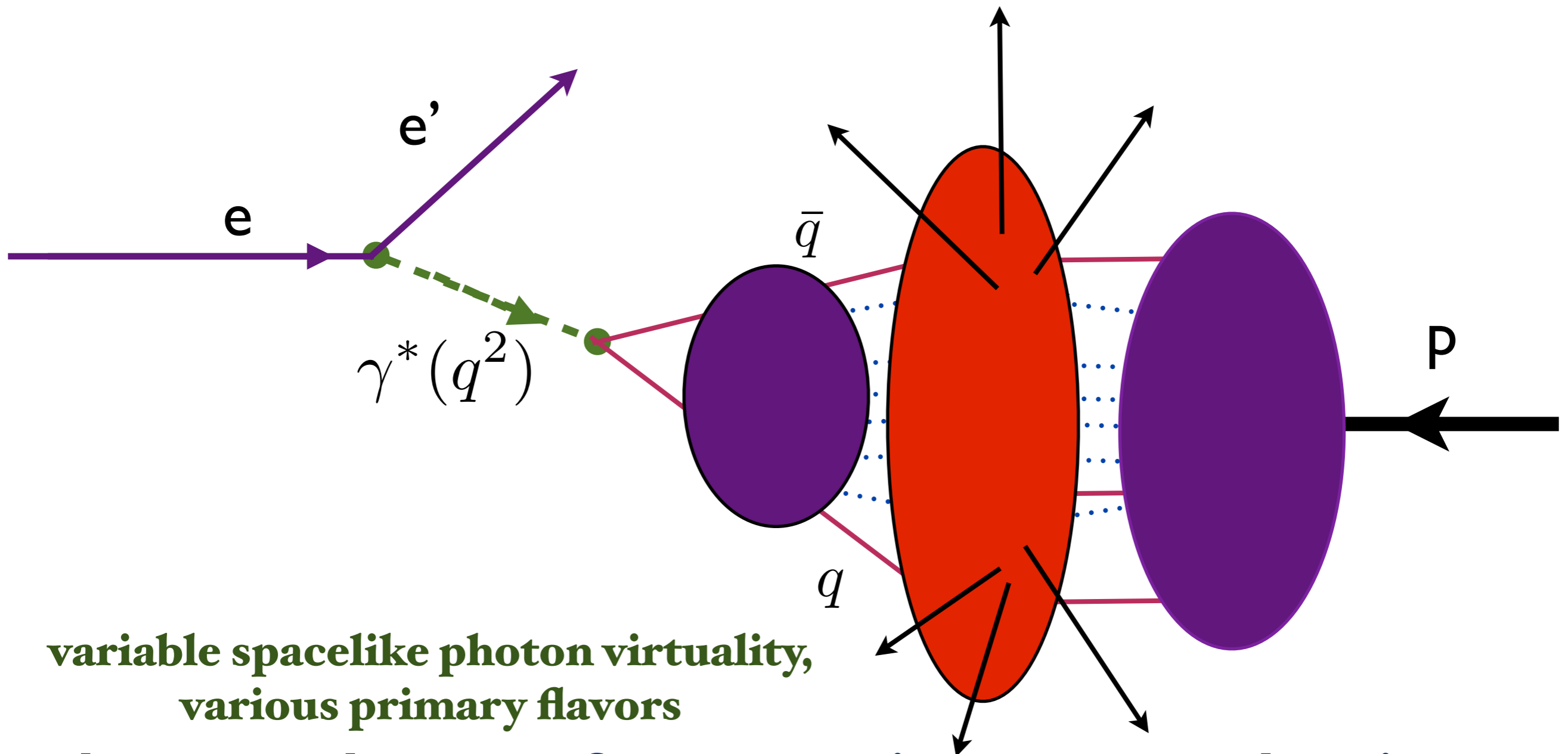
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LHeC: Virtual Photon-Proton Collider

Perspective from the e-p collider frame



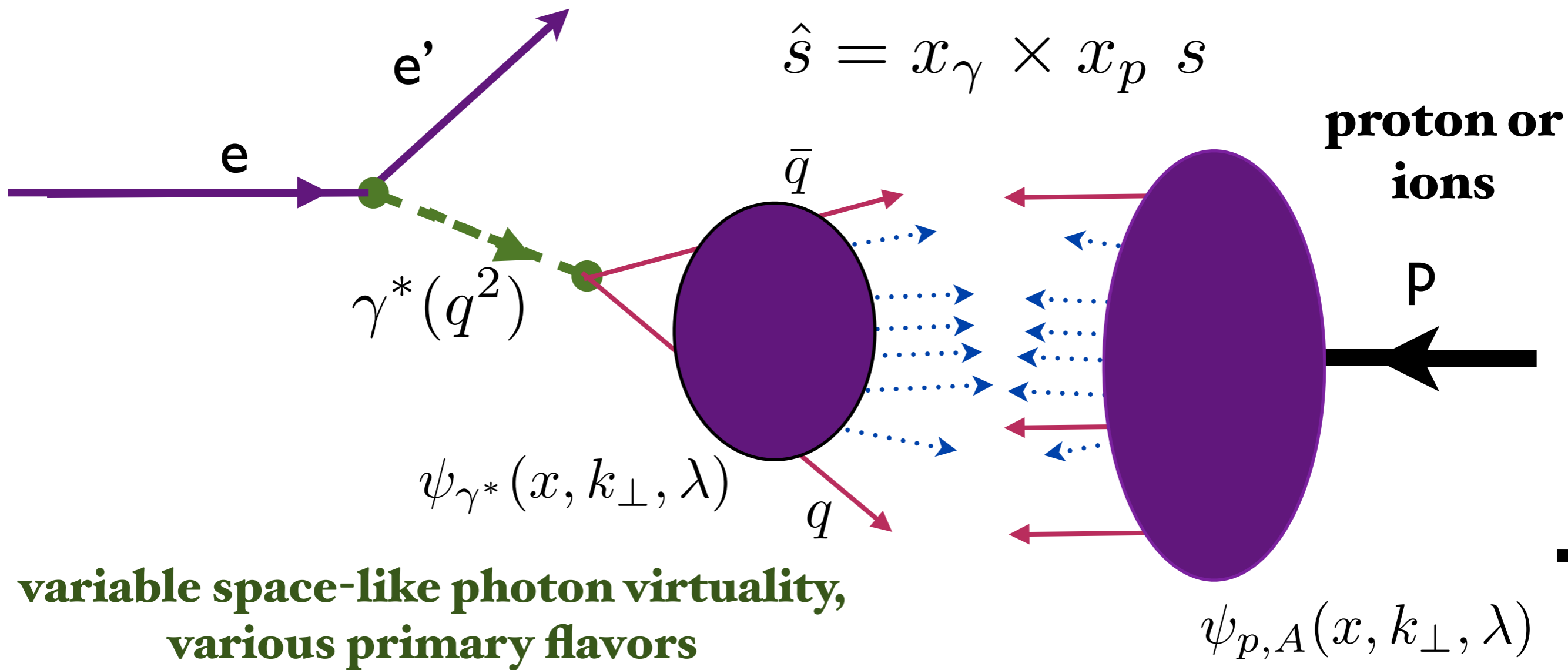
**variable spacelike photon virtuality,
various primary flavors**

photon and proton fragmentation vs. central regions

Saturation, nuclear shadowing, antishadowing

Electron-Ion Colliders: Virtual Photon-Ion Collider

Perspective from the e-p collider frame



$\bar{q}q$ plane aligned with lepton scattering plane $\sim \cos^2\phi$
Front-surface dynamics: shadowing/antishadowing

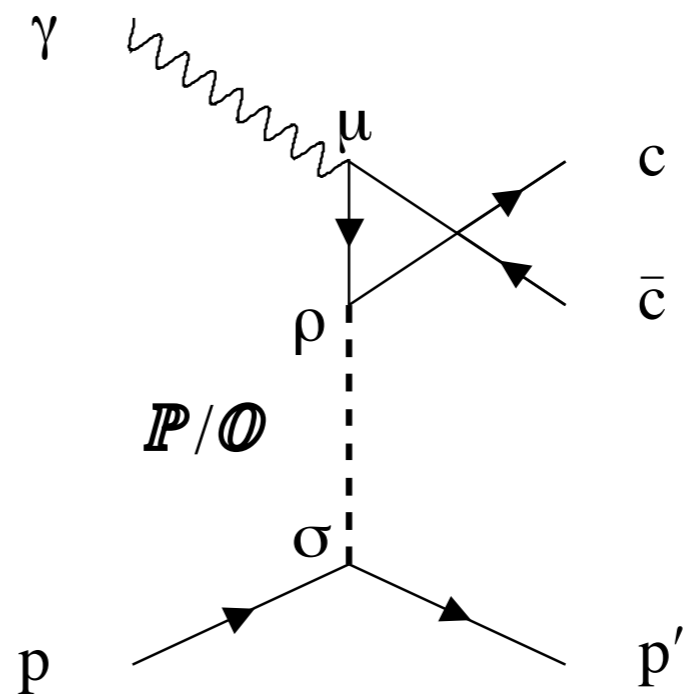
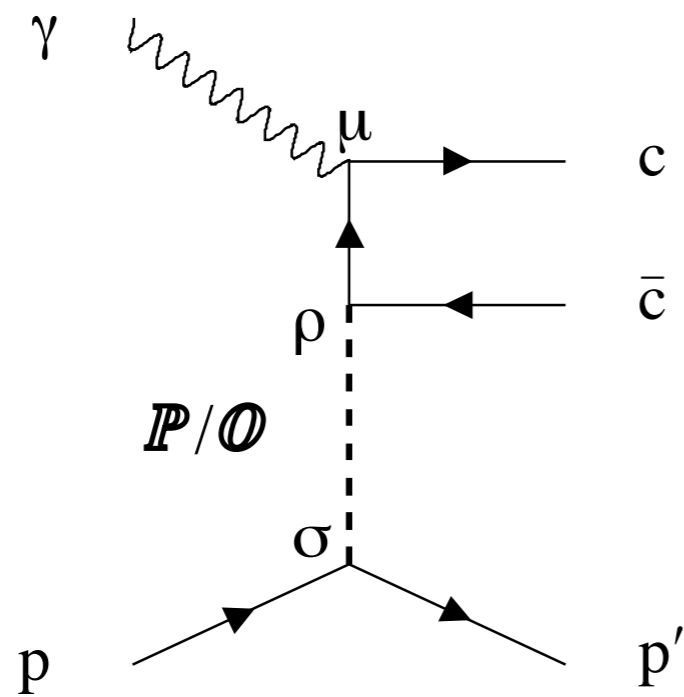
Novel QCD Physics at an Electron-Ion Collider

- **Control Collisions of Flux Tubes and Ridge Phenomena**
- **Study Flavor-Dependence of Anti-Shadowing**
- **Heavy Quarks at Large x ; Exotic States**
- **Direct, color-transparent hard subprocesses and the baryon anomaly**
- **Tri-Jet Production and the proton's LFWF**
- **Odderon-Pomeron Interference**
- **Digluon-initiated subprocesses and anomalous nuclear dependence of quarkonium production**
- **Factorization-Breaking Lensing Corrections**

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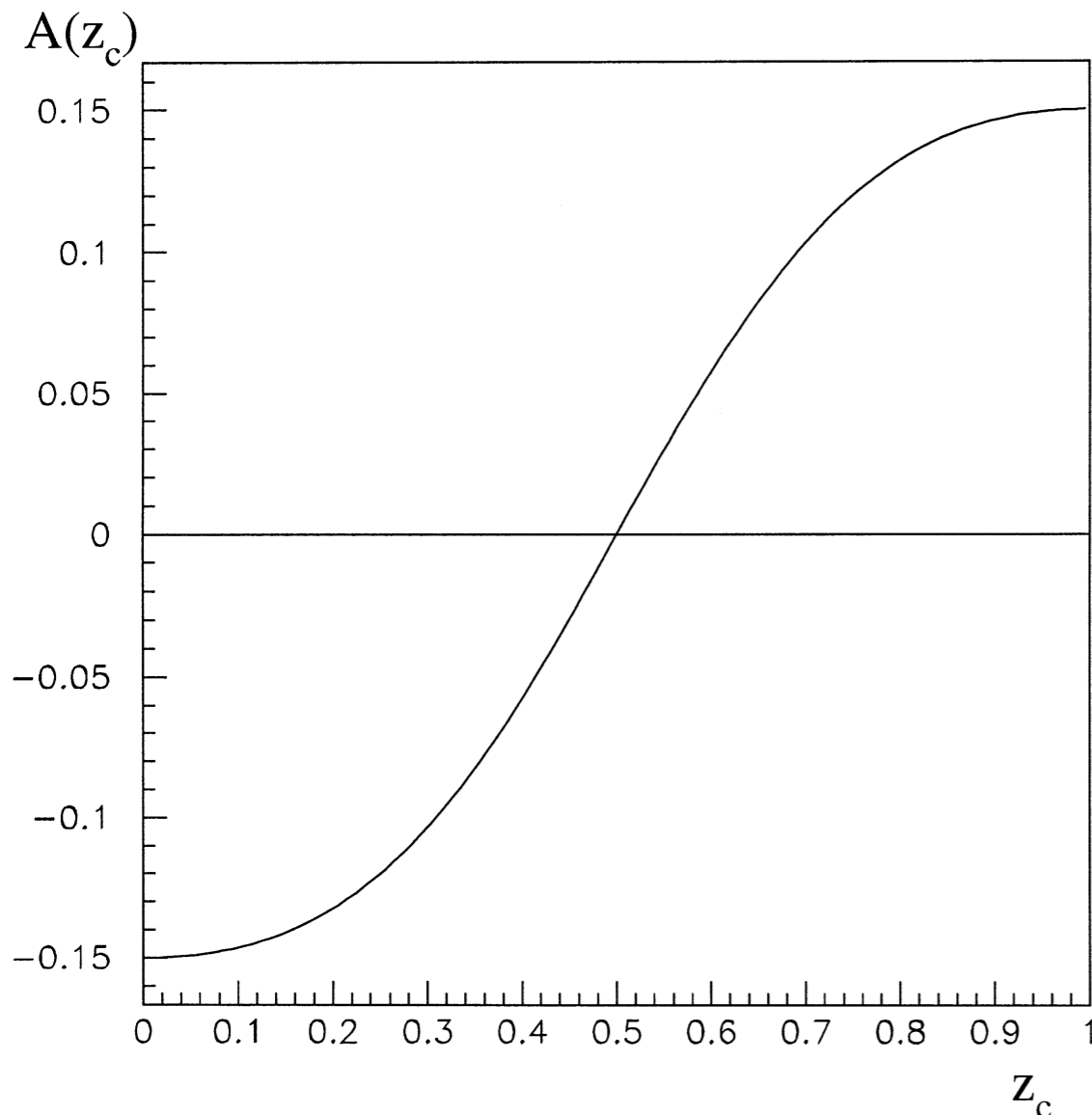


New Perspectives for Hadron Physics



$$\gamma^* p \rightarrow c\bar{c}p$$

Odderon-Pomeron Interference!



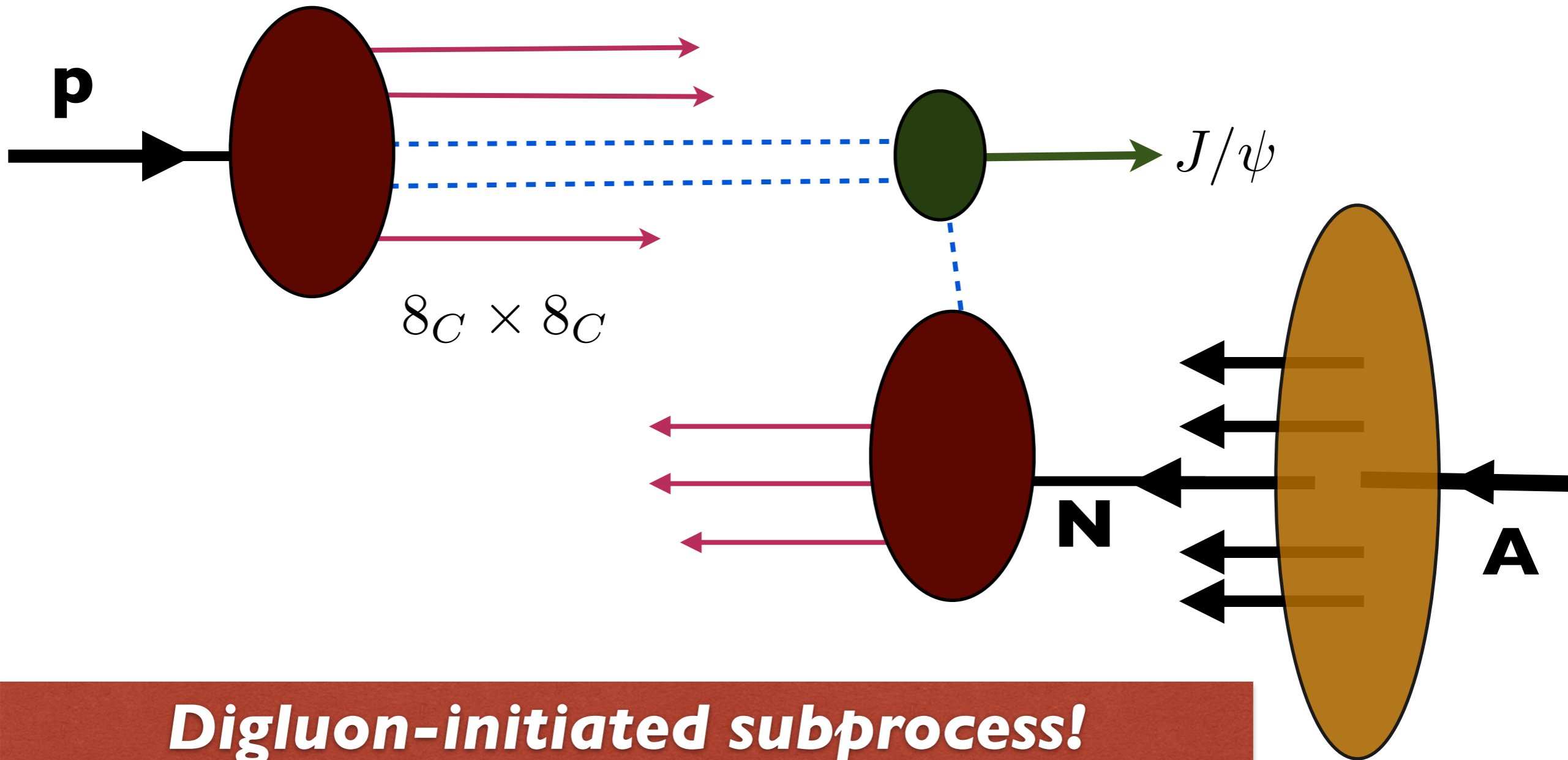
$$\mathcal{A}(t \simeq 0, M_X^2, z_c) \simeq 0.45 \left(\frac{s_{\gamma p}}{M_X^2} \right)^{-0.25} \frac{2z_c - 1}{z_c^2 + (1 - z_c)^2}$$

Measure charm asymmetry in photon fragmentation region

Merino, Rathsmann, sjb

$$pA \rightarrow J/\psi X$$

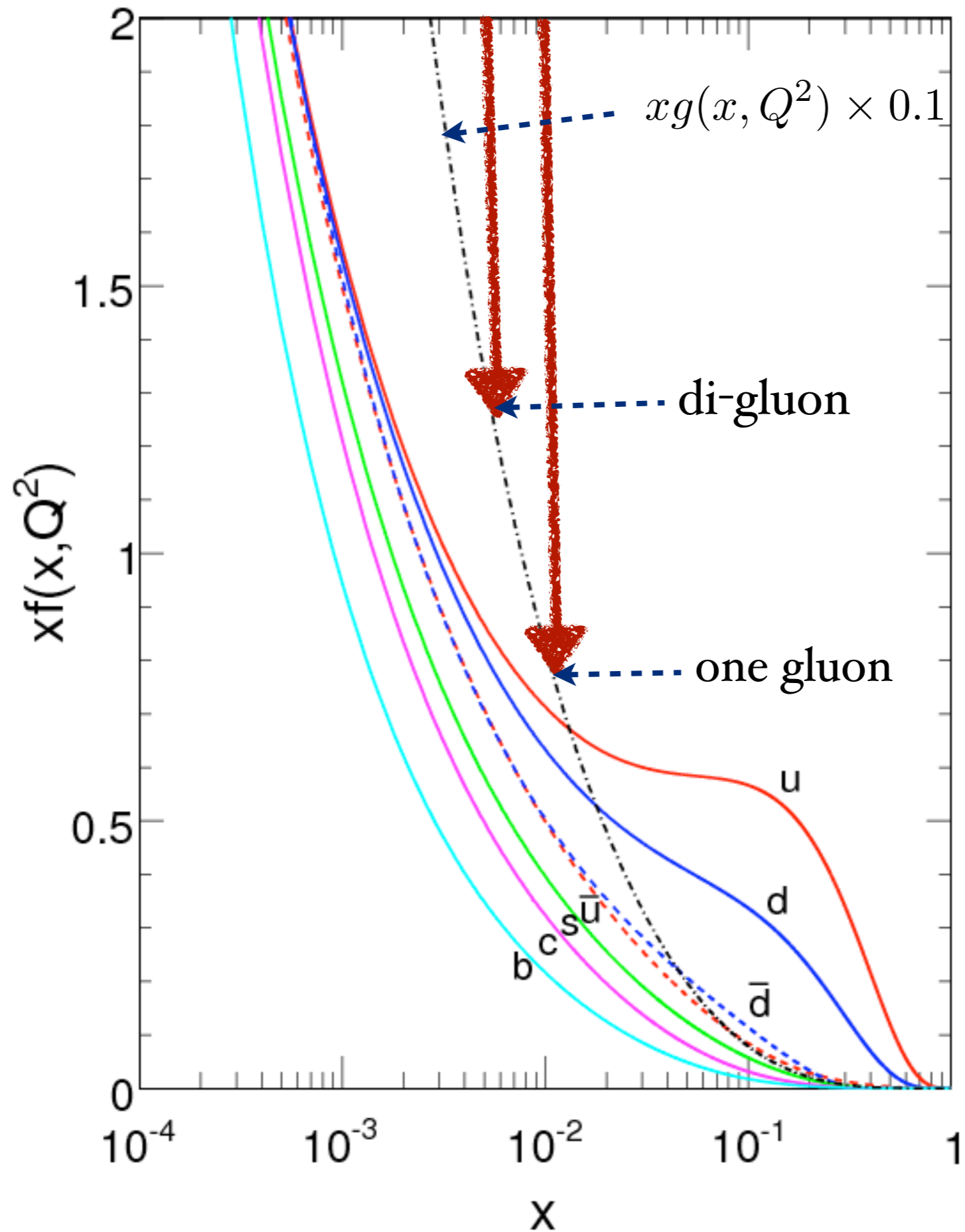
$$(gg)_{8_C} + g_{8_C} \rightarrow J/\psi$$



Digluon-initiated subprocess!

Higher-Twist but can dominate at forward rapidity, small p_T

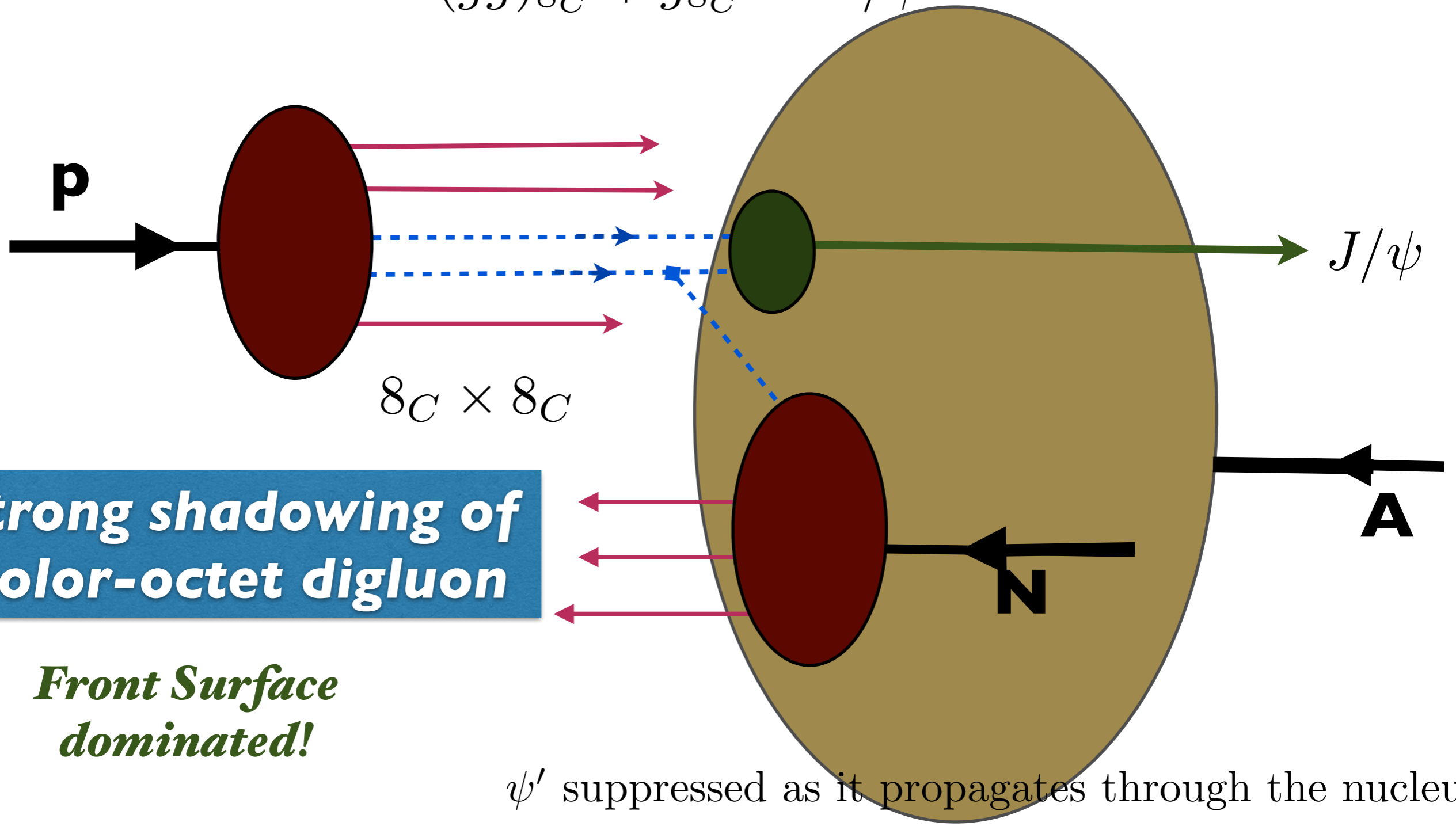
Two gluons at $g(0.005) \sim \frac{13}{0.005} = 2600$ vs. one gluon at $g(0.01) \sim \frac{8}{0.01} = 800$



Forward rapidity $y \sim 4$

$$pA \rightarrow J/\psi X$$

$$(gg)_{8_C} + g_{8_C} \rightarrow J/\psi$$



Strong shadowing of color-octet digluon

Front Surface dominated!

Crossing: Diffractive & pomeron exchange

ψ' suppressed as it propagates through the nucleus

Digluon-initiated subprocess!

Exact frame-independent formulation of nonperturbative QCD!

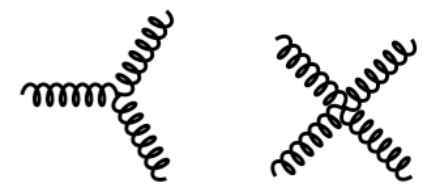
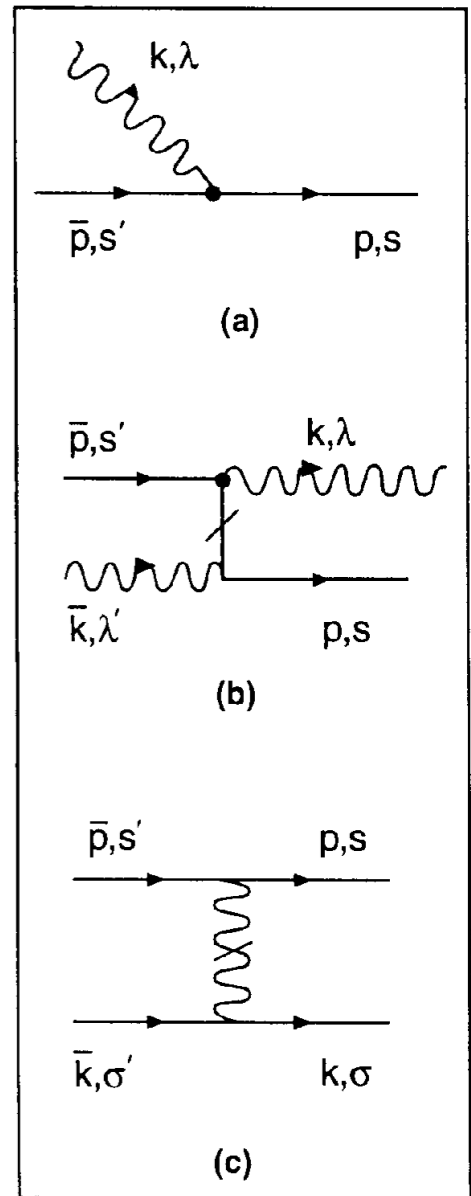
$$L^{QCD} \rightarrow H_{LF}^{QCD}$$

$$H_{LF}^{QCD} = \sum_i \left[\frac{m^2 + k_{\perp}^2}{x} \right]_i + H_{LF}^{int}$$

H_{LF}^{int} : Matrix in Fock Space

$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

$$|p, J_z\rangle = \sum_{n=3} \psi_n(x_i, \vec{k}_{\perp i}, \lambda_i) |n; x_i, \vec{k}_{\perp i}, \lambda_i\rangle$$



H_{LF}^{int}

Eigenvalues and Eigensolutions give Hadronic Spectrum and Light-Front wavefunctions

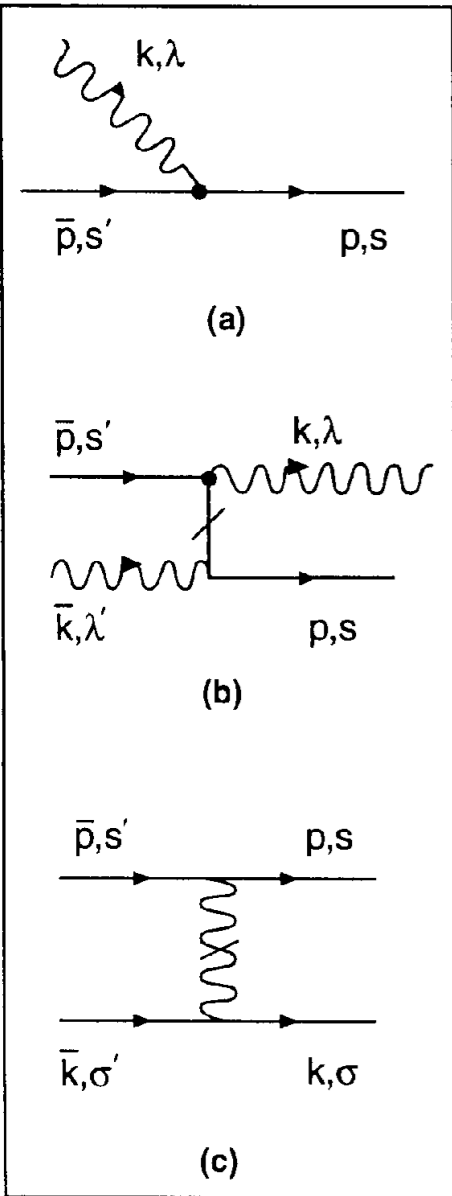
LFWFs: Off-shell in P- and invariant mass

Light-Front QCD
Heisenberg Equation

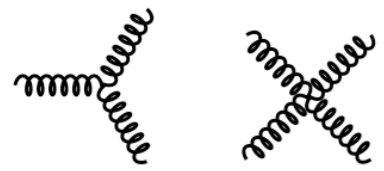
$$H_{LC}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

DLCQ: Solve QCD(1+1) for any quark mass and flavors

Hornbostel, Pauli, sjb



n	Sector	1 q \bar{q}	2 gg	3 q \bar{q} g	4 q \bar{q} q \bar{q}	5 gg g	6 q \bar{q} gg	7 q \bar{q} q \bar{q} g	8 q \bar{q} q \bar{q} q \bar{q}	9 gg gg	10 q \bar{q} gg g	11 q \bar{q} q \bar{q} gg	12 q \bar{q} q \bar{q} q \bar{q} g	13 q \bar{q} q \bar{q} q \bar{q} q \bar{q}
1	q \bar{q}				
2	gg			
3	q \bar{q} g							
4	q \bar{q} q \bar{q}	
5	gg g
6	q \bar{q} gg								.				.	.
7	q \bar{q} q \bar{q} g
8	q \bar{q} q \bar{q} q \bar{q}			
9	gg gg
10	q \bar{q} gg g
11	q \bar{q} q \bar{q} gg
12	q \bar{q} q \bar{q} q \bar{q} g			
13	q \bar{q} q \bar{q} q \bar{q} q \bar{q}		



Minkowski space; frame-independent; no fermion doubling; no ghosts
trivial vacuum

$$|p, S_z\rangle = \sum_{n=3} \Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i) |n; \vec{k}_{\perp i}, \lambda_i\rangle$$

sum over states with $n=3, 4, \dots$ constituents

The Light Front Fock State Wavefunctions

$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

are boost invariant; they are independent of the hadron's energy and momentum P^μ .

