

# Parton Distributions from Proton to nuclei

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# Outline

- \* Parton distributions and QCD evolution equations;
- \* A new approach to parton distributions in the proton and nuclei;
- \* Quark saturation in low  $Q^2$  range.

# Parton distributions and QCD evolution equations

- \* Quark and gluon distributions are important knowledge of nucleon and nuclei.
- \* Worldwide proposing EIC will complete our knowledge of sea quark, gluon and nuclear media modification.

MEIC, eRHIC, EIC@HIAF...

# Parton distributions and QCD evolution equations

\* “QCD is successful and we believe QCD”

DGLAP  $x > 10^{-3}, Q^2 > 1 \sim 4 \text{ GeV}^2$

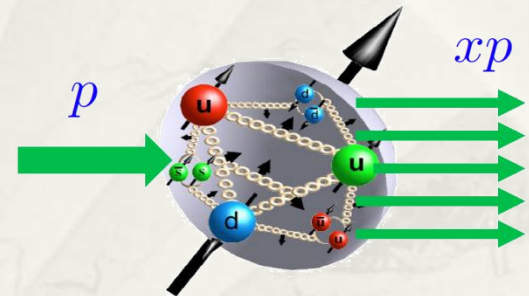
GLR-MQ  $x < \sim 10^{-2}$

BFKL  $x < \sim 10^{-4}$

BK  $x < \sim 10^{-4}$

JIMWLK  $x < \sim 10^{-4}, Q^2 > 1 \text{ GeV}^2$

ZSR  $x > 10^{-4}, Q^2 > 0.06 \text{ GeV}^2$



QCD evolution equations



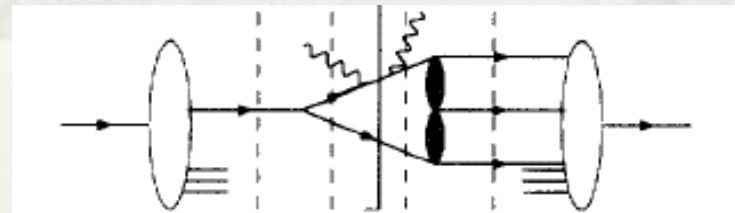
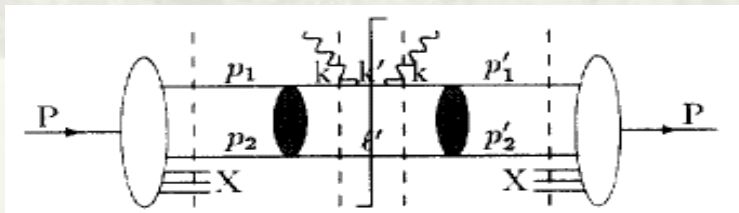
# Parton distributions and QCD evolution equations

GLR-MQ-ZSR: DGLAP + Leading high twist corrections

W. Zhu, Nucl. Phys. B551, 245 (1999);

W. Zhu and J.H.Ruan, Nucl. Phys. B559, 378 (1999);

W. Zhu, Z.Q. Shen and J.H. Ruan, Nucl. Phys. B692 ,417.(2004)



# GLR-MQ-ZSR

For valence quark  $Q^2 \frac{dx f_{v_j}(x, Q^2)}{dQ^2}$

$$\begin{aligned}
 &= \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} P_{qq}(z) x f_{v_j}(y, Q^2) \\
 &\quad - \frac{\alpha_s(Q^2)}{2\pi} x f_{v_j}(x, Q^2) \int_0^1 dz P_{qq}(z) \\
 &\quad - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{qg \rightarrow qg}(z) [y f_g(y, Q^2) y f_{v_j}(y, Q^2)] \\
 &\quad + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{qg \rightarrow qg}(z) [y f_g(y, Q^2) y f_{v_j}(y, Q^2)] \\
 &\quad - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{qq \rightarrow qq}(z) [y \Sigma_{sea}(y, Q^2) y f_{v_j}(y, Q^2)] \\
 &\quad + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{qq \rightarrow qq}(z) [y \Sigma_{sea}(y, Q^2) y f_{v_j}(y, Q^2)],
 \end{aligned}$$



# GLR-MQ-ZSR

$$\begin{aligned}
 & Q^2 \frac{dx f_{\bar{q}_i}(x, Q^2)}{dQ^2} \\
 = & \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} P_{qq}(z) x f_{\bar{q}_i}(y, Q^2) \\
 & - \frac{\alpha_s(Q^2)}{2\pi} x f_{\bar{q}_i}(x, Q^2) \int_0^1 dz P_{qq}(z) \\
 & + \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} P_{qg}(z) x f_g(y, Q^2) \\
 & - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{gq \rightarrow q\bar{q}}(z) [y f_g(y, Q^2)]^2 \\
 & + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{gq \rightarrow q\bar{q}}(z) [y f_g(y, Q^2)]^2 \\
 & - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{q\bar{q} \rightarrow q\bar{q}}(z) y f_{q_i}(x, Q^2) y f_{\bar{q}_i}(y, Q^2) \\
 & + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{q\bar{q} \rightarrow q\bar{q}}(z) y f_{q_i}(x, Q^2) y f_{\bar{q}_i}(y, Q^2) \\
 & - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{qq \rightarrow qq}(z) y [\Sigma(x, Q^2) - f_{q_i}(x, Q^2)] y f_{\bar{q}_i}(y, Q^2) \\
 & + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{qq \rightarrow qq}(z) y [\Sigma(x, Q^2) - f_{q_i}(x, Q^2)] y f_{\bar{q}_i}(y, Q^2) \\
 & - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{qg \rightarrow qg}(z) [y f_g(y, Q^2) y f_{\bar{q}_i}(y, Q^2)] \\
 & + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{qg \rightarrow qg}(z) [y f_g(y, Q^2) y f_{\bar{q}_i}(y, Q^2)],
 \end{aligned}$$

For sea quark

# GLR-MQ-ZSR

$$\begin{aligned}
 & Q^2 \frac{dx f_g(x, Q^2)}{dQ^2} \\
 &= \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} P_{gq}(z) x \Sigma(y, Q^2) \\
 &+ \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dy}{y} P_{gg}(z) x f_g(y, Q^2) \\
 &- f \frac{\alpha_s(Q^2)}{2\pi} x f_g(x, Q^2) \int_0^1 dz P_{qg}(z) \\
 &- \frac{1}{2} \frac{\alpha_s(Q^2)}{2\pi} x f_g(x, Q^2) \int_0^1 dz P_{gg}(z) \\
 &- \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{gq \rightarrow gq}(z) [y f_g(y, Q^2)]^2 \\
 &+ \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{gq \rightarrow gq}(z) [y f_g(y, Q^2)]^2
 \end{aligned}$$

$$\begin{aligned}
 & - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{q\bar{q} \rightarrow gq}(z) \sum_{i=1}^f [y f_{\bar{q}_i}(y, Q^2)]^2 \\
 & + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{q\bar{q} \rightarrow gq}(z) \sum_{i=1}^f [y f_{\bar{q}_i}(y, Q^2)]^2 \\
 & - \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_x^{1/2} \frac{dy}{y} x P_{qg \rightarrow qg}(z) [y \Sigma(y, Q^2) y f_g(y, Q^2)] \\
 & + \frac{\alpha_s^2(Q^2)}{4\pi R^2 Q^2} \int_{x/2}^x \frac{dy}{y} x P_{qg \rightarrow qg}(z) [y \Sigma(y, Q^2) y f_g(y, Q^2)],
 \end{aligned}$$

For gluon



# Parton distributions in the proton

Proton(global analysis)

GRV 23 parameters+DGLAP

CTEQ 36 parameters+DGLAP

nuclei HKM +16 parameters

EPS +24 parameters

A large uncertainty ( $>100\%$ ) in nuclear gluon density. We need to pin it down.

# Parton distributions in the proton

- \* GLR-MQ-ZSR + natural input

We have only 4 free parameters to give PDF predictions.

And first realize three valence quark input as the non-perturbative input.

# Parton distributions in the proton

Natural input:  
 $\mu^2 = 0.064 \text{ GeV}^2$

$$x f_{v_u}(x, \mu^2) = A_u x^{B_u} (1-x)^{C_u},$$

$$x f_{v_d}(x, \mu^2) = A_d x^{B_d} (1-x)^{C_d},$$

$$f_g(x, \mu^2) = 0, \quad f_{q_i}(x, \mu^2) = f_{\bar{q}_i}(x, \mu^2) = 0.$$

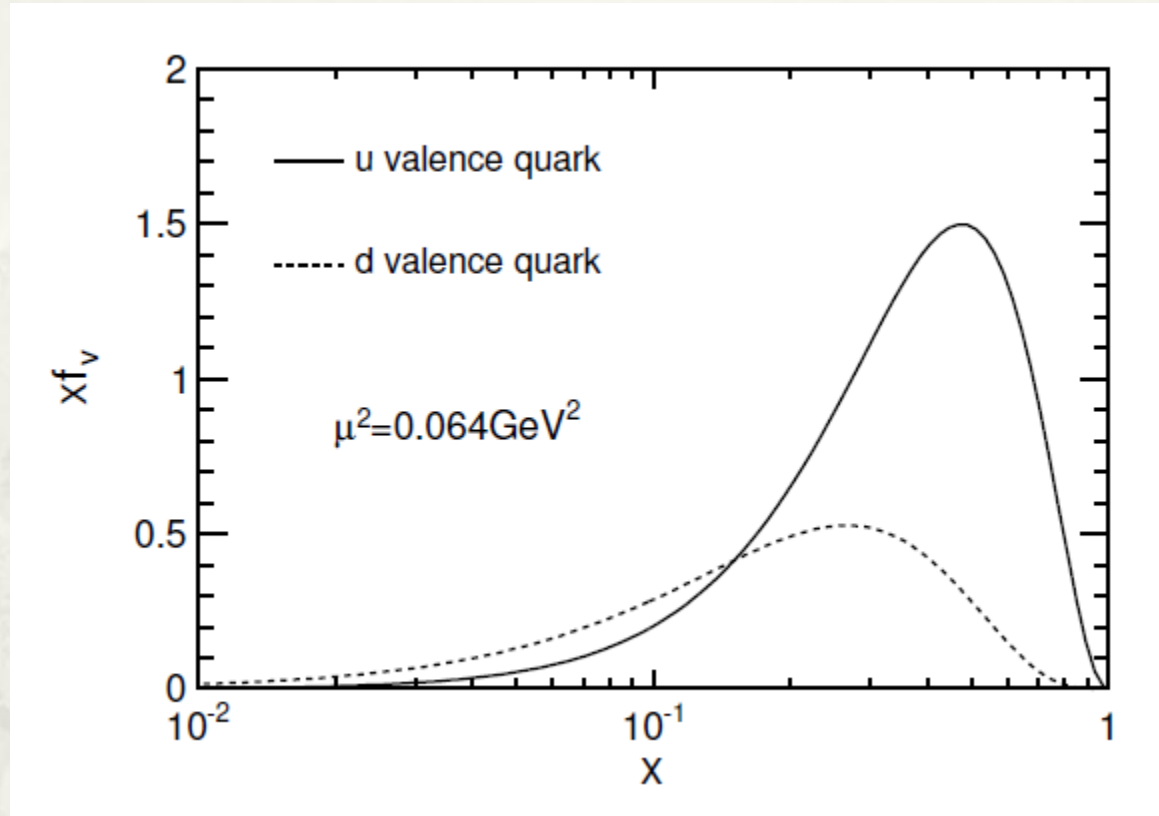
$$\int_0^1 dx x [f_{v_u}(x, \mu^2) + f_{v_d}(x, \mu^2)] = 1.$$

