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JLab, 24 May 2007  
*on behalf of the COMPASS Collaboration*

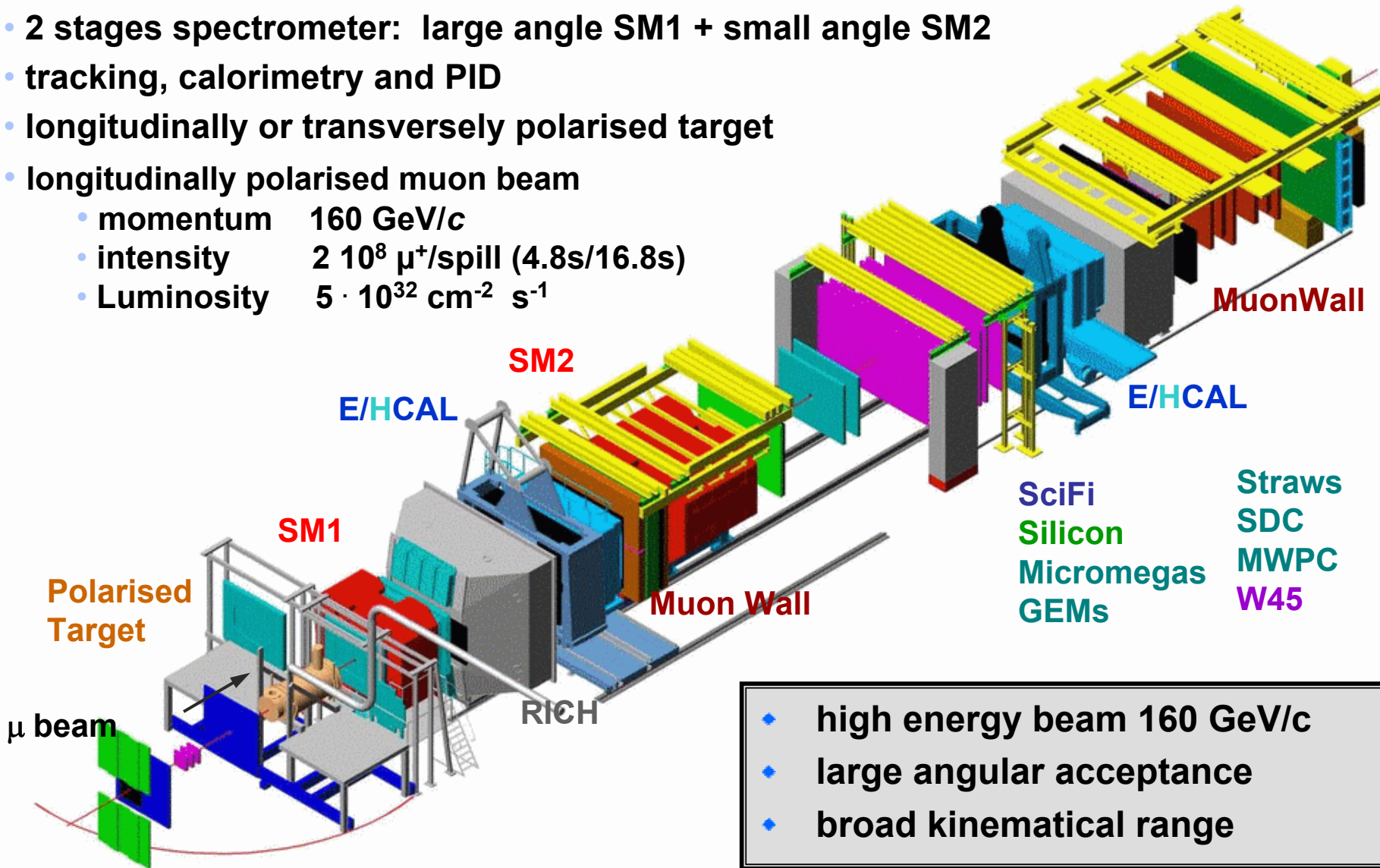
## Outline

- ◆ The COMPASS experiment
- ◆ Results on:
  - ◆ Collins/Sivers asymmetries  
positive and negative hadrons  
 $\pi^\pm, K^\pm$
  - ◆ two identified hadron asymmetries  
 $\pi\pi, \pi K, KK$
  - ◆ other Transv. Mom. Dependent (TMD) asymmetries