The light-cone momentum fraction

$$x_i = \frac{k_i^+}{p^+} = \frac{k_i^0 + k_i^z}{P^0 + P^z}$$

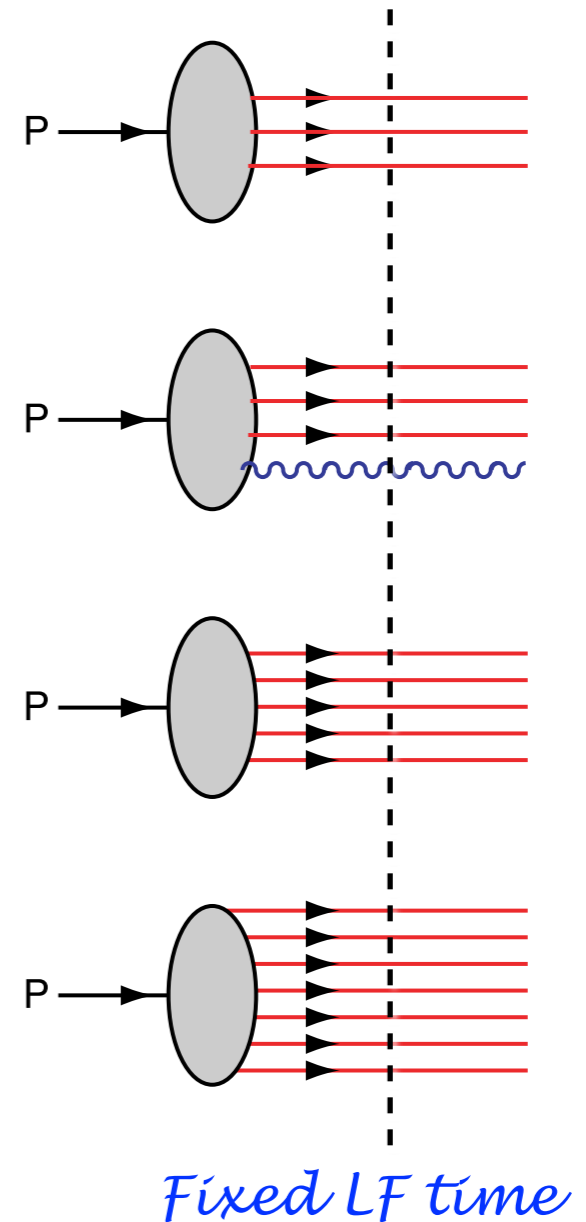
are boost invariant.

$$\sum_i^n k_i^+ = P^+, \quad \sum_i^n x_i = 1, \quad \sum_i^n \vec{k}_i^\perp = \vec{0}^\perp.$$

Intrinsic heavy quarks
 $s(x), c(x), b(x)$ at high x !

$$\bar{s}(x) \neq s(x)$$

$$\bar{u}(x) \neq \bar{d}(x)$$

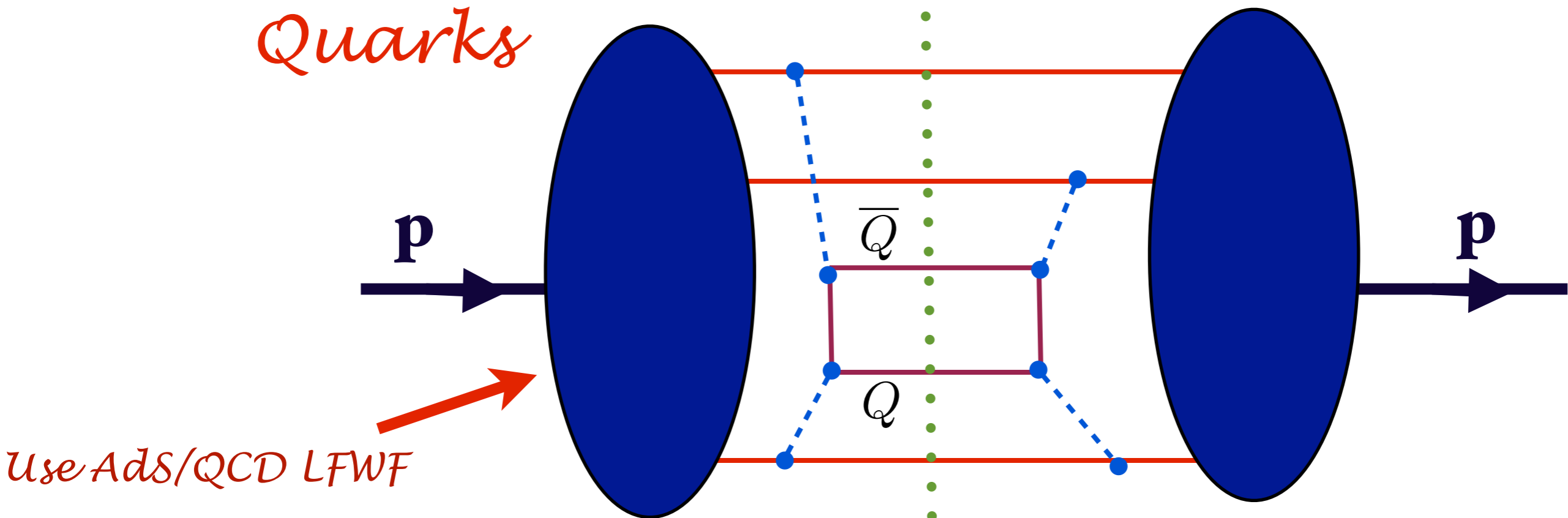


Hidden Color

*Proton Self Energy
Intrinsic Heavy
Quarks*

Fixed LF time

$$x_Q \propto (m_Q^2 + k_{\perp}^2)^{1/2}$$



Use AdS/QCD LFWF

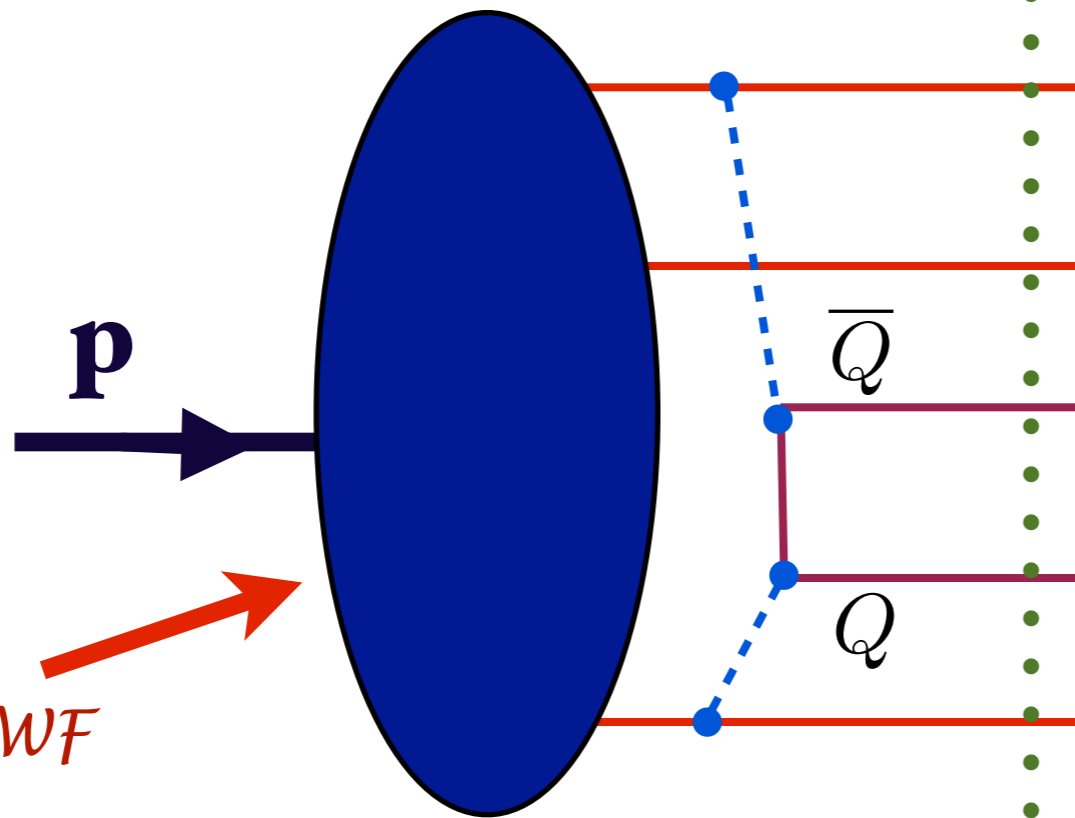
Probability (QED) $\propto \frac{1}{M_{\ell}^4}$

Probability (QCD) $\propto \frac{1}{M_Q^2}$

**Collins, Ellis, Gunion, Mueller, sjb
Polyakov, et al.**

*Proton 5-quark Fock State:
Intrinsic Heavy Quarks*

Fixed LF time



*QCD predicts
Intrinsic Heavy
Quarks at high x .*

Minimal off-shellness

Use AdS/QCD LFWF

$$x_Q \propto (m_Q^2 + k_{\perp}^2)^{1/2}$$

Probability (QED) $\propto \frac{1}{M_{\ell}^4}$

Probability (QCD) $\propto \frac{1}{M_Q^2}$

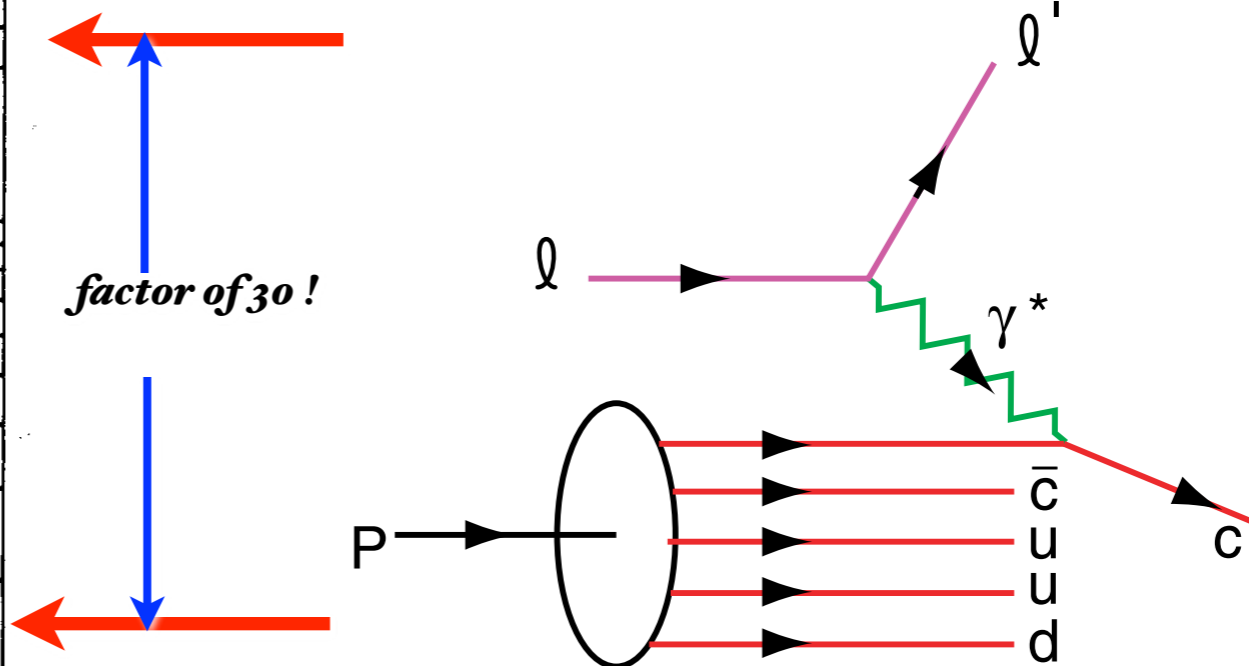
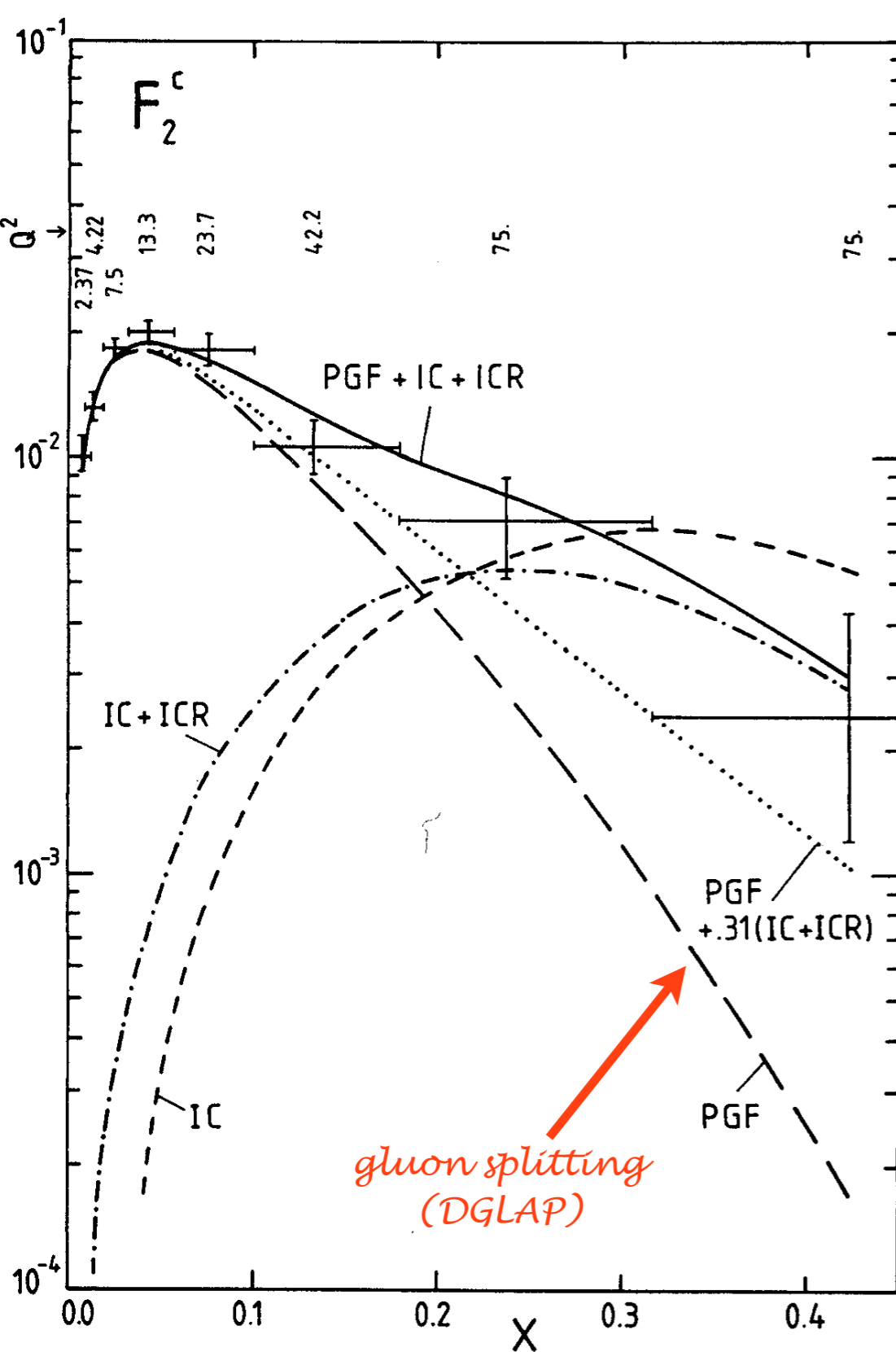
**Collins, Ellis, Gunion, Mueller, sjb
Polyakov, et al.**

Measurement of Charm Structure Function

J. J. Aubert et al. [European Muon Collaboration], "Production Of Charmed Particles In 250-GeV Mu+ - Iron Interactions," Nucl. Phys. B 213, 31 (1983).

First Evidence for Intrinsic Charm

Hoyer, Peterson, Sakai, sjb



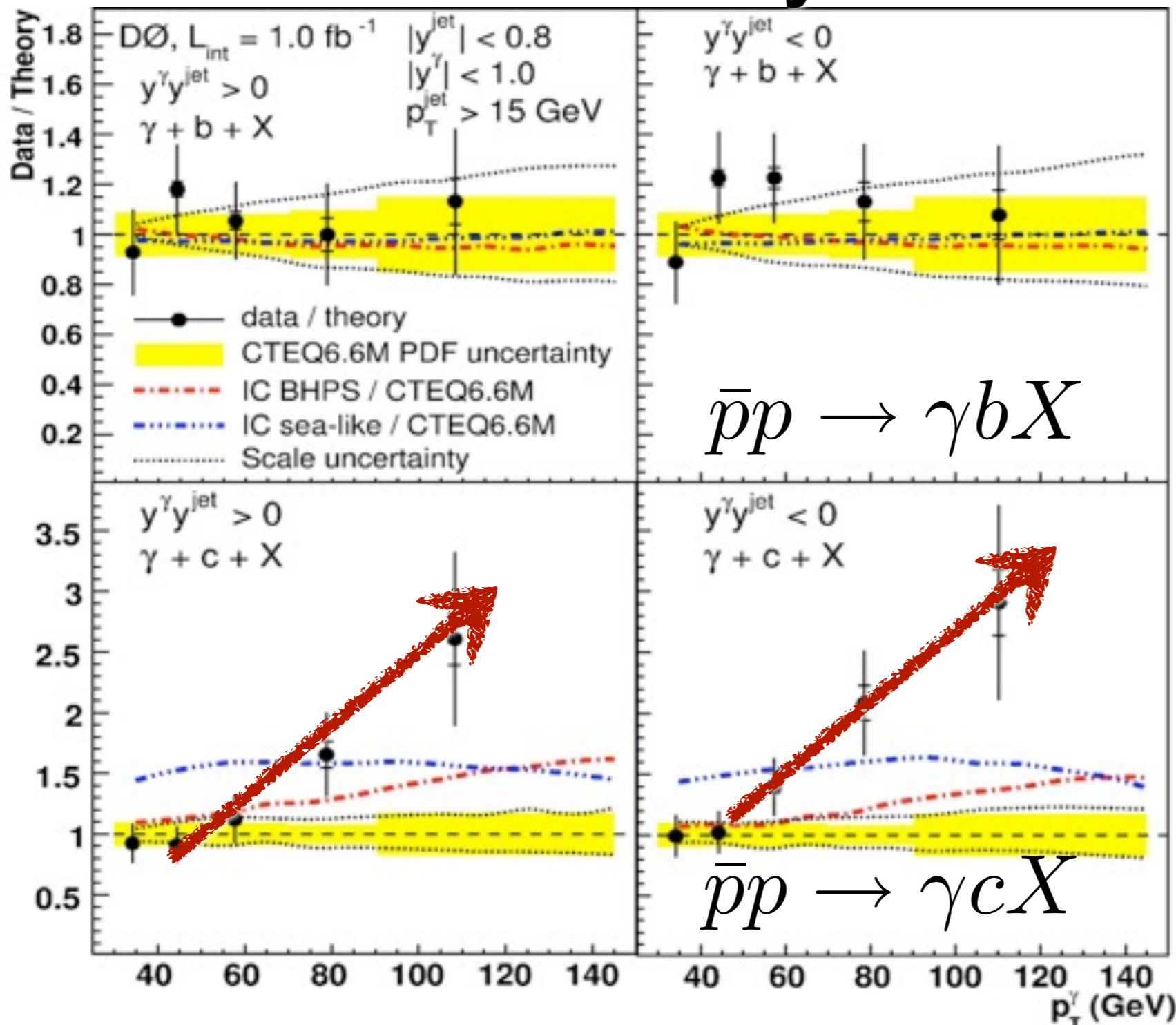
DGLAP / Photon-Gluon Fusion: factor of 30 too small

Two Components (separate evolution):

$$c(x, Q^2) = c(x, Q^2)_{\text{extrinsic}} + c(x, Q^2)_{\text{intrinsic}}$$

Measurement of $\gamma + b + X$ and $\gamma + c + X$ Production Cross Sections
in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ TeV

Data/Theory



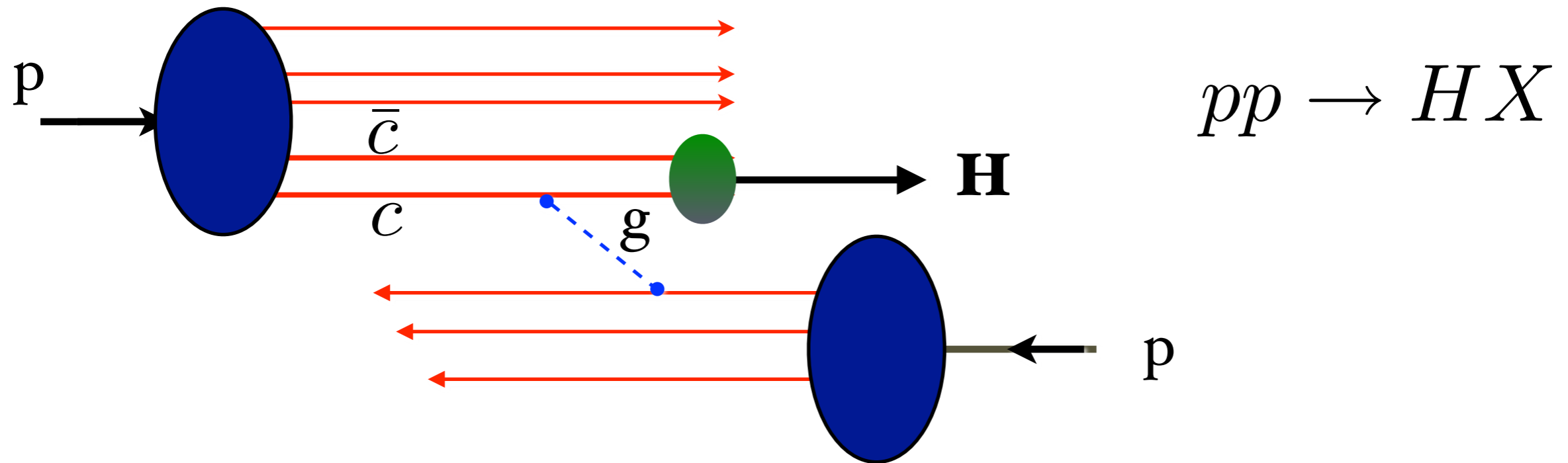
$$\frac{\Delta\sigma(\bar{p}p \rightarrow \gamma c X)}{\Delta\sigma(\bar{p}p \rightarrow \gamma b X)}$$

**Ratio insensitive
to gluon PDF,
scales**

**Signal for significant
IC
at $x > 0.1$**

*Consistent with EMC measurement of charm
structure function at high x*

*Intrinsic Charm Mechanism for Inclusive
High- x_F Higgs Production*



Also: intrinsic strangeness, bottom, top

Higgs can have > 80% of Proton Momentum!

New production mechanism for Higgs

AFTER: Higgs production at threshold!

JLab: Charm production near threshold!

● **LF wavefunctions play the role of Schrödinger wavefunctions in Atomic Physics**

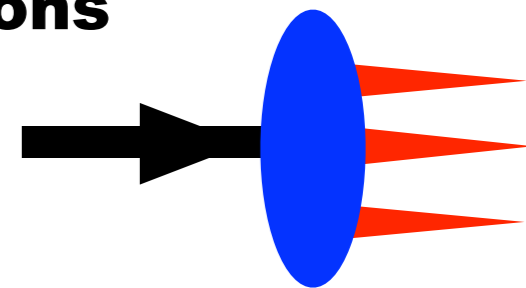
● **LFWFs=Hadron Eigensolutions: Direct Connection to QCD Lagrangian**

● **Relativistic, frame-independent: no boosts, no disc contraction, Melosh built into LF spinors**

● **Hadronic observables computed from LFWFs: Form factors, Structure Functions, Distribution Amplitudes, GPDs, TMDs, Weak Decays, modulo 'lensing' from ISIs, FSIs**

● **Cannot compute current matrix elements using instant form from eigensolutions alone -- need to include vacuum currents!**

● **Hadron Physics without LFWFs is like Biology without DNA!**



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$



Need a First Approximation to QCD

*Comparable in simplicity to
Schrödinger Theory in Atomic Physics*

Relativistic, Frame-Independent, Color-Confining

H_{QED}

QED atoms: positronium and muonium

Coupled Fock states

$$(H_0 + H_{int}) |\Psi\rangle = E |\Psi\rangle$$

$$\left[-\frac{\Delta^2}{2m_{red}} + V_{eff}(\vec{S}, \vec{r})\right] \psi(\vec{r}) = E \psi(\vec{r})$$

Effective two-particle equation

Includes Lamb Shift, quantum corrections

$$\left[-\frac{1}{2m_{red}} \frac{d^2}{dr^2} + \frac{1}{2m_{red}} \frac{\ell(\ell+1)}{r^2} + V_{eff}(r, S, \ell)\right] \psi(r) = E \psi(r)$$

Spherical Basis r, θ, ϕ

$$V_{eff} \rightarrow V_C(r) = -\frac{\alpha}{r}$$

Semiclassical first approximation to QED

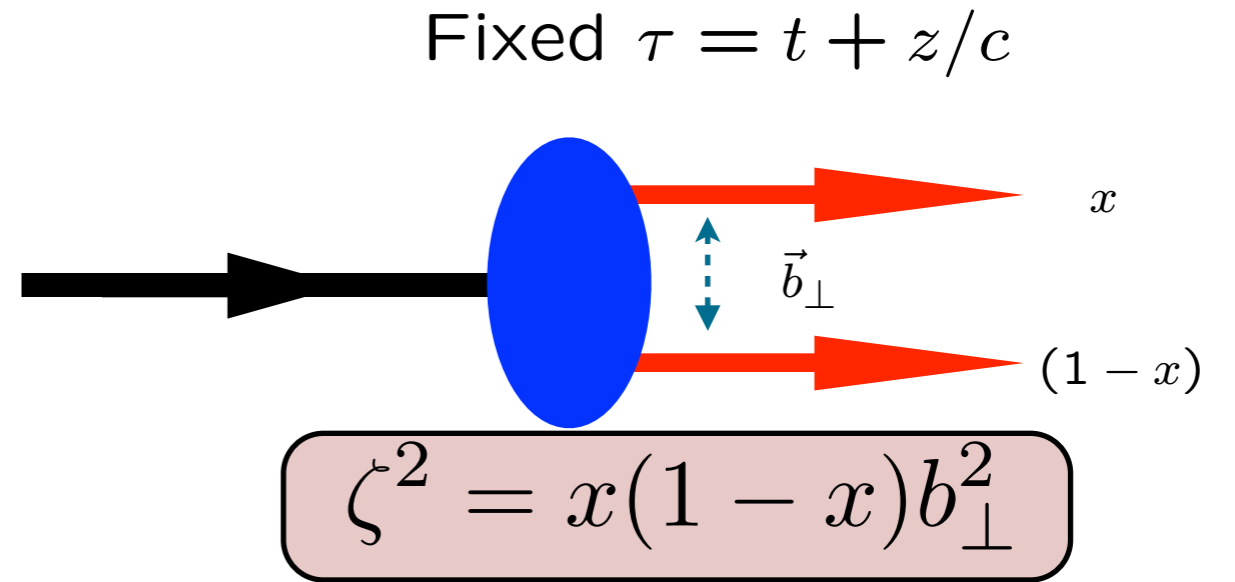
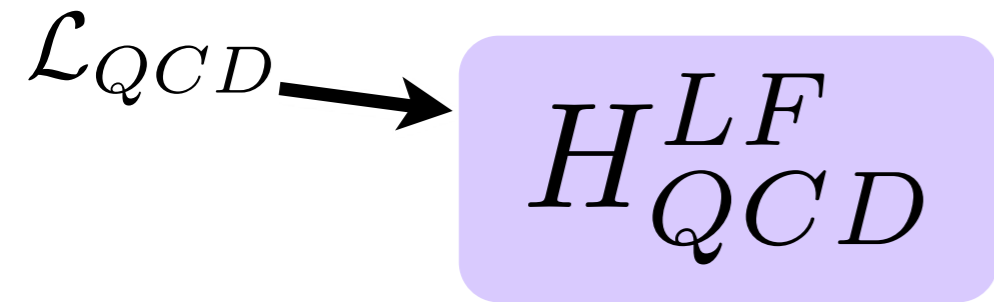


Coulomb potential

Bohr Spectrum

Schrödinger Eq.

Light-Front QCD



$$(H_{LF}^0 + H_{LF}^I) |\Psi\rangle = M^2 |\Psi\rangle$$

Coupled Fock states

Eliminate higher Fock states and retarded interactions

$$\left[\frac{\vec{k}_{\perp}^2 + m^2}{x(1-x)} + V_{\text{eff}}^{LF} \right] \psi_{LF}(x, \vec{k}_{\perp}) = M^2 \psi_{LF}(x, \vec{k}_{\perp})$$

Effective two-particle equation

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$

Azimuthal Basis

$$\zeta, \phi$$

$$m_q = 0$$

AdS/QCD:

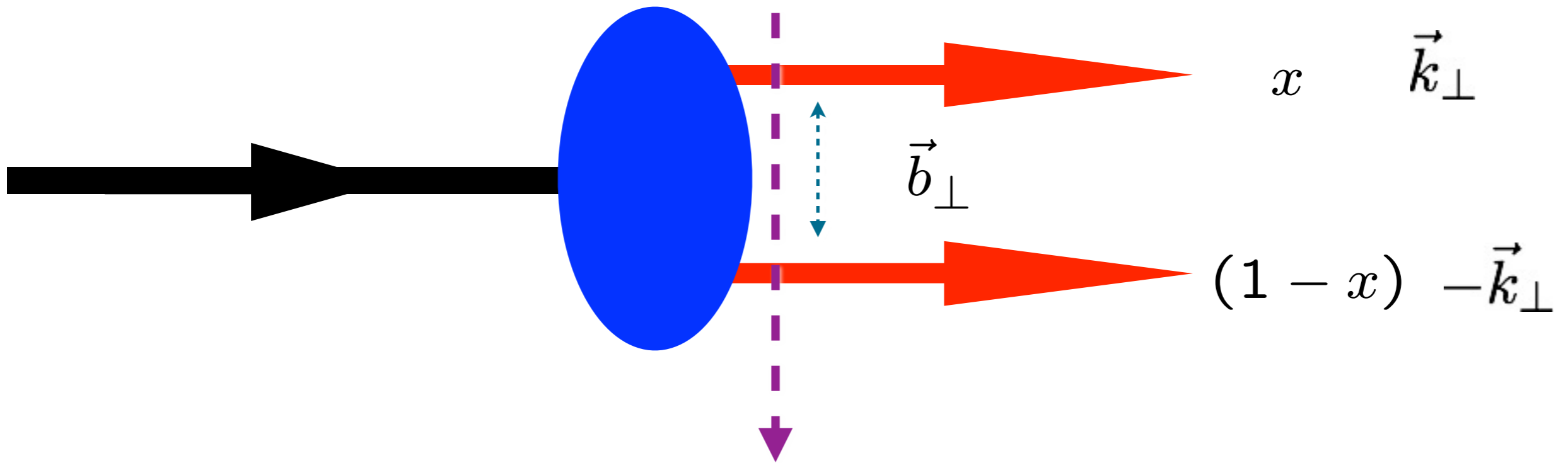
$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

Confining AdS/QCD potential!

Semiclassical first approximation to QCD

Sums an infinite # diagrams

Fixed $\tau = t + z/c$



$$\zeta^2 \equiv b_{\perp}^2 x(1-x)$$

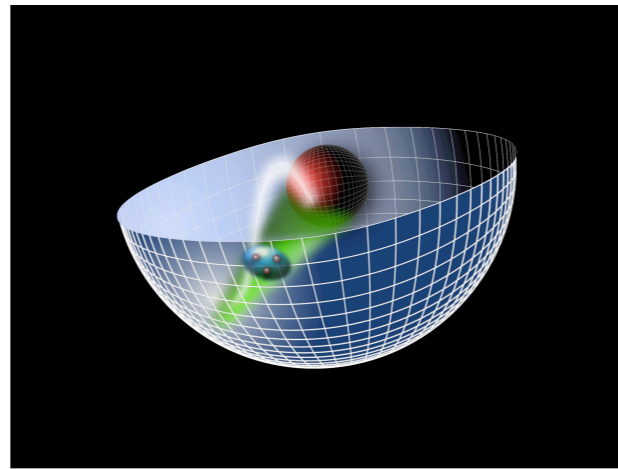
Invariant transverse separation

$$\zeta^2 \text{ conjugate to } \frac{k_{\perp}^2}{x(1-x)} = (p_q + p_{\bar{q}})^2 = \mathcal{M}_{q+\bar{q}}^2$$

$$\int dk^- \Psi_{BS}(P, k) \rightarrow \psi_{LF}(x, \vec{k}_{\perp})$$

*AdS/QCD
Soft-Wall Model*

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)b_{\perp}^2.$$

Light-Front Holography

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2(L + S - 1)$$

$$\kappa \simeq 0.6 \text{ GeV}$$

Confinement scale:

$$1/\kappa \simeq 1/3 \text{ fm}$$

***Unique
Confinement Potential!***

*Preserves Conformal Symmetry
of the action*

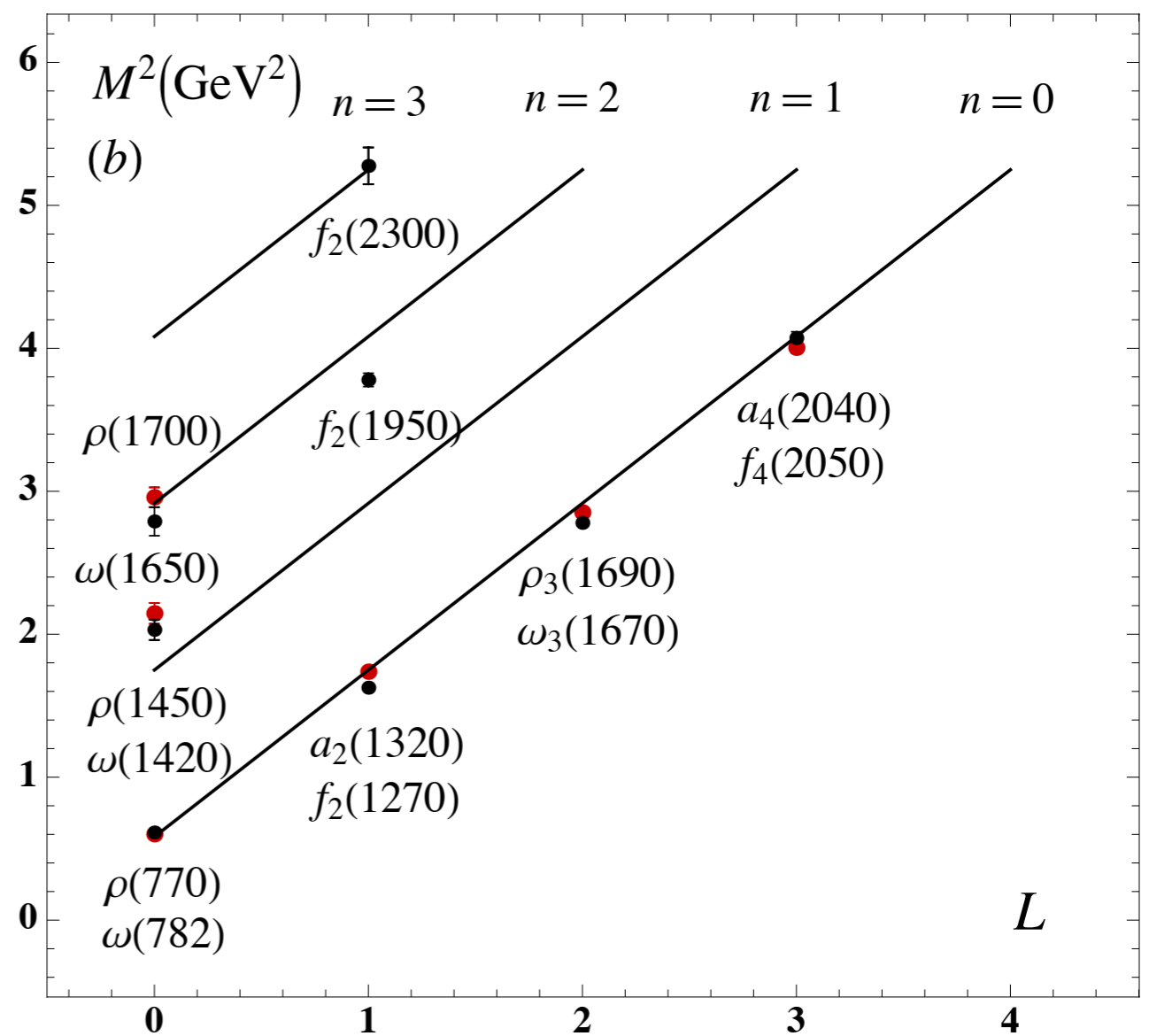
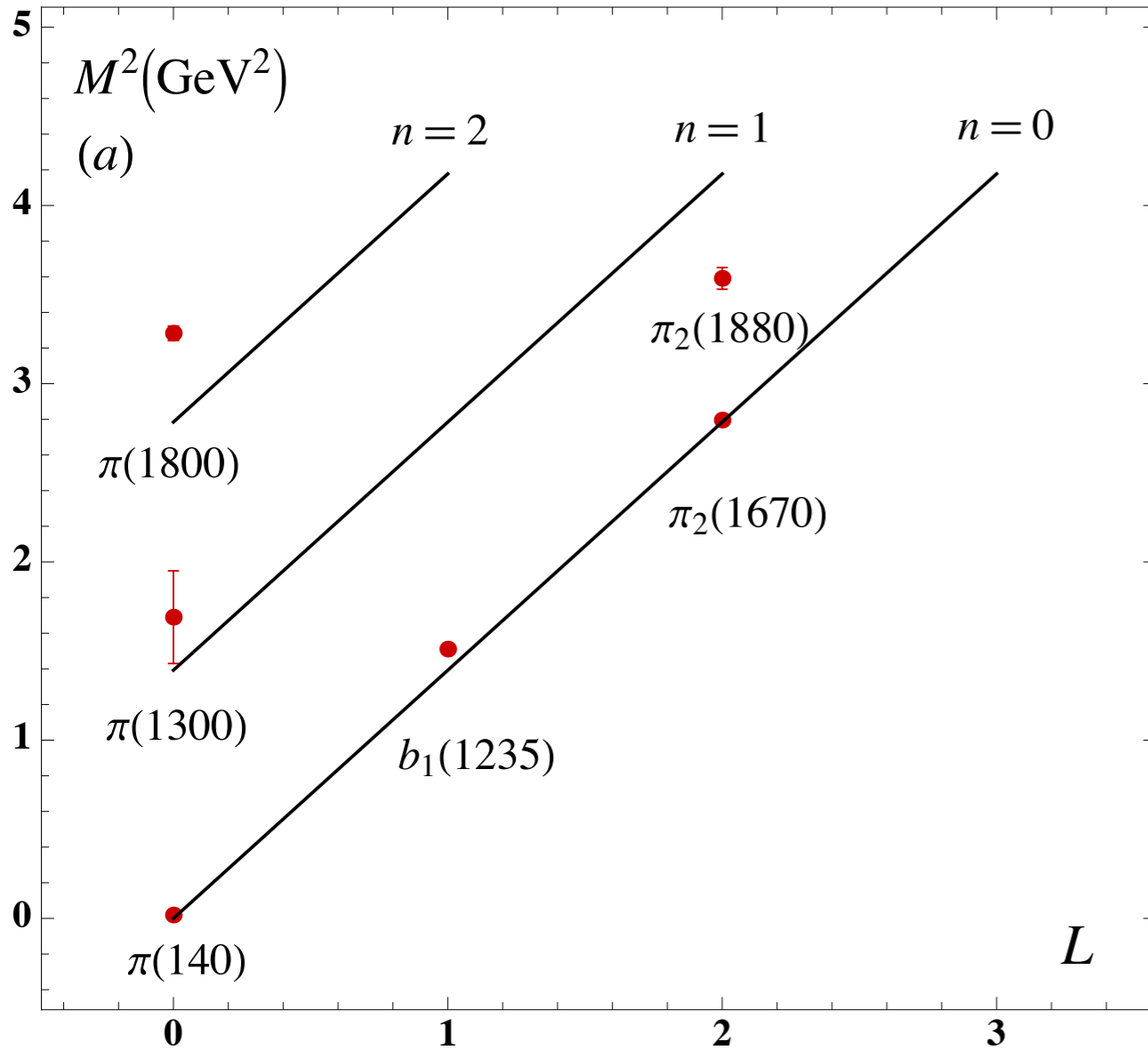
● **de Alfaro, Fubini, Furlan:**