$$\int_0^1 dx f_{v_u}(x, \mu^2) = 2, \quad \int_0^1 dx f_{v_d}(x, \mu^2) = 1.$$

Another free parameter  
in GLR-MQ-ZSR equation.

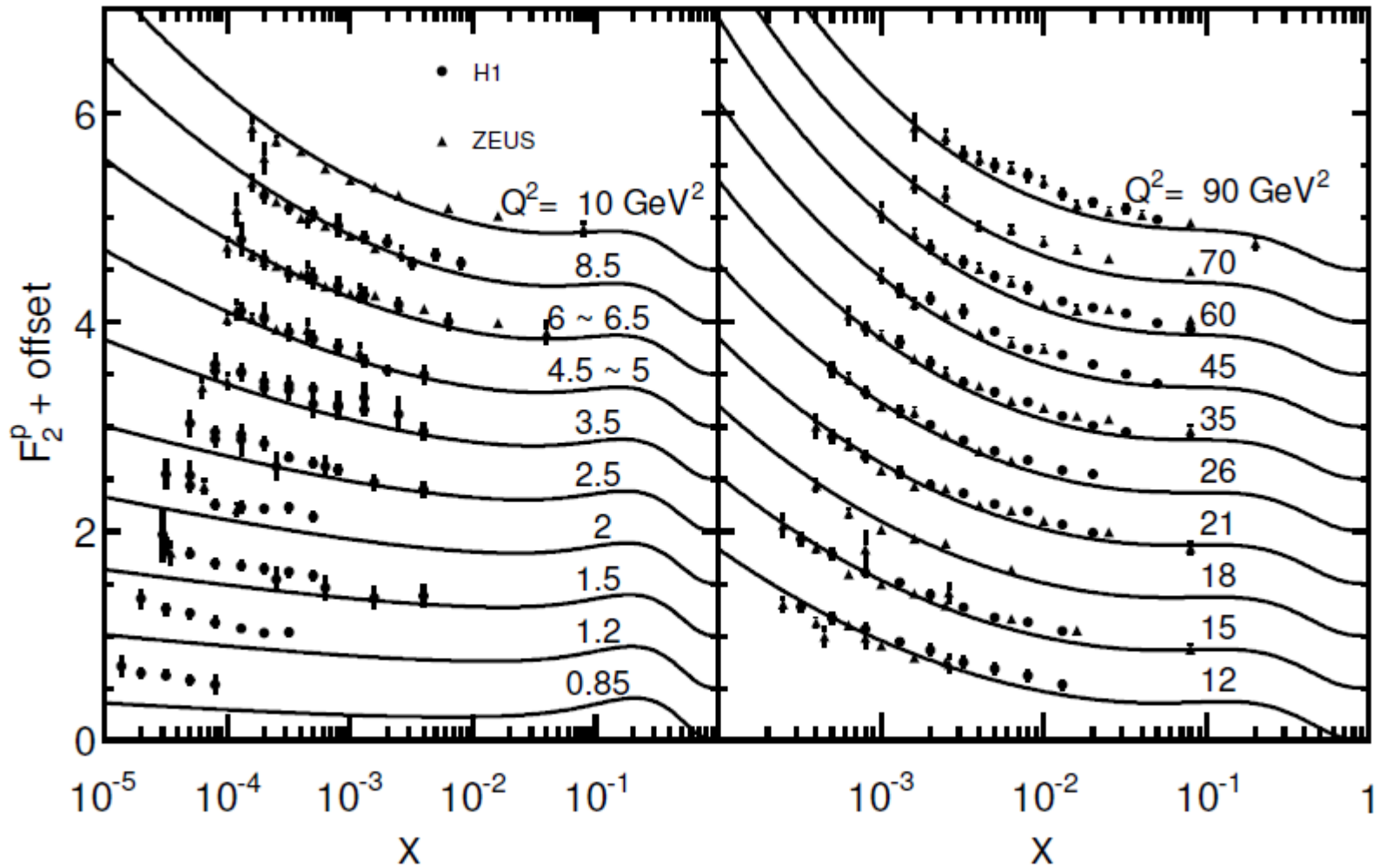
$$R = 4.17 \text{ GeV}^{-1}$$

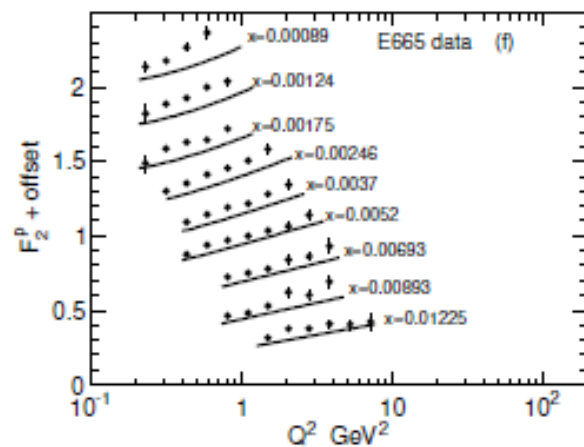
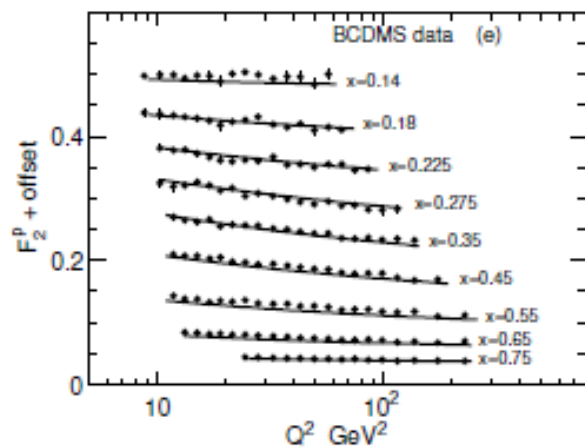
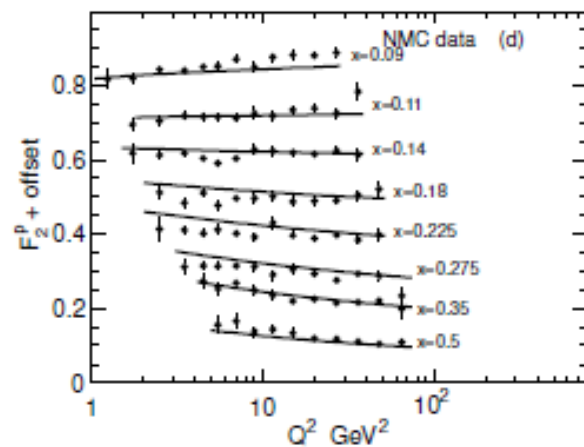
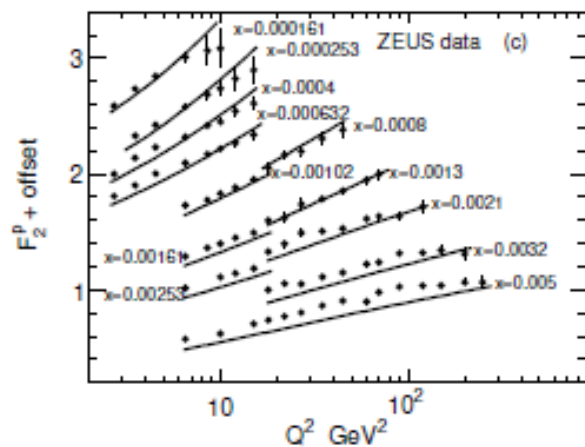
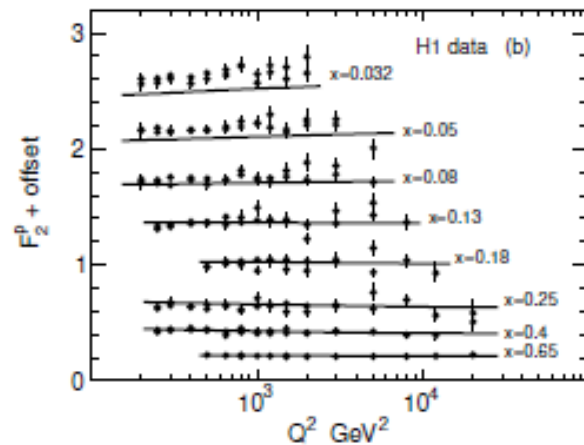
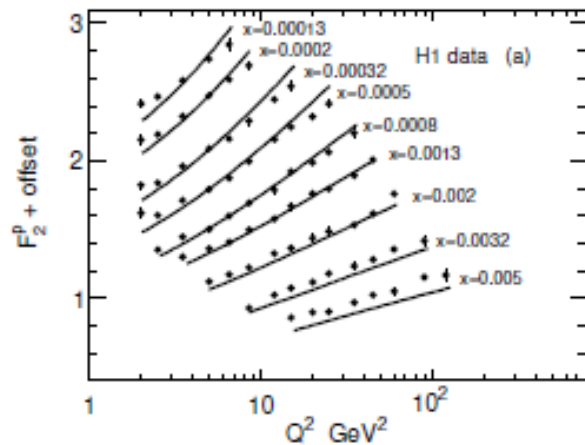
# Parton distributions in the proton



The initial input for GLR-MQ-ZSR evolution.