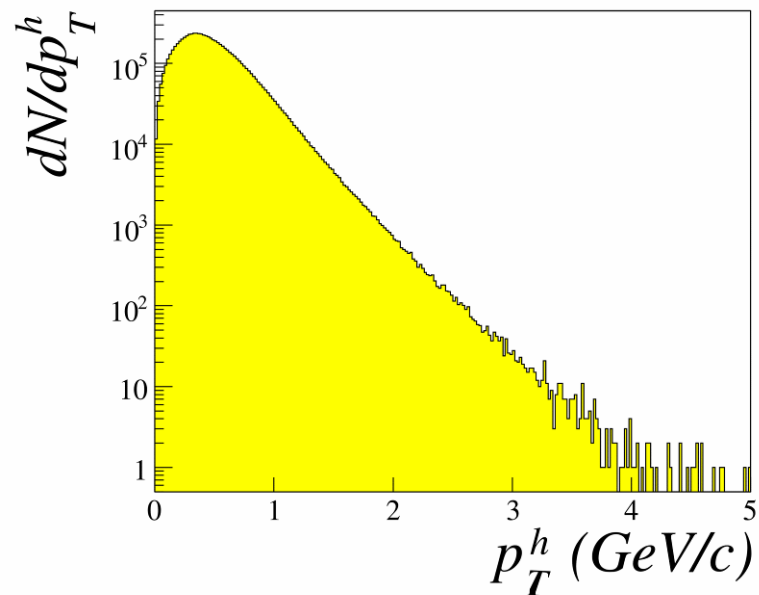
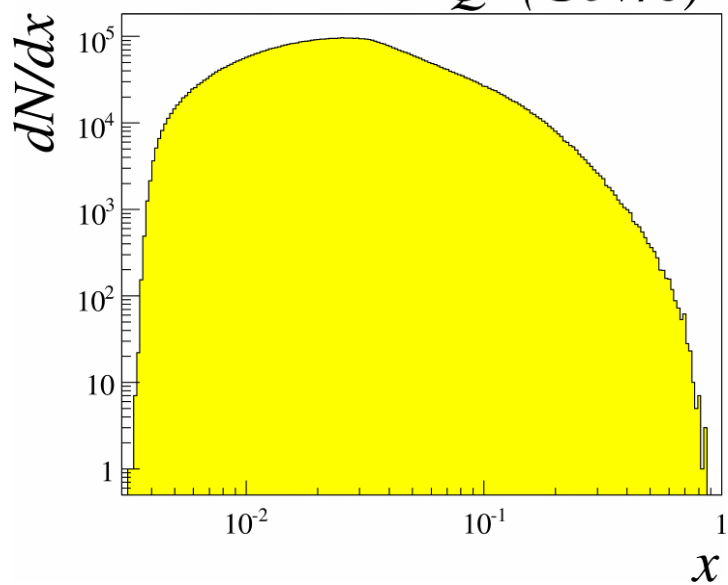
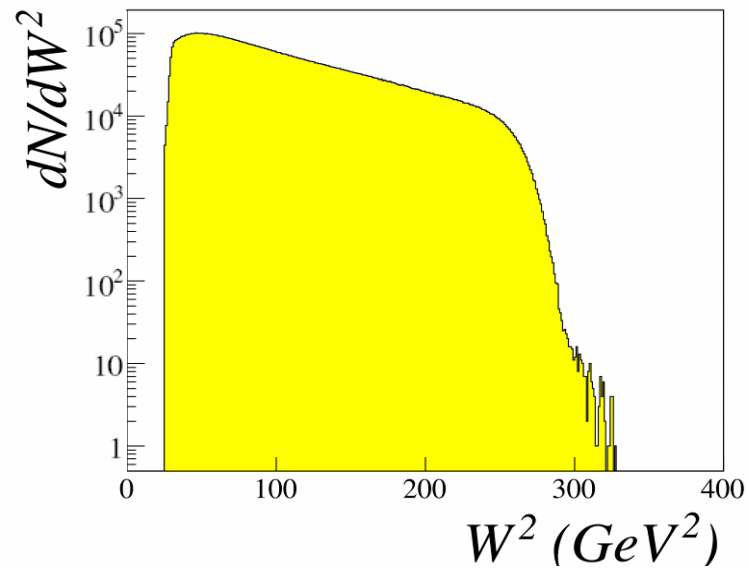
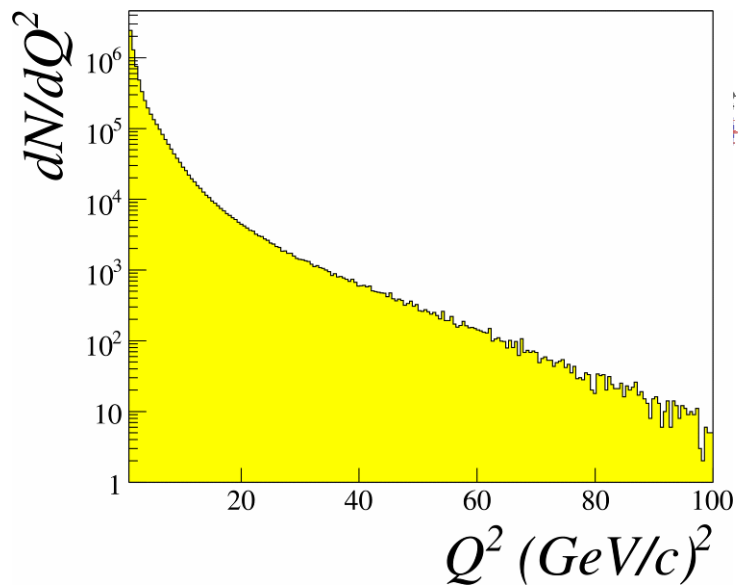
# COMPASS spectrometer