● **Fubini, Rabinovici:**

**Scale can appear in Hamiltonian and EQM
without affecting conformal invariance of action!**

$$m_u = m_d = 0$$

Preview



$$M^2(n, L, S) = 4\kappa^2(n + L + S/2)$$

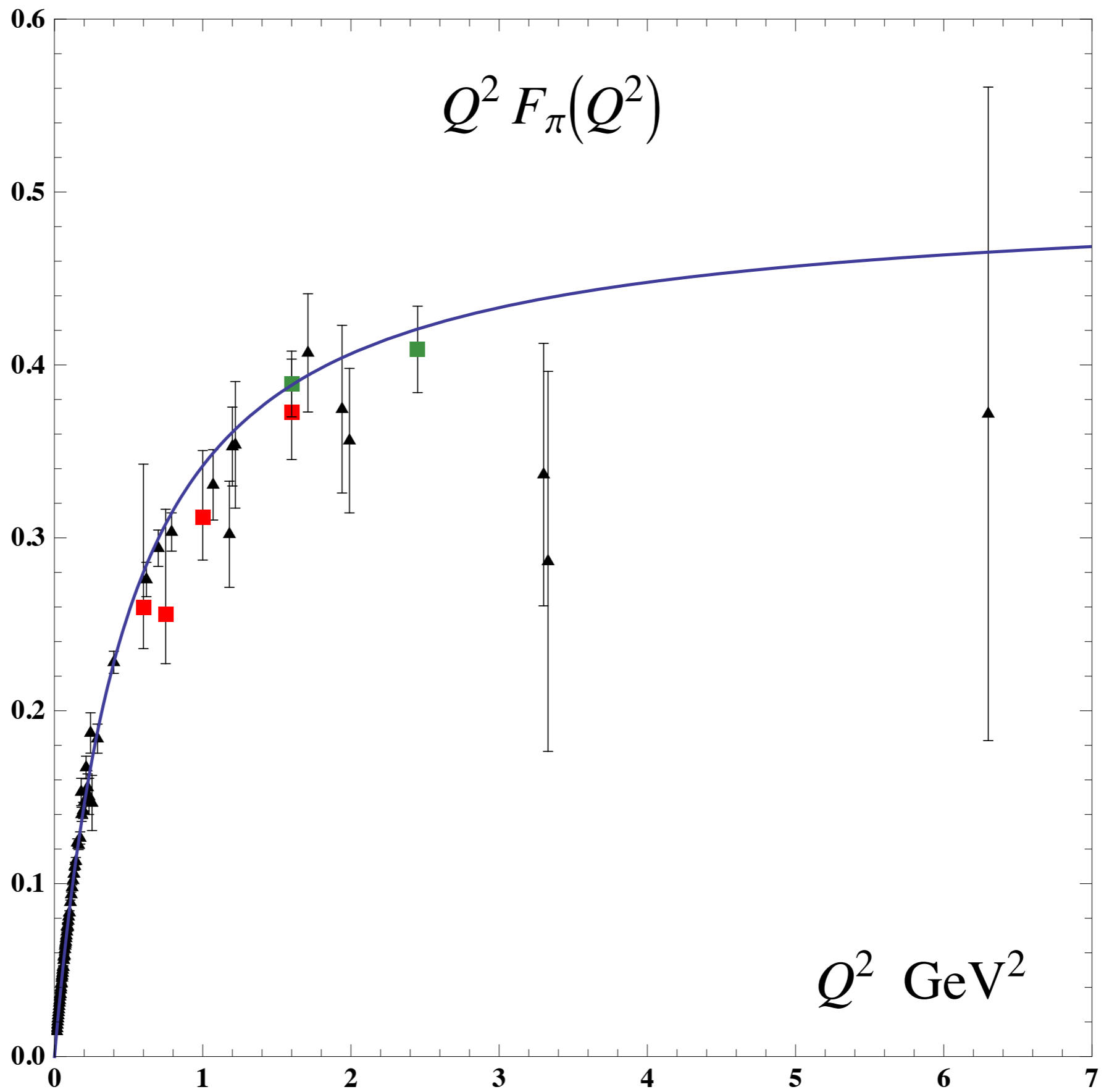
Kunshan
August 3, 2015

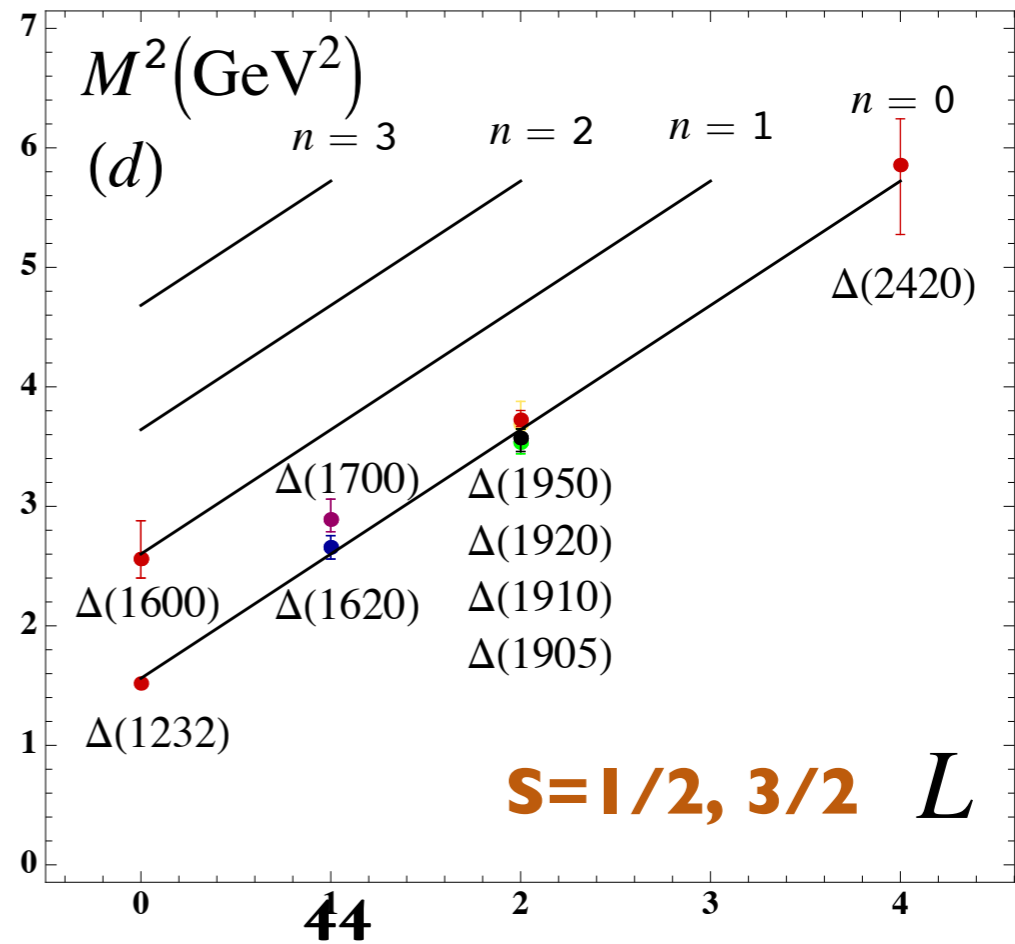
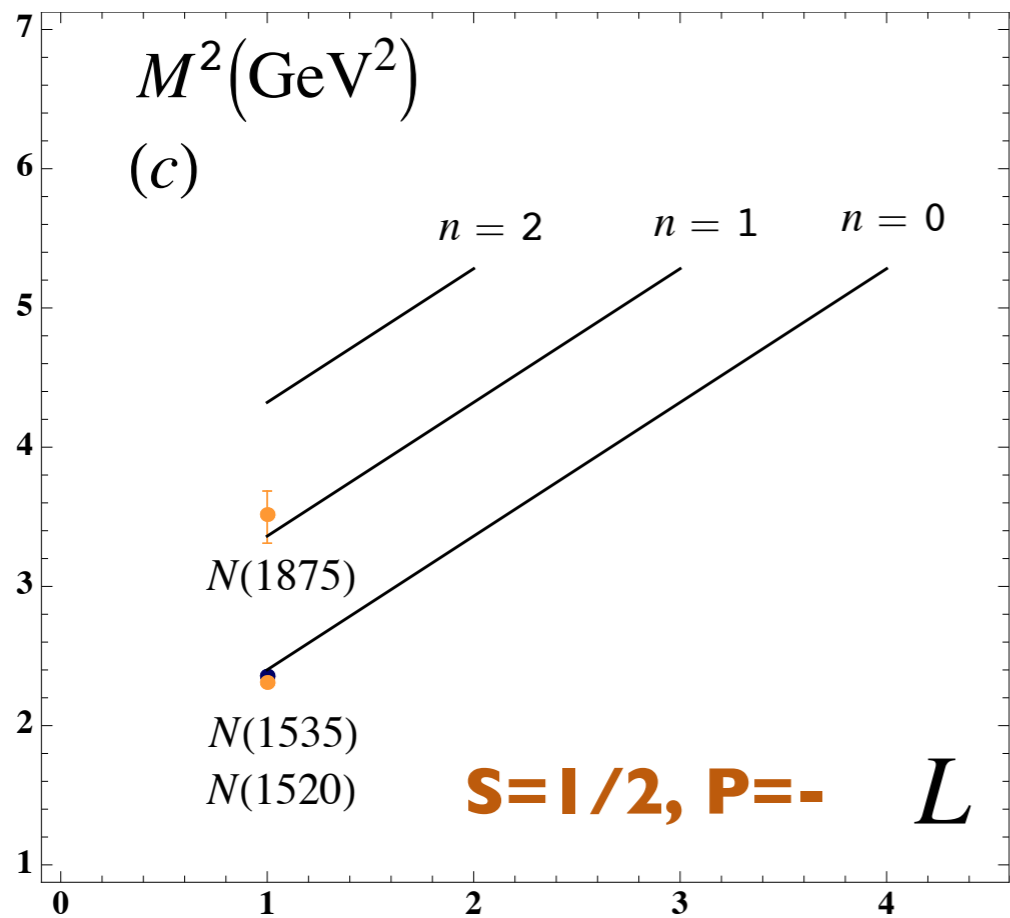
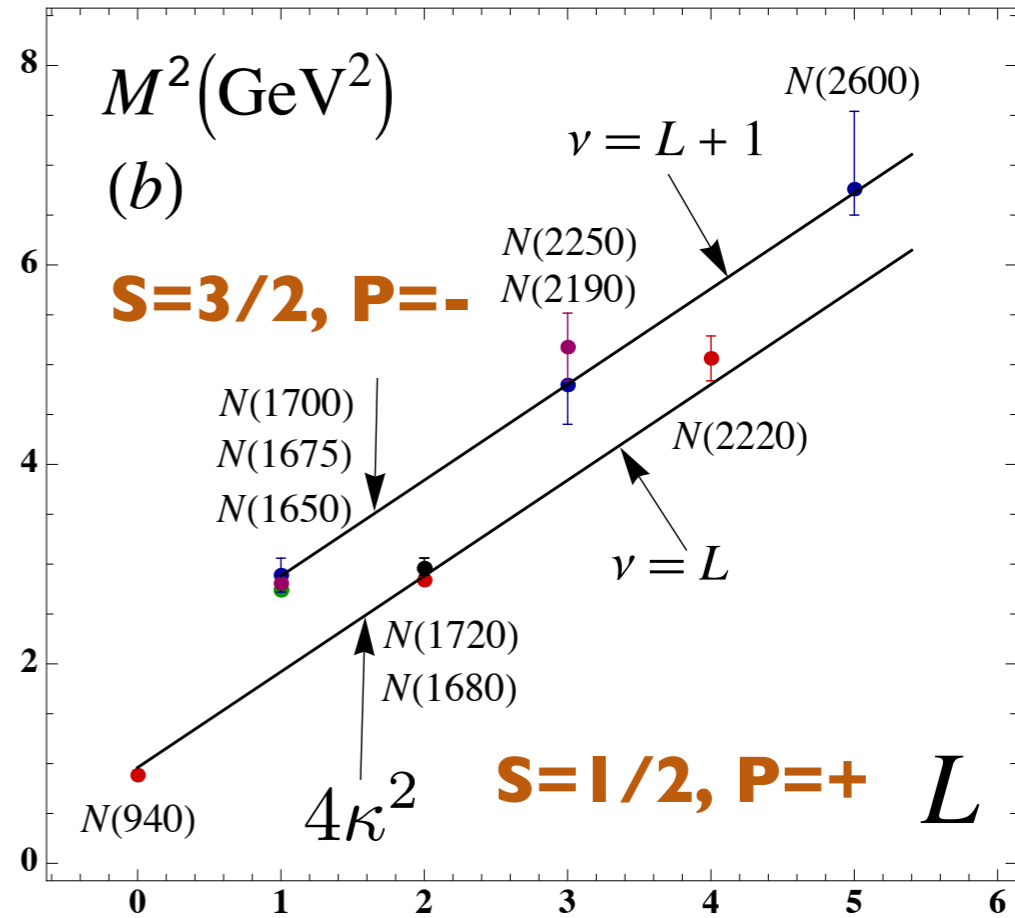
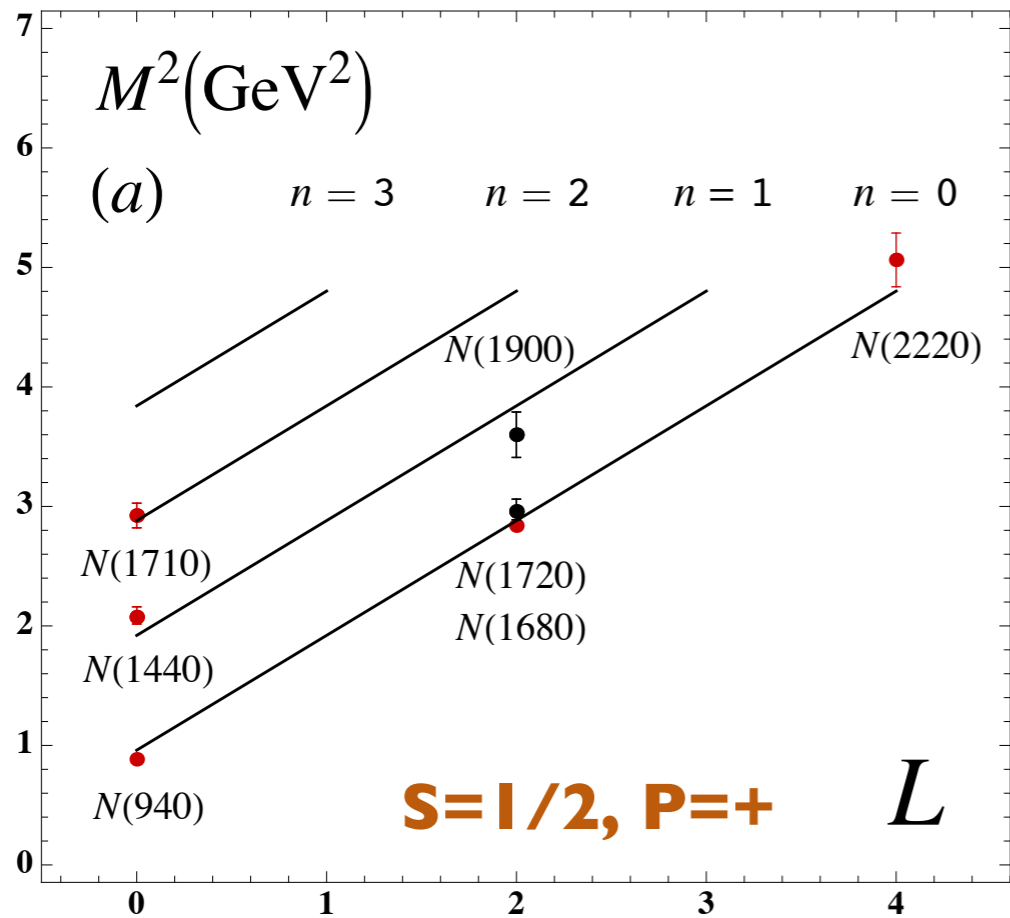


New Perspectives for Hadron Physics

Stan Brodsky



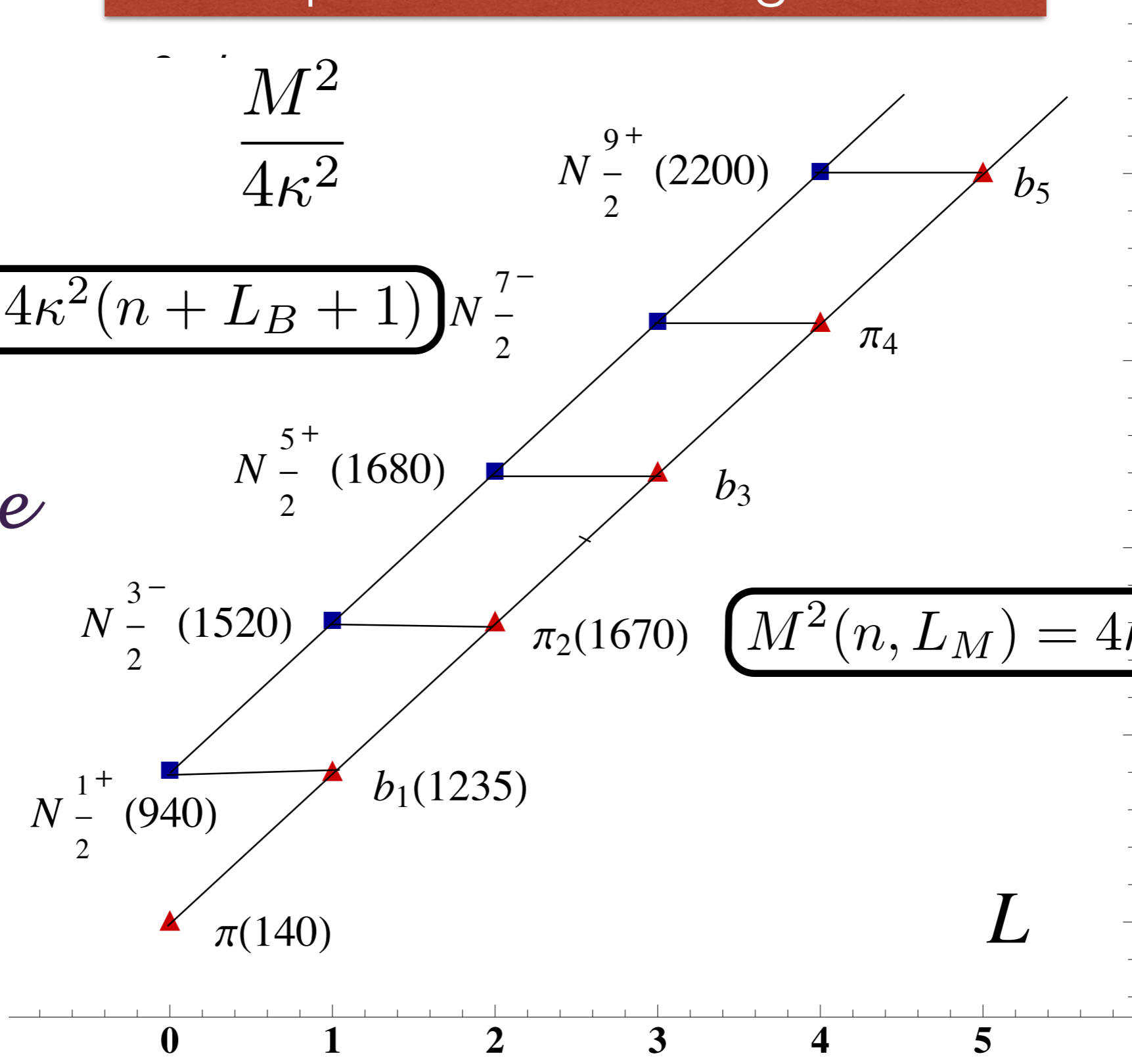




Superconformal Algebra

$$M^2(n, L_B) = 4\kappa^2(n + L_B + 1)$$

Same slope

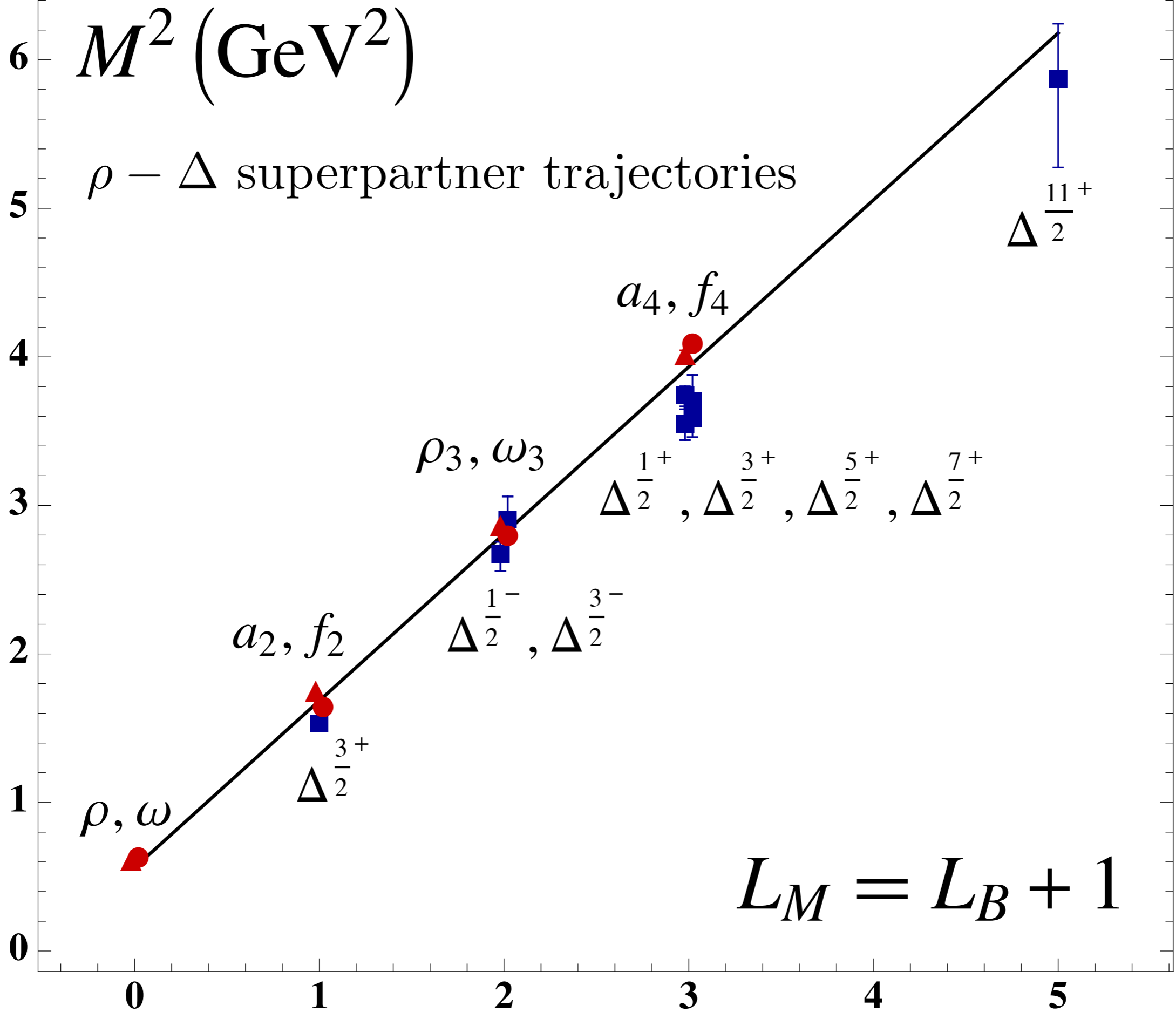


$$M^2(n, L_M) = 4\kappa^2(n + L_M)$$

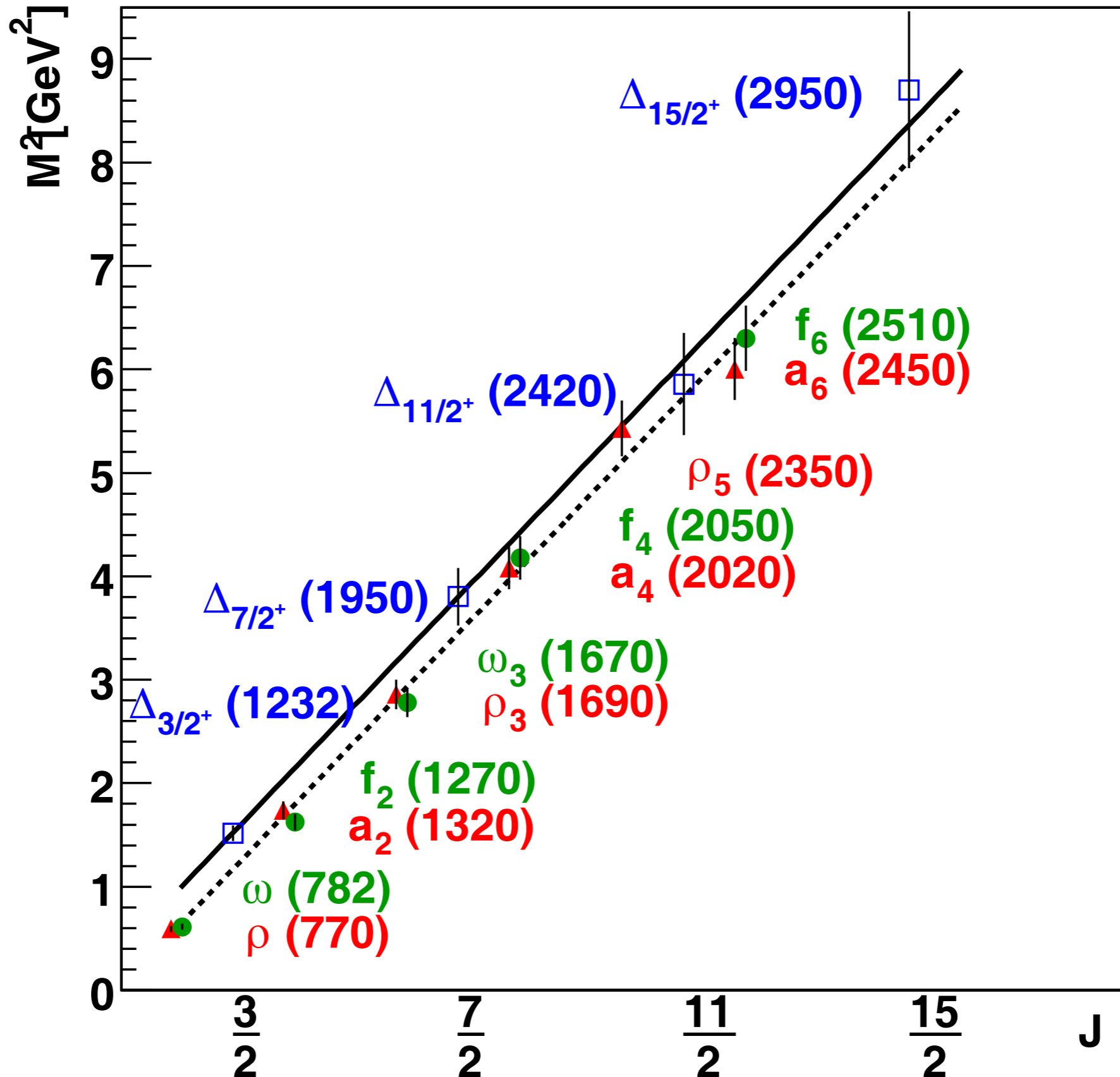
**Meson-Baryon
Mass Degeneracy
for $L_M=L_B+1$**

M^2 (GeV²)

$\rho - \Delta$ superpartner trajectories

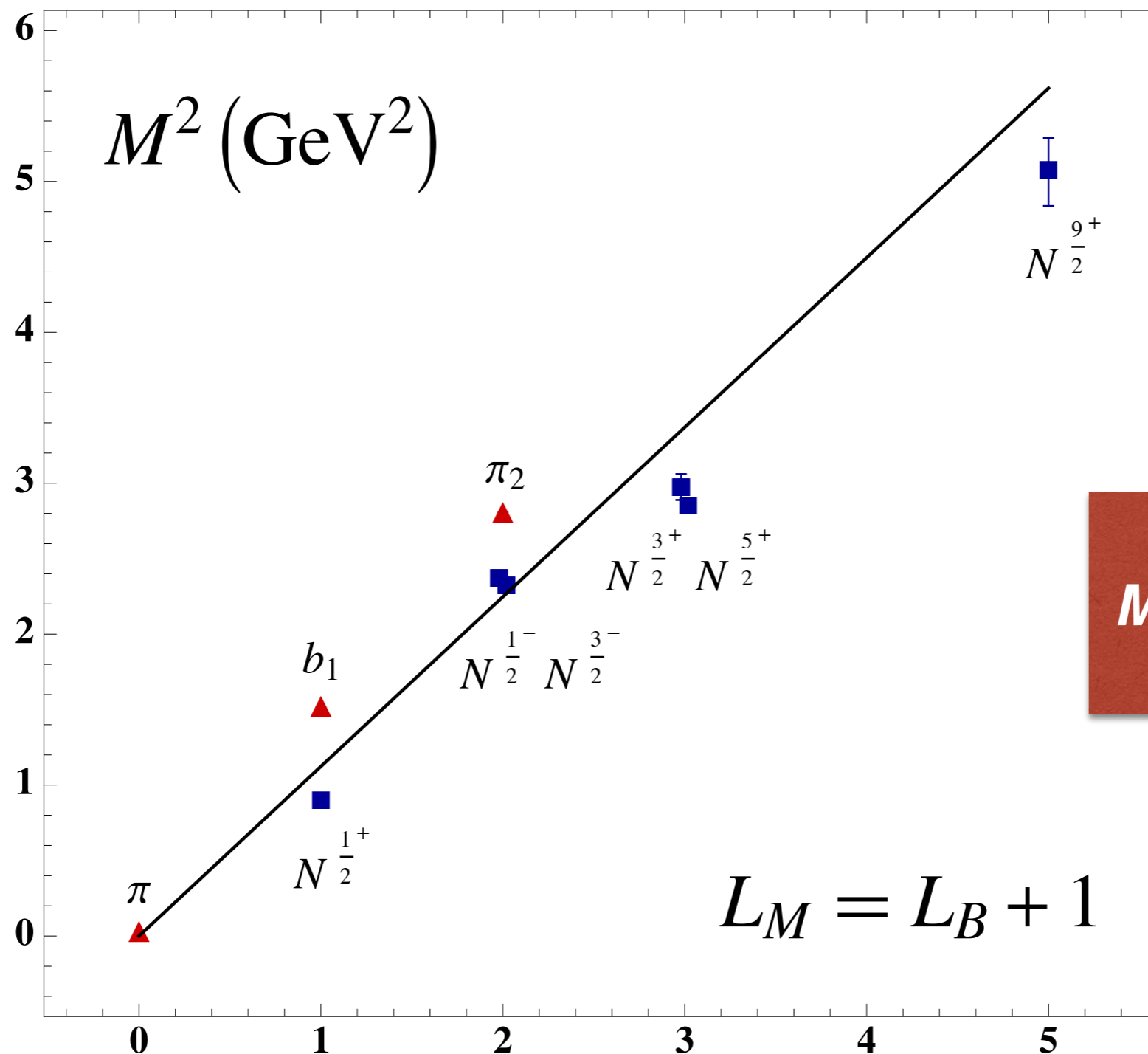


$$L_M = L_B + 1$$



The leading Regge trajectory: Δ resonances with maximal J in a given mass range. Also shown is the Regge trajectory for mesons with $|L+S|$.

Superconformal AdS Light-Front Holographic QCD (LFHQCD): Identical meson and baryon spectra!



$S=0, I=I$ Meson is superpartner of $S=1/2, I=I$ Baryon

LF Holography

Baryon Equation

$$\left(-\partial_{\zeta}^2 + \kappa^4 \zeta^2 + 2\kappa^2(L_B + 1) + \frac{4L_B^2 - 1}{4\zeta^2} \right) \psi_J^+ = M^2 \psi_J^+ \quad \text{G}_{22}$$

$$\left(-\partial_{\zeta}^2 + \kappa^4 \zeta^2 + 2\kappa^2 L_B + \frac{4(L_B + 1)^2 - 1}{4\zeta^2} \right) \psi_J^- = M^2 \psi_J^- \quad \text{G}_{11}$$

$$M^2(n, L_B) = 4\kappa^2(n + L_B + 1) \quad \mathbf{S=1/2, P=+}$$

both chiralities

Meson Equation

$$\left(-\partial_{\zeta}^2 + \kappa^4 \zeta^2 + 2\kappa^2(J - 1) + \frac{4L_M^2 - 1}{4\zeta^2} \right) \phi_J = M^2 \phi_J \quad \text{G}_{11}$$

$$M^2(n, L_M) = 4\kappa^2(n + L_M) \quad \text{Same } \kappa !$$

S=0, I=I Meson is superpartner of S=1/2, I=I Baryon

Meson-Baryon Degeneracy for $L_M=L_B+1$

Some Features of AdS/QCD

- **Regge spectroscopy—same slope in n, L for mesons,**
- **Chiral features for $m_q=0$: $m_\pi=0$, chiral-invariant proton**
- **Hadronic LFWFs**
- **Counting Rules**
- **Connection between hadron masses and $\Lambda_{\overline{MS}}$**

Superconformal AdS Light-Front Holographic QCD (LFHQCD)

Meson-Baryon Mass Degeneracy for $L_M=L_B+1$



$$m_\rho = \sqrt{2}\kappa$$

$$m_p = 2\kappa_1$$

Deur, de Tèramond, sjb

All-Scale QCD Coupling

$$\frac{\alpha_{g_1}^s(Q^2)}{\pi}$$

$$e^{-\frac{Q^2}{4\kappa^2}}$$

Expt:
 $\Lambda_{\overline{MS}} = 0.341 \pm 0.024 \text{ GeV}$

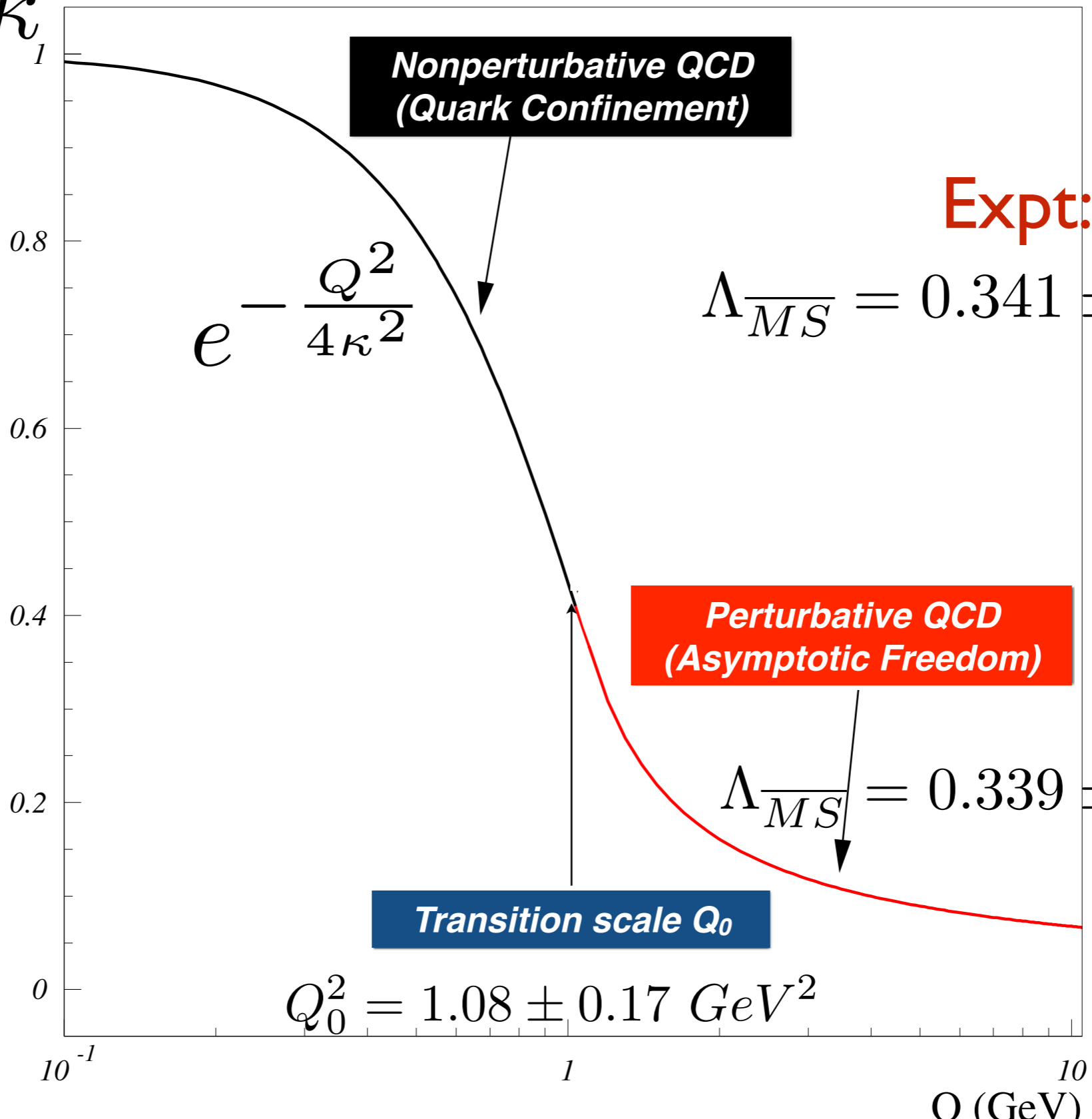
**Perturbative QCD
(Asymptotic Freedom)**

$$\Lambda_{\overline{MS}} = 0.339 \pm 0.016 \text{ GeV}$$

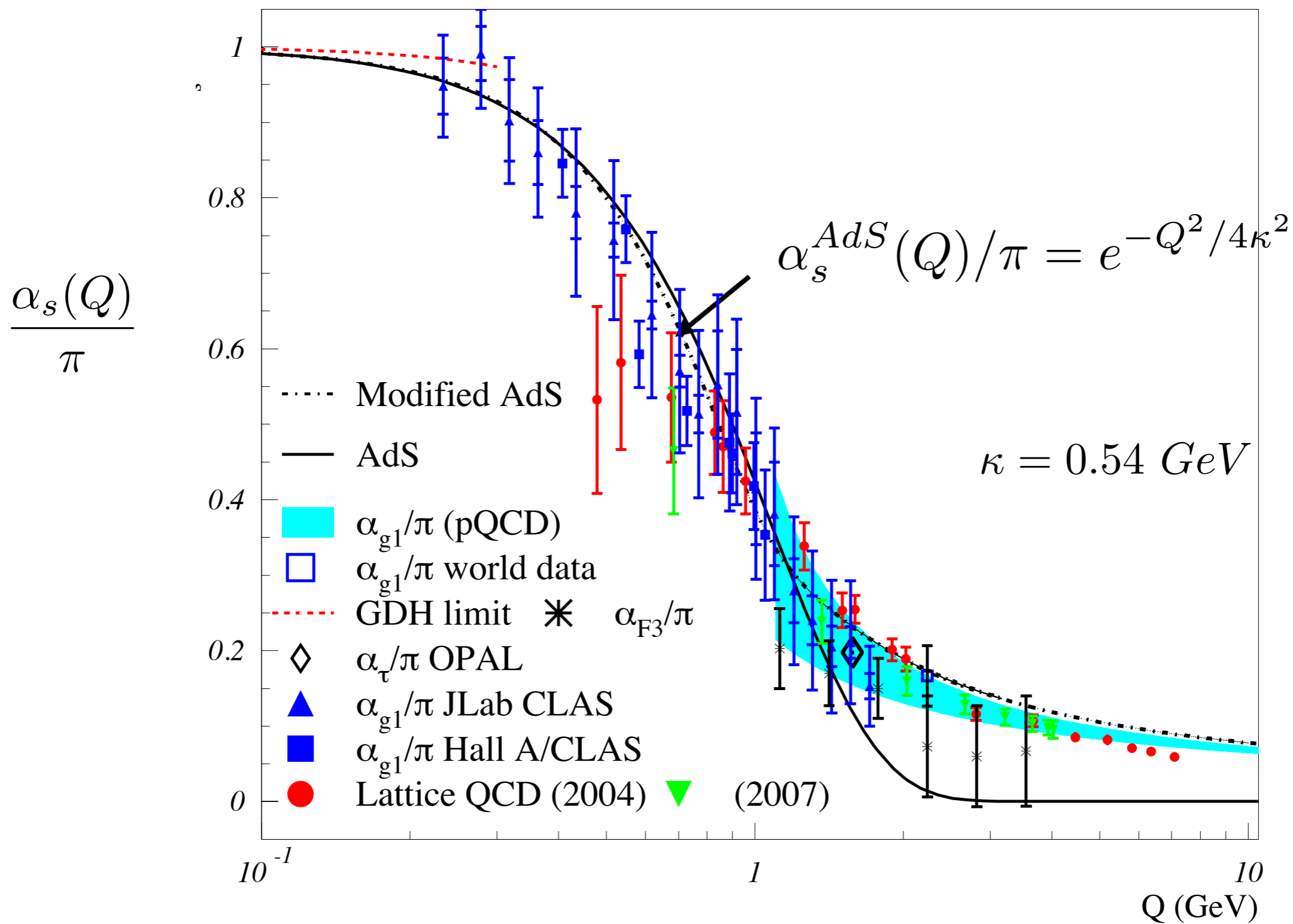
Transition scale Q_0

$$Q_0^2 = 1.08 \pm 0.17 \text{ GeV}^2$$

$$\lambda \equiv \kappa^2$$

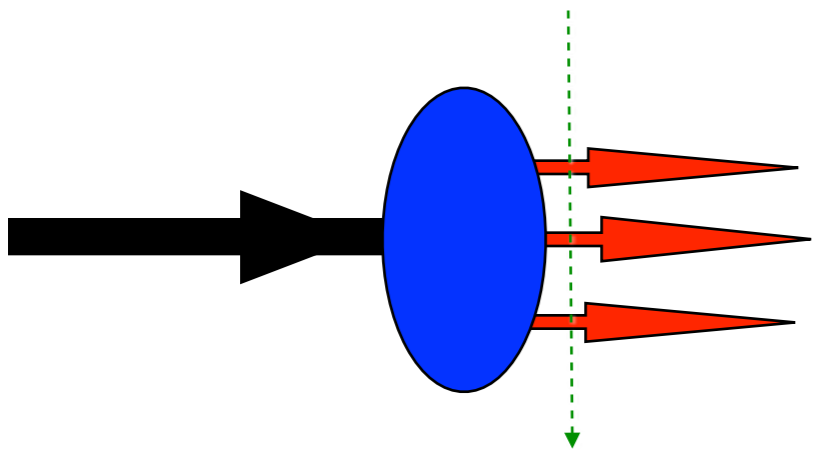


Analytic, defined at all scales, IR Fixed Point



$$e^\varphi = e^{+\kappa^2 z^2}$$

Dirac's Front Form: Fixed $\tau = t + z/c$



$$\psi(x_i, \vec{k}_{\perp i}, \lambda_i)$$

$$x_i = \frac{k_i^+}{P^+}$$

Invariant under boosts. Independent of P^μ

$$H_{LF}^{QCD} |\psi\rangle = M^2 |\psi\rangle$$

Direct connection to QCD Lagrangian

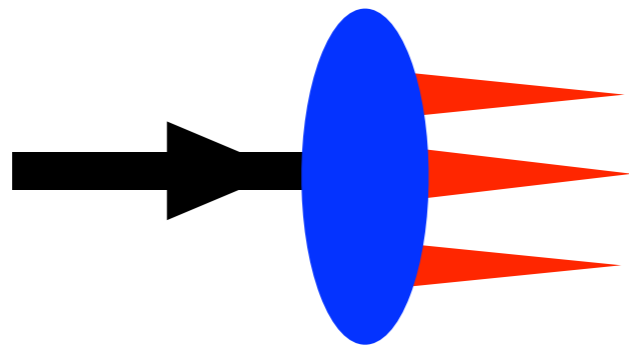
Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space