# Parton distributions in the proton

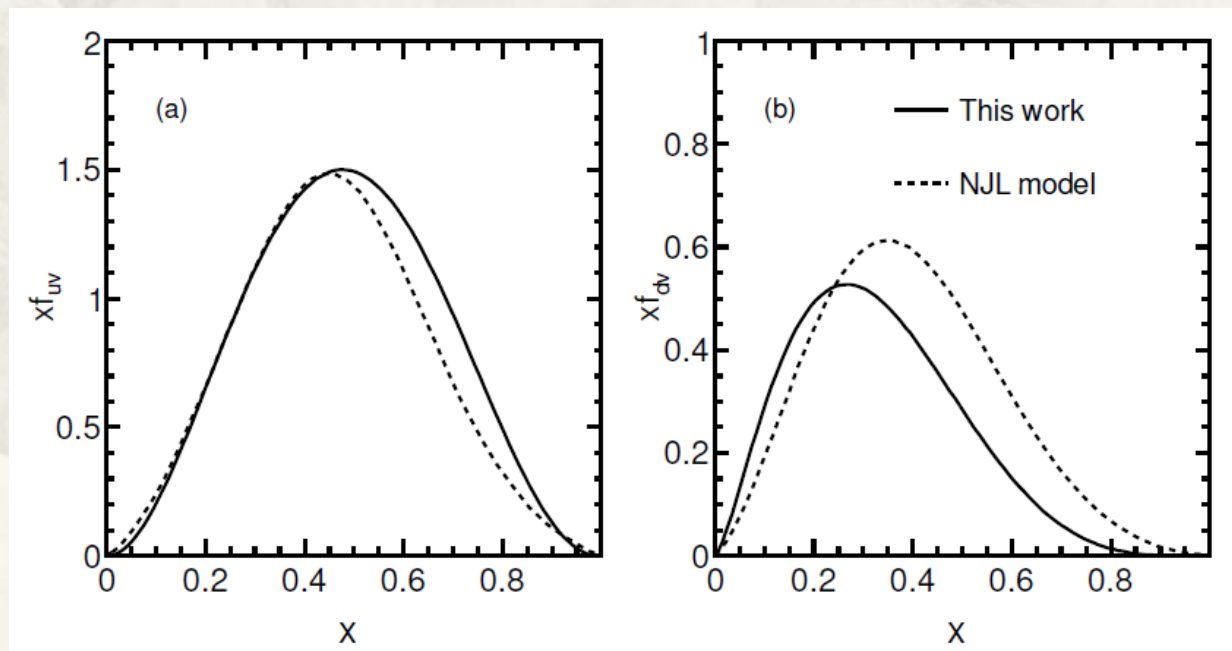






# Parton distributions in the proton

- \* GLR-MQ-ZSR offers a bridge to connect quark models of hadron and various non-perturbative effects on them with measured structure functions at high scale  $Q^2$ .



# Parton distributions in the nuclei

- \* Add 2 free parameters for nuclear media effect. Fermi motion is considered.

Shadowing  
enhancement

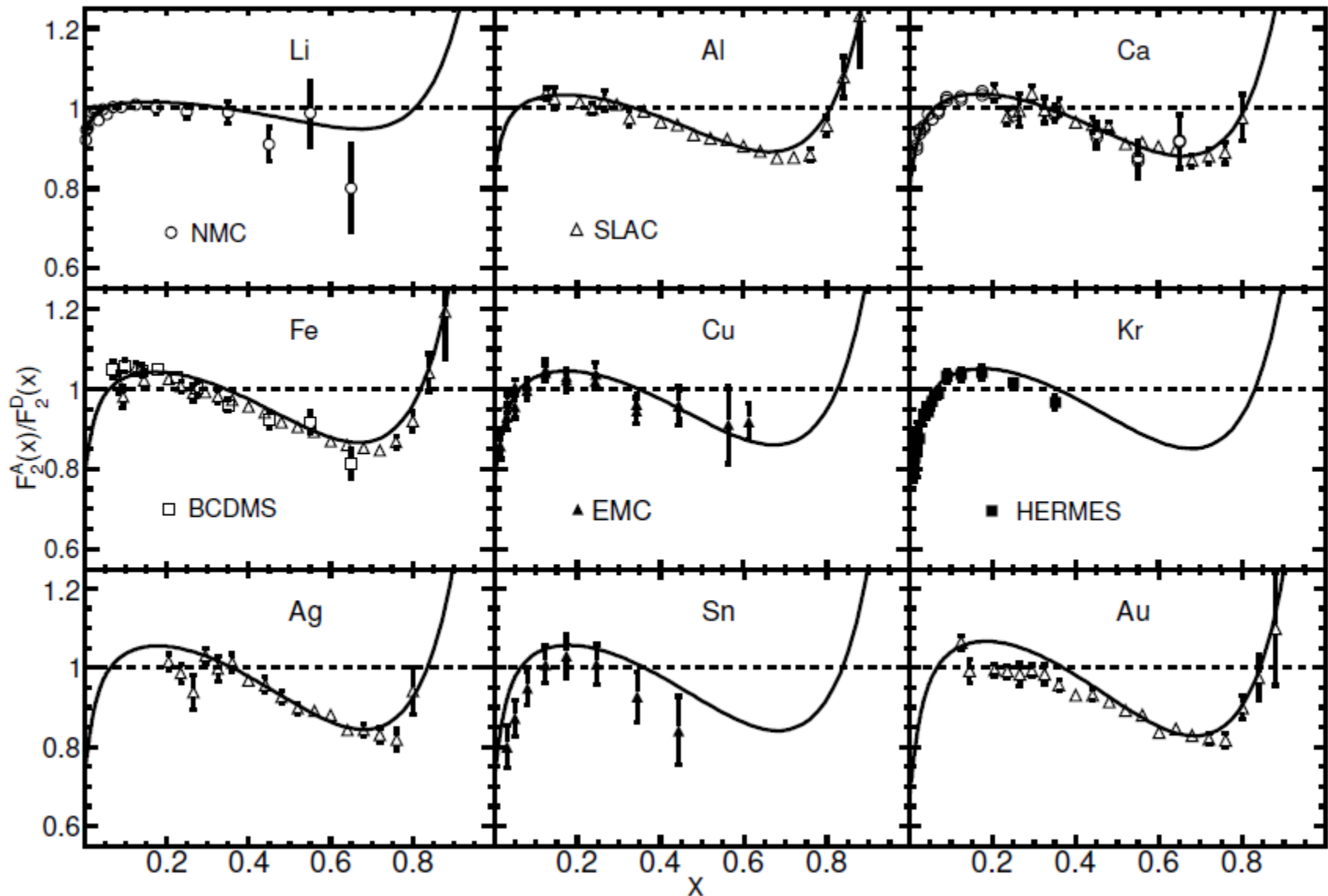
$$A_{eff} = 1 + \beta(A^{1/3} - 1)$$

Nucleon swelling

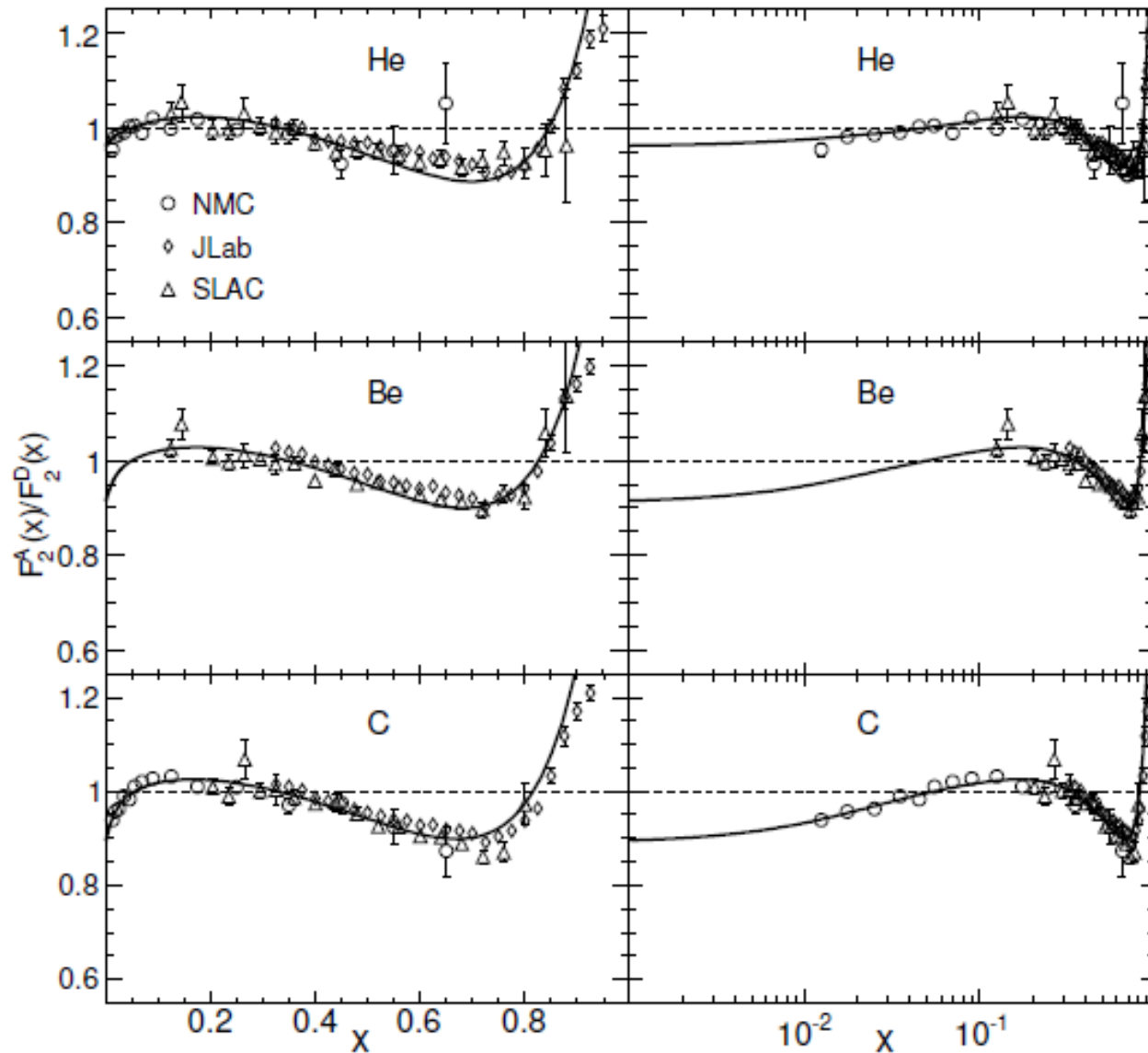
$$R_p/R_A \equiv \frac{1}{1 + \delta_A}$$