- 2 stages spectrometer: large angle SM1 + small angle SM2
- tracking, calorimetry and PID
- longitudinally or transversely polarised target
- longitudinally polarised muon beam
  - momentum 160 GeV/c
  - intensity  $2 \cdot 10^8 \mu^+/\text{spill}$  (4.8s/16.8s)
  - Luminosity  $5 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

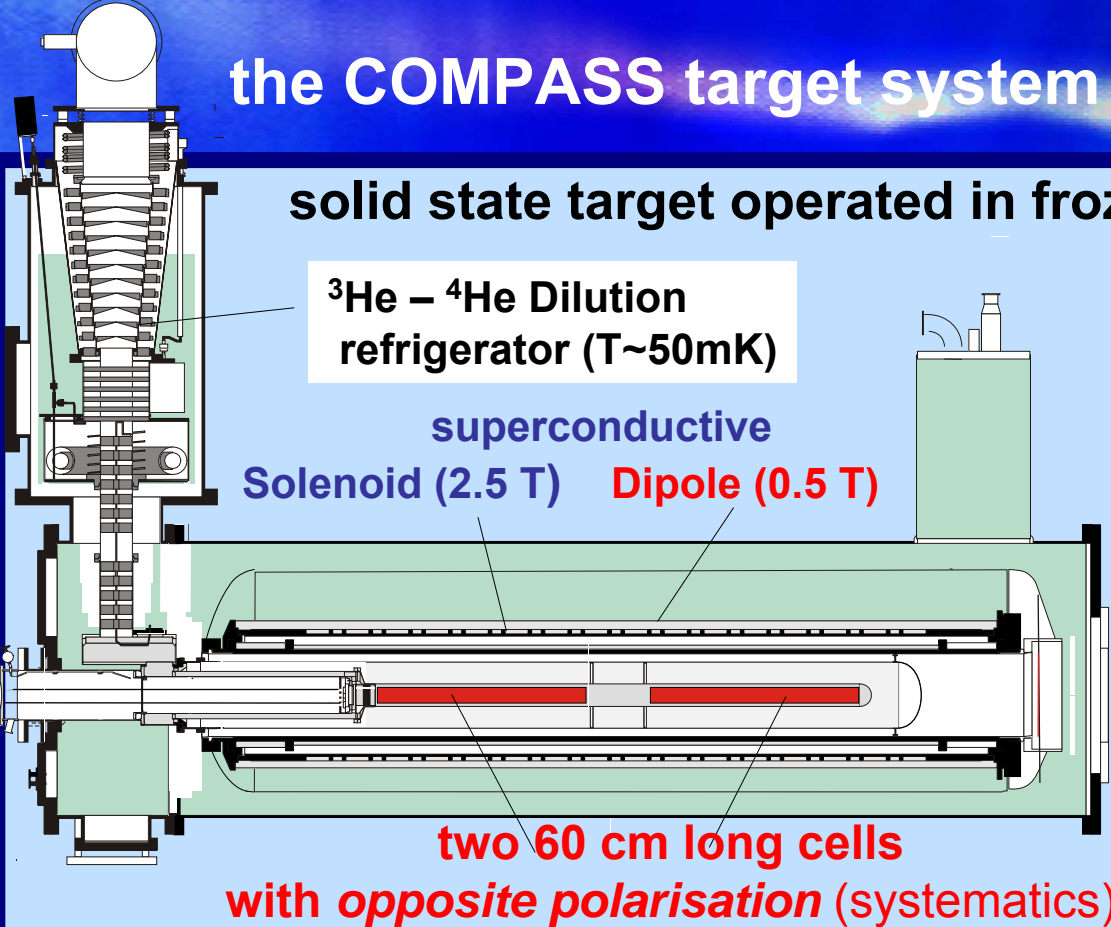


# SIDIS kinematics

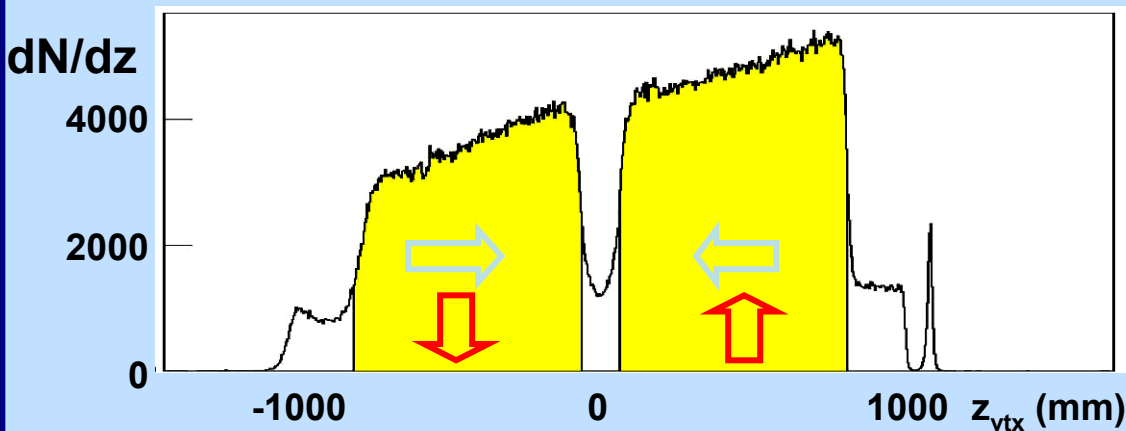


# the COMPASS target system

solid state target operated in frozen spin mode



2002-2004:  $^6\text{LiD}$   
 dilution factor  $f = 0.38$   
 polarization  $P_T = 50\%$   
 ~20% of the time transversely polarised



during data taking with transverse polarization

- dipole field always  $\uparrow$
- polarization reversal in the 2 cells after  $\sim 5$  days

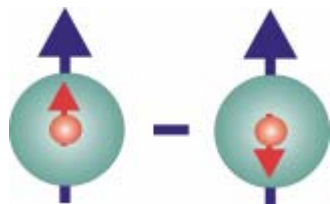
# transversity DF

$$\Delta_T q(x) = q^{\uparrow\uparrow}(x) - q^{\uparrow\downarrow}(x)$$

$h_1^q(x)$ ,

$\delta q(x)$ ,

$\delta_T q(x)$



$q = u_v, d_v, q_{sea}$