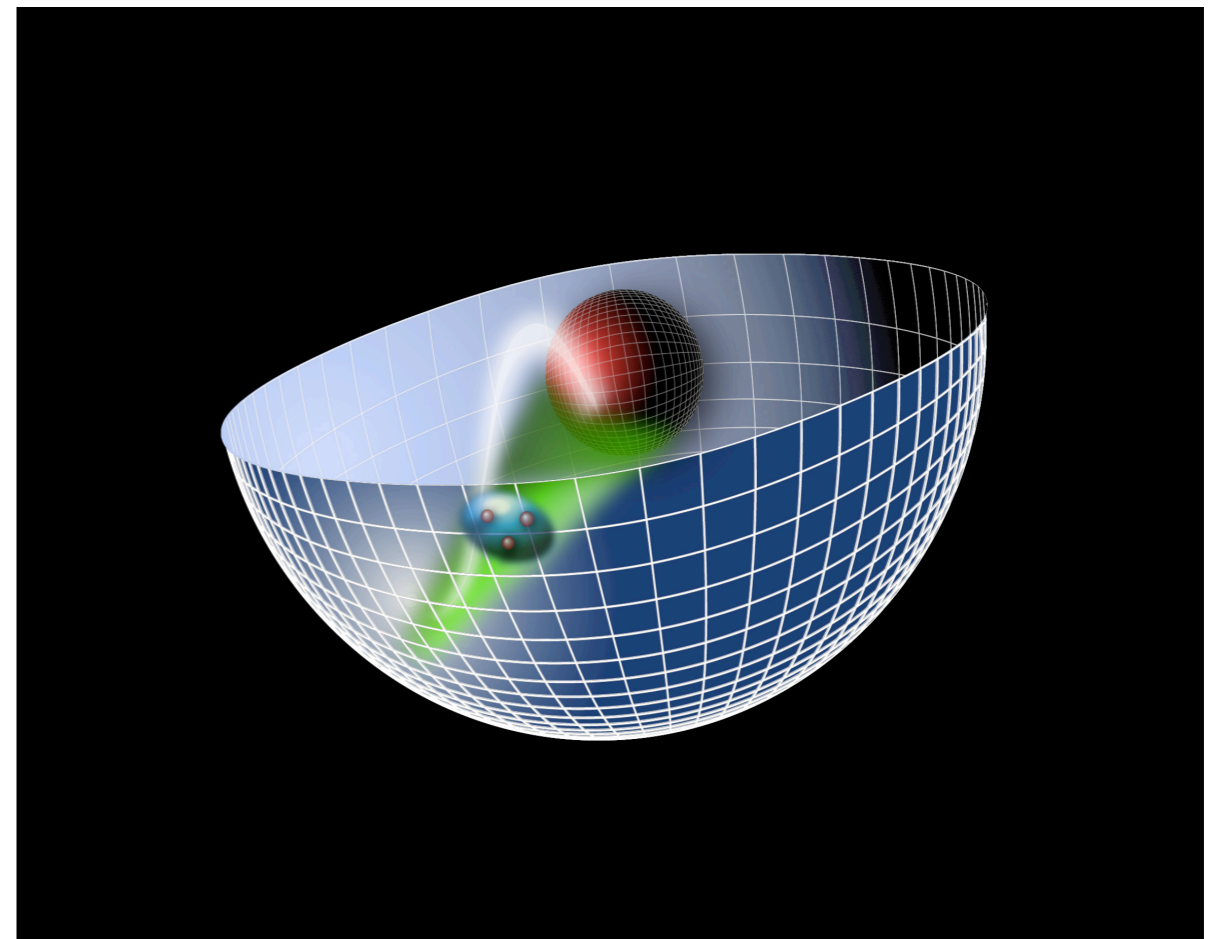
Light-Front Holography and Non-Perturbative QCD

Goal:
**Use AdS/QCD duality to construct
a first approximation to QCD**

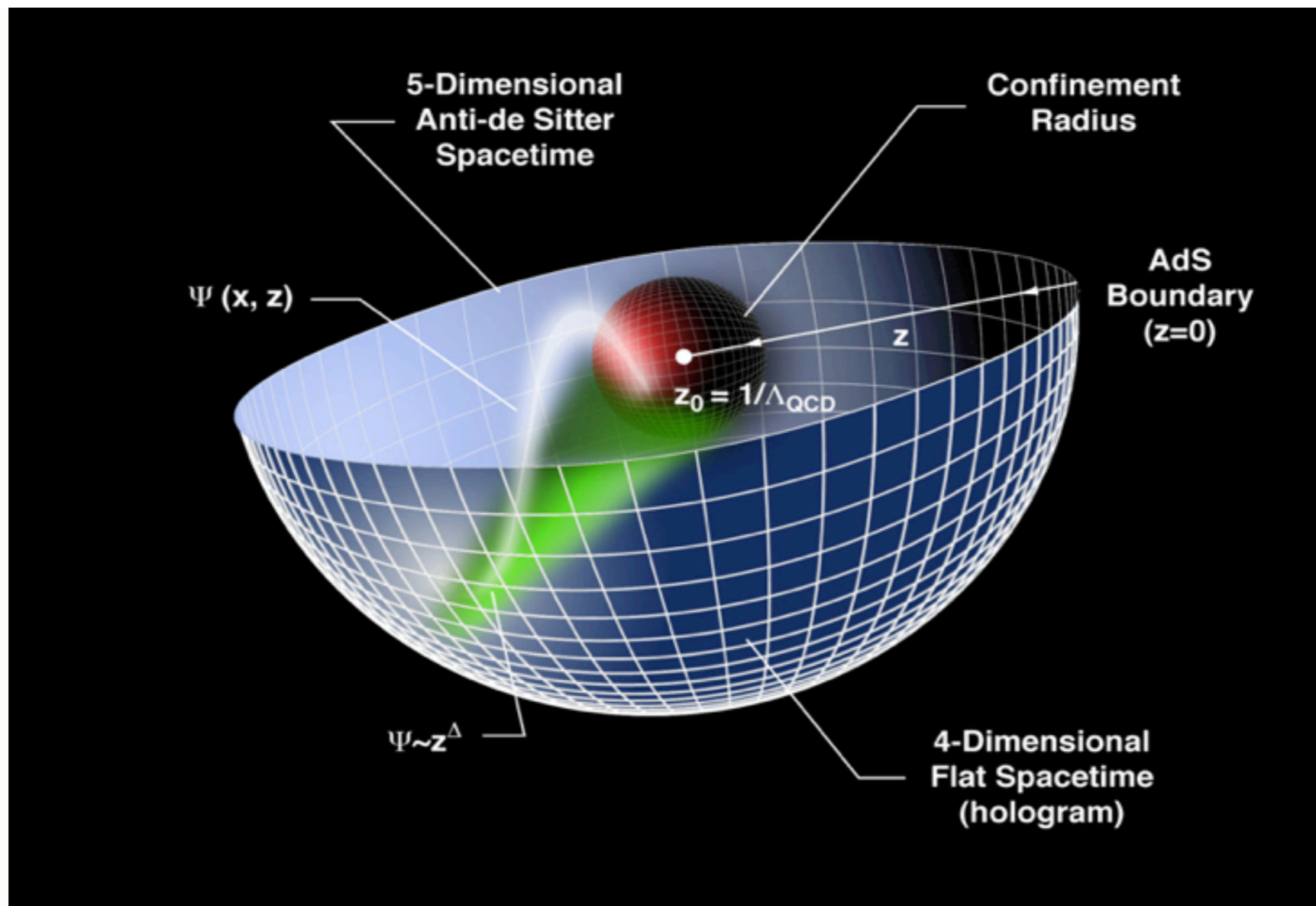
*Hadron Spectrum
Light-Front Wavefunctions,
Form Factors, DVCS, etc*



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$



in collaboration with Guy de Teramond and H. Guenter Dosch



Changes in physical length scale mapped to evolution in the 5th dimension z

8-2007
8685A14

- Truncated AdS/CFT (Hard-Wall) model: cut-off at $z_0 = 1/\Lambda_{\text{QCD}}$ breaks conformal invariance and allows the introduction of the QCD scale (Hard-Wall Model) **Polchinski and Strassler (2001)**.
- Smooth cutoff: introduction of a background dilaton field $\varphi(z)$ – usual linear Regge dependence can be obtained (Soft-Wall Model) **Karch, Katz, Son and Stephanov (2006)**.

Kunshan
August 3, 2015



New Perspectives for Hadron Physics


Stan Brodsky

SLAC
NATIONAL ACCELERATOR LABORATORY

AdS/CFT

- Isomorphism of $SO(4, 2)$ of conformal QCD with the group of isometries of AdS space

$$ds^2 = \frac{R^2}{z^2} (\eta_{\mu\nu} dx^\mu dx^\nu - dz^2),$$

invariant measure 

$x^\mu \rightarrow \lambda x^\mu, z \rightarrow \lambda z$, maps scale transformations into the holographic coordinate z .

- AdS mode in z is the extension of the hadron wf into the fifth dimension.
- Different values of z correspond to different scales at which the hadron is examined.

$$x^2 \rightarrow \lambda^2 x^2, \quad z \rightarrow \lambda z.$$

$x^2 = x_\mu x^\mu$: invariant separation between quarks

- The AdS boundary at $z \rightarrow 0$ correspond to the $Q \rightarrow \infty$, UV zero separation limit.

Dilaton-Modified AdS/QCD

$$ds^2 = e^{\varphi(z)} \frac{R^2}{z^2} (\eta_{\mu\nu} x^\mu x^\nu - dz^2)$$

- **Soft-wall dilaton profile breaks conformal invariance** $e^{\varphi(z)} = e^{+\kappa^2 z^2}$
- **Color Confinement**
- **Introduces confinement scale** κ
- **Uses AdS₅ as template for conformal theory**



$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$

Positive-sign dilaton

• Dosch, de Teramond, sjb

AdS Soft-Wall Schrodinger Equation for bound state of two scalar constituents:

$$\left[-\frac{d^2}{dz^2} - \frac{1 - 4L^2}{4z^2} + U(z) \right] \Phi(z) = \mathcal{M}^2 \Phi(z)$$

$$U(z) = \kappa^4 z^2 + 2\kappa^2 (L + S - 1)$$

Derived from variation of Action for Dilaton-Modified AdS₅

Identical to Light-Front Bound State Equation!

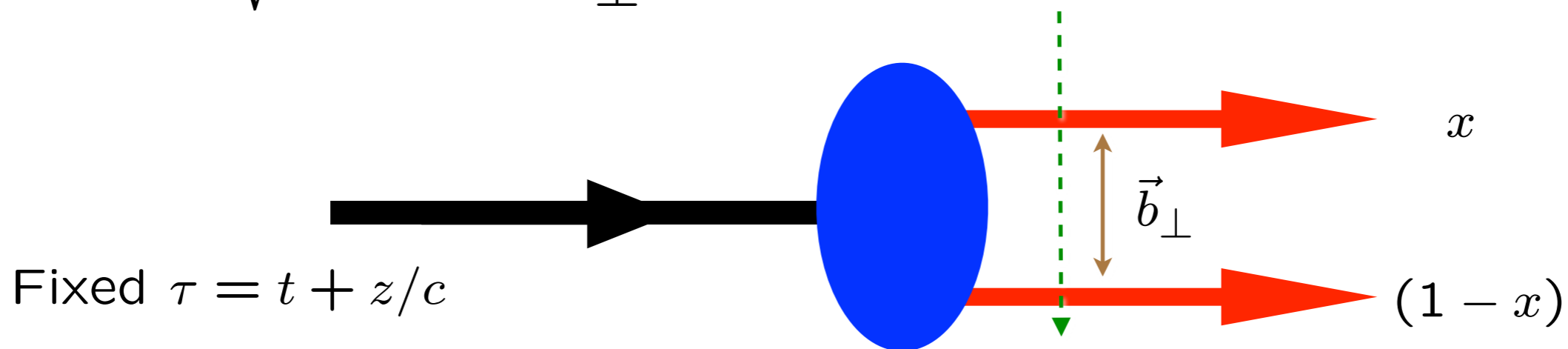
$$z \quad \longleftrightarrow \quad \zeta = \sqrt{x(1-x)\vec{b}_\perp^2}$$

$$LF(3+1) \longleftrightarrow AdS_5$$

Light-Front Holographic Dictionary

$$\psi(x, \vec{b}_\perp) \longleftrightarrow \phi(z)$$

$$\zeta = \sqrt{x(1-x)} \vec{b}_\perp^2 \longleftrightarrow z$$



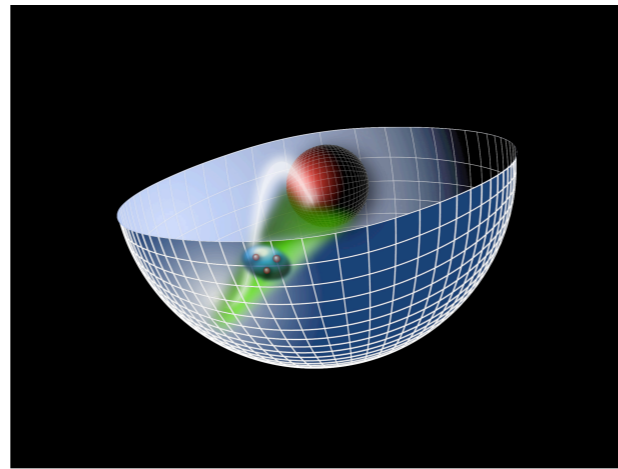
$$\psi(x, \zeta) = \sqrt{x(1-x)} \zeta^{-1/2} \phi(\zeta)$$

$$(\mu R)^2 = L^2 - (J - 2)^2$$

Light-Front Holography: Unique mapping derived from equality of LF and AdS formula for EM and gravitational current matrix elements and identical equations of motion

*AdS/QCD
Soft-Wall Model*

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)b_{\perp}^2.$$

Light-Front Holography

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2(L + S - 1)$$

$$\kappa \simeq 0.6 \text{ GeV}$$

Confinement scale:

$$1/\kappa \simeq 1/3 \text{ fm}$$

***Unique
Confinement Potential!***

*Preserves Conformal Symmetry
of the action*

● **de Alfaro, Fubini, Furlan:**

● **Fubini, Rabinovici:**

**Scale can appear in Hamiltonian and EQM
without affecting conformal invariance of action!**

Meson Spectrum in Soft Wall Model

Pion: Negative term for $J=0$ cancels positive terms from LFKE and potential



- Effective potential: $U(\zeta^2) = \kappa^4 \zeta^2 + 2\kappa^2(J - 1)$

- LF WE

$$\left(-\frac{d^2}{d\zeta^2} - \frac{1 - 4L^2}{4\zeta^2} + \kappa^4 \zeta^2 + 2\kappa^2(J - 1) \right) \phi_J(\zeta) = M^2 \phi_J(\zeta)$$

- Normalized eigenfunctions $\langle \phi | \phi \rangle = \int d\zeta \phi^2(z)^2 = 1$

$$\phi_{n,L}(\zeta) = \kappa^{1+L} \sqrt{\frac{2n!}{(n+L)!}} \zeta^{1/2+L} e^{-\kappa^2 \zeta^2 / 2} L_n^L(\kappa^2 \zeta^2)$$

- Eigenvalues

$$\mathcal{M}_{n,J,L}^2 = 4\kappa^2 \left(n + \frac{J+L}{2} \right)$$

G. de Teramond, H. G. Dosch, sjb

Quark separation increases with L

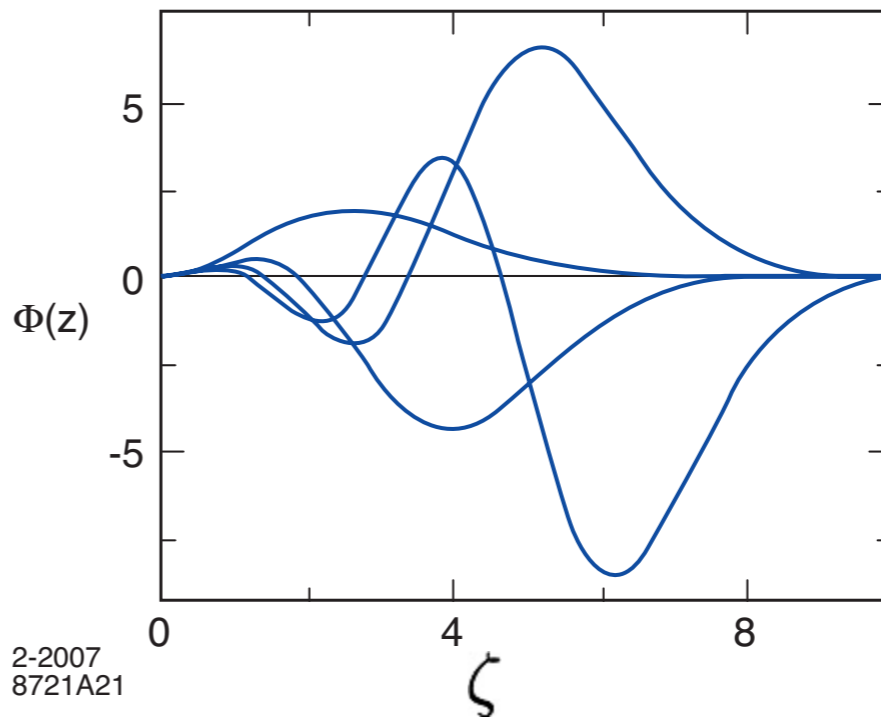
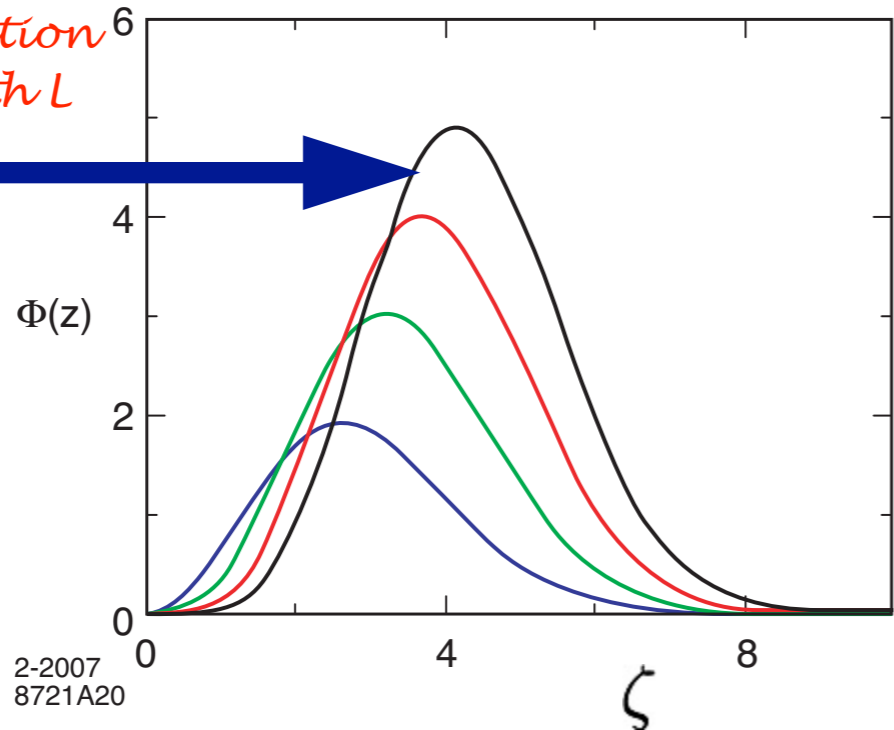
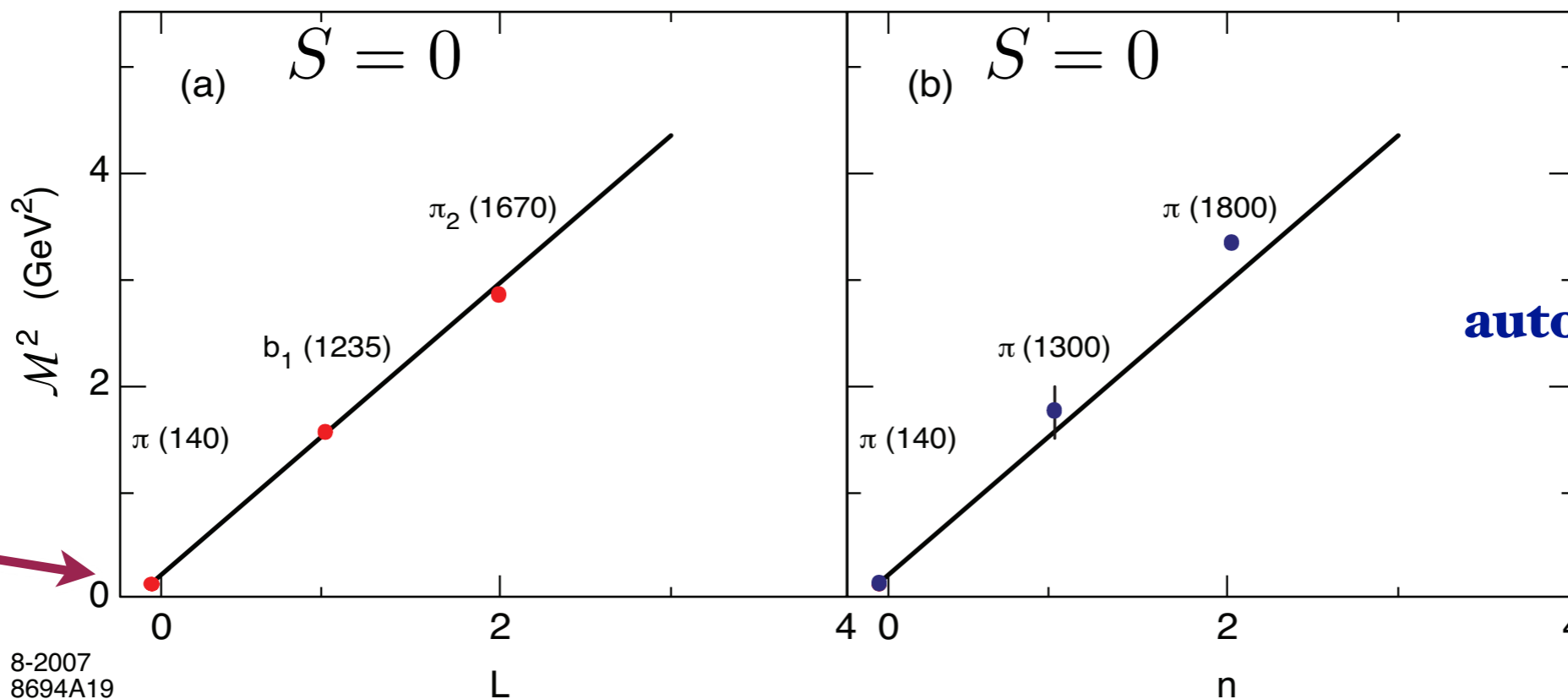


Fig: Orbital and radial AdS modes in the soft wall model for $\kappa = 0.6$ GeV .

Same slope in n and L !

Soft Wall Model



Pion has zero mass!

Pion mass automatically zero!

$$m_q = 0$$

Light meson orbital (a) and radial (b) spectrum for $\kappa = 0.6$ GeV.

- $J = L + S, I = 1$ meson families

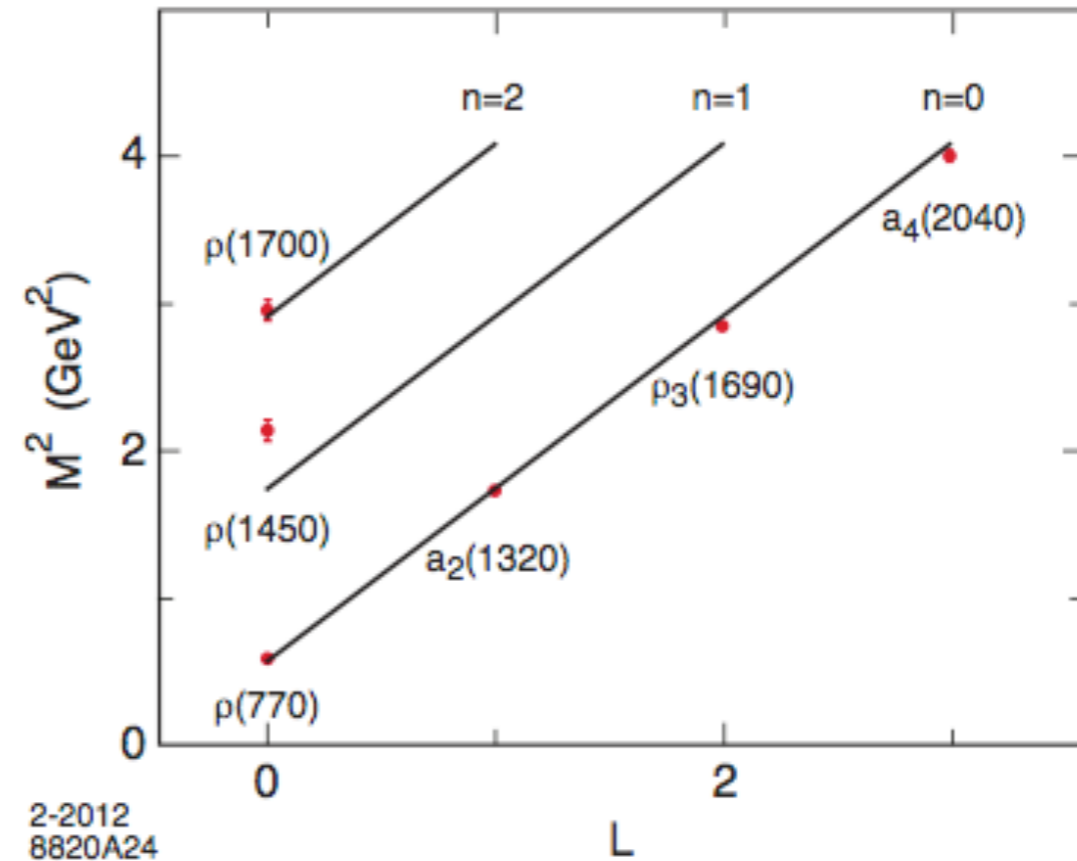
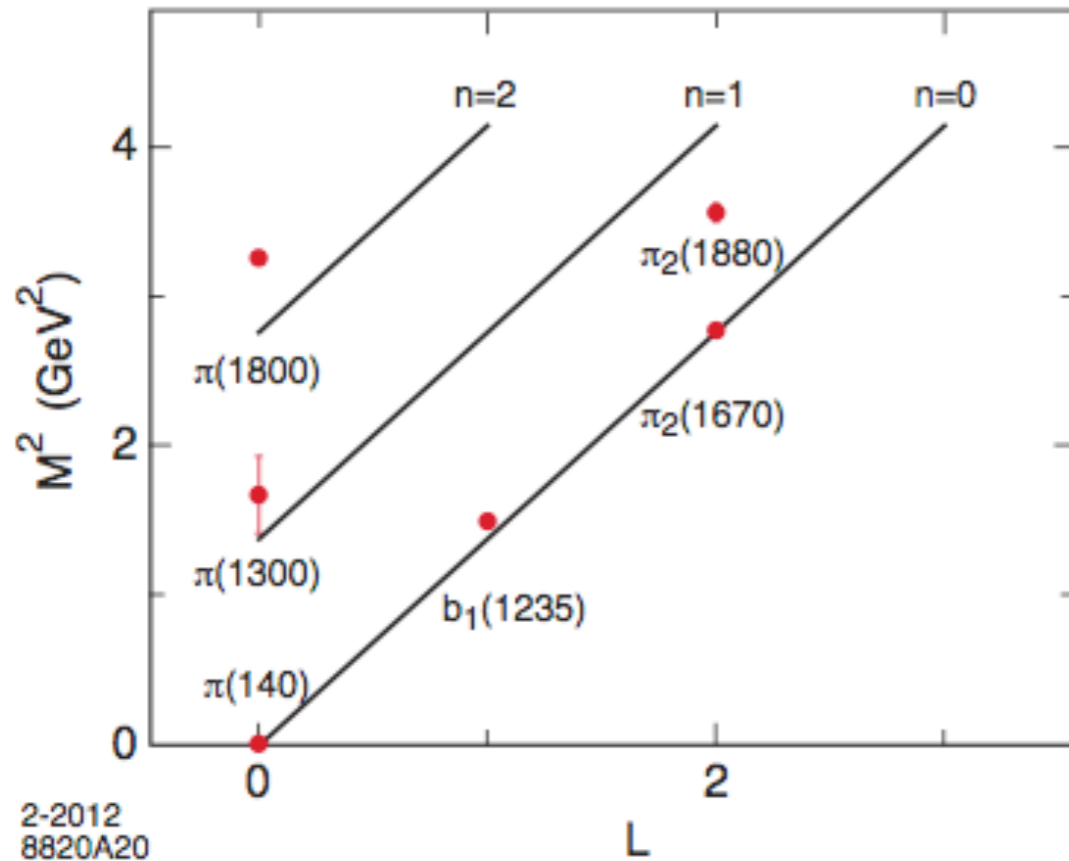
$$\mathcal{M}_{n,L,S}^2 = 4\kappa^2 (n + L + S/2)$$

$$\begin{aligned} 4\kappa^2 &\text{ for } \Delta n = 1 \\ 4\kappa^2 &\text{ for } \Delta L = 1 \\ 2\kappa^2 &\text{ for } \Delta S = 1 \end{aligned}$$

$$m_q = 0$$

Massless pion in Chiral Limit!

Same slope in n and L !



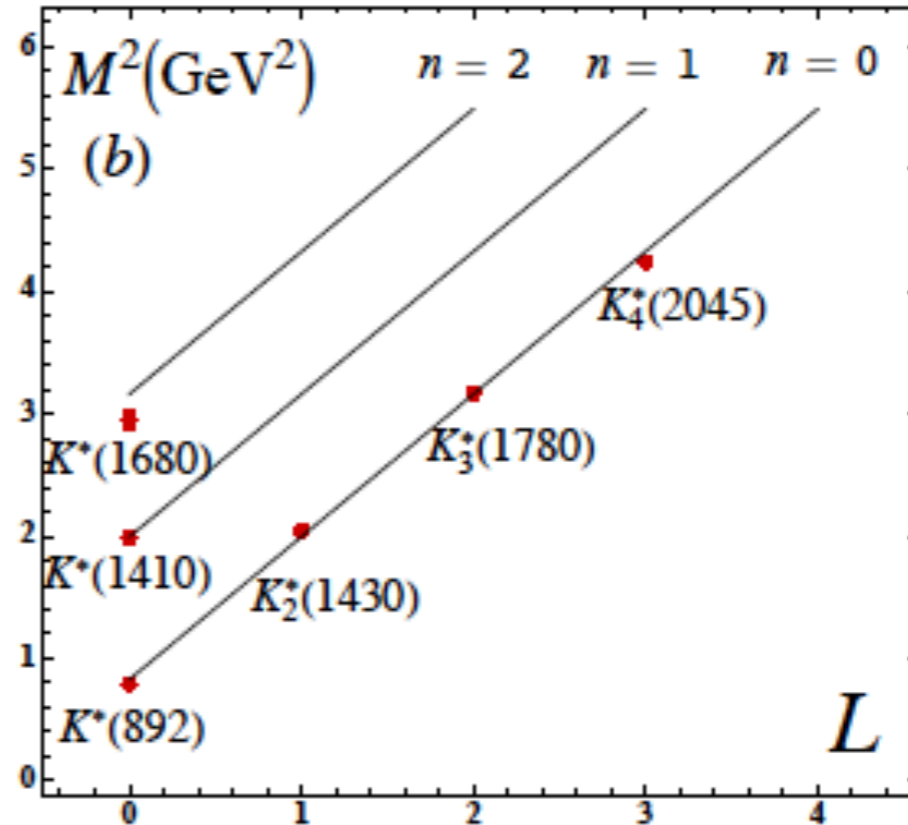
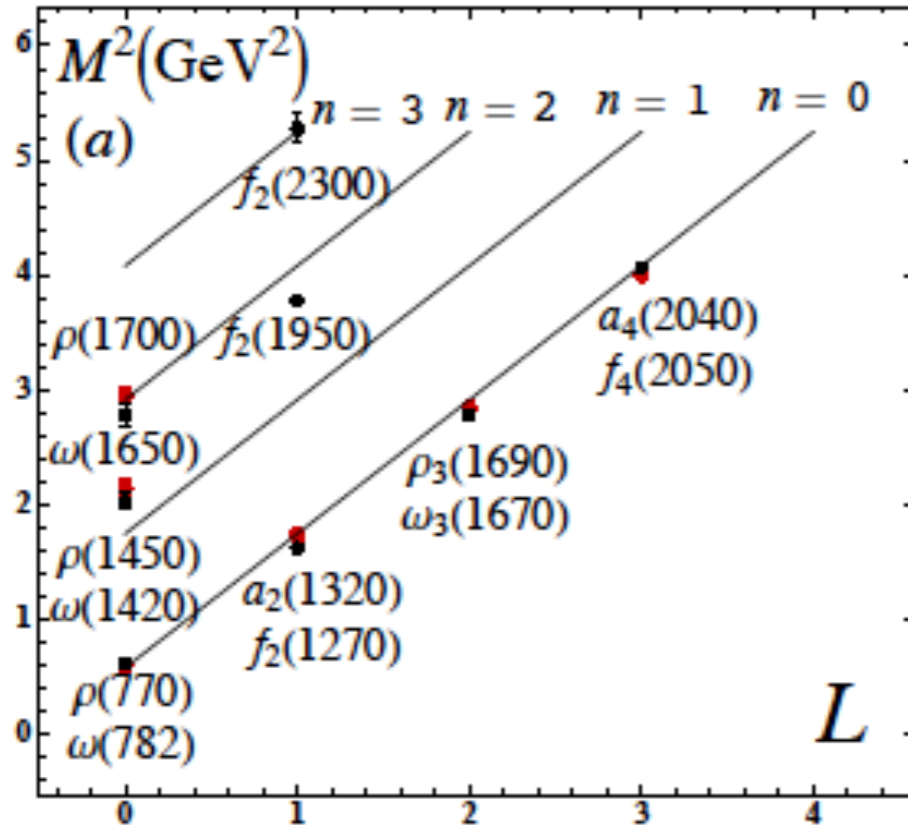
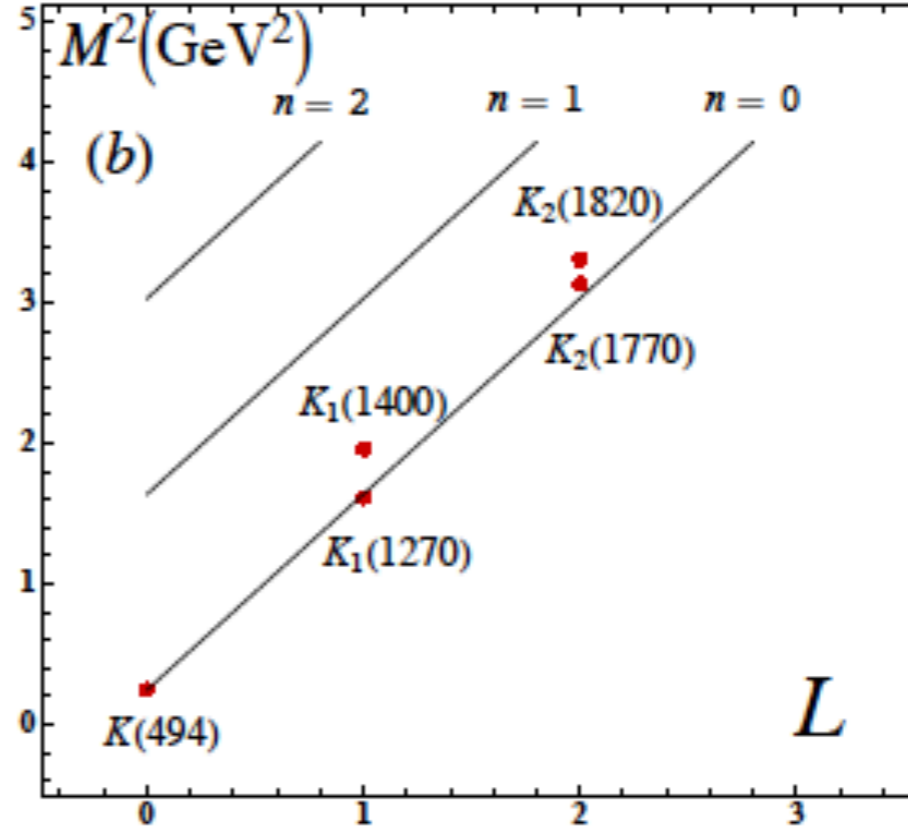
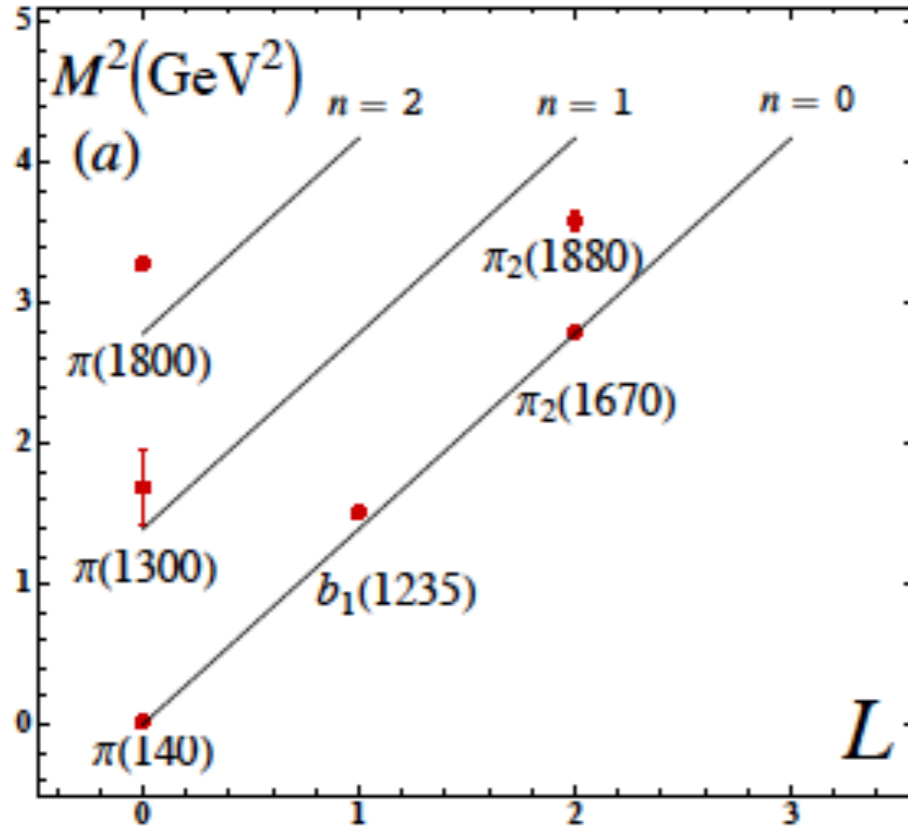
$I=1$ orbital and radial excitations for the π ($\kappa = 0.59$ GeV) and the ρ -meson families ($\kappa = 0.54$ GeV)