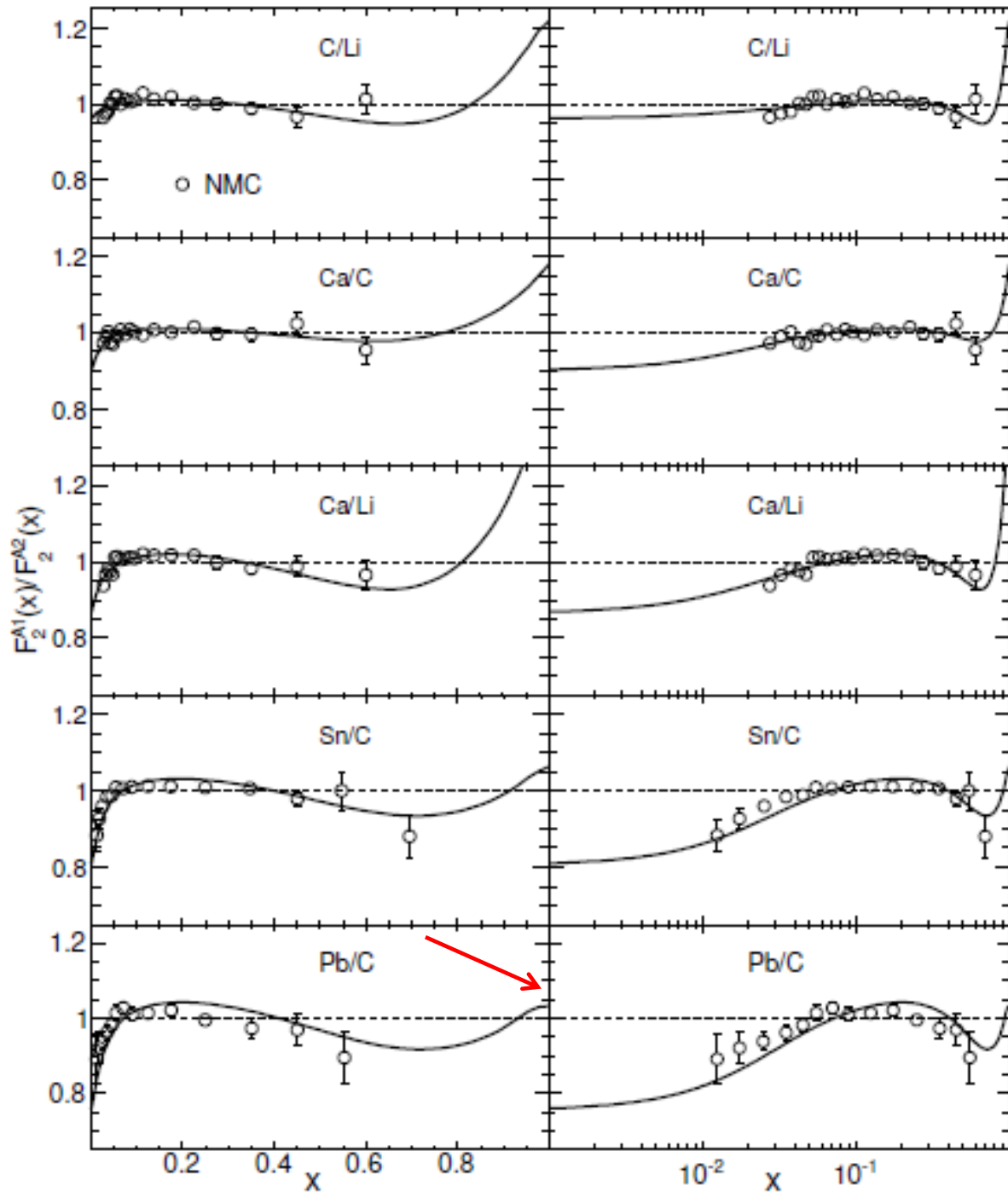
$$\delta_A = [1 - P_s(A)]\delta_0 + P_s(A)\delta_0/4$$

# Parton distributions in the nuclei



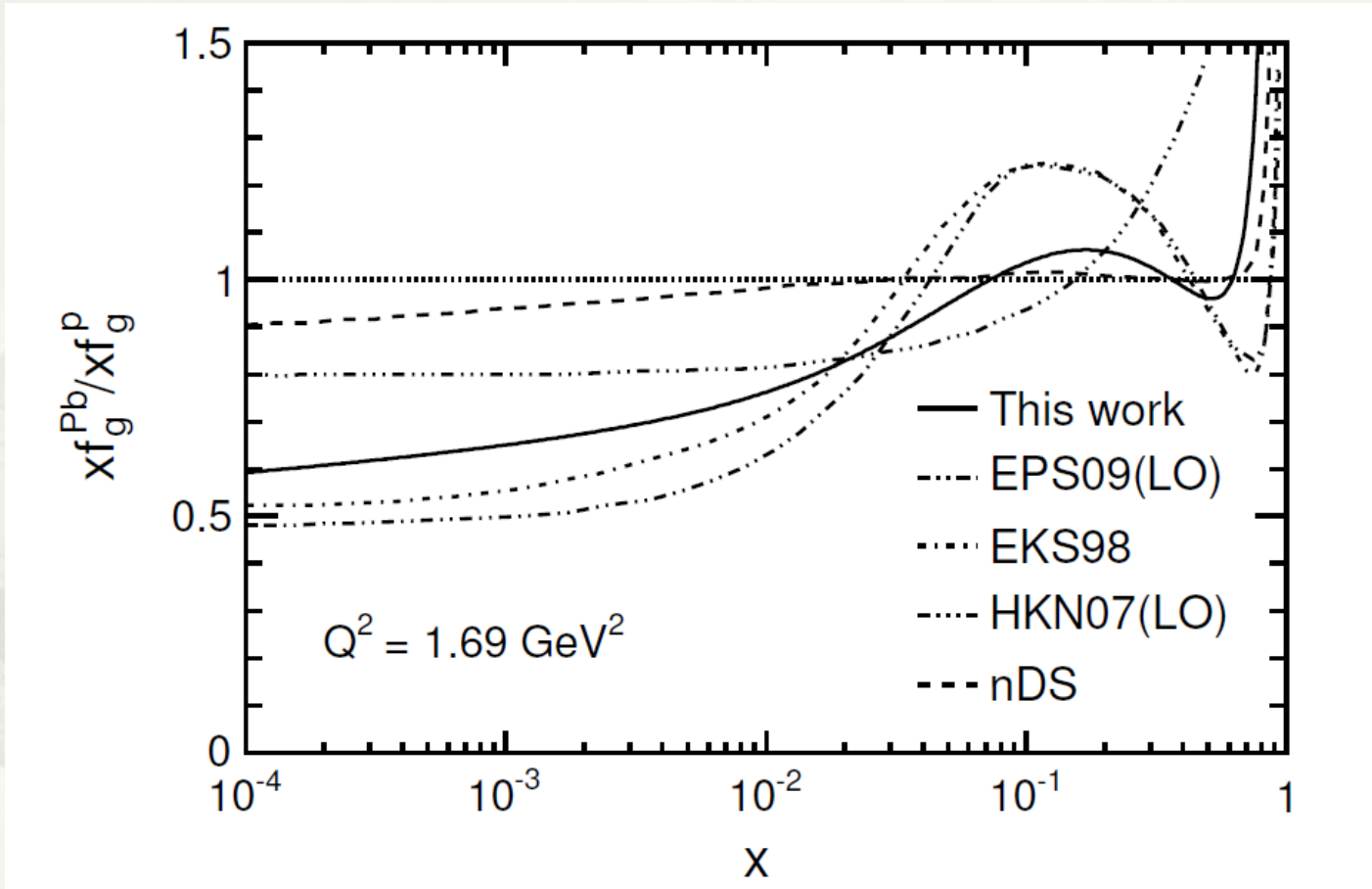
# Parton distributions in the nuclei





Weak short range correlation caused by fermi motion.