quark with **spin** parallel to the nucleon spin in a transversely polarised nucleon

Properties:

- probes the relativistic nature of **quark dynamics**
- **no contribution from the gluons** → simple  $Q^2$  evolution
- Positivity: Soffer bound.....  $2|\Delta_T q| \leq q + \Delta q$  *Soffer, PRL 74 (1995)*
- first moments: tensor charge.....  $\Delta_T q \equiv \int dx \Delta_T q(x)$
- sum rule for transverse spin  
in Parton Model framework.....  $\frac{1}{2} = \frac{1}{2} \sum \Delta_T q + L_q + L_g$   
*Bakker, Leader, Trueman, PRD 70 (04)*
- it is related to **GPD's**
- is **chiral-odd**: decouples from inclusive DIS



the Transversity DF is chiral-odd:

→ survives only by the product with another chiral-odd function

can be measured in SIDIS on a transversely polarised target via  
“quark polarimetry”

$\ell N^\uparrow \rightarrow \ell' h X$	Collins Asymmetry (Collins FF)
$\ell N^\uparrow \rightarrow \ell' hh X$	Two hadrons asymmetry (Interference FF)
$\ell N^\uparrow \rightarrow \ell' \Lambda X$	$\Lambda$ polarization (FF of $q^\uparrow \rightarrow \Lambda$ )