- Triplet splitting for the $I = 1, L = 1, J = 0, 1, 2$, vector meson a -states

$$\mathcal{M}_{a_2(1320)} > \mathcal{M}_{a_1(1260)} > \mathcal{M}_{a_0(980)}$$

Mass ratio of the ρ and the a_1 mesons: coincides with Weinberg sum rules

$$M^2 = M_0^2 + \left\langle X \left| \frac{m_q^2}{x} \right| X \right\rangle + \left\langle X \left| \frac{m_{\bar{q}}^2}{1-x} \right| X \right\rangle$$



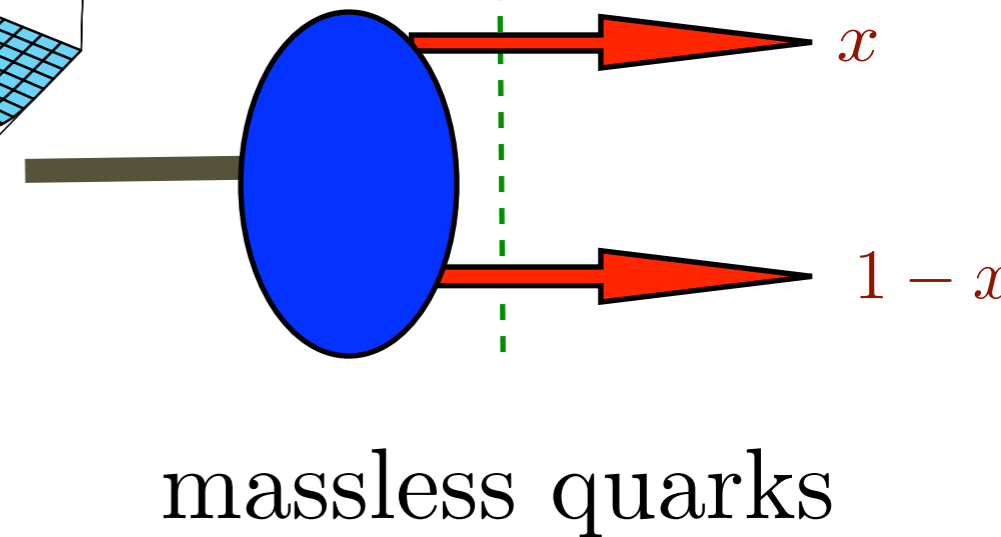
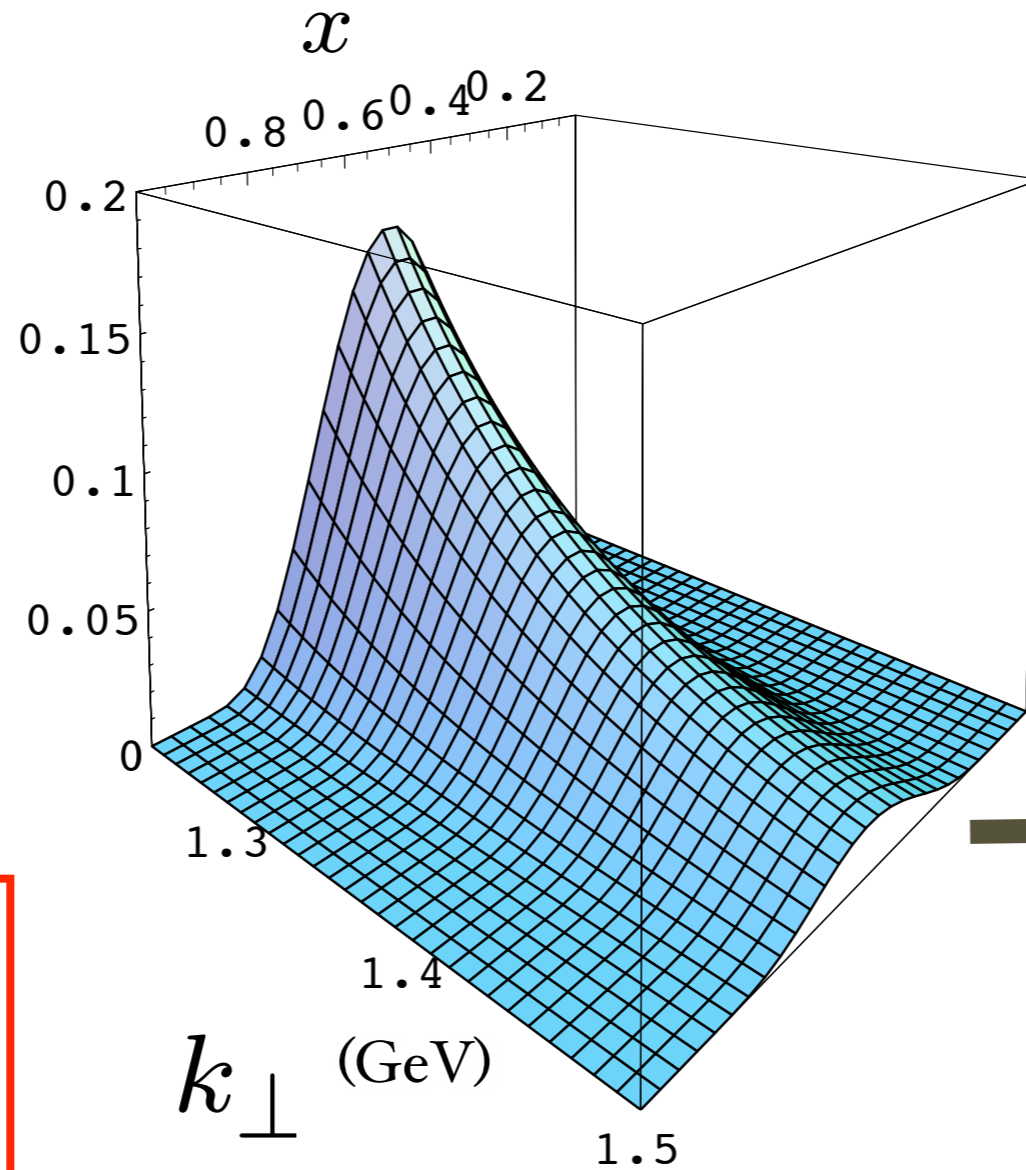
Prediction from AdS/QCD: Meson LFWF

$$e^{\varphi(z)} = e^{+\kappa^2 z}$$

de Teramond,
Cao, sjb

“Soft Wall”
model

$$\psi_M(x, k_{\perp}^2)$$



Note coupling

$$k_{\perp}^2, x$$

$$\psi_M(x, k_{\perp}) = \frac{4\pi}{\kappa \sqrt{x(1-x)}} e^{-\frac{k_{\perp}^2}{2\kappa^2 x(1-x)}}$$

$$\phi_{\pi}(x) = \frac{4}{\sqrt{3}\pi} f_{\pi} \sqrt{x(1-x)}$$

$$f_{\pi} = \sqrt{P_{q\bar{q}}} \frac{\sqrt{3}}{8} \kappa = 92.4 \text{ MeV}$$

Same as DSE!

Provides Connection of Confinement to Hadron Structure

AdS/QCD Holographic Wave Function for the ρ Meson and Diffractive ρ Meson Electroproduction

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(Received 5 April 2012; published 20 August 2012)*

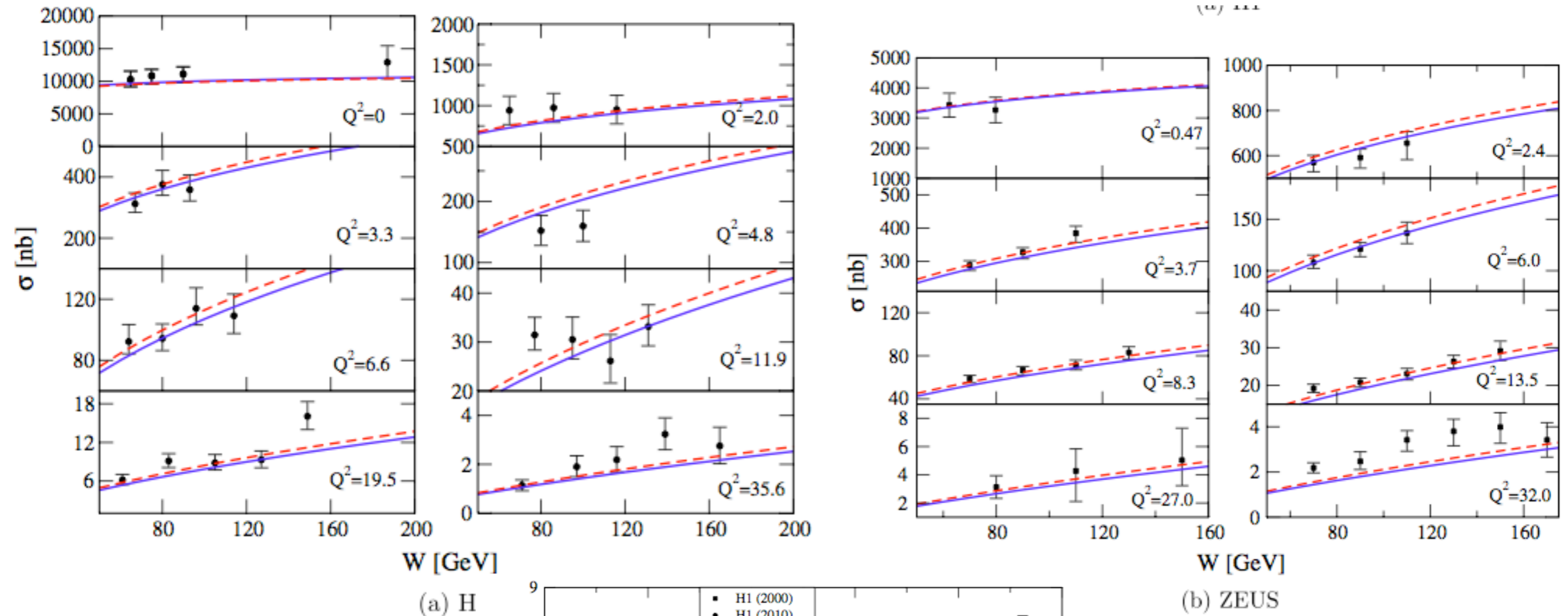
We show that anti-de Sitter/quantum chromodynamics generates predictions for the rate of diffractive ρ -meson electroproduction that are in agreement with data collected at the Hadron Electron Ring Accelerator electron-proton collider.

$$\psi_M(x, k_\perp) = \frac{4\pi}{\kappa \sqrt{x(1-x)}} e^{-\frac{k_\perp^2}{2\kappa^2 x(1-x)}}$$

**See also Ferreira
and Dosch**

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$

AdS/QCD Holographic Wave Function for the ρ Meson and Diffractive ρ Meson Electroproduction

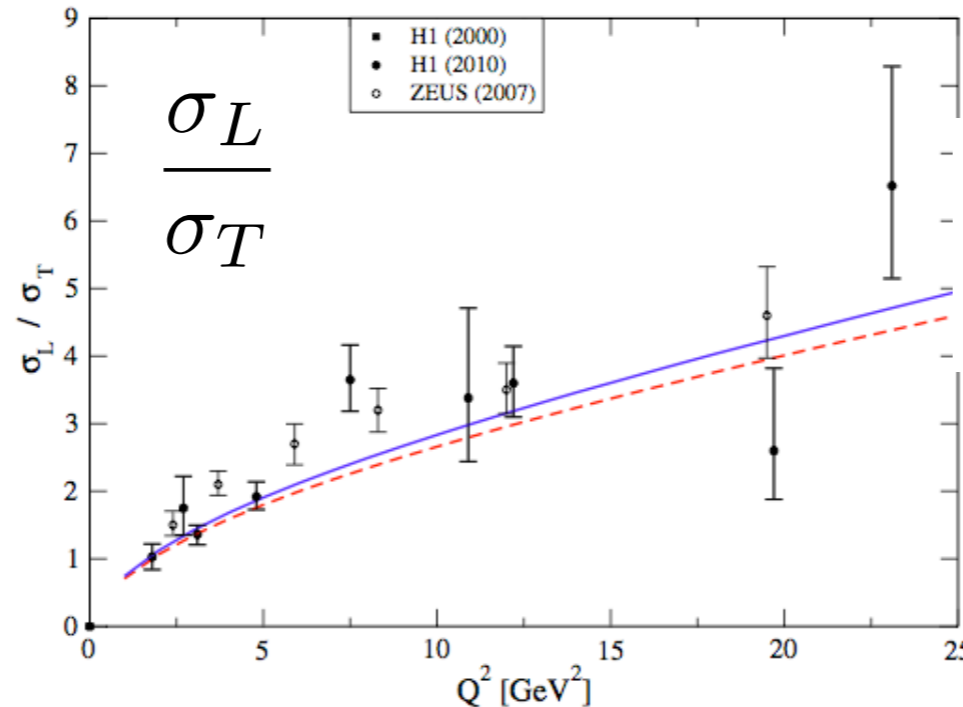


(a) H

(b) ZEUS

**J. R. Forshaw,
R. Sandapen**

$$\gamma^* p \rightarrow \rho^0 p'$$



$$\tilde{\phi}(x, k) \propto \frac{1}{\sqrt{x(1-x)}} \exp\left(-\frac{M_{q\bar{q}}^2}{2\kappa^2}\right)$$

**See also Ferreira
and Dosch**

Uniqueness

de Tèramond, Dosch, sjb

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2(L + S - 1) \quad e^{\varphi(z)} = e^{+\kappa^2 z^2}$$

- **ζ^2 confinement potential and dilaton profile unique!**
- **Linear Regge trajectories in n and L : same slope!**
- **Massless pion in chiral limit! No vacuum condensate!**
- **Conformally invariant action for massless quarks retained despite mass scale**
- **Same principle, equation of motion as de Alfaro, Furlan, Fubini, Conformal Invariance in Quantum Mechanics Nuovo Cim. A34 (1976) 569**

Hadron Form Factors from AdS/QCD

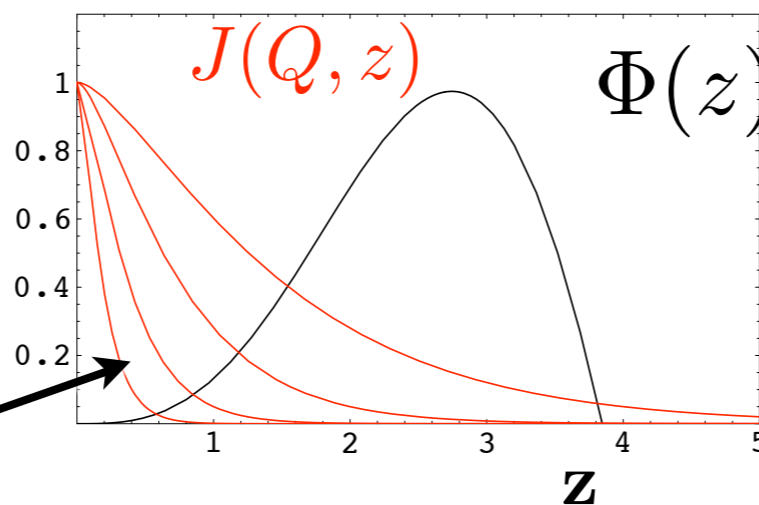
Propagation of external perturbation suppressed inside AdS.

$$J(Q, z) = zQ K_1(zQ)$$

$$F(Q^2)_{I \rightarrow F} = \int \frac{dz}{z^3} \Phi_F(z) J(Q, z) \Phi_I(z)$$

High Q^2
from
small $z \sim 1/Q$

high Q^2



**Polchinski, Strassler
de Teramond, sjb**

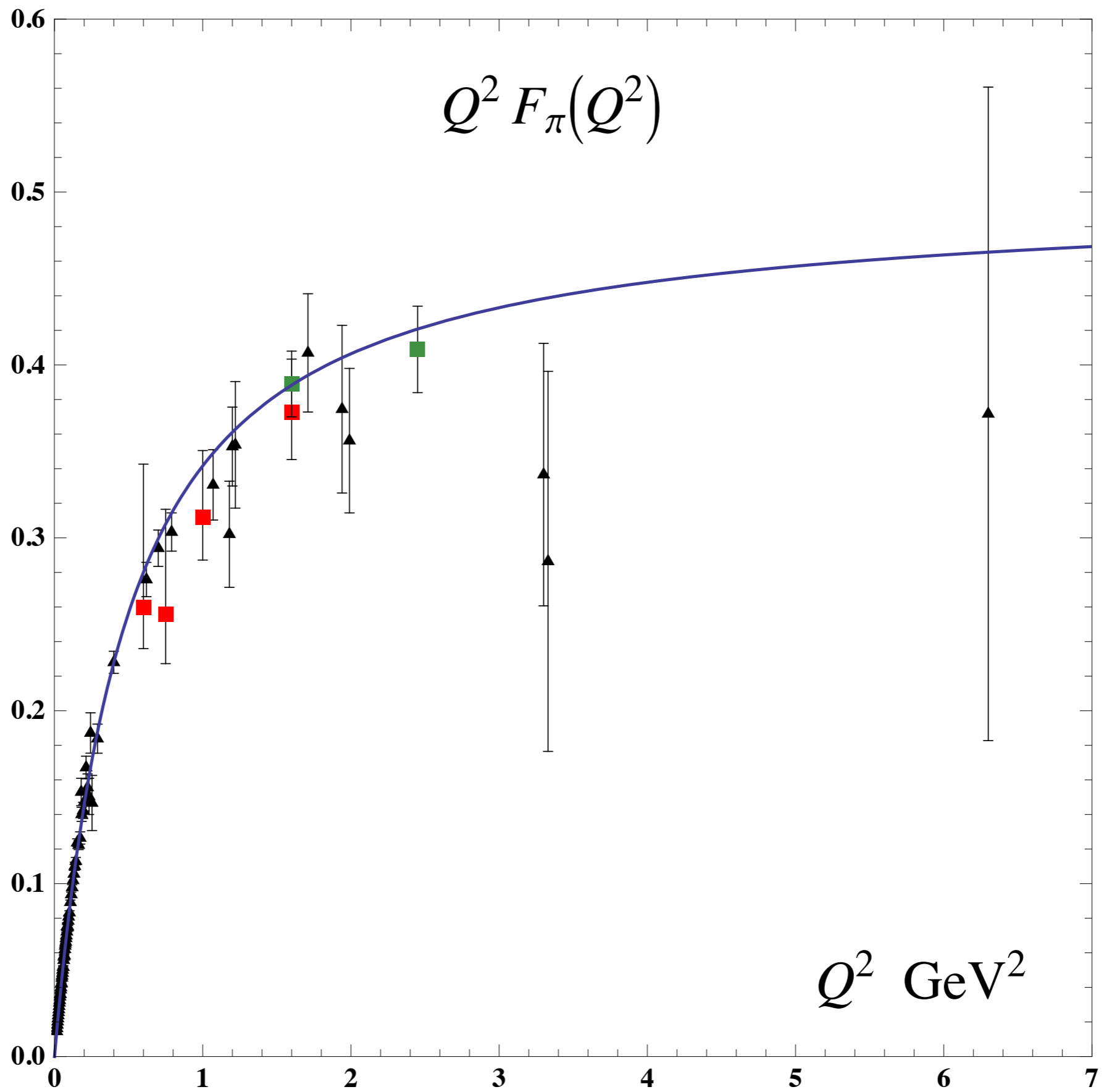
Consider a specific AdS mode $\Phi^{(n)}$ dual to an n partonic Fock state $|n\rangle$. At small z , $\Phi^{(n)}$ scales as $\Phi^{(n)} \sim z^{\Delta_n}$. Thus:

$$F(Q^2) \rightarrow \left[\frac{1}{Q^2} \right]^{\tau-1},$$

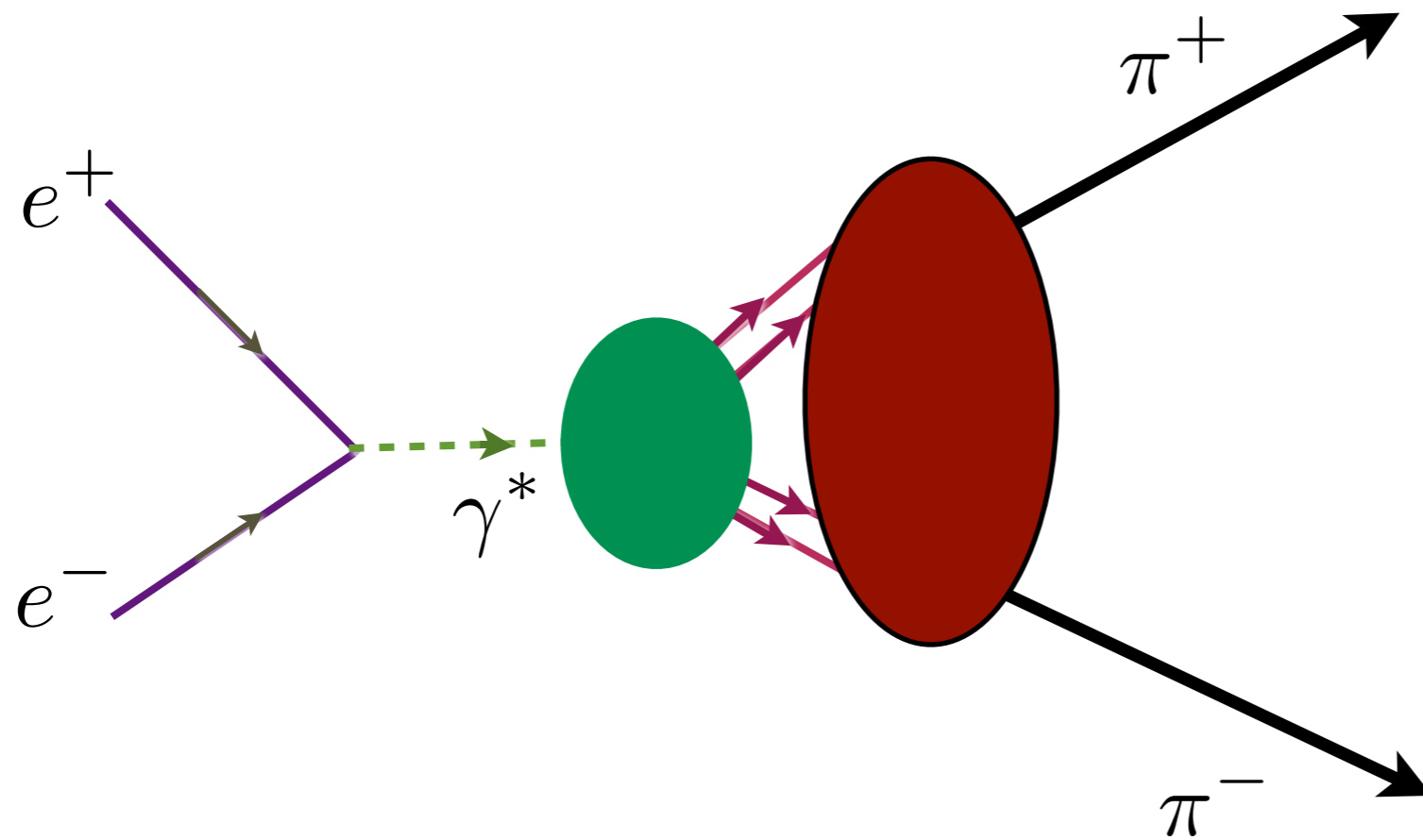
**Dimensional Quark Counting Rules:
General result from
AdS/CFT and Conformal Invariance**

where $\tau = \Delta_n - \sigma_n$, $\sigma_n = \sum_{i=1}^n \sigma_i$.

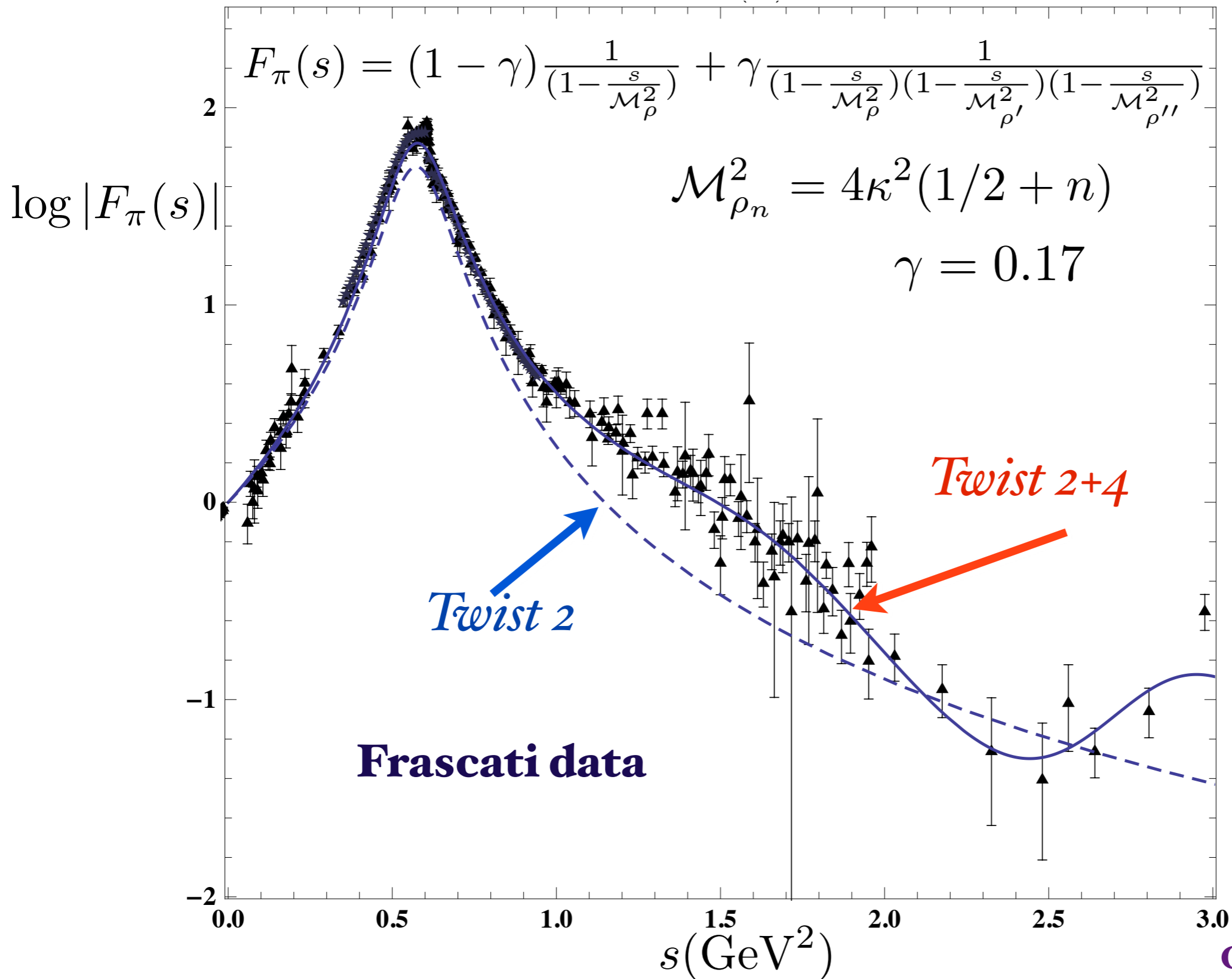
Twist $\tau = n + L$



Dressed soft-wall current brings in higher Fock states and more vector meson poles



Timelike Pion Form Factor from AdS/QCD and Light-Front Holography



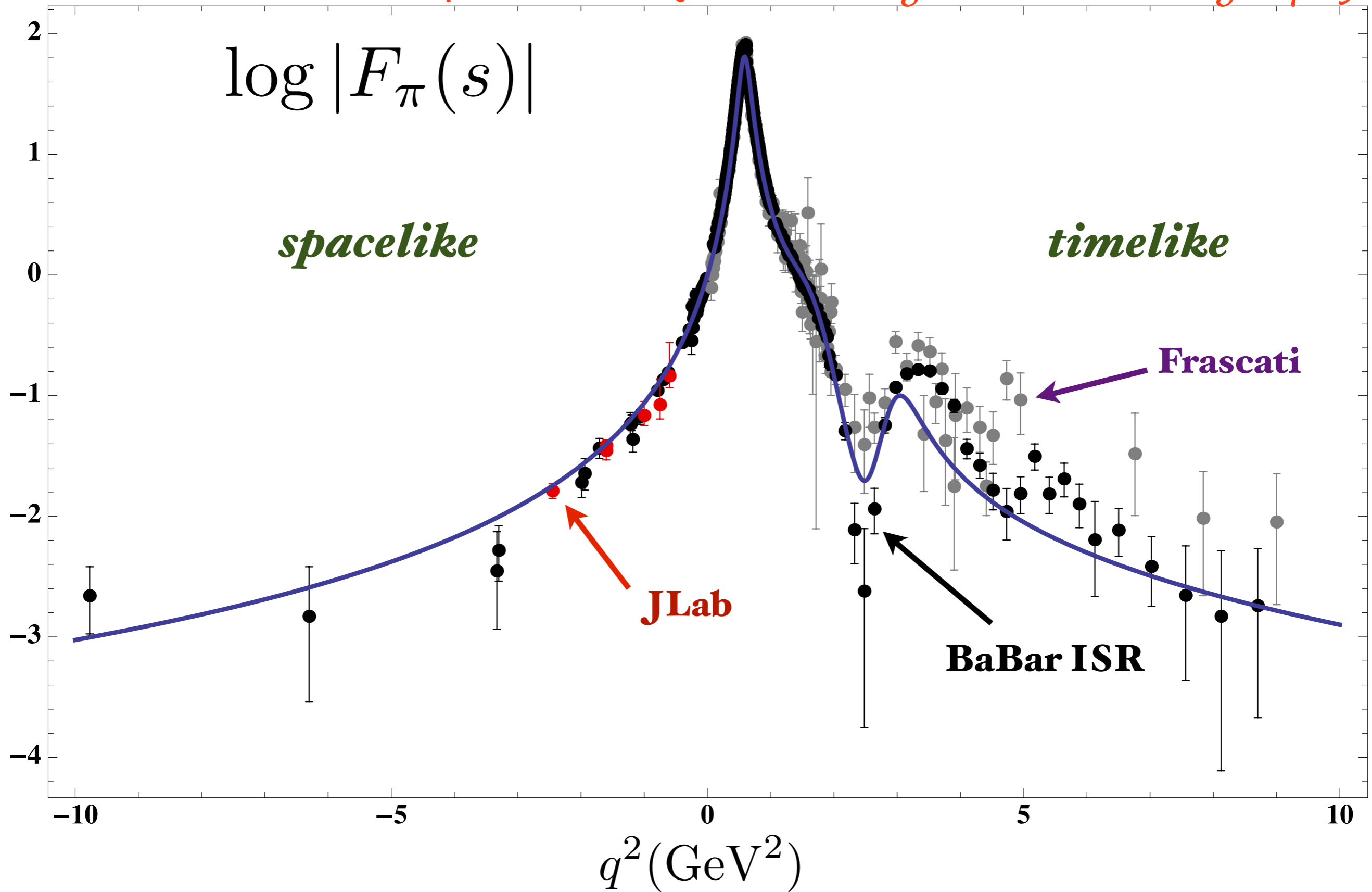
Prescription for Timelike poles :

$$\frac{1}{s - M^2 + i\sqrt{s}\Gamma}$$

14% four-quark probability

G. de Teramond & sjb

Pion Form Factor from AdS/QCD and Light-Front Holography



Remarkable Features of Light-Front Schrödinger Equation

- **Relativistic, frame-independent**
- **QCD scale appears - unique LF potential**
- **Reproduces spectroscopy and dynamics of light-quark hadrons with one parameter**
- **Zero-mass pion for zero mass quarks!**
- **Regge slope same for n and L -- not usual HO**
- **Splitting in L persists to high mass -- contradicts conventional wisdom based on breakdown of chiral symmetry**
- **Phenomenology: LFWFs, Form factors, electroproduction**
- **Extension to heavy quarks**

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2 (L + S - 1)$$

Kunshan
August 3, 2015



New Perspectives for Hadron Physics

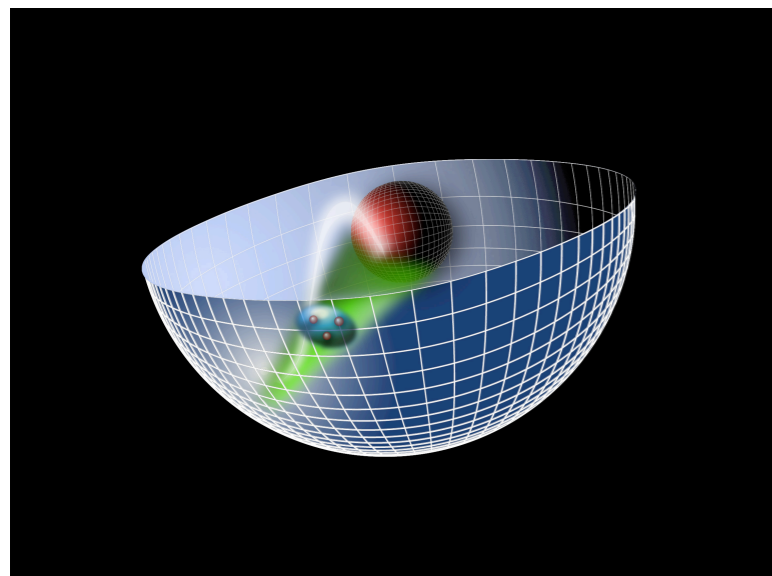
Stan Brodsky

SLAC
NATIONAL ACCELERATOR LABORATORY

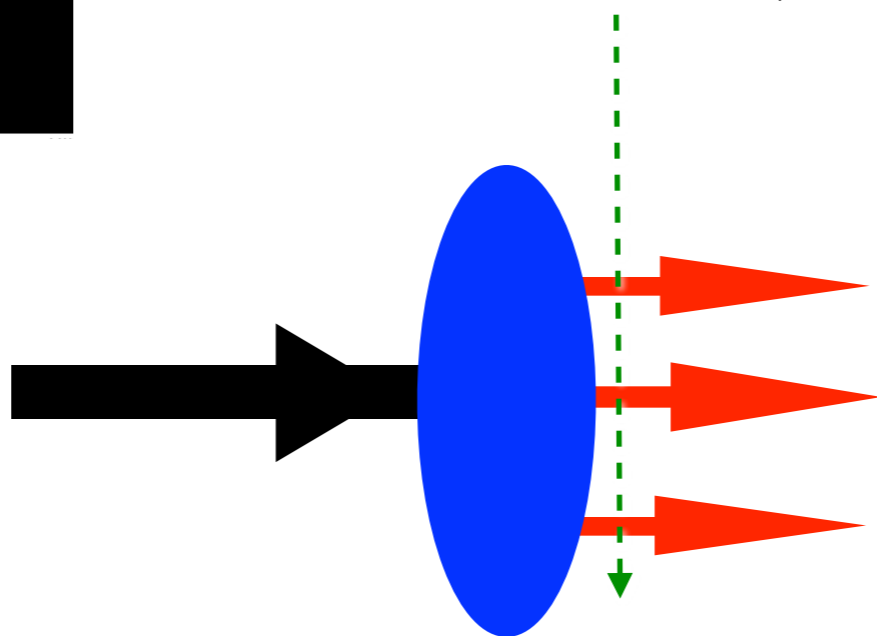
$$\phi(z)$$

AdS₅: Conformal Template for QCD

- *Light-Front Holography*

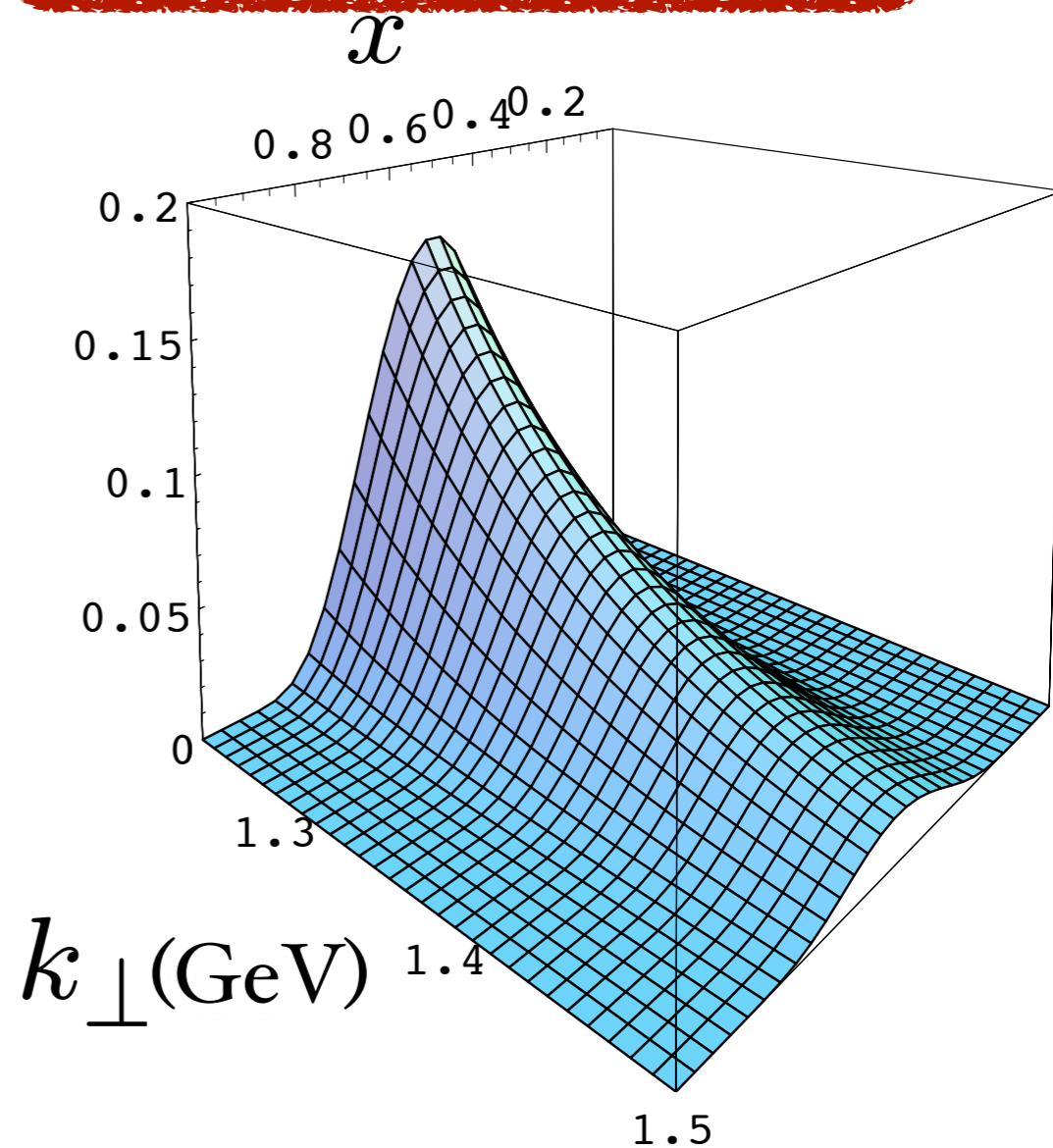


Fixed $\tau = t + z/c$



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

Duality of AdS₅ with LF Hamiltonian Theory



- *Light Front Wavefunctions:*

**Light-Front Schrödinger Equation
Spectroscopy and Dynamics**

QCD Lagrangian

$$\mathcal{L}_{QCD} = -\frac{1}{4} \text{Tr}(G^{\mu\nu} G_{\mu\nu}) + \sum_{f=1}^{n_f} i\bar{\Psi}_f D_\mu \gamma^\mu \Psi_f + \sum_{f=1}^{n_f} \cancel{m_f} \bar{\Psi}_f \Psi_f$$

$$iD^\mu = i\partial^\mu - gA^\mu \quad G^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu - g[A^\mu, A^\nu]$$

Classical Chiral Lagrangian is Conformally Invariant

Where does the QCD Mass Scale Λ_{QCD} come from?