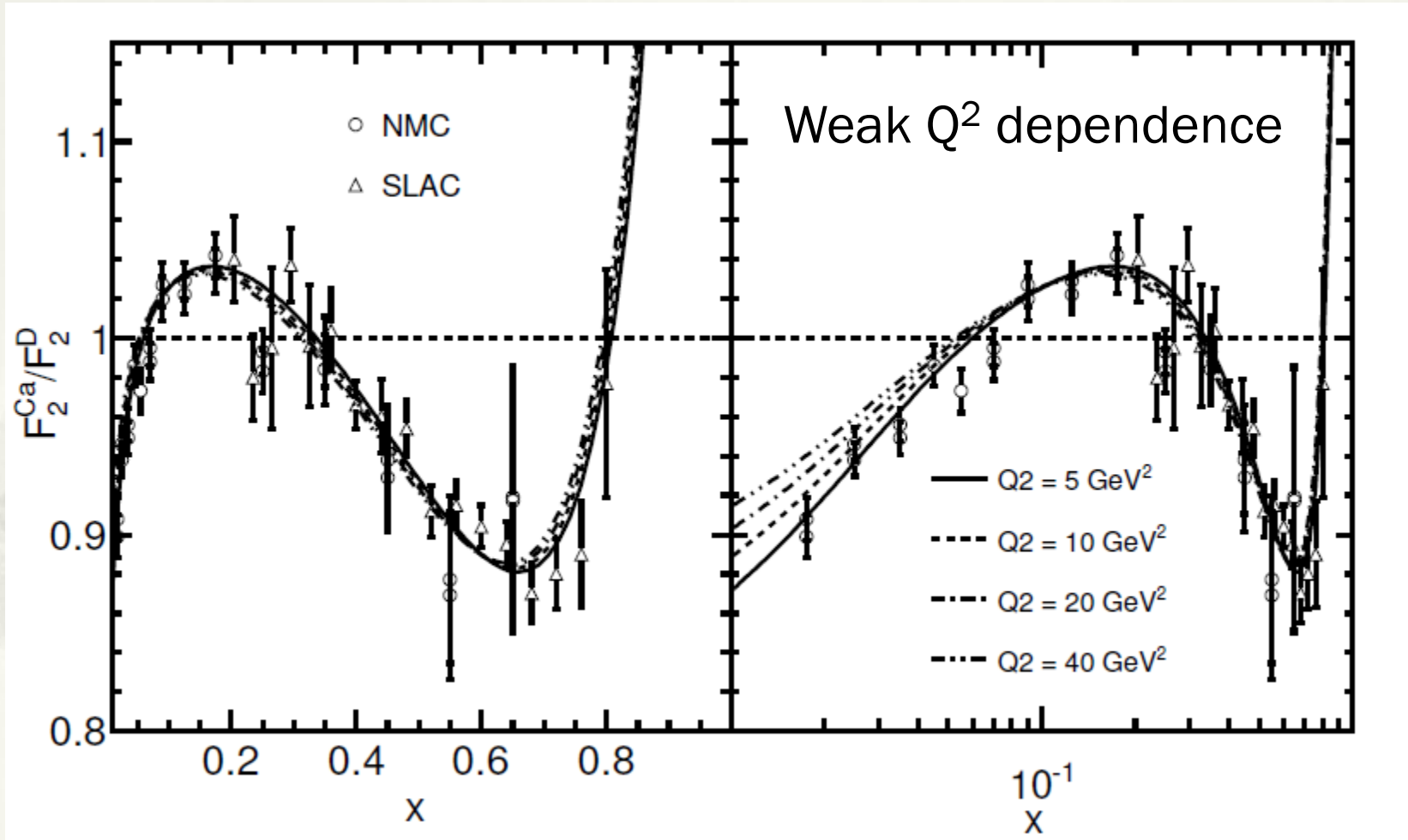
# Parton distributions in the nuclei



The gluon is dynamically determined. The uncertainty is small.



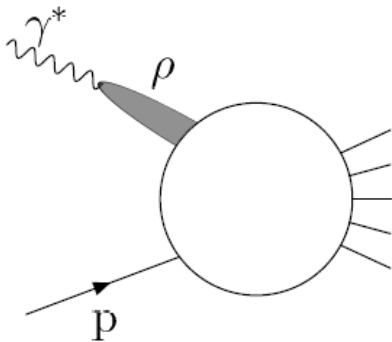
# Parton distributions in the nuclei



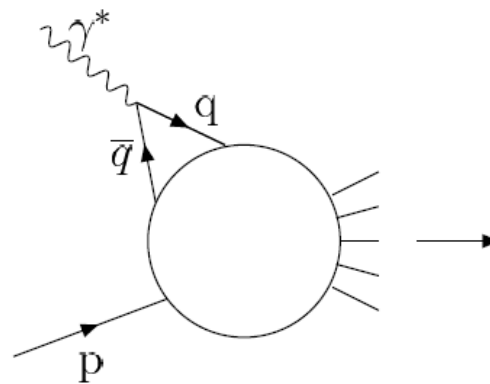
# Quark saturation in low $Q^2$ range

- \* Add 1 parameter for hadron part in low  $Q^2$  range.

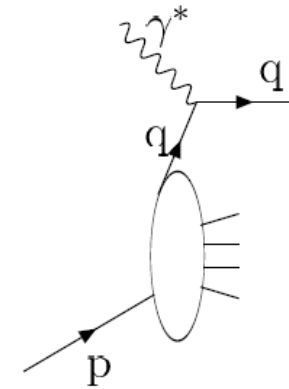
$$F_2^p(x, Q^2) = F_2^{\text{parton}}(x, Q^2) + F_2^{\text{hadron}}(x, Q^2)$$

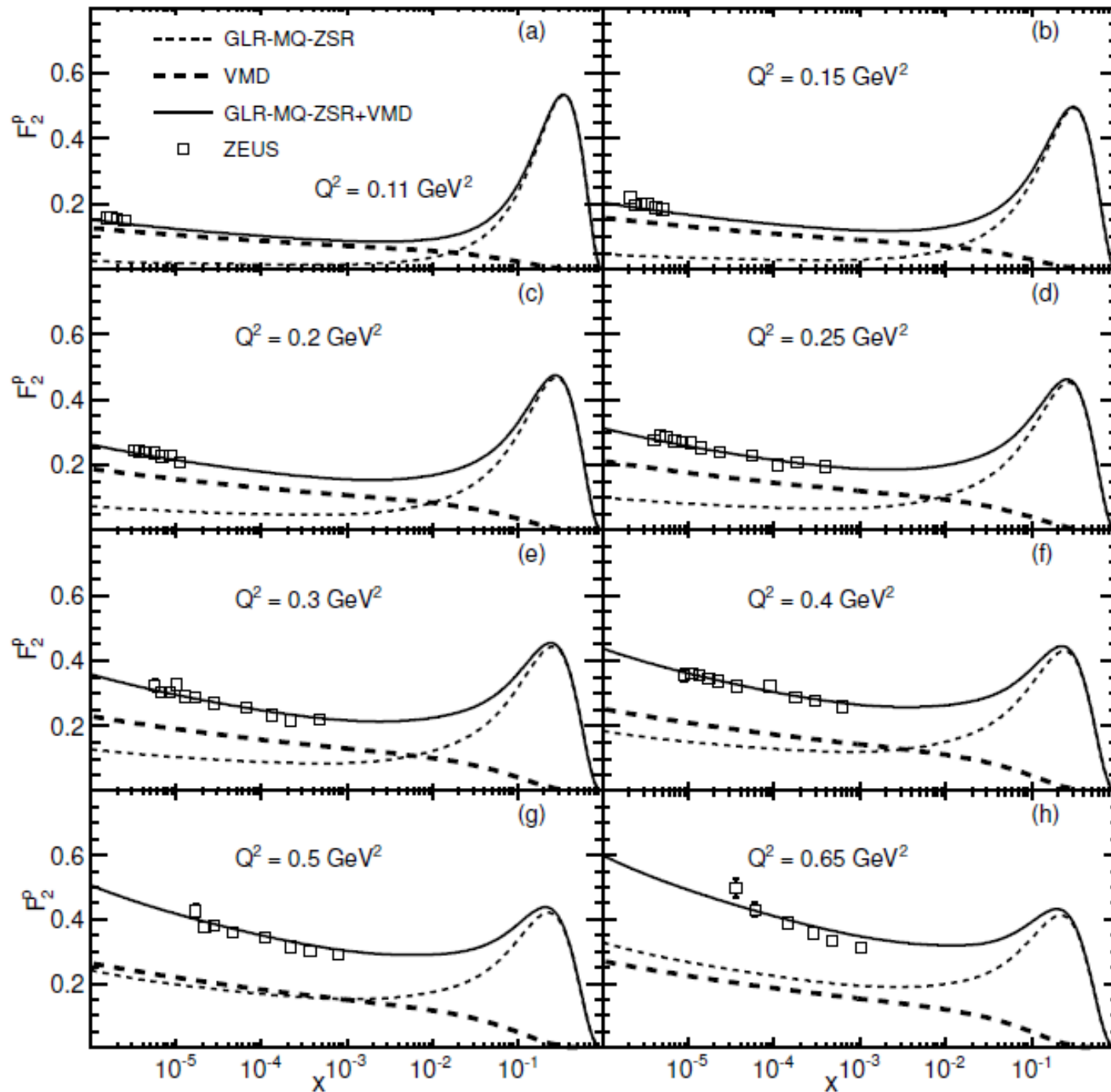


(a)