# single hadron asymmetries

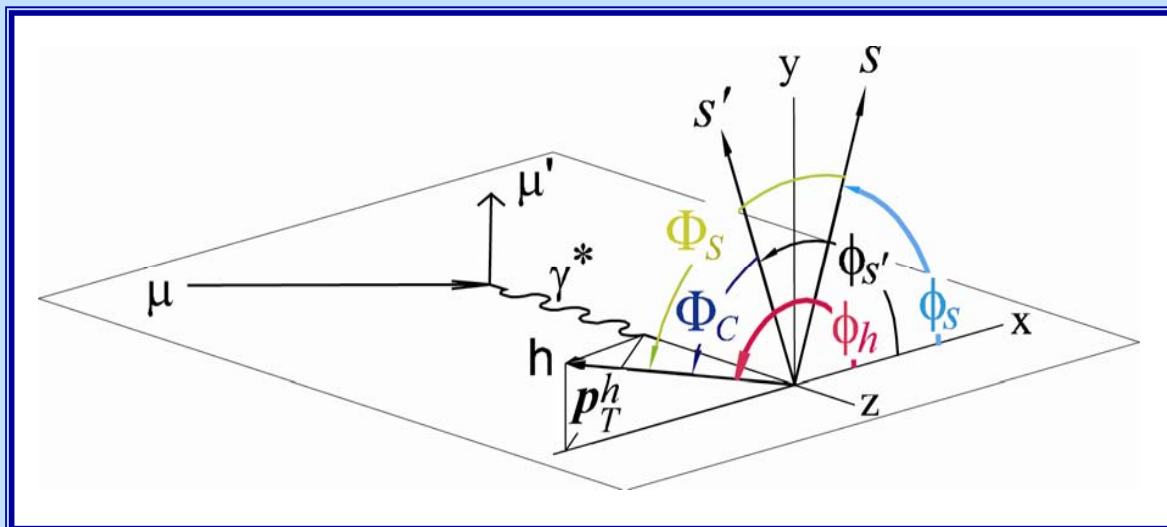
Collins and Sivers terms in SIDIS cross sections depend on different combination of angles:

Collins angle

$$\begin{aligned}\Phi_C &= \phi_h - \phi_{s'} \\ &= \phi_h + \phi_s - \pi\end{aligned}$$

Sivers angle

$$\Phi_S = \phi_h - \phi_s$$



$\phi_h$  azimuthal angle of the hadron

$\phi_s$  azimuthal angle of the transverse spin of the initial quark

$\phi_{s'}$  azimuthal angle of the transverse spin of the fragmenting quark

$$\phi_{s'} = \pi - \phi_s \text{ (spin flip)}$$

$\Phi_C$  and  $\Phi_S$  are independent angles  $\rightarrow$  independent extraction of the asymmetries

## Collins effect:

a quark moving horizontally and polarized upward (downward) prefers to emit the leading meson to the left (right) side of the jet (quark direction).

i.e.

the fragmentation function of a transversely polarized quark has a spin dependent part

$$D_q^h(z, \vec{p}_T^h) = D_q^h(z, p_T^h) + \Delta_T^0 D_q^h(z, p_T^h) \times \sin(\varphi_h - \varphi_{s'})$$

And the resulting measured asymmetry:

$$\mathbf{N}_h^\pm(\Phi_C) = \mathbf{N}_h^0 \cdot \left\{ 1 \pm \mathbf{A}_C^h \cdot \sin\Phi_C \right\} \quad \Phi_C = \phi_h + \phi_s - \pi$$

$$\mathbf{A}_{\text{Coll}} = \frac{\mathbf{A}_C^h}{\mathbf{f} \cdot \mathbf{P}_T \cdot \mathbf{D}} = \frac{\sum_q e_q^2 \cdot \Delta_T^0 \cdot \Delta_T^0 D_q^h}{\sum_q e_q^2 \cdot \mathbf{q} \cdot \mathbf{D}_q^h}$$



DIS cuts:

- ◆  $Q^2 > 1$
- ◆  $0.1 < y < 0.9$
- ◆  $W > 5 \text{ GeV}/c$

All hadron selection:

- ◆  $z > 0.20$
- ◆  $p_t > 0.1$

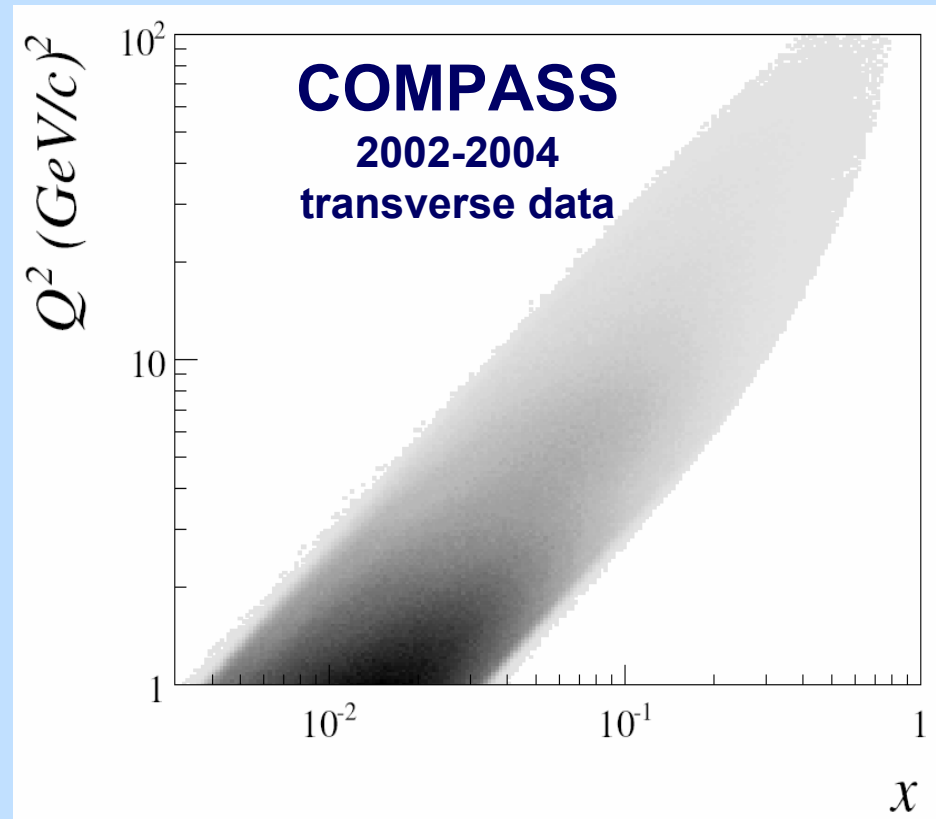
Plus for leading hadron:

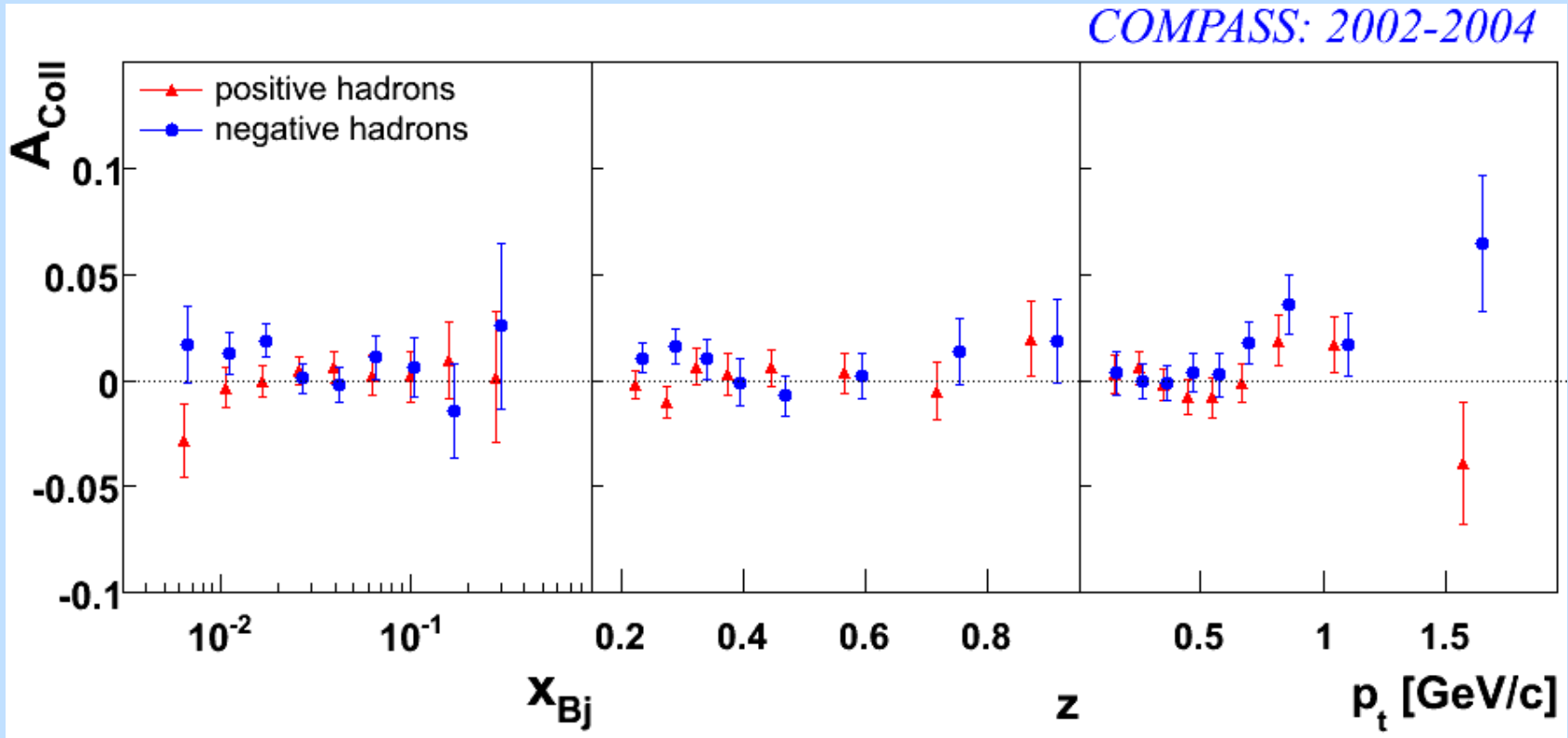
- ◆  $z_1 > 0.25$
- ◆ No signals in the CALOs from neutral particles with  $z > z_1$