How does color confinement arise?

● **de Alfaro, Fubini, Furlan:**

Scale can appear in Hamiltonian and EQM without affecting conformal invariance of action!

Unique confinement potential!

$$G|\psi(\tau)\rangle = i\frac{\partial}{\partial\tau}|\psi(\tau)\rangle$$

$$G = uH + vD + wK$$

New term

$$G = H_\tau = \frac{1}{2}\left(-\frac{d^2}{dx^2} + \frac{g}{x^2} + \frac{4uw - v^2}{4}x^2\right)$$

Retains conformal invariance of action despite mass scale!

$$4uw - v^2 = \kappa^4 = [M]^4$$

Identical to LF Hamiltonian with unique potential and dilaton!

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta)\right]\psi(\zeta) = \mathcal{M}^2\psi(\zeta)$$

$$U(\zeta) = \kappa^4\zeta^2 + 2\kappa^2(L + S - 1)$$

What determines the QCD mass scale Λ_{QCD} ?

- Mass scale does not appear in the QCD Lagrangian (massless quarks)
- Dimensional Transmutation? Requires external constraint such as $\alpha_s(M_Z)$
- dAFF: Confinement Scale κ appears spontaneously via the Hamiltonian: $G = uH + vD + wK \quad 4uw - v^2 = \kappa^4 = [M]^4$
- The confinement scale regulates infrared divergences, connects Λ_{QCD} to the confinement scale κ
- Only dimensionless mass ratios (and M times R) predicted
- Mass and time units [GeV] and [sec] from physics external to QCD
- New feature: bounded frame-independent relative time between constituents

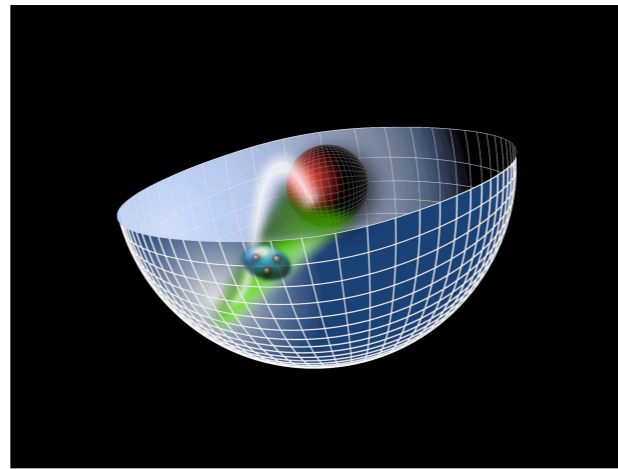


dAFF: New Time Variable

$$\tau = \frac{2}{\sqrt{4uw - v^2}} \arctan \left(\frac{2tw + v}{\sqrt{4uw - v^2}} \right),$$

- **Identify with difference of LF time $\Delta x^+ / P^+$ between constituents**
- **Finite range**
- **Measure in Double-Parton Processes**





*AdS/QCD
Soft-Wall Model*

Light-Front Holography

$$\zeta^2 = x(1-x)b_{\perp}^2.$$

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2(L + S - 1)$$

$$\kappa \simeq 0.6 \text{ GeV}$$

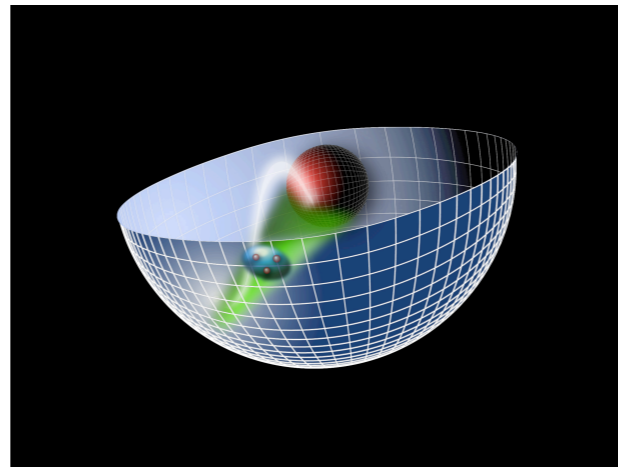
Confinement scale:

$$1/\kappa \simeq 1/3 \text{ fm}$$

***Unique
Confinement Potential!
Conformal Symmetry
of the action***

● **de Alfaro, Fubini, Furlan:**

**Scale can appear in Hamiltonian and EQM
without affecting conformal invariance of action!**



*AdS/QCD
Soft-Wall Model*

Light-Front Holography

Semi-Classical Approximation to QCD

Relativistic, frame-independent

Unique color-confining potential

Zero mass pion for massless quarks

Regge trajectories with equal slopes in n and L

Light-Front Wavefunctions

Light-Front Schrödinger Equation

*Conformal Symmetry
of the action*

1+1

$$\{\psi, \psi^+\} = 1$$

*two anti-commuting
fermionic operators*

$$\psi = \frac{1}{2}(\sigma_1 - i\sigma_2), \quad \psi^+ = \frac{1}{2}(\sigma_1 + i\sigma_2)$$

Realization as Pauli Matrices

$$Q = \psi^+[-\partial_x + W(x)], \quad Q^+ = \psi[\partial_x + W(x)], \quad W(x) = \frac{f}{x}$$

(Conformal)

$$S = \psi^+ x, \quad S^+ = \psi x$$

Introduce new spinor operators

$$\{Q, Q^+\} = 2H, \quad \{S, S^+\} = 2K$$

$$Q \simeq \sqrt{H}, \quad S \simeq \sqrt{K}$$

$$\{Q, Q\} = \{Q^+, Q^+\} = 0, \quad [Q, H] = [Q^+, H] = 0$$

Superconformal Algebra

Baryon Equation

Consider $R_w = Q + wS$; w : dimensions of mass squared

$$G = \{R_w, R_w^+\} = 2H + 2w^2K + 2wfI - 2wB \quad 2B = \sigma_3$$

Retains Conformal Invariance of Action

Fubini and Rabinovici

New Extended Hamiltonian G is diagonal:

$$G_{11} = \left(-\partial_x^2 + w^2x^2 + 2wf - w + \frac{4(f + \frac{1}{2})^2 - 1}{4x^2} \right)$$

$$G_{22} = \left(-\partial_x^2 + w^2x^2 + 2wf + w + \frac{4(f - \frac{1}{2})^2 - 1}{4x^2} \right)$$

Identify $f - \frac{1}{2} = L_B$, $w = \kappa^2$

Eigenvalue of G : $M^2(n, L) = 4\kappa^2(n + L_B + 1)$ 83

LF Holography

Baryon Equation

$$x \rightarrow \zeta$$

$$\left(-\frac{d^2}{d\zeta^2} + \lambda_B^2 \zeta^2 + 2\lambda_B(L_B + 1) + \frac{4L_B^2 - 1}{4\zeta^2} \right) \psi_J^+ = M^2 \psi_J^+, \quad \text{G}_{22}$$
$$\left(-\frac{d^2}{d\zeta^2} + \lambda_B^2 \zeta^2 + 2\lambda_B L_B + \frac{4(L_B + 1)^2 - 1}{4\zeta^2} \right) \psi_J^- = M^2 \psi_J^-. \quad \text{G}_{11}$$

$$M_B^2(N, L_B) = 4\lambda_B(n + L_B + 1)$$

S=1/2, P=+

both chiralities

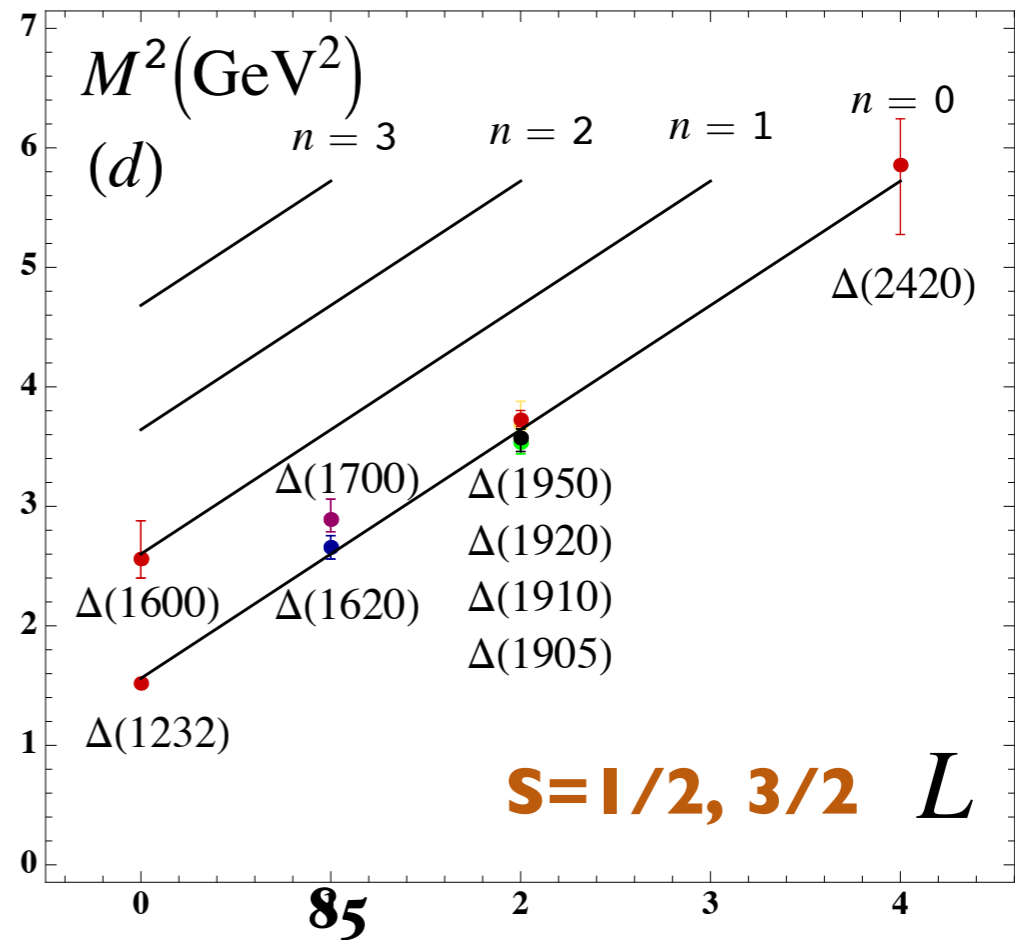
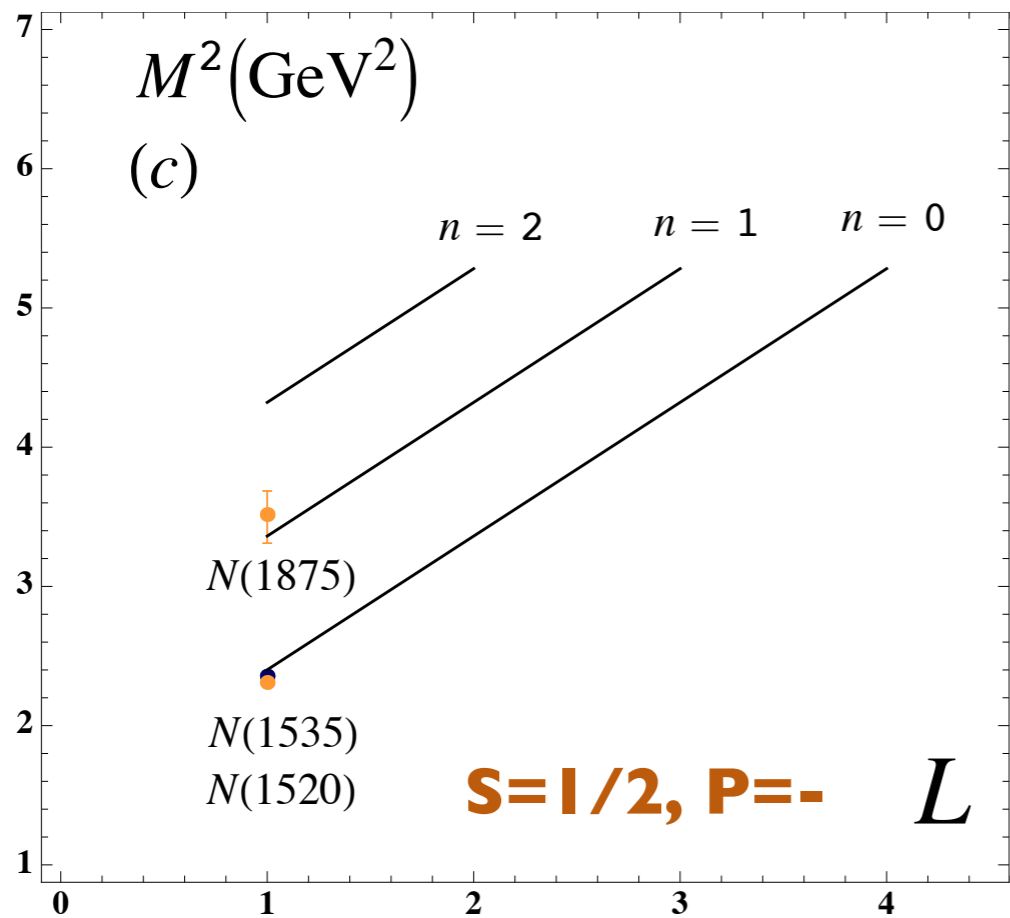
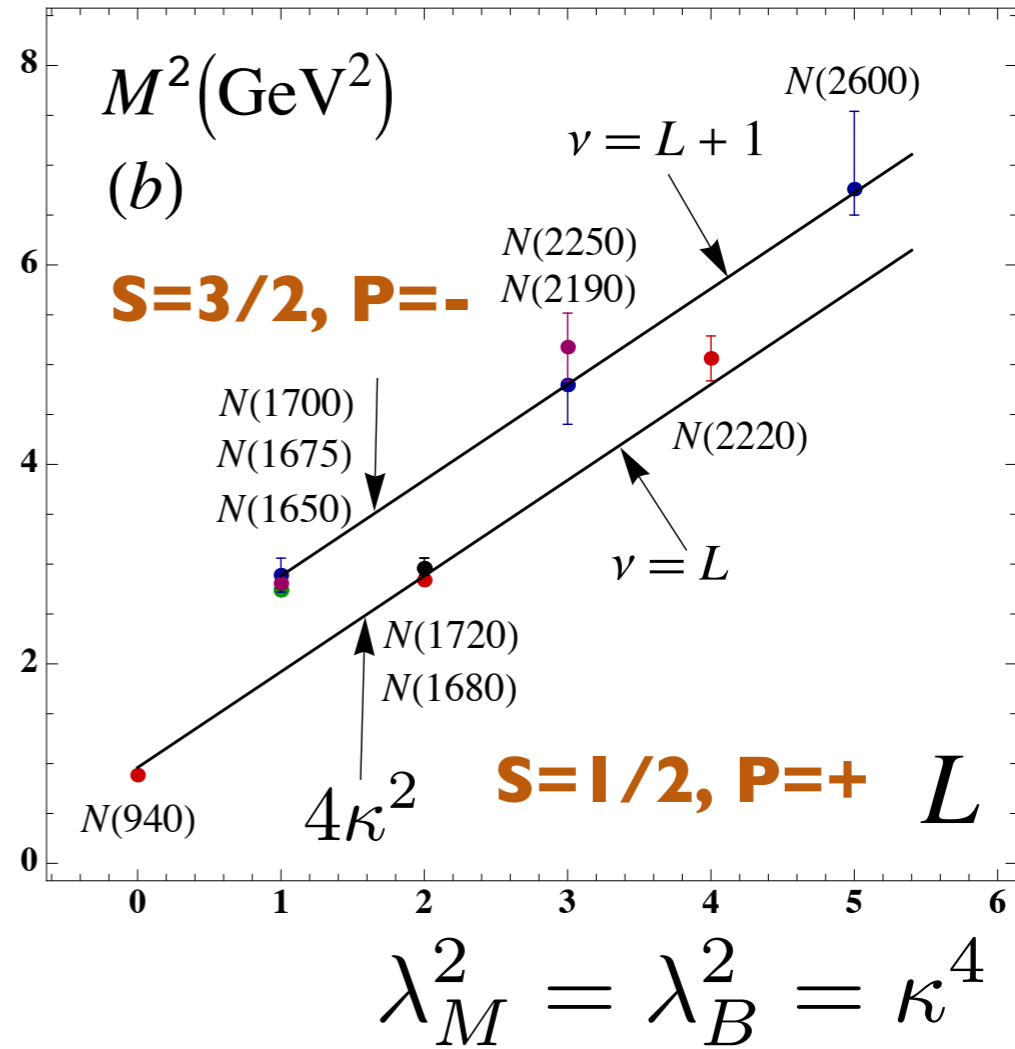
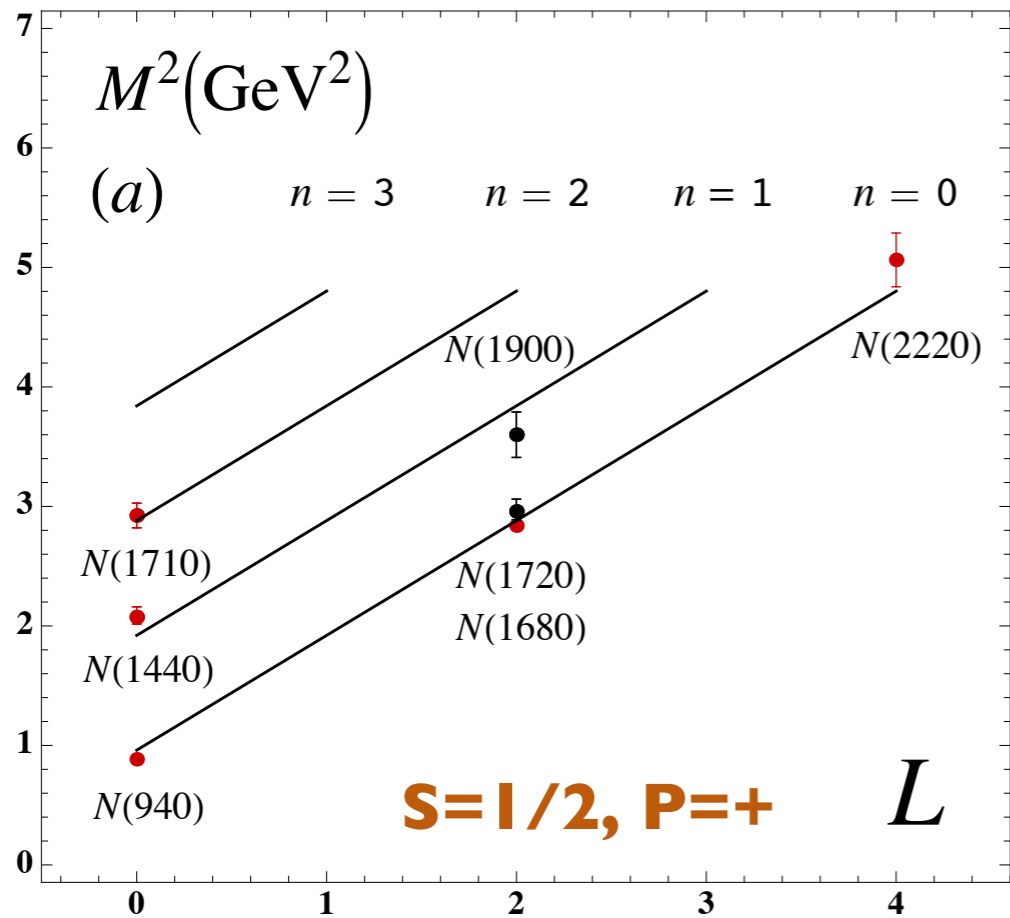
Meson Equation

$$\left(-\frac{d^2}{d\zeta^2} + \lambda_M^2 \zeta^2 + 2\lambda_M(J - 1) + \frac{4L_M^2 - 1}{4\zeta^2} \right) \phi_J = M^2 \phi_j \quad \text{G}_{11}$$

$$M_M^2(N, L_M, S = 0) = 4\lambda_M(n + L_M)$$

S=0, I=I Meson is superpartner of S=1/2, I=I Baryon

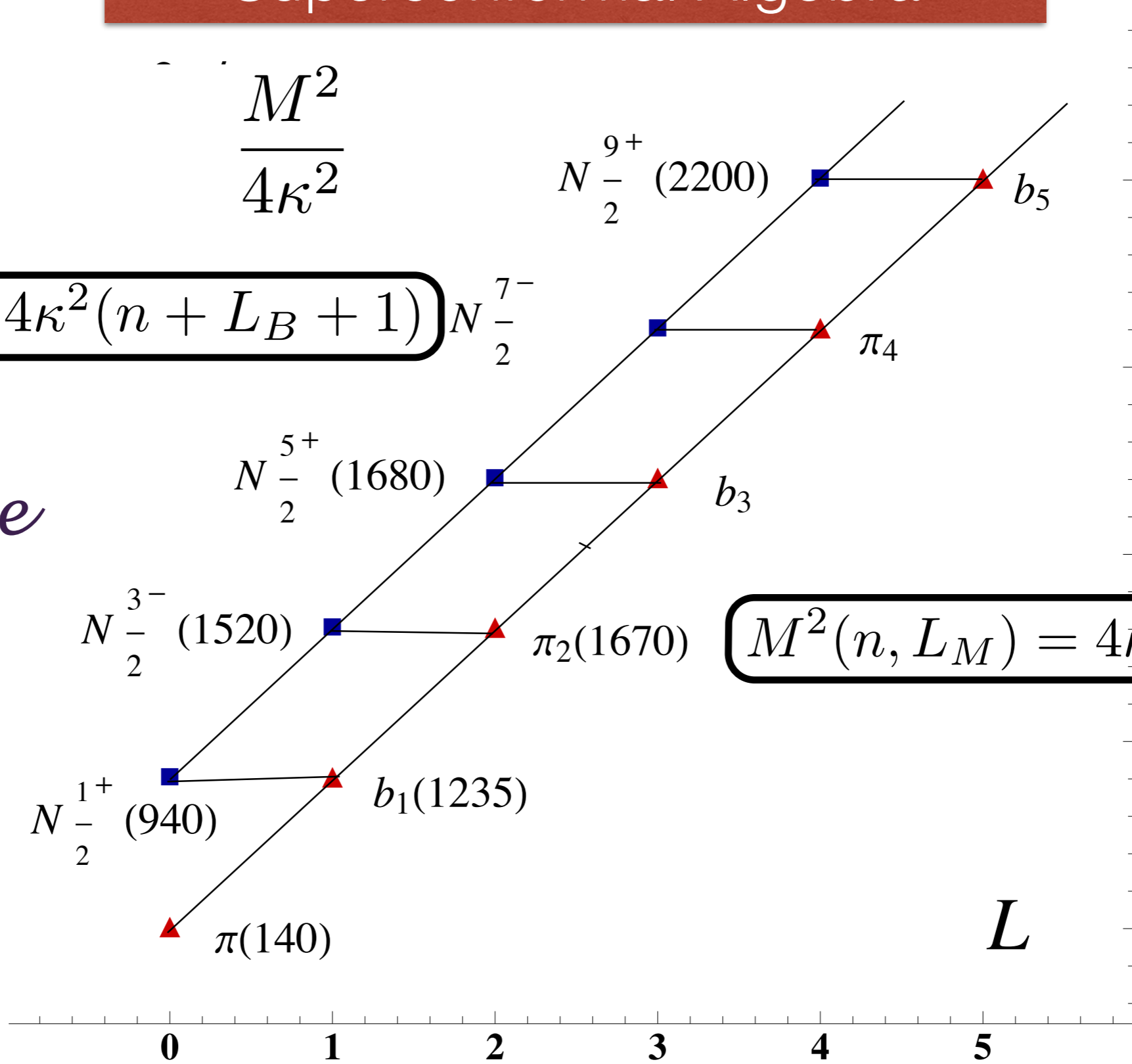
Meson-Baryon Degeneracy for $L_M=L_B+1$ $\lambda_M^2 = \lambda_B^2 = \kappa^4$



Superconformal Algebra

$$M^2(n, L_B) = 4\kappa^2(n + L_B + 1)$$

Same slope



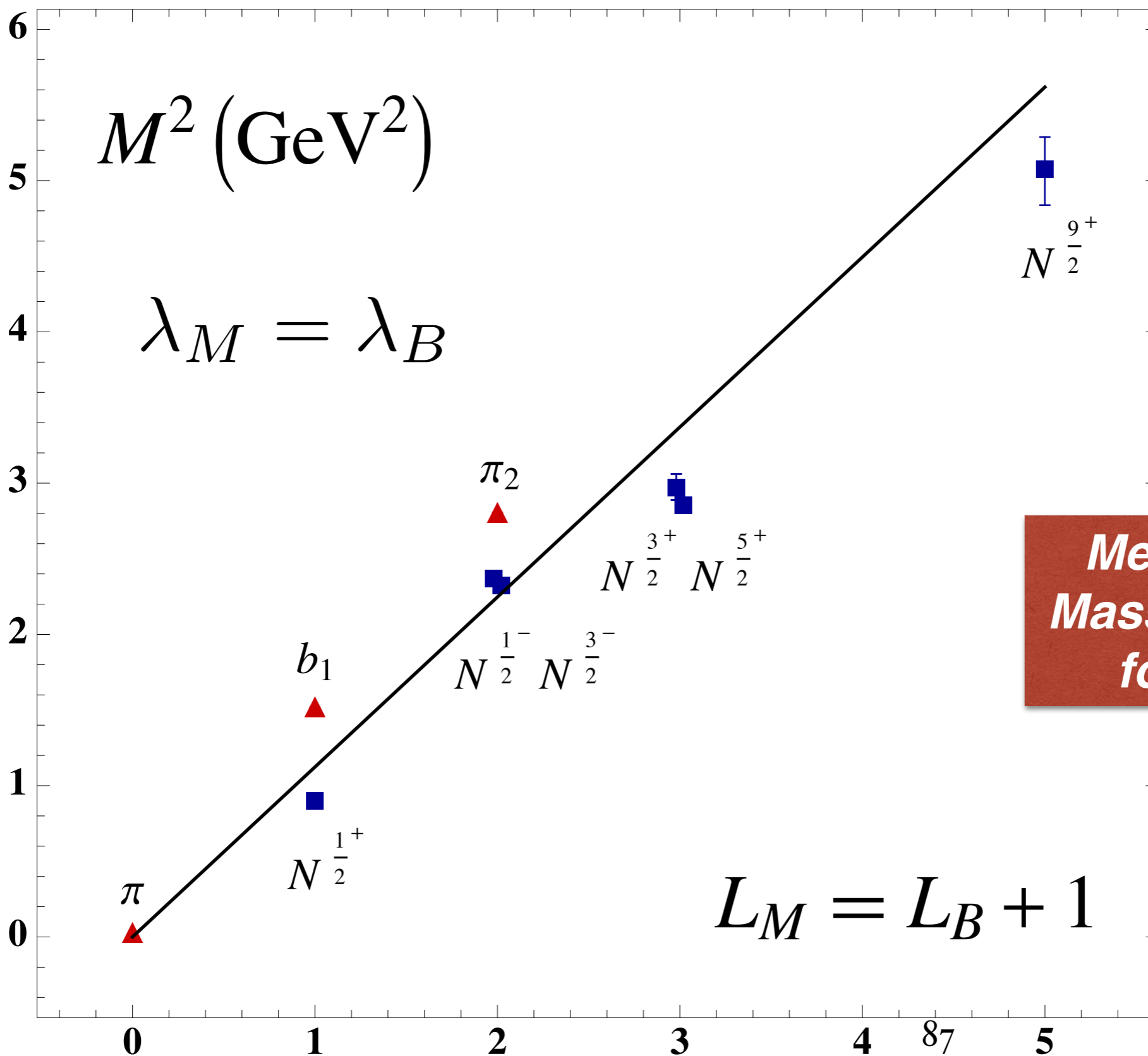
$$M^2(n, L_M) = 4\kappa^2(n + L_M)$$

**Meson-Baryon
Mass Degeneracy
for $L_M=L_B+1$**

$$\lambda_M^2 = \lambda_B^2 = \kappa^4$$

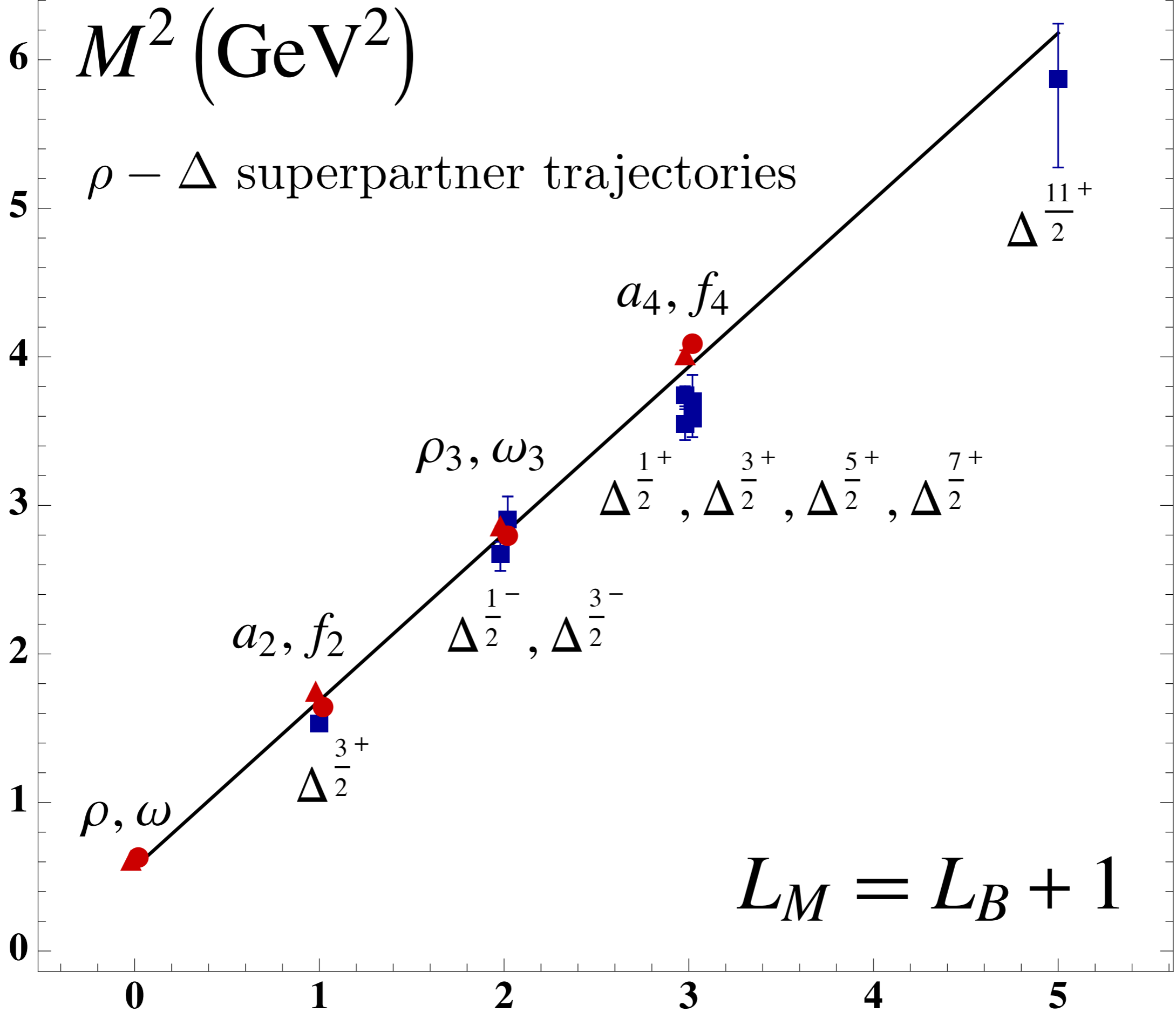
**Superconformal AdS Light-Front Holographic
QCD (LFHQCD):
Identical meson and baryon spectra!**

$$\lambda = \kappa^2$$



M^2 (GeV²)

$\rho - \Delta$ superpartner trajectories



Chiral Features of Soft-Wall AdS/QCD Model

- **Boost Invariant**
- **Trivial LF vacuum! No condensate, but consistent with GMOR**
- **Massless Pion**
- **Hadron Eigenstates (even the pion) have LF Fock components of different L^z**
- **Proton: equal probability** $S^z = +1/2, L^z = 0; S^z = -1/2, L^z = +1$

$$J^z = +1/2 : \langle L^z \rangle = 1/2, \langle S_q^z \rangle = 0$$

- **Self-Dual Massive Eigenstates: Proton is its own chiral partner.**
- **Label State by minimum L as in Atomic Physics**
- **Minimum L dominates at short distances**
- **AdS/QCD Dictionary: Match to Interpolating Operator Twist at $z=0$.**

No mass-degenerate parity partners!