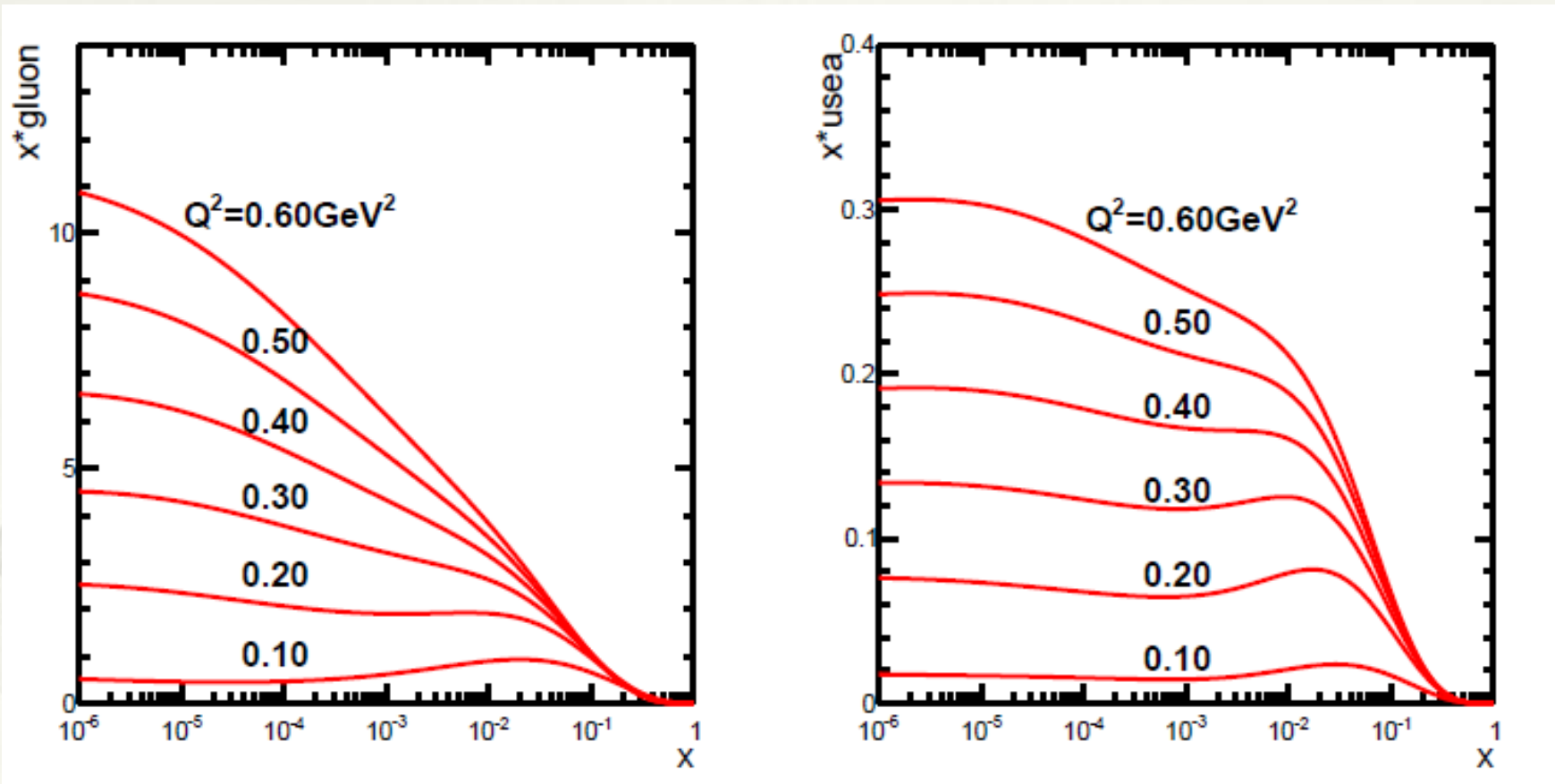
(b)





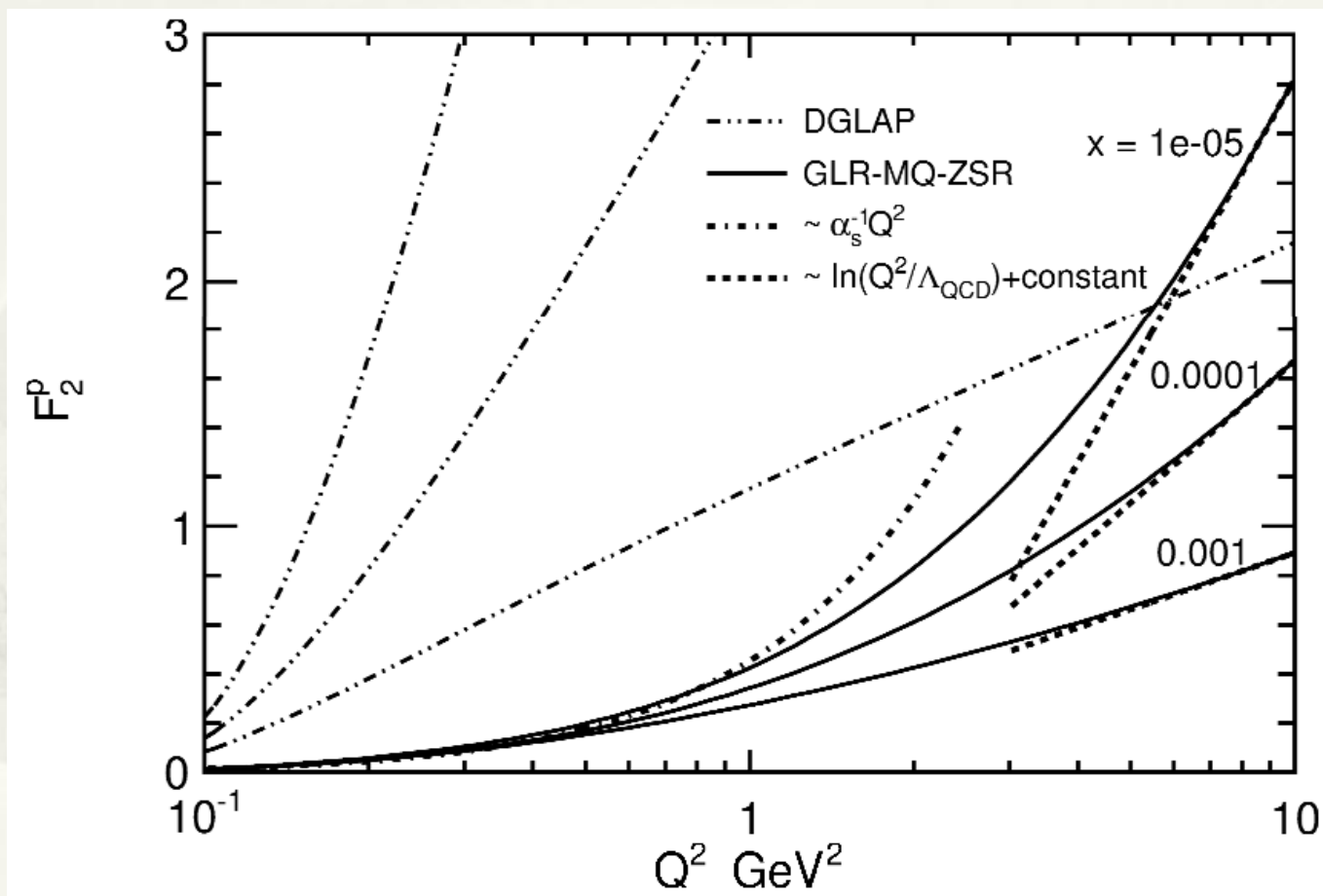
Fit good to experimental data.

# Quark saturation in low $Q^2$ range

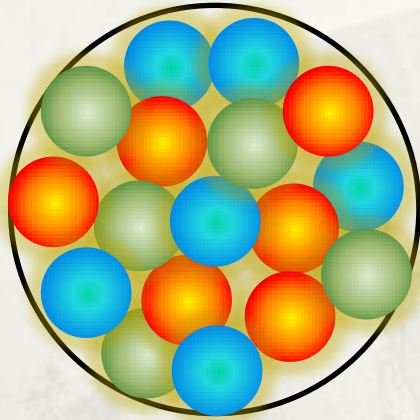


Quark distributions exhibit a flattish form at small  $x$ , while gluon distribution does not.

# Quark saturation in low $Q^2$ range



# Quark saturation in low $Q^2$ range



GLR-MQ-ZSR saturation,  
 $Q_s < 1$  GeV, quark saturation.

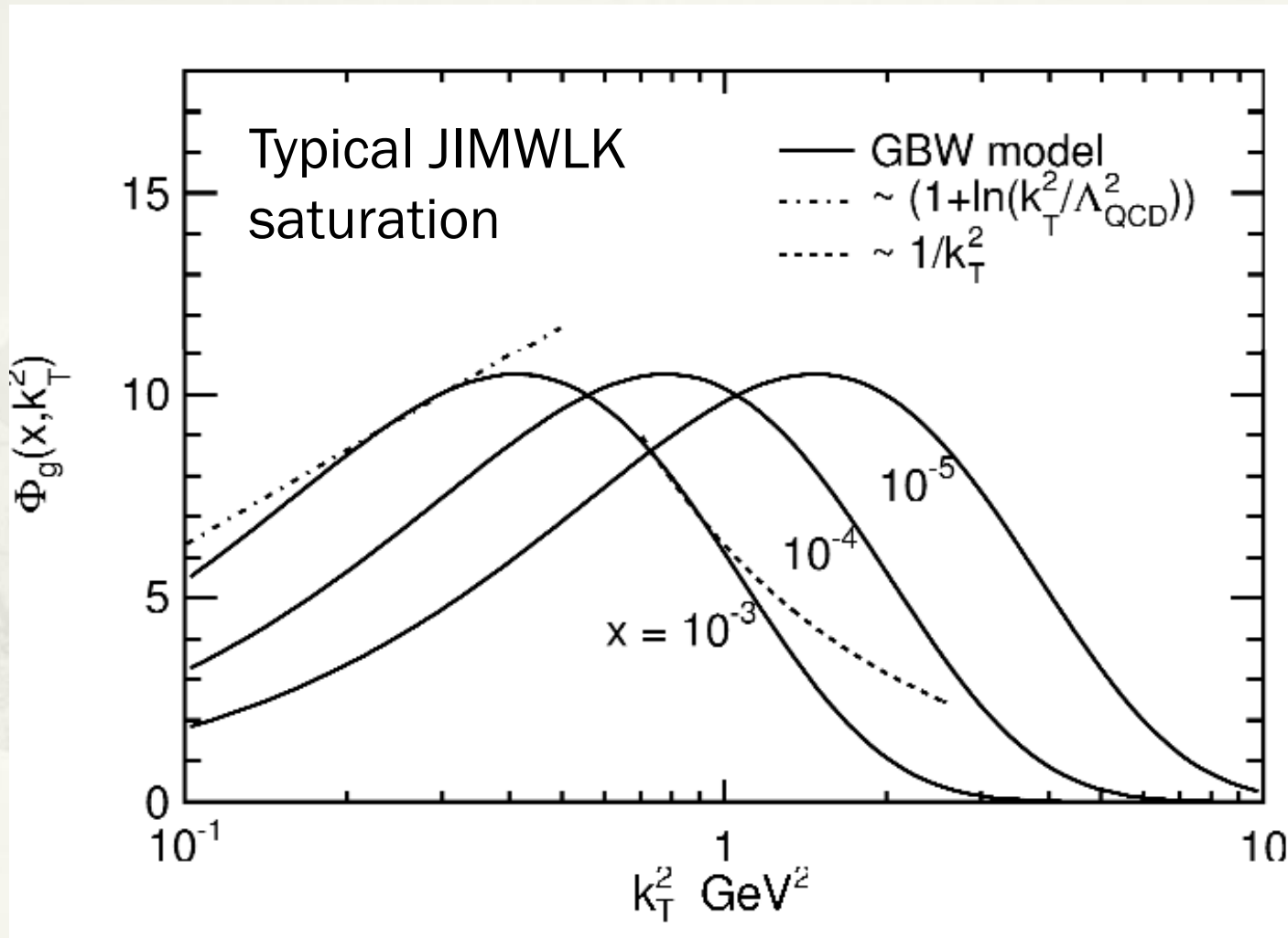


JIMWLK saturation,  
 $Q_s > 1$  GeV, gluon saturation.