Statistics 2002 - 2004:

$8.5 \cdot 10^6$  positive hadrons

$7.0 \cdot 10^6$  negative hadrons

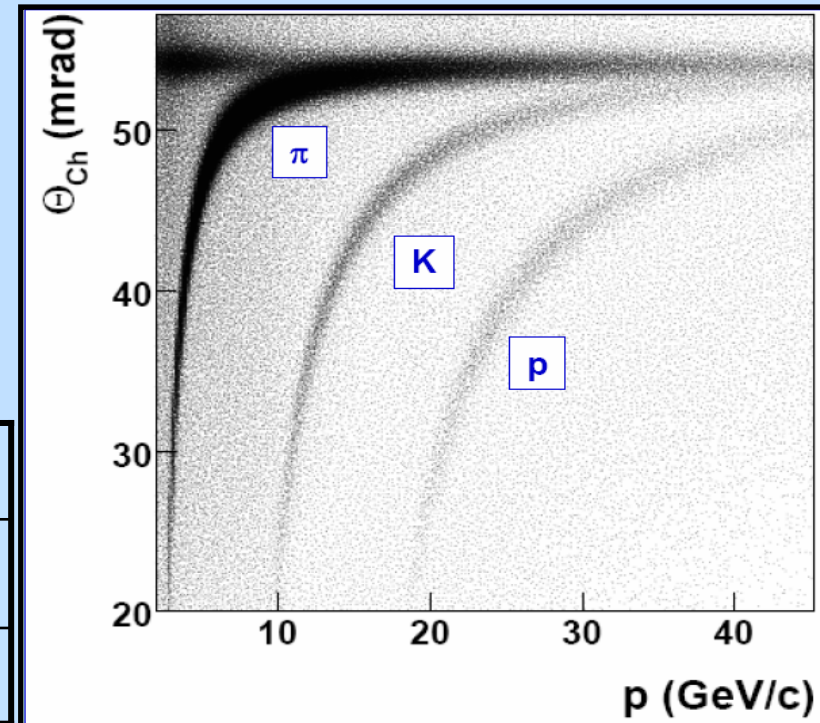




- only statistical errors shown ( $\sim 1\%$ ),
- systematic errors considerably smaller
- **small asymmetries compatible with 0 for both + and - hadrons**