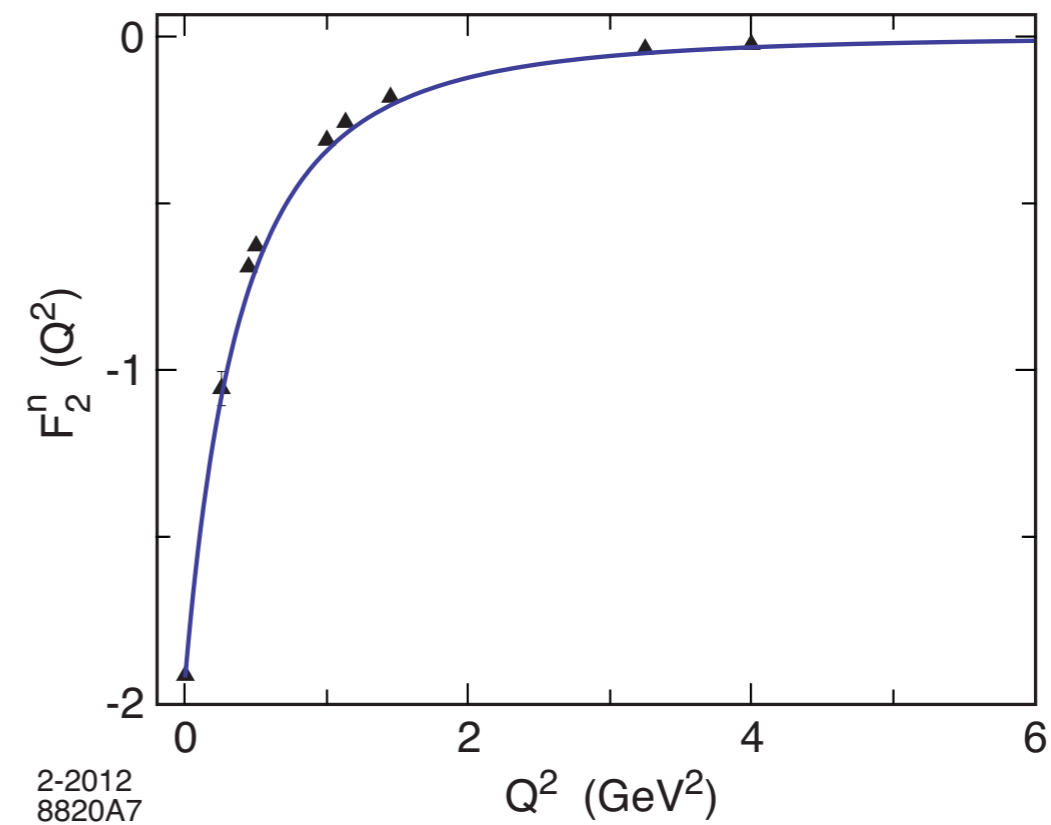
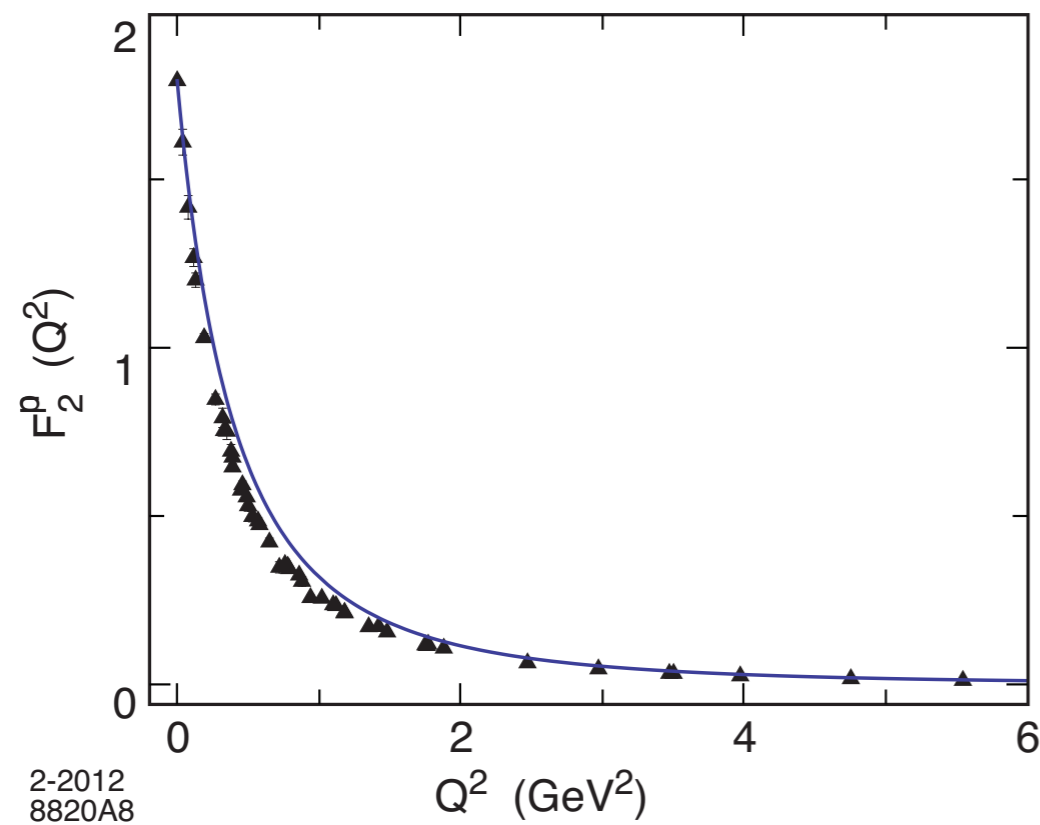
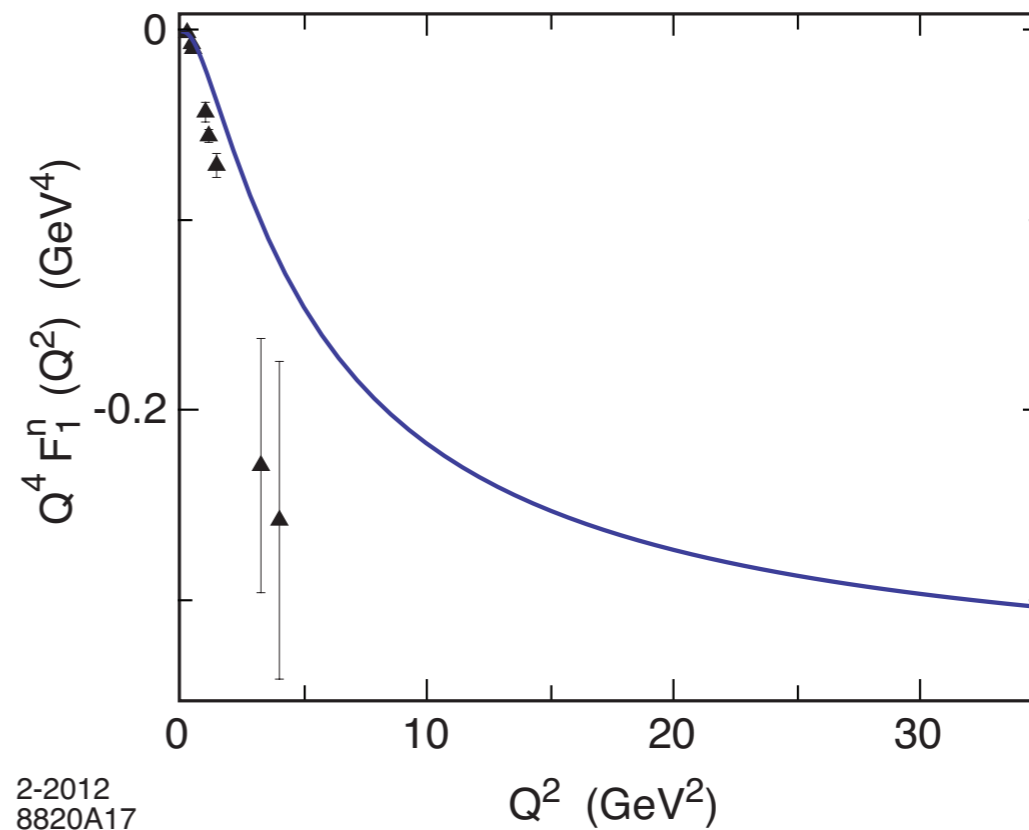
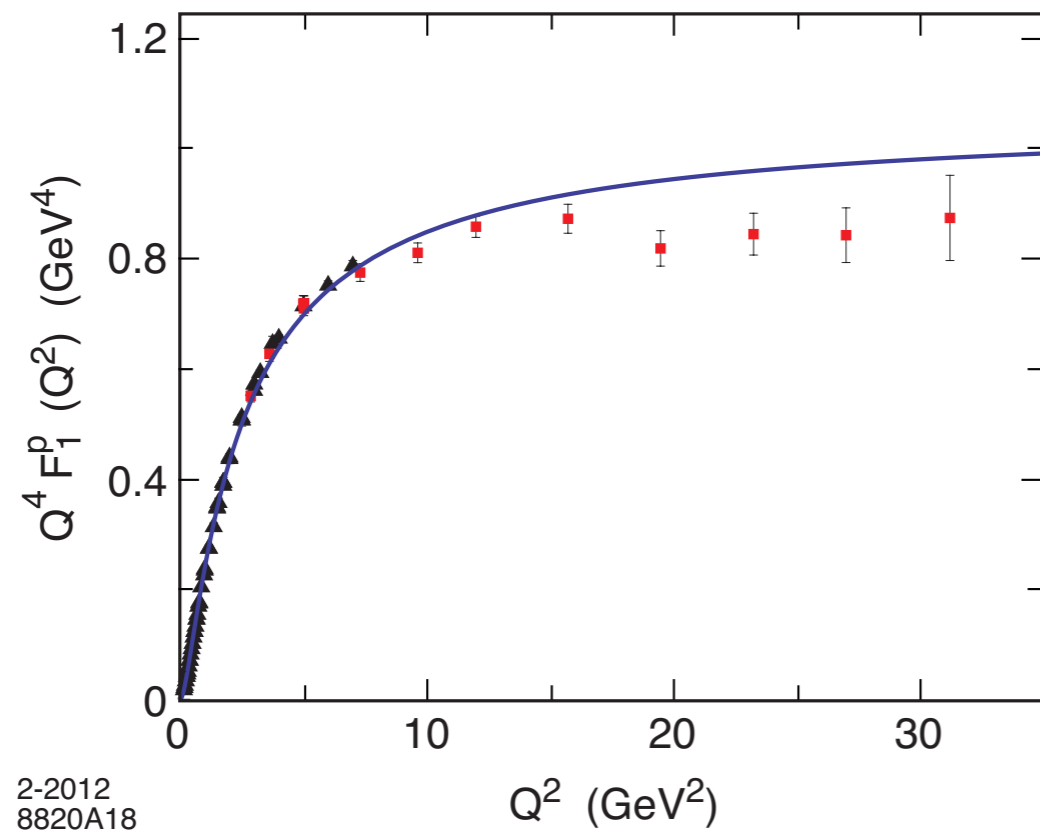
Features of Supersymmetric Equations

- $J = L + S$ baryon simultaneously satisfies both equations of G with L , $L + 1$ for same mass eigenvalue
- $J^z = L^z + 1/2 = (L^z + 1) - 1/2$ $S^z = \pm 1/2$
- Baryon spin carried by quark orbital angular momentum: $\langle J^z \rangle = L^z + 1/2$
- Mass-degenerate meson “superpartner” with $L_M = L_B + 1$. *“Shifted meson-baryon Duality”*

Meson and baryon have same κ !

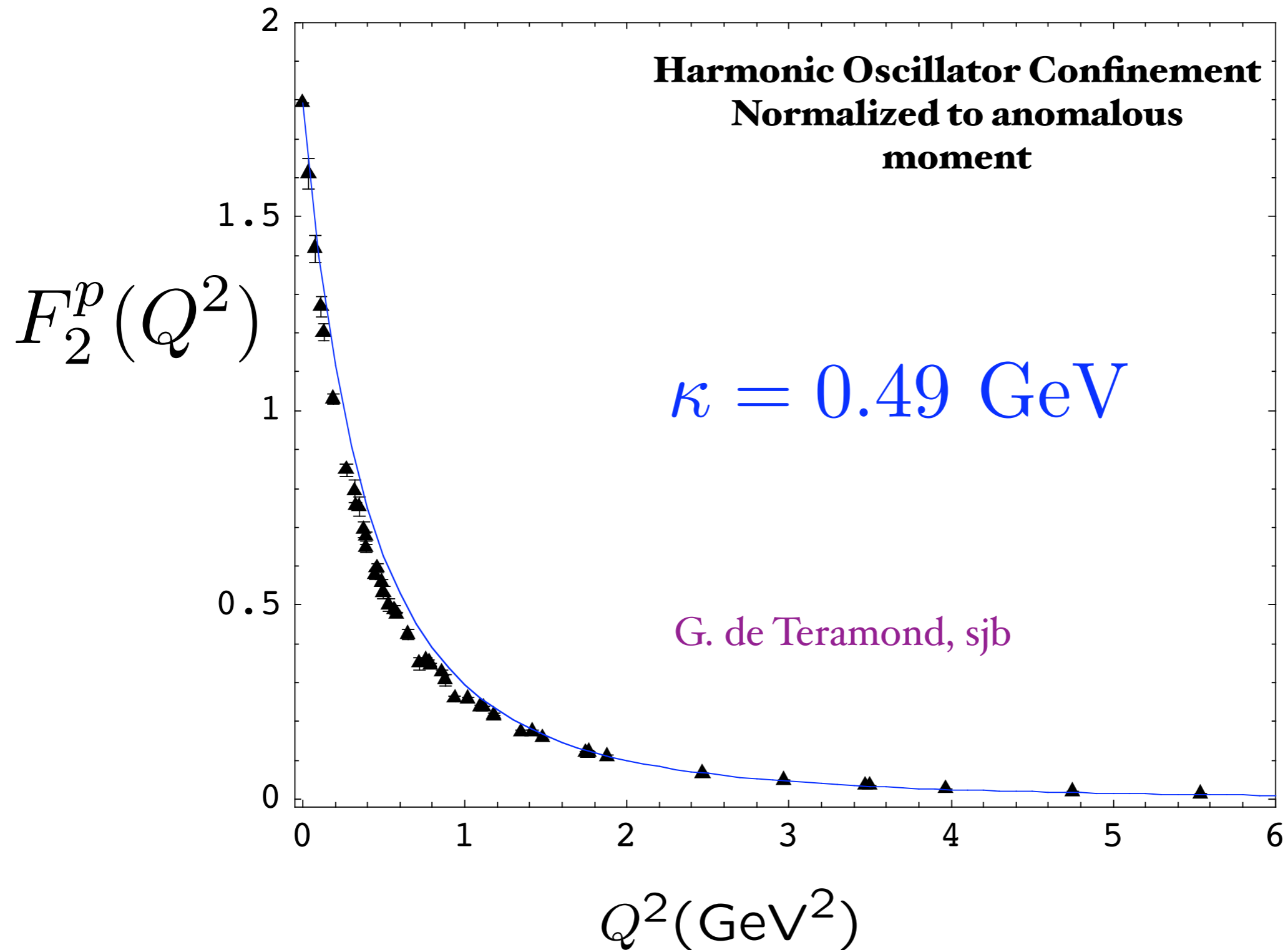


Using $SU(6)$ flavor symmetry and normalization to static quantities



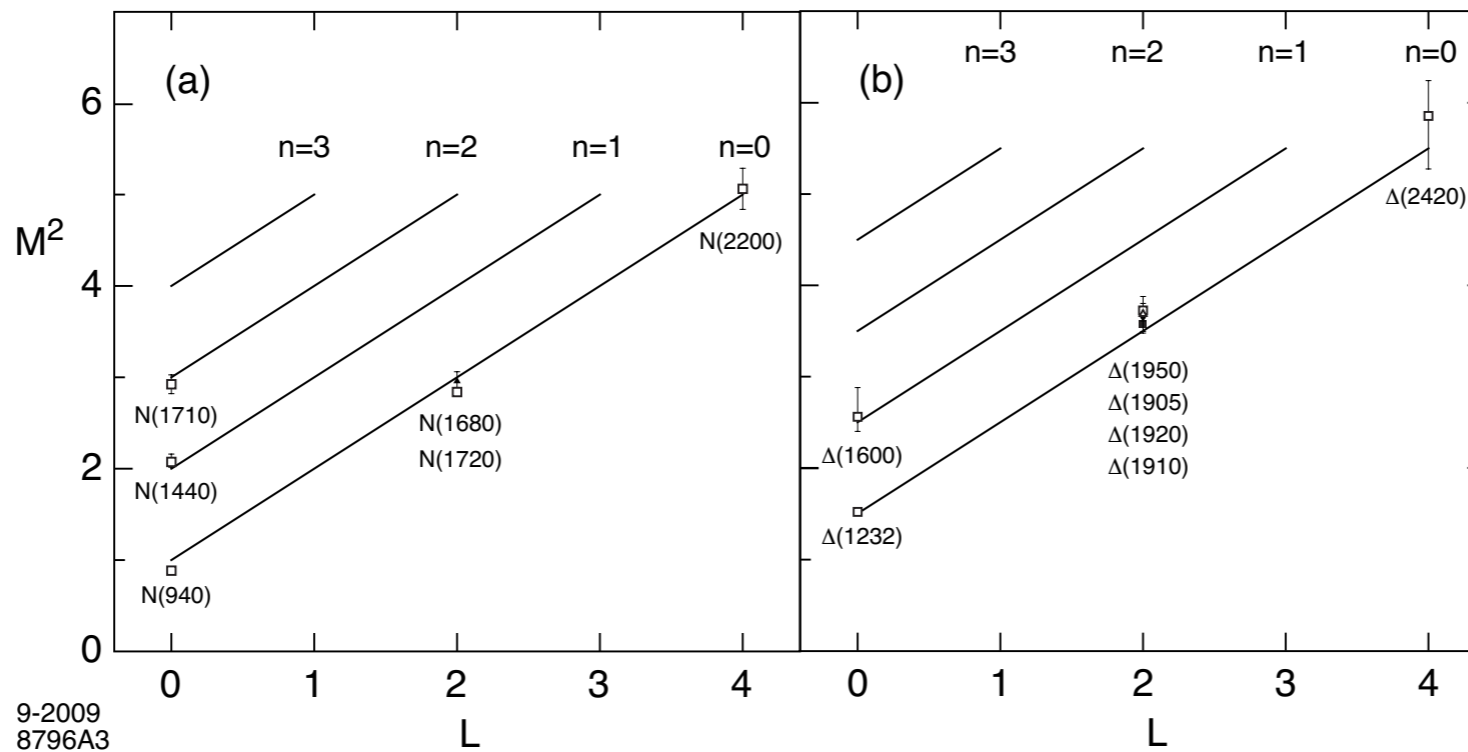
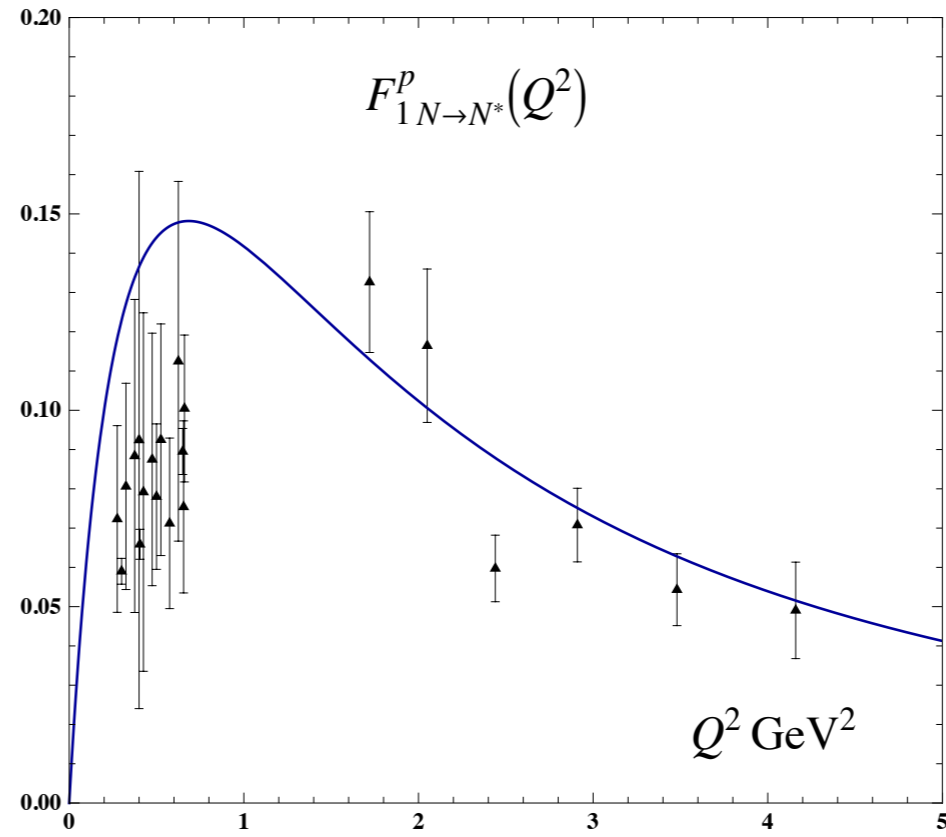
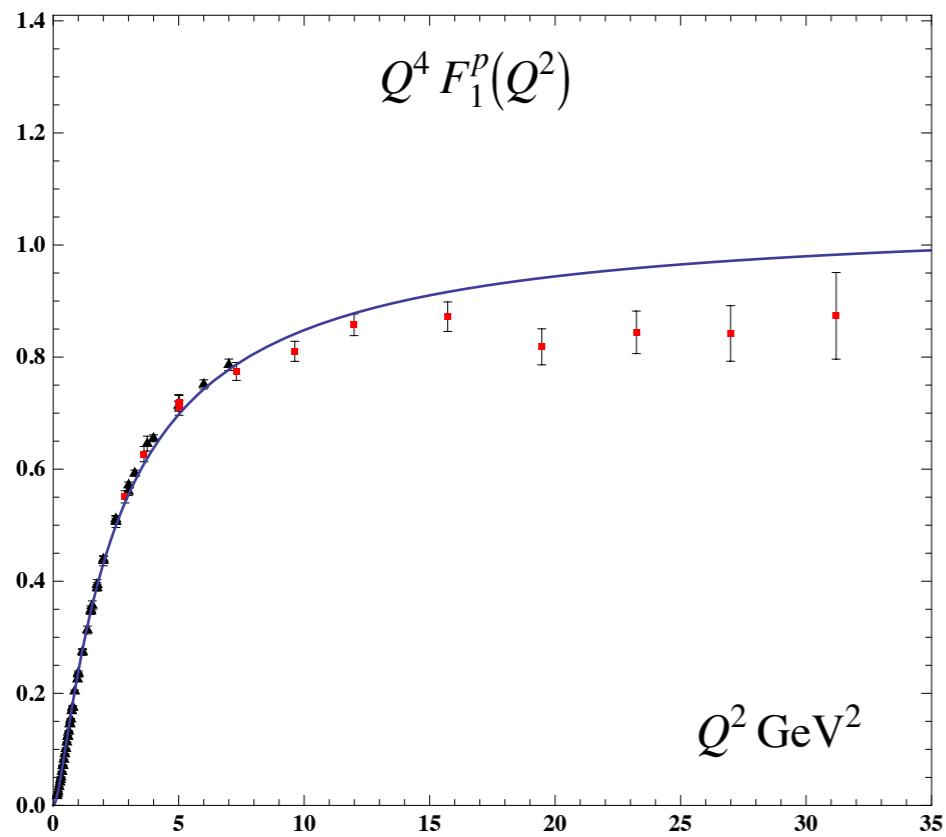
Spacelike Pauli Form Factor

From overlap of $L = 1$ and $L = 0$ LFWFs



Excited Baryons in Holographic QCD

G. de Teramond & sjb



9-2009
8796A3

Nucleon Transition Form Factors

- Compute spin non-flip EM transition $N(940) \rightarrow N^*(1440)$: $\Psi_+^{n=0,L=0} \rightarrow \Psi_+^{n=1,L=0}$
- Transition form factor

$$F_{1N \rightarrow N^*}^p(Q^2) = R^4 \int \frac{dz}{z^4} \Psi_+^{n=1,L=0}(z) V(Q, z) \Psi_+^{n=0,L=0}(z)$$

- Orthonormality of Laguerre functions $(F_{1N \rightarrow N^*}^p(0) = 0, \quad V(Q=0, z) = 1)$

$$R^4 \int \frac{dz}{z^4} \Psi_+^{n',L}(z) \Psi_+^{n,L}(z) = \delta_{n,n'}$$

- Find

$$F_{1N \rightarrow N^*}^p(Q^2) = \frac{2\sqrt{2}}{3} \frac{\frac{Q^2}{M_P^2}}{\left(1 + \frac{Q^2}{\mathcal{M}_\rho^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho'}^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho''}^2}\right)}$$

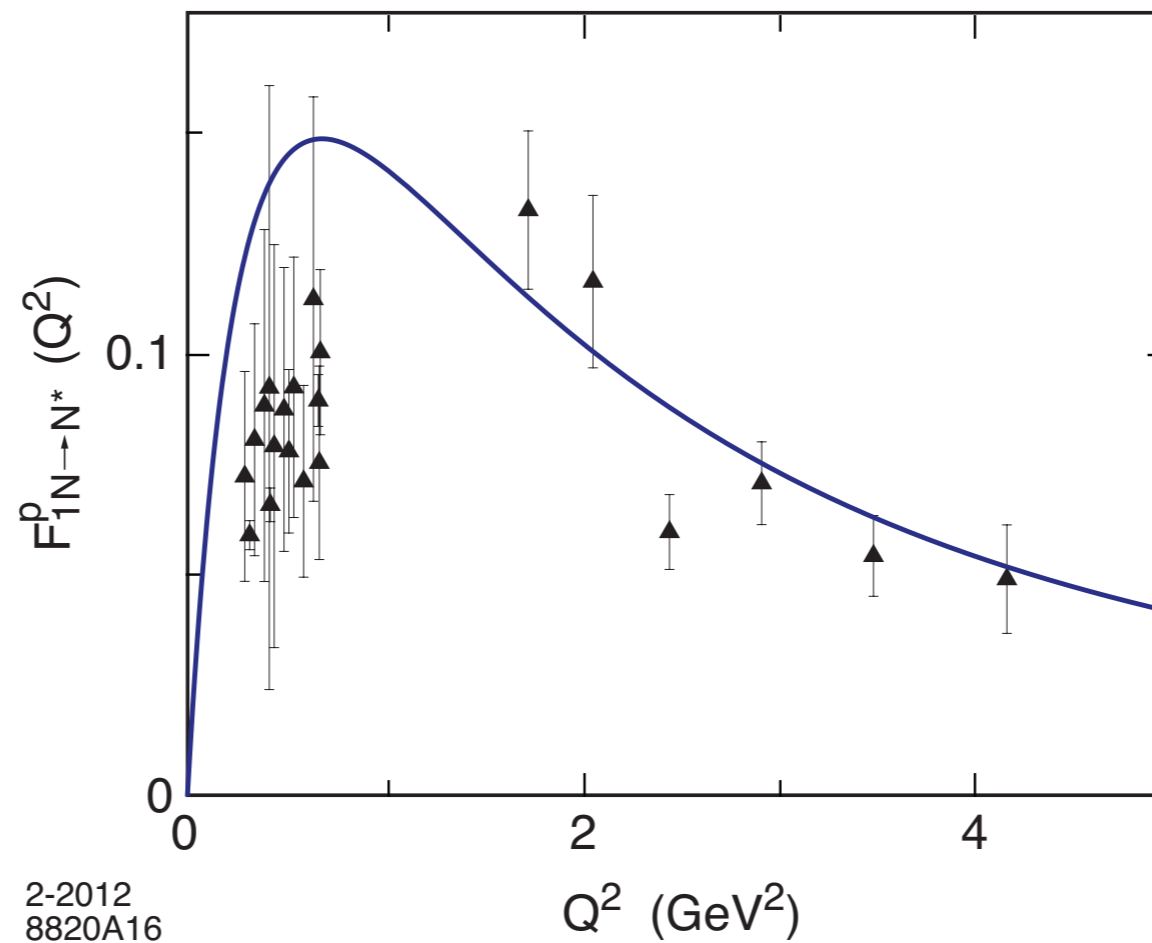
with $\mathcal{M}_{\rho_n}^2 \rightarrow 4\kappa^2(n + 1/2)$

de Teramond, sjb

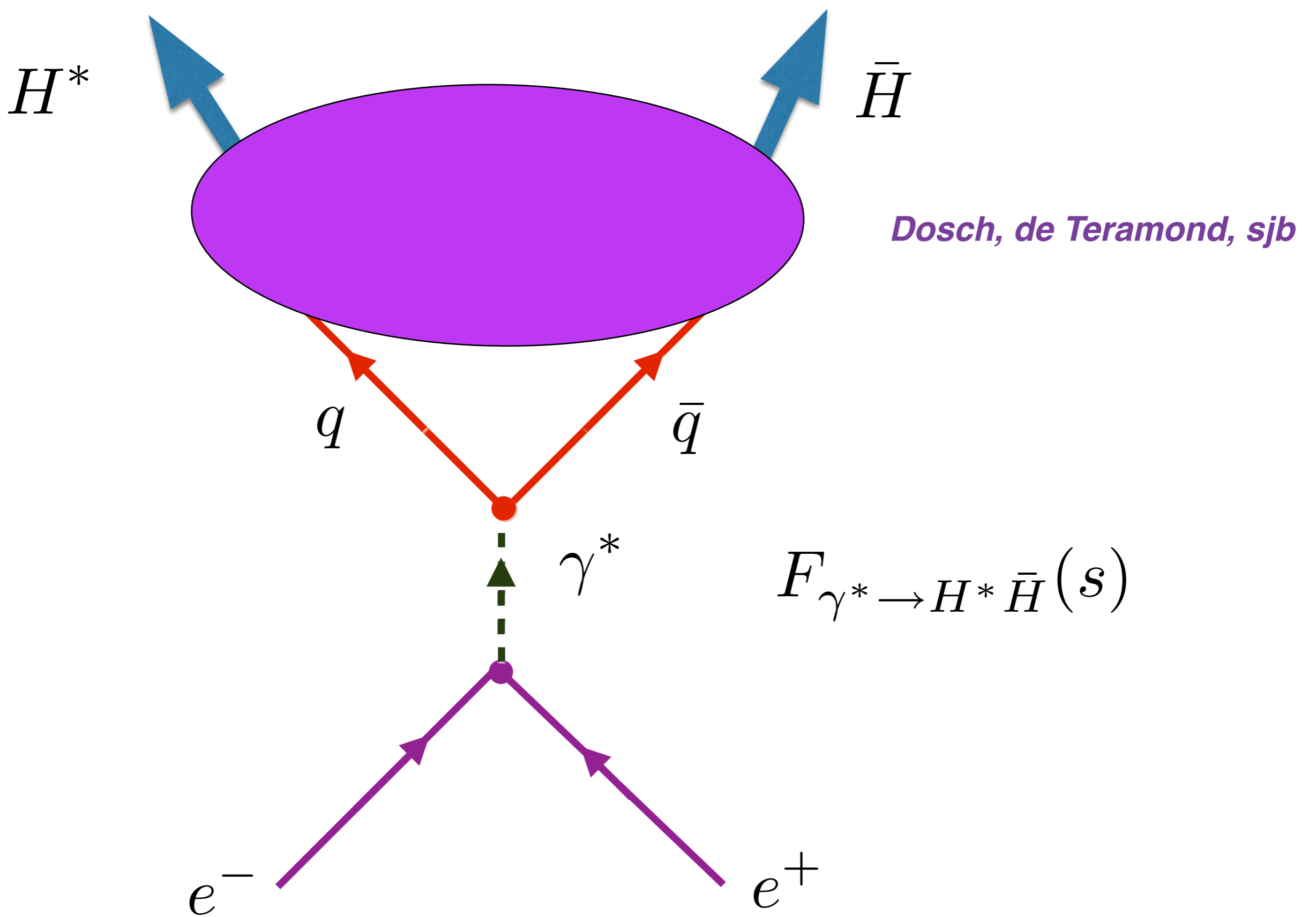
Consistent with counting rule, twist 3

Nucleon Transition Form Factors

$$F_{1N \rightarrow N^*}^p(Q^2) = \frac{\sqrt{2}}{3} \frac{\frac{Q^2}{\mathcal{M}_\rho^2}}{\left(1 + \frac{Q^2}{\mathcal{M}_\rho^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho'}^2}\right) \left(1 + \frac{Q^2}{\mathcal{M}_{\rho''}^2}\right)}.$$



Proton transition form factor to the first radial excited state. Data from JLab



Prediction from Super Conformal AdS/QCD:
 Same Form Factors for $H=M$ and $H=B$ if $L_M=L_B+1$

Bjorken sum rule defines effective charge

$$\alpha_{g1}(Q^2)$$

$$\int_0^1 dx [g_1^{ep}(x, Q^2) - g_1^{en}(x, Q^2)] \equiv \frac{g_a}{6} \left[1 - \frac{\alpha_{g1}(Q^2)}{\pi} \right]$$

- **Can be used as standard QCD coupling**
- **Well measured**
- **Asymptotic freedom at large Q^2**
- **Computable at large Q^2 in any pQCD scheme**
- **Universal β_0, β_1**



Running Coupling from Modified AdS/QCD

Deur, de Teramond, sjb

- Consider five-dim gauge fields propagating in AdS_5 space in dilaton background $\varphi(z) = \kappa^2 z^2$

$$S = -\frac{1}{4} \int d^4x dz \sqrt{g} e^{\varphi(z)} \frac{1}{g_5^2} G^2$$

- Flow equation

$$\frac{1}{g_5^2(z)} = e^{\varphi(z)} \frac{1}{g_5^2(0)} \quad \text{or} \quad g_5^2(z) = e^{-\kappa^2 z^2} g_5^2(0)$$

where the coupling $g_5(z)$ incorporates the non-conformal dynamics of confinement

- YM coupling $\alpha_s(\zeta) = g_{YM}^2(\zeta)/4\pi$ is the five dim coupling up to a factor: $g_5(z) \rightarrow g_{YM}(\zeta)$
- Coupling measured at momentum scale Q

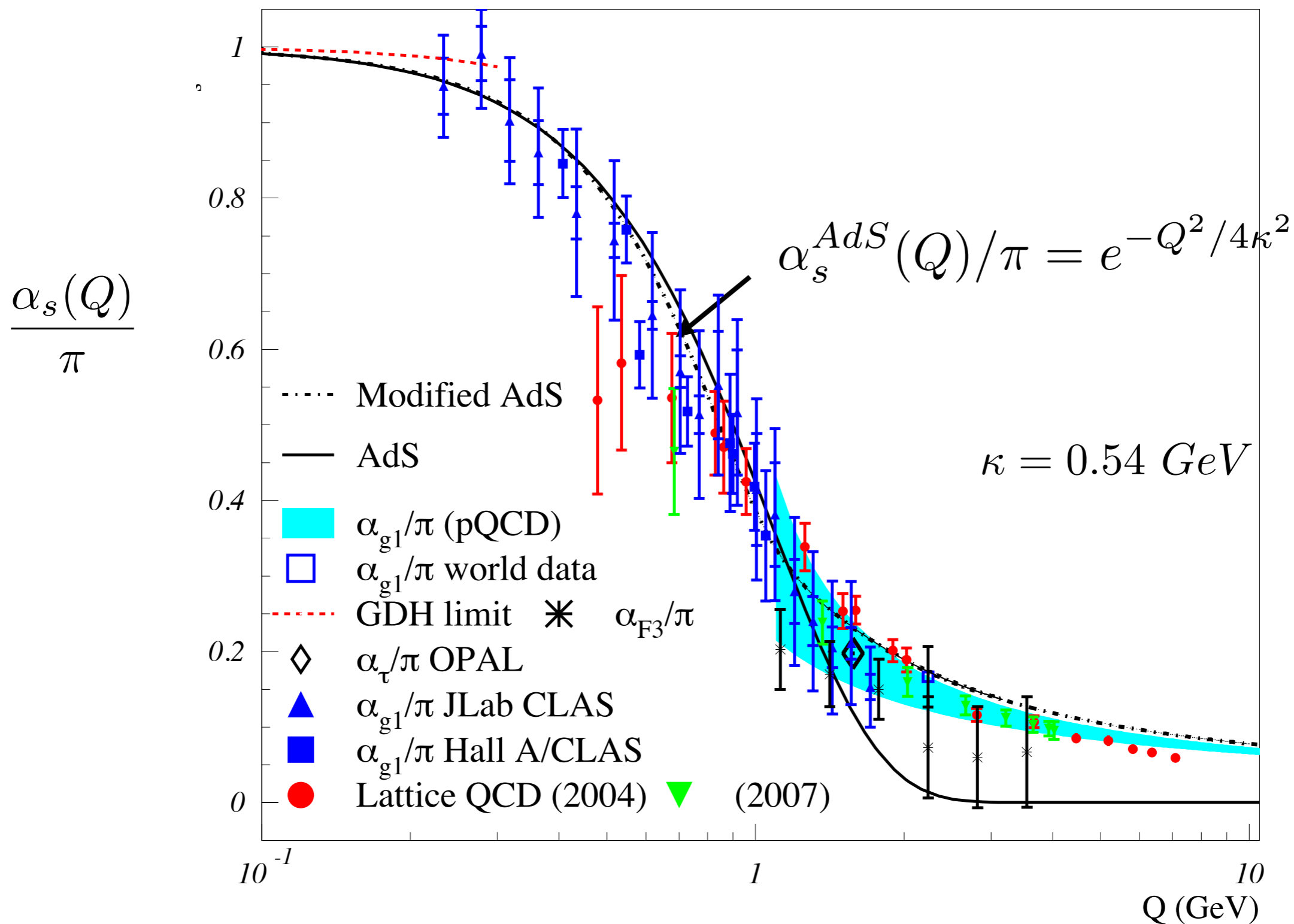
$$\alpha_s^{AdS}(Q) \sim \int_0^\infty \zeta d\zeta J_0(\zeta Q) \alpha_s^{AdS}(\zeta)$$

- Solution

$$\alpha_s^{AdS}(Q^2) = \alpha_s^{AdS}(0) e^{-Q^2/4\kappa^2}.$$

where the coupling α_s^{AdS} incorporates the non-conformal dynamics of confinement

Analytic, defined at all scales, IR Fixed Point



AdS/QCD dilaton captures the higher twist corrections to effective charges for $Q < 1 \text{ GeV}$

$$e^\varphi = e^{+\kappa^2 z^2}$$

Deur, de Teramond, sjb
99

$$m_\rho = \sqrt{2}\kappa$$

$$m_p = 2\kappa_1$$

Deur, de Teramond, sjb

All-Scale QCD Coupling

$$\frac{\alpha_{g_1}^s(Q^2)}{\pi}$$

$$e^{-\frac{Q^2}{4\kappa^2}}$$

Expt:
 $\Lambda_{\overline{MS}} = 0.341 \pm 0.024 \text{ GeV}$

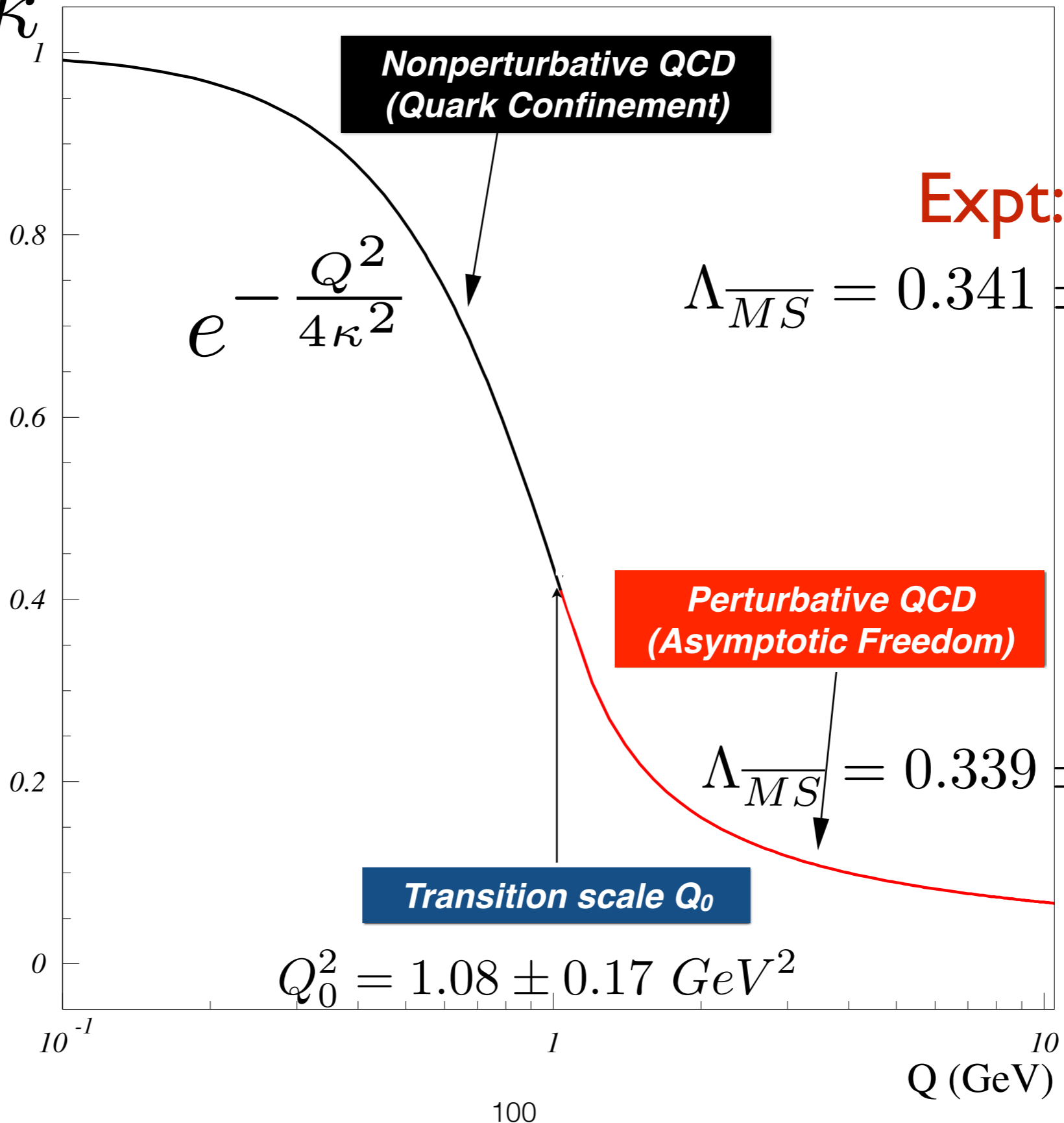
**Perturbative QCD
(Asymptotic Freedom)**

$$\Lambda_{\overline{MS}} = 0.339 \pm 0.016 \text{ GeV}$$

Transition scale Q_0

$$Q_0^2 = 1.08 \pm 0.17 \text{ GeV}^2$$

$$\lambda \equiv \kappa^2$$



Interpretation of Mass Scale \mathcal{K}

- Does not affect conformal symmetry of QCD action
- Self-consistent regularization of IR divergences
- Determines all mass and length scales for zero quark mass
- Compute scheme-dependent $\Lambda_{\overline{MS}}$ determined in terms of \mathcal{K}
- Value of \mathcal{K} itself not determined -- place holder
- Need external constraint such as f_π