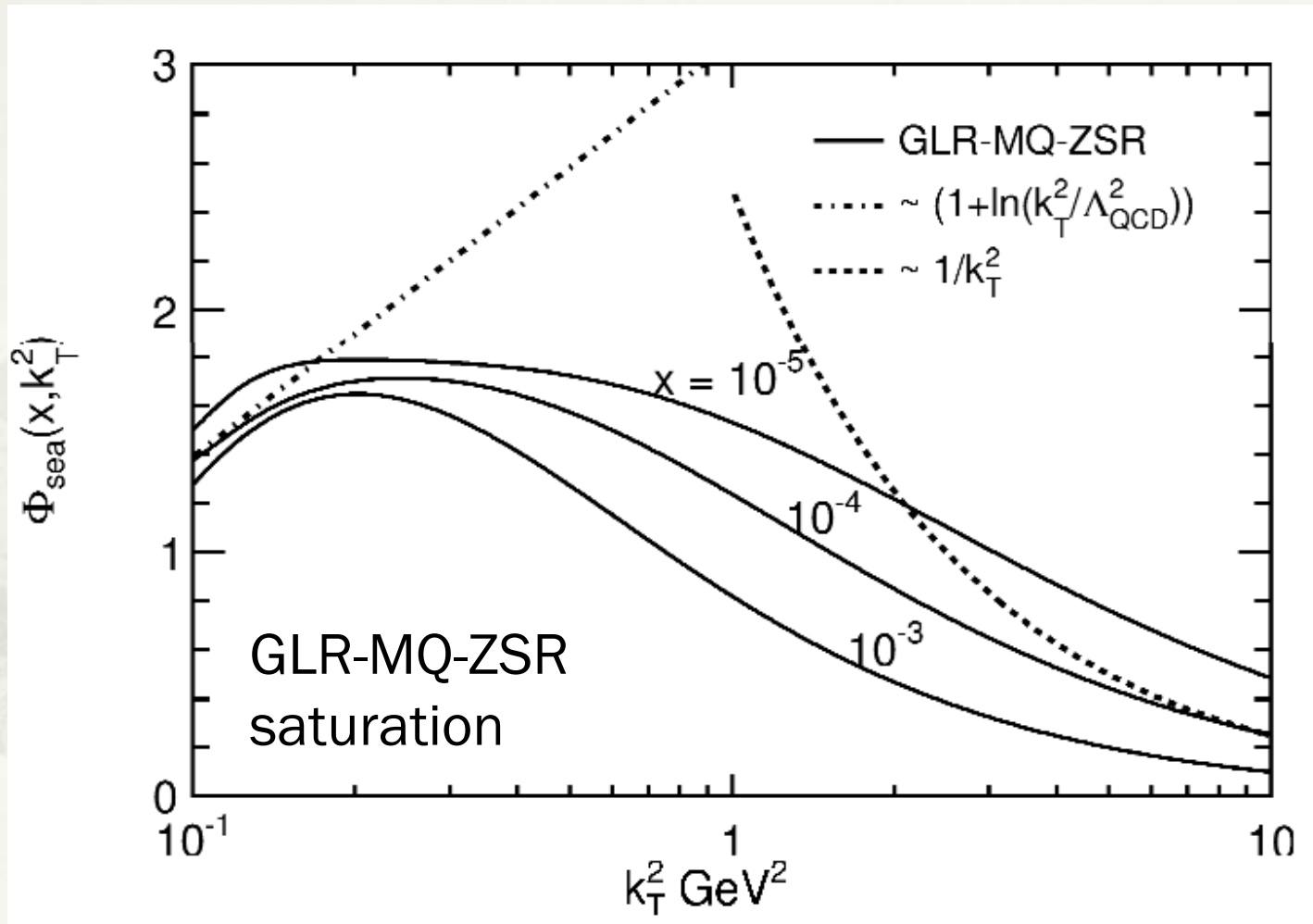
**GLR-MQ-ZRS saturation has similar properties  
as JIMWLK saturation**



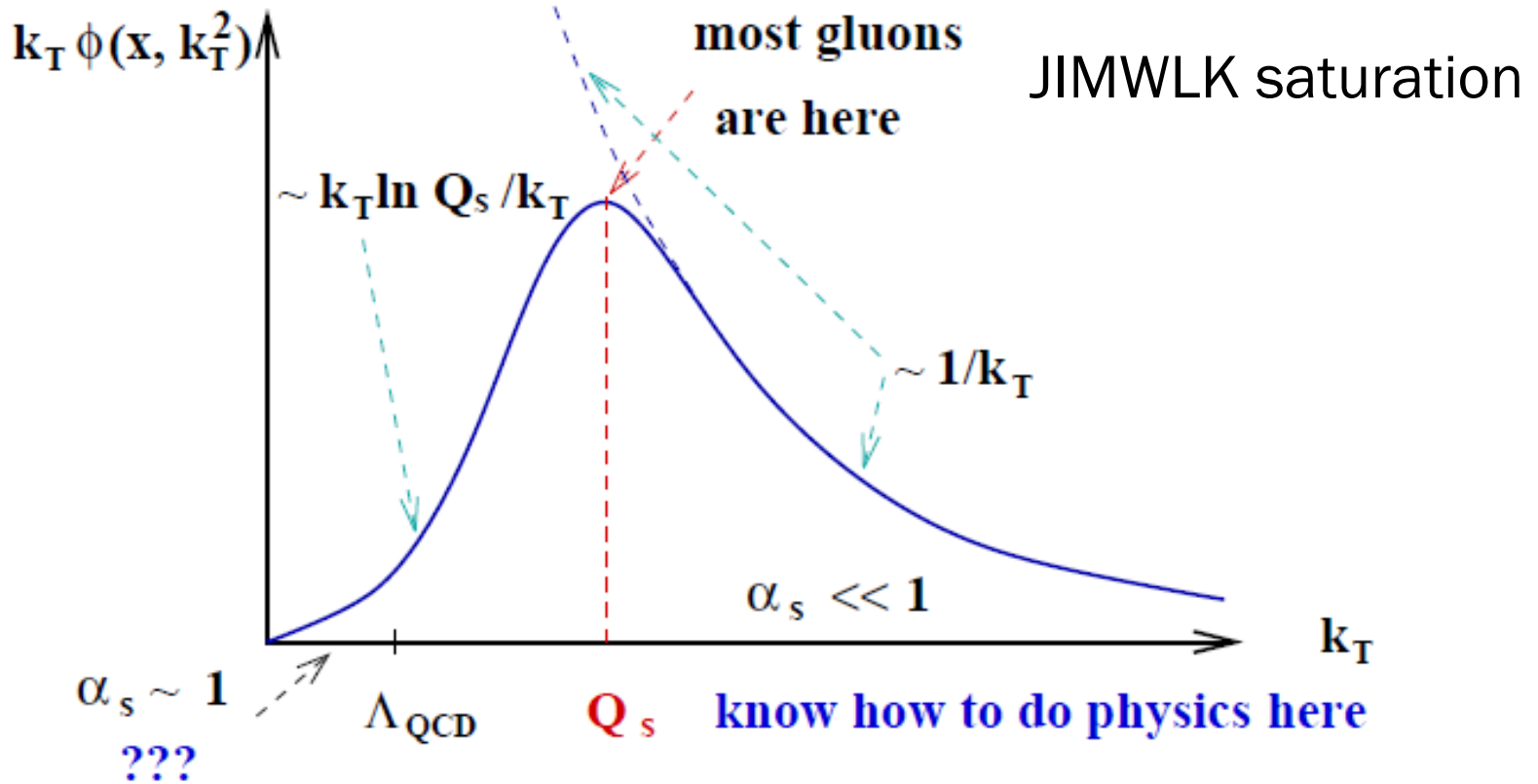
# Quark saturation in low $Q^2$ range



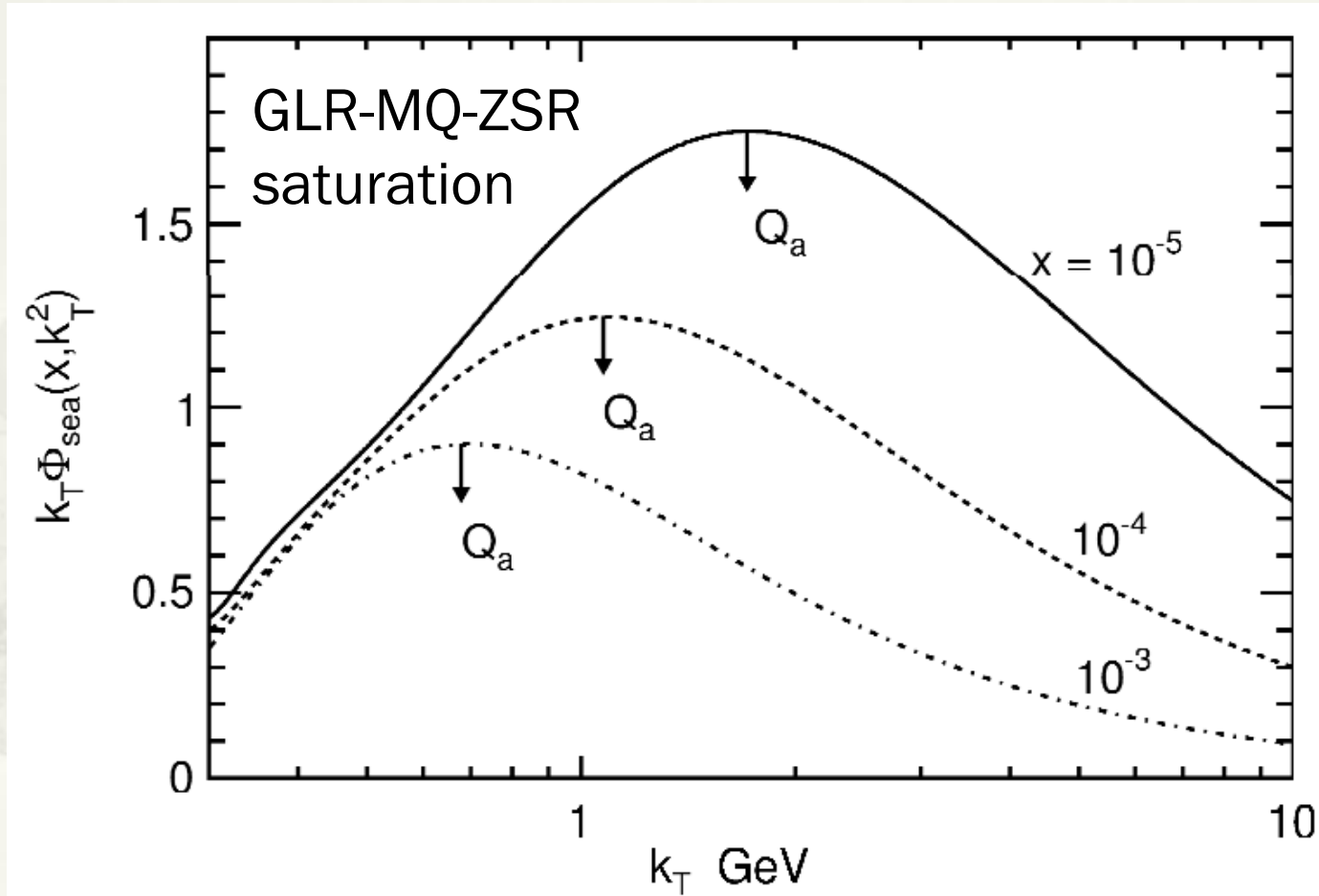
# Quark saturation in low $Q^2$ range



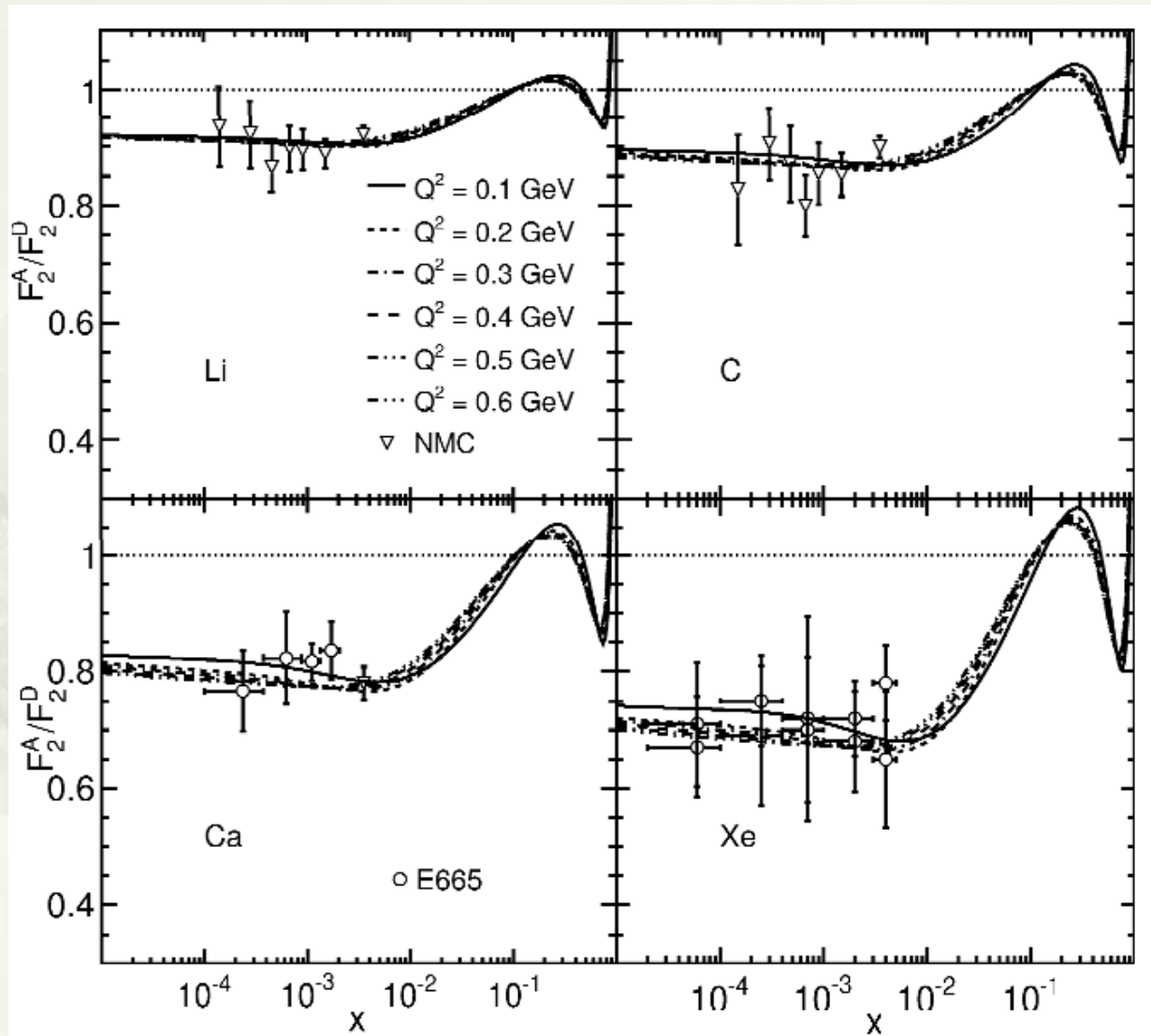
# Gluon saturation



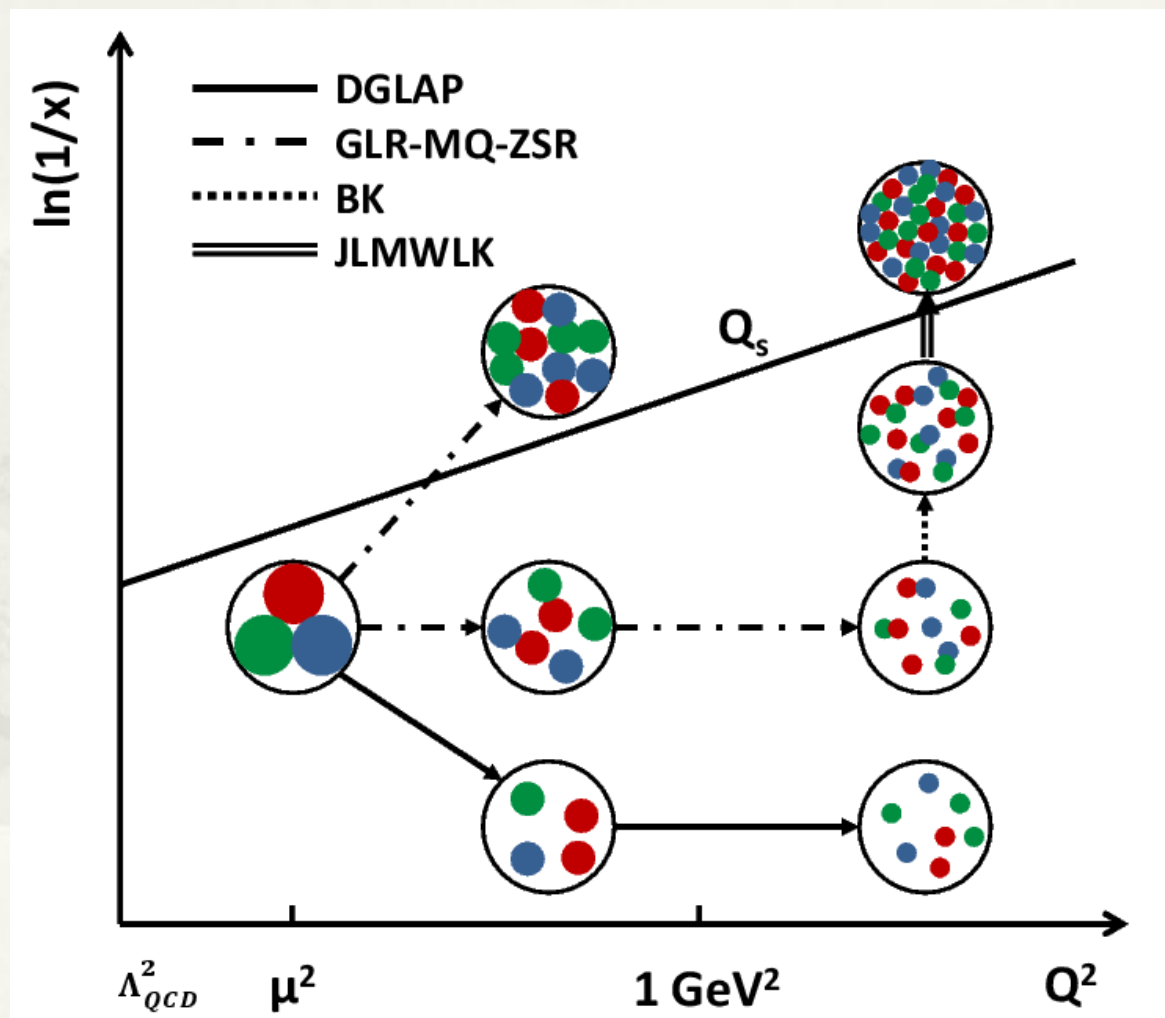
# Quark saturation in low $Q^2$ range



# Quark saturation in low $Q^2$ range



# Evolution of parton distributions in different kinematic region





# Next step

- \* PDF database for proton
- \* nPDF database for nuclei
- \* Spin-dependent PDF
- \* PDF for pion and kaon
- \* SRC
- \* Antishadowing
- \* TMD.....

# The end

Thank you!