AdS/QCD and Light-Front Holography

$$\mathcal{M}_{n,J,L}^2 = 4\kappa^2 \left(n + \frac{J+L}{2} \right)$$

- **Zero mass pion for $m_q = 0$ ($n=J=L=0$)**
- **Regge trajectories: equal slope in n and L**
- **Form Factors at high Q^2 : Dimensional counting**
 $[Q^2]^{n-1} F(Q^2) \rightarrow \text{const}$
- **Space-like and Time-like Meson and Baryon Form Factors**
- **Running Coupling for NPQCD** $\alpha_s(Q^2) \propto e^{-\frac{Q^2}{4\kappa^2}}$
- **Meson Distribution Amplitude** $\phi_\pi(x) \propto f_\pi \sqrt{x(1-x)}$



Features of AdS/QCD

- **Color confining potential $\kappa^4 \zeta^2$ and universal mass scale from dilaton**
$$e^{\phi(z)} = e^{\kappa^2 z^2} \quad \alpha_s(Q^2) \propto \exp -Q^2/4\kappa^2$$
- **Dimensional transmutation** $\Lambda_{\overline{MS}} \leftrightarrow \kappa \leftrightarrow m_H$
- **Chiral Action remains conformally invariant despite mass scale** *DAFF*
- **Light-Front Holography: Duality of AdS and frame-independent LF QCD**
- **Reproduces observed Regge spectroscopy — same slope in n, L, and J for mesons and baryons**
- **Massless pion for massless quark**
- **Supersymmetric meson-baryon dynamics and spectroscopy:**
 $L_M = L_B + I$
- **Dynamics: LFWFs, Form Factors, GPDs**

*Superconformal Algebra
Fubini and Rabinovici*

An analytic first approximation to QCD

AdS/QCD + Light-Front Holography

- **As Simple as Schrödinger Theory in Atomic Physics**
- **LF radial variable ζ conjugate to invariant mass squared**
- **Relativistic, Frame-Independent, Color-Confining**
- **Unique confining potential!**
- **QCD Coupling at all scales: Essential for Gauge Link phenomena**
- **Hadron Spectroscopy and Dynamics from one parameter**
- **Wave Functions, Form Factors, Hadronic Observables, Constituent Counting Rules**
- **Insight into QCD Condensates: Zero cosmological constant!**
- **Systematically improvable with DLCQ-BLFQ Methods**

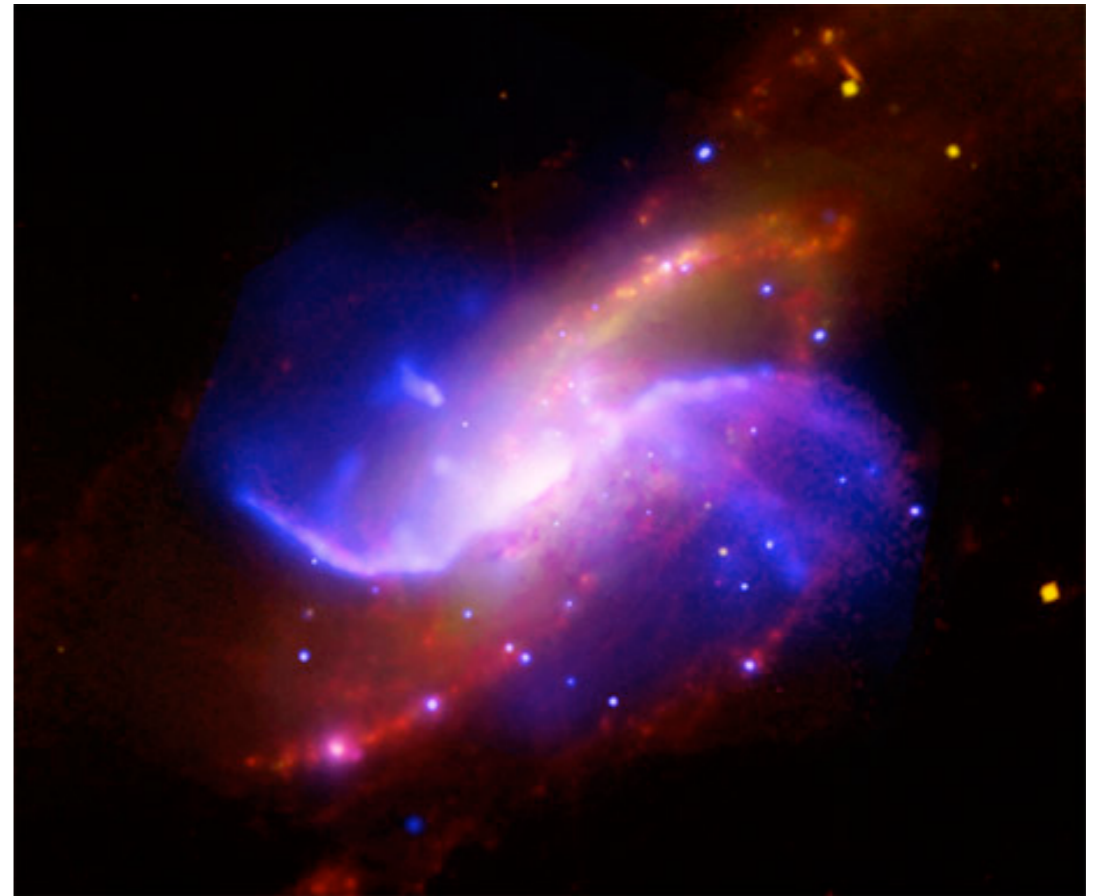


Future Directions for AdS/QCD

- **Hadronization at the Amplitude Level**
- **Diffraction dissociation of pion and proton to jets**
- **Identify the factorization Scale for ERBL, DGLAP evolution: Q_0**
- **Compute Tetraquark Spectroscopy Sequentially**
- **Update $SU(6)$ spin-flavor symmetry**
- **Heavy Quark States: Supersymmetry, not conformal**
- **Compute higher Fock states; e.g. Intrinsic Heavy Quarks**
- **Nuclear States — Hidden Color**
- **Basis LF Quantization**

*We view the universe
as light reaches us
along the light-front
at fixed*

$$\tau = t + z/c$$



Front Form Vacuum Describes the Empty, Causal Universe

“One of the gravest puzzles of theoretical physics”

DARK ENERGY AND THE COSMOLOGICAL CONSTANT PARADOX

A. ZEE

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Kavil Institute for Theoretical Physics, University of California,
Santa Barbara, CA 93106, USA
zee@kitp.ucsb.edu*

$$(\Omega_{\Lambda})_{QCD} \sim 10^{45}$$

$$(\Omega_{\Lambda})_{EW} \sim 10^{56}$$

$$\Omega_{\Lambda} = 0.76(\text{expt})$$

Extraordinary conflict between the conventional definition of the vacuum in quantum field theory and cosmology

Elements of the solution:

(A) Light-Front Quantization: causal, frame-independent vacuum

(B) New understanding of QCD “Condensates”

(C) Higgs Light-Front Zero Mode

Two Definitions of Vacuum State

Instant Form: Lowest Energy Eigenstate of Instant-Form Hamiltonian

$$H|\psi_0\rangle = E_0|\psi_0\rangle, E_0 = \min\{E_i\}$$

*Eigenstate defined at one time t over all space;
Acausal! Frame-Dependent*

Front Form: Lowest Invariant Mass Eigenstate of Light-Front Hamiltonian

$$H_{LF}|\psi_0\rangle_{LF} = M_0^2|\psi_0\rangle_{LF}, M_0^2 = 0.$$

*Frame-independent eigenstate at fixed LF time $\tau = t+z/c$
within causal horizon*

Frame-independent description of the causal physical universe!

Light-Front vacuum can simulate empty universe

Shrock, Tandy, Roberts, sjb

- **Independent of observer frame**
- **Causal**
- **Lowest invariant mass state $M=0$.**
- **Trivial up to $k^+=0$ zero modes-- already normal-ordering**
- **Higgs theory consistent with trivial LF vacuum (Srivastava, sjb)**
- **QCD and AdS/QCD: “In-hadron” condensates (Maris, Tandy Roberts) -- GMOR satisfied.**
- **QED vacuum; no loops**
- **Zero cosmological constant from QED, QCD**

Kunshan
August 3, 2015



New Perspectives for Hadron Physics

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Standard Model on the Light-Front

- Same phenomenological predictions
- Higgs field has three components
- Real part creates Higgs particle
- Imaginary part (Goldstone) become longitudinal components of W, Z
- *Higgs VEV of instant form becomes $k^+=0$ LF zero mode!*
- Analogous to a background static classical Zeeman or Stark Fields
- Zero contribution to T^{μ}_{μ} ; zero coupling to gravity



Goals

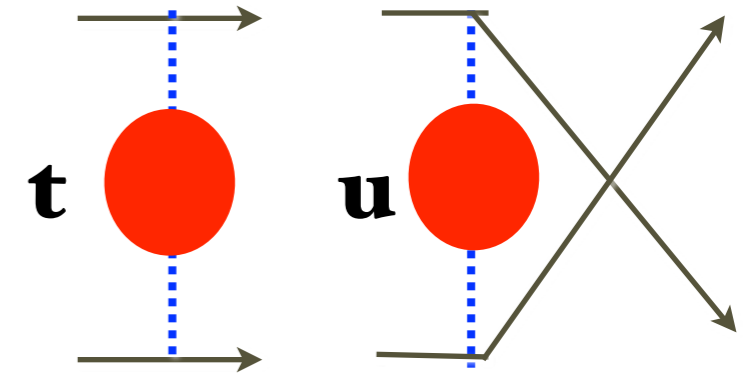
- **Test QCD to maximum precision at the LHC**
- **Maximize sensitivity to new physics**
- **High precision determination of fundamental parameters**
- **Determine renormalization scales without ambiguity**
- **Eliminate scheme dependence**

Predictions for physical observables cannot depend on theoretical conventions such as the renormalization scheme

Electron-Electron Scattering in QED

$$\mathcal{M}_{ee \rightarrow ee}(++;++) = \frac{8\pi s}{t} \alpha(t) + \frac{8\pi s}{u} \alpha(u)$$

- **Two separate physical scales: $t, u =$ photon virtuality**
- **Gauge Invariant. Dressed photon propagator**
- **Sums all vacuum polarization, non-zero beta terms into running coupling. This is the purpose of the running coupling!**
- **If one chooses a different initial scale, one must sum an infinite number of graphs -- but always recover same result!**
- **Number of active leptons correctly set**
- **Analytic: reproduces correct behavior at lepton mass thresholds**
- **No renormalization scale ambiguity!**



δ -Renormalization Scheme (\mathcal{R}_δ scheme)

In dim. reg. $1/\epsilon$ poles come in powers of [Bollini & Gambiagi, 't Hooft & Veltman, '72]

$$\ln \frac{\mu^2}{\Lambda^2} + \frac{1}{\epsilon} + c$$

In the **modified minimal subtraction** scheme ($\overline{\text{MS}}$) one subtracts together with the pole a constant [Bardeen, Buras, Duke, Muta (1978) on DIS results]:

$$\ln(4\pi) - \gamma_E$$

This corresponds to a shift in the scale:

$$\mu_{\overline{\text{MS}}}^2 = \mu^2 \exp(\ln 4\pi - \gamma_E)$$

A finite subtraction from infinity is arbitrary. *Let's make use of this!*

Subtract an arbitrary constant and keep it in your calculation: \mathcal{R}_δ -scheme

M. Mojaza, Xing-Gang Wu, sjb

$$\ln(4\pi) - \gamma_E - \delta,$$

$$\mu_\delta^2 = \mu_{\overline{\text{MS}}}^2 \exp(-\delta) = \mu^2 \exp(\ln 4\pi - \gamma_E - \delta)$$

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Exposing the Renormalization Scheme Dependence

Observable in the \mathcal{R}_δ -scheme:

$$\rho_\delta(Q^2) = r_0 + r_1 a(\mu) + [r_2 + \beta_0 r_1 \delta] a(\mu)^2 + [r_3 + \beta_1 r_1 \delta + 2\beta_0 r_2 \delta + \beta_0^2 r_1 \delta^2] a(\mu)^3 + \dots$$

$$\mathcal{R}_0 = \overline{\text{MS}}, \quad \mathcal{R}_{\ln 4\pi - \gamma_E} = \text{MS} \quad \mu^2 = \mu_{\overline{\text{MS}}}^2 \exp(\ln 4\pi - \gamma_E), \quad \mu_{\delta_2}^2 = \mu_{\delta_1}^2 \exp(\delta_2 - \delta_1)$$

Note the divergent 'renormalon series' $n! \beta^n \alpha_s^n$

Renormalization Scheme Equation

$$\frac{d\rho}{d\delta} = -\beta(a) \frac{d\rho}{da} \stackrel{!}{=} 0 \quad \longrightarrow \text{PMC}$$

$$\rho_\delta(Q^2) = r_0 + r_1 a_1(\mu_1) + (r_2 + \beta_0 r_1 \delta_1) a_2(\mu_2)^2 + [r_3 + \beta_1 r_1 \delta_1 + 2\beta_0 r_2 \delta_2 + \beta_0^2 r_1 \delta_1^2] a_3(\mu_3)^3$$

The $\delta_k^p a^n$ -term indicates the term associated to a diagram with $1/\epsilon^{n-k}$ divergence for any p . Grouping the different δ_k -terms, one recovers in the $N_c \rightarrow 0$ Abelian limit the dressed skeleton expansion.



Special Degeneracy in PQCD

There is nothing special about a particular value for δ , thus for any δ

$$\rho(Q^2) = r_{0,0} + r_{1,0}a(Q) + [r_{2,0} + \beta_0 \underline{r_{2,1}}]a(Q)^2 + [r_{3,0} + \beta_1 \underline{r_{2,1}} + 2\beta_0 \underline{r_{3,1}} + \beta_0^2 \underline{r_{3,2}}]a(Q)^3 \\ + [r_{4,0} + \beta_2 \underline{r_{2,1}} + 2\beta_1 \underline{r_{3,1}} + \frac{5}{2}\beta_1 \beta_0 \underline{r_{3,2}} + 3\beta_0 r_{4,1} + 3\beta_0^2 r_{4,2} + \beta_0^3 r_{4,3}]a(Q)^4$$

According to the **principal of maximum conformality** we must set the scales such to absorb all 'renormalon-terms', i.e. **non-conformal terms**

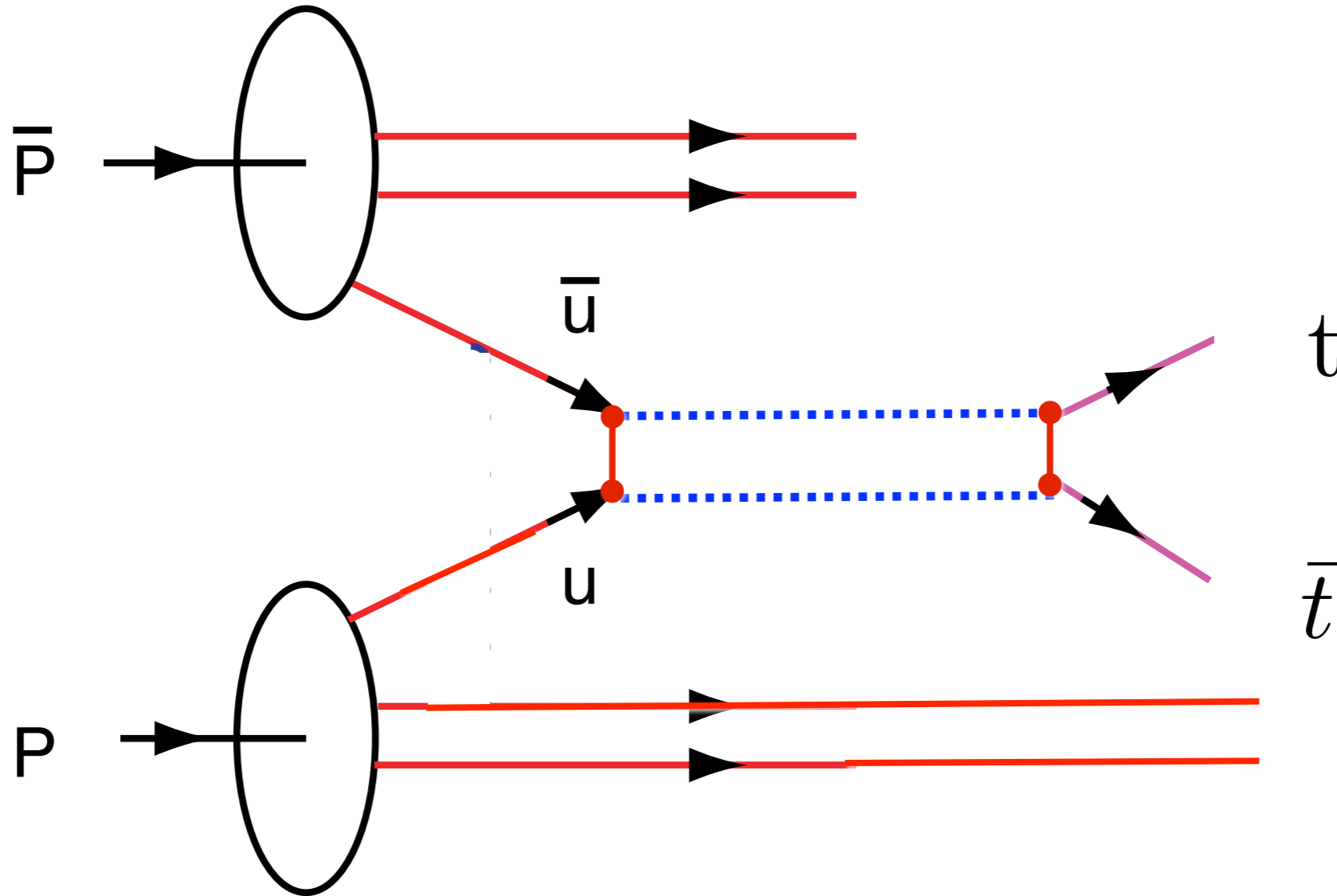
$$\rho(Q^2) = r_{0,0} + r_{1,0}a(Q) + (\beta_0 a(Q)^2 + \beta_1 a(Q)^3 + \beta_2 a(Q)^4 + \dots) \underline{r_{2,1}} \\ + (\beta_0^2 a(Q)^3 + \frac{5}{2}\beta_1 \beta_0 a(Q)^4 + \dots) \underline{r_{3,2}} + (\beta_0^3 + \dots) r_{4,3} \\ + r_{2,0}a(Q)^2 + 2a(Q)(\beta_0 a(Q)^2 + \beta_1 a(Q)^3 + \dots) \underline{r_{3,1}} \\ + \dots$$

$$r_{1,0}a(Q_1) = r_{1,0}a(Q) - \beta(a)r_{2,1} + \frac{1}{2}\beta(a)\frac{\partial\beta}{\partial a}r_{3,2} + \dots + \frac{(-1)^n}{n!} \frac{d^{n-1}\beta}{(d \ln \mu^2)^{n-1}} r_{n+1,n}$$

$$r_{2,0}a(Q_2)^2 = r_{2,0}a(Q)^2 - 2a(Q)\beta(a)r_{3,1} + \dots$$



Implications for the $\bar{p}p \rightarrow t\bar{t}X$ asymmetry at the Tevatron



Interferes with Born term.

Small value of renormalization scale increases asymmetry, just as in QED

Xing-Gang Wu, sjb

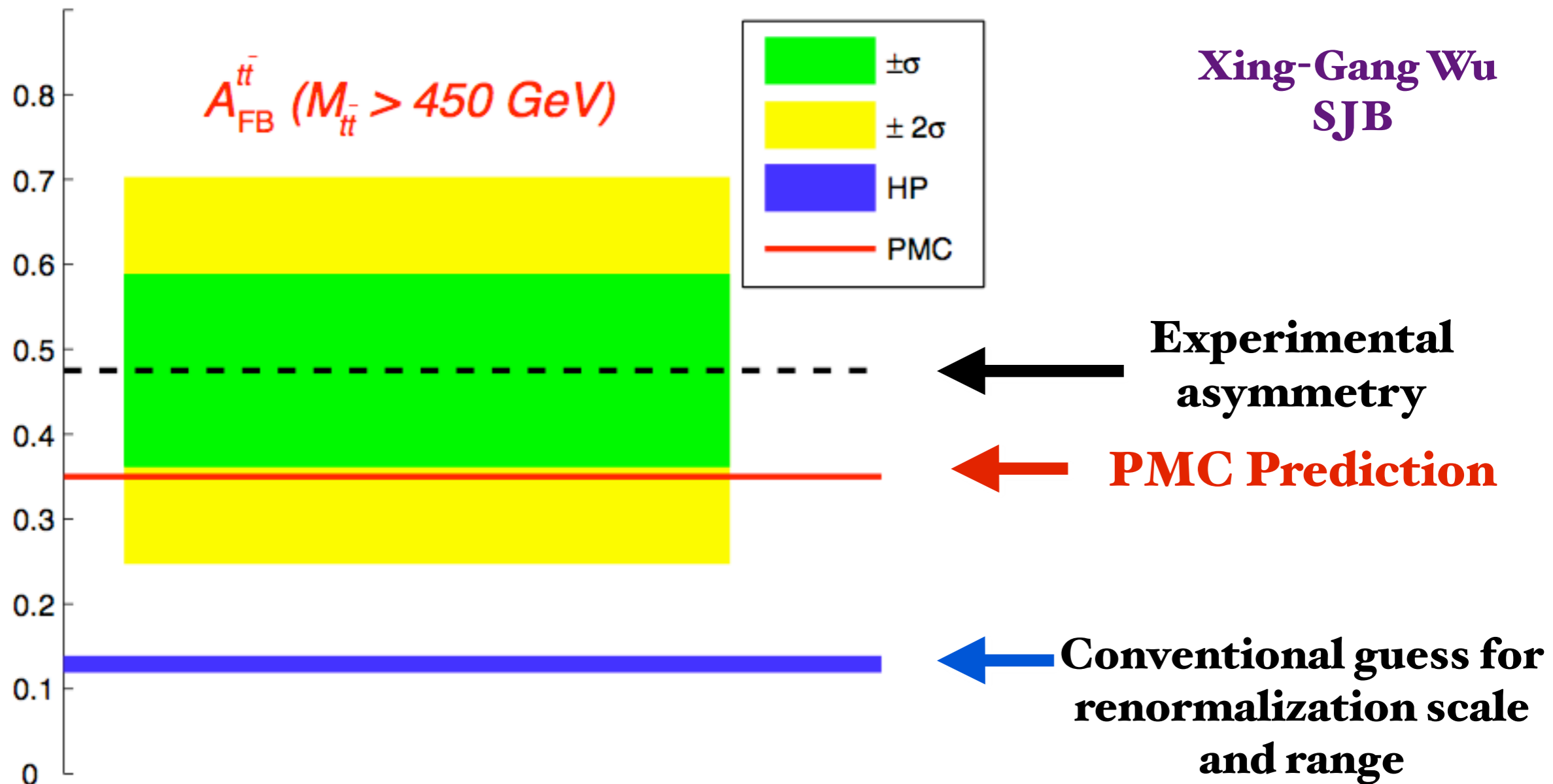
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**Kunshan
August 3, 2015**



New Perspectives for Hadron Physics

*The Renormalization Scale Ambiguity for Top-Pair Production
Eliminated Using the 'Principle of Maximum Conformality' (PMC)*

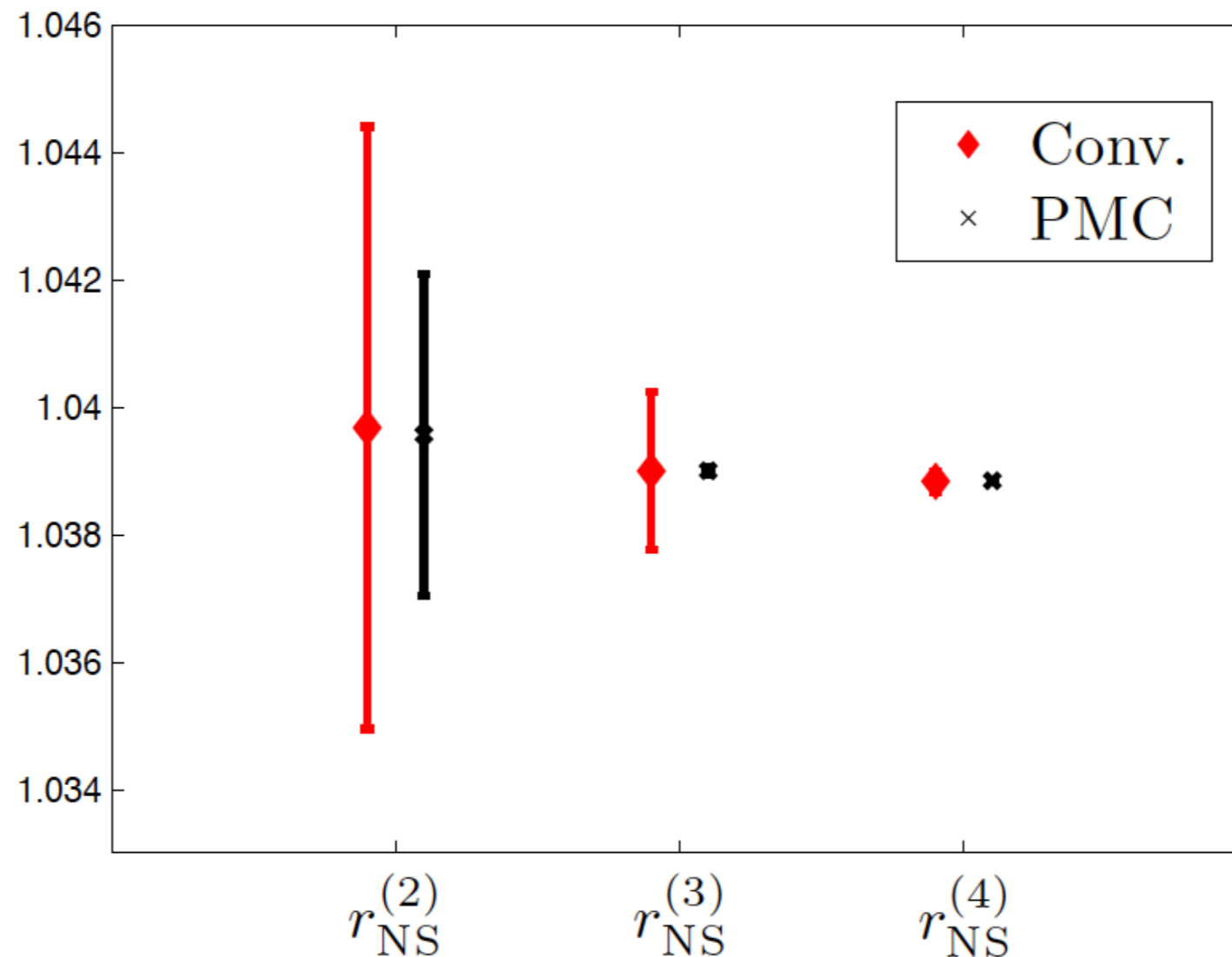


Top quark forward-backward asymmetry predicted by pQCD NNLO within 1σ of CDF/D0 measurements using PMC/BLM scale setting

Reanalysis of the Higher Order Perturbative QCD corrections to Hadronic Z Decays using the Principle of Maximum Conformality

S-Q Wang, X-G Wu, sjb

P.A. Baikov, K.G. Chetyrkin, J.H. Kuhn, and J. Rittinger,
Phys. Rev. Lett. 108, 222003 (2012).



The values of $r_{\text{NS}}^{(n)} = 1 + \sum_{i=1}^n C_i^{\text{NS}} a_s^i$ and their errors $\pm |C_n^{\text{NS}} a_s^n|_{\text{MAX}}$. The diamonds and the crosses are for conventional (Conv.) and PMC scale settings, respectively. The central values assume the initial scale choice $\mu_r^{\text{init}} = M_Z$.

What is PMC ?

Principle of Maximum Conformality

Choose renormalization scheme; e.g. $\alpha_s^R(\mu_R^{\text{init}})$

Choose μ_R^{init} ; arbitrary initial renormalization scale

Identify $\{\beta_i^R\}$ – terms using n_f – terms through the PMC – BLM correspondence principle

order-by-order

Shift scale of α_s to μ_R^{PMC} to eliminate $\{\beta_i^R\}$ – terms

Conformal Series

Result is independent of μ_R^{init} and scheme at fixed order

*Xing-Gang Wu, Martin Mojaza
Leonardo di Giustino, SJB*

PMC-BLM – one

Phys. Rev. Lett. **109**, 042002 (2012)

R_δ -scheme – two

Phys. Rev. Lett. **110**, 192001 (2013)

Eliminate β -terms

n_f dependence of pQCD series does not uniquely identify the β terms

Features of BLM/PMC

- **Predictions are scheme-independent**
- **Matches conformal series**
- **Commensurate Scale Relations between observables: Generalized Crewther Relation (Kataev, Lu, Rathsmann, sjb)**
- **No $n!$ Renormalon growth**
- **New scale at each order; n_F determined at each order**
- **Multiple Physical Scales Incorporated**
- **Rigorous: Satisfies all Renormalization Group Principles**
- **Realistic Estimate of Higher-Order Terms**
- **Eliminates unnecessary theory error**

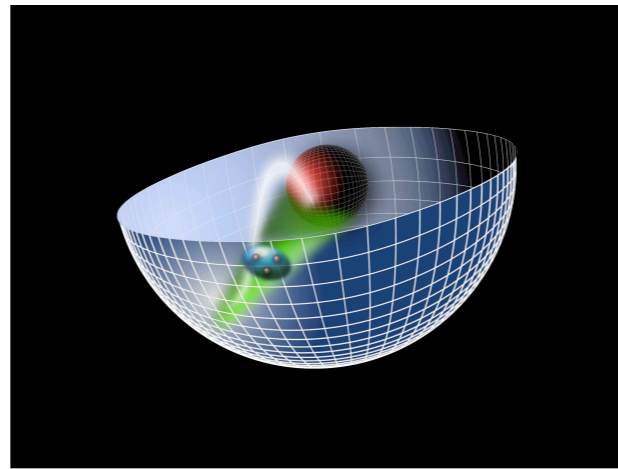
QCD Myths

- **Anti-Shadowing is Universal**
- **ISI and FSI are higher twist effects and universal**
- **High transverse momentum hadrons arise only from jet fragmentation -- baryon anomaly!**
- **Heavy quarks only from gluon splitting**
- **Renormalization scale cannot be fixed**
- **QCD condensates are vacuum effects**
- **QCD gives 10^{42} to the cosmological constant**



*AdS/QCD
Soft-Wall Model*

$$e^{\varphi(z)} = e^{+\kappa^2 z^2}$$



$$\zeta^2 = x(1-x)b_{\perp}^2.$$

Light-Front Holography

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1-4L^2}{4\zeta^2} + U(\zeta) \right] \psi(\zeta) = \mathcal{M}^2 \psi(\zeta)$$



Light-Front Schrödinger Equation

$$U(\zeta) = \kappa^4 \zeta^2 + 2\kappa^2(L + S - 1)$$

$$\kappa \simeq 0.6 \text{ GeV}$$

Confinement scale:

$$1/\kappa \simeq 1/3 \text{ fm}$$

***Unique
Confinement Potential!***

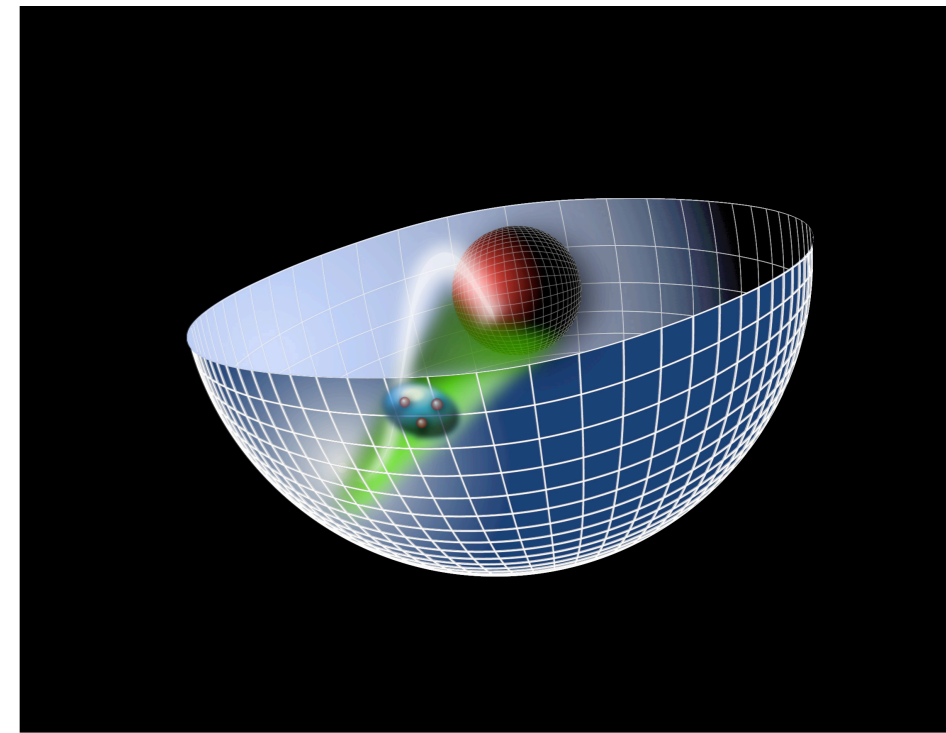
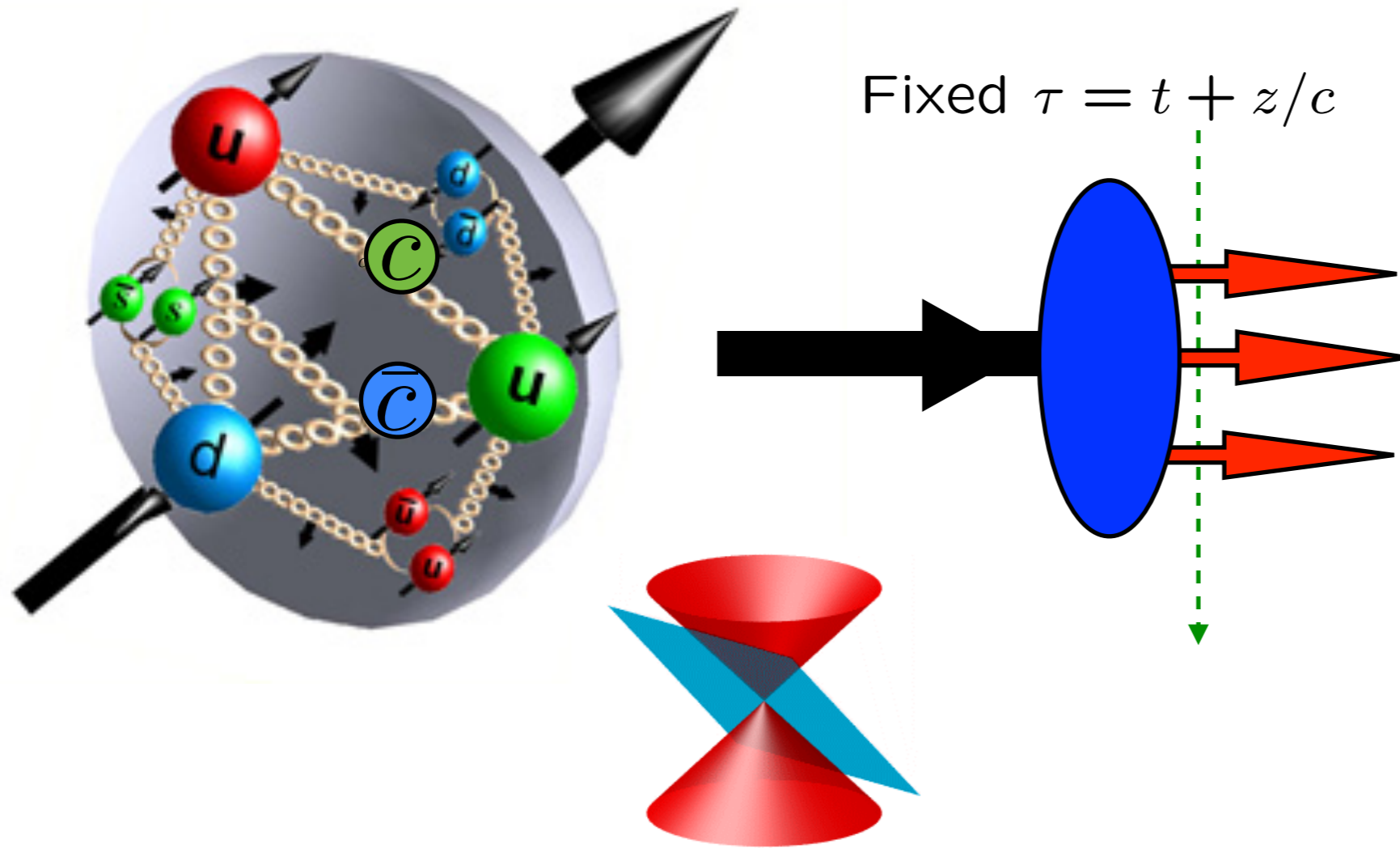
*Preserves Conformal Symmetry
of the action*

● **de Alfaro, Fubini, Furlan:**

● **Fubini, Rabinovici:**

**Scale can appear in Hamiltonian and EQM
without affecting conformal invariance of action!**

New Perspectives for Hadron Physics



*Kunshan, China
August 3, 2015*

Stan Brodsky