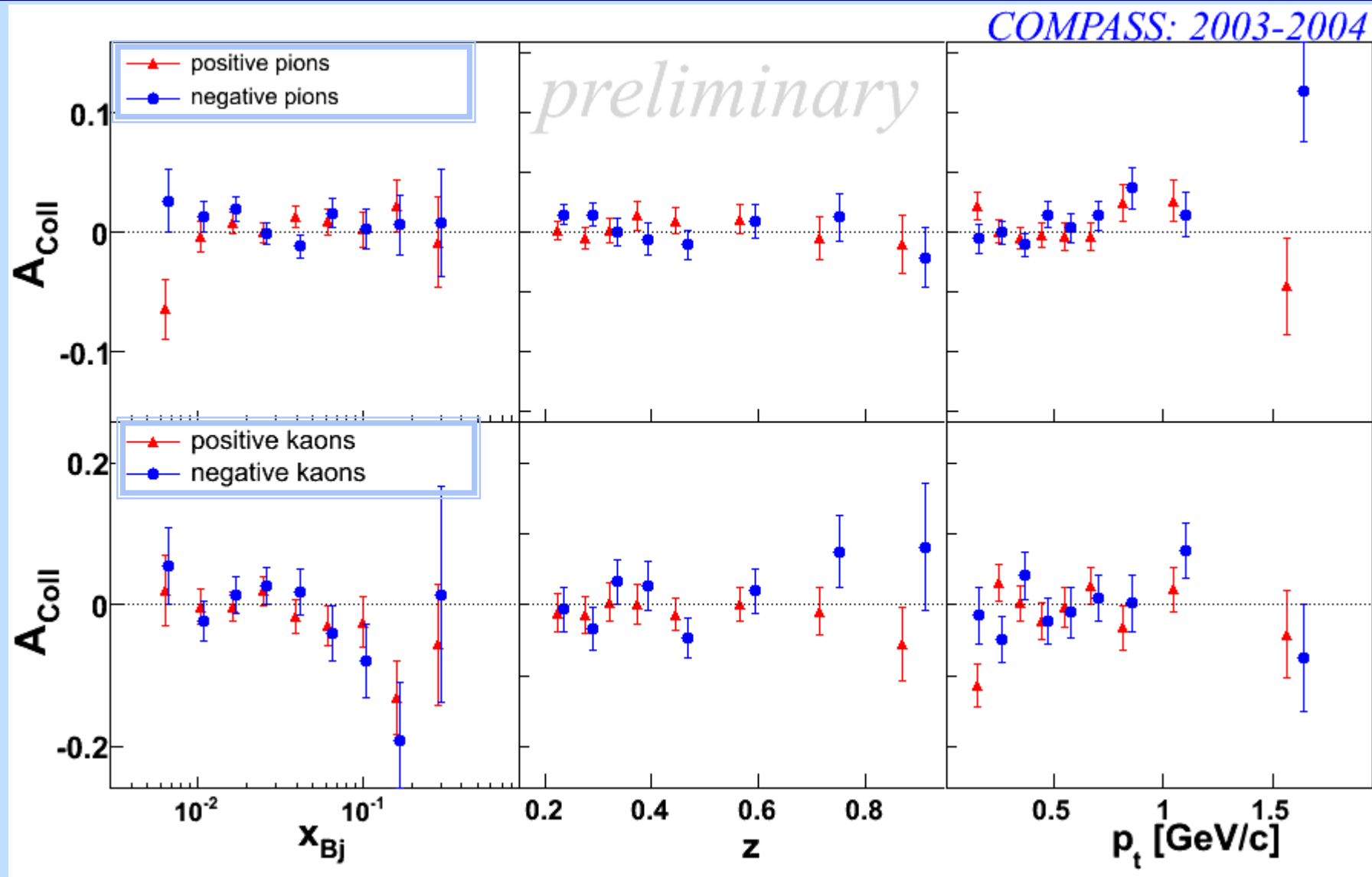
[[NP B765 \(2007\) 31-70](#)]

- Hadron identification is based on **RICH response**: several studies performed on the stability in time of the detector.
- Cherenkov thresholds:  $\pi \sim 3 \text{ GeV}/c$   
 $K \sim 9 \text{ GeV}/c$   
 $p \sim 17 \text{ GeV}/c$
- $2 \sigma$   $\pi/K$  separation at  $43 \text{ GeV}/c$
- In the leading hadron sample:
  - ~76% pions
  - ~12% kaons



	positive	negative
leading $\pi$	3.4M	2.8M
leading $K$	0.7M	0.4M

# COLLINS asymmetries 2003-2004 data



# naïve interpretation (parton model, valence region)

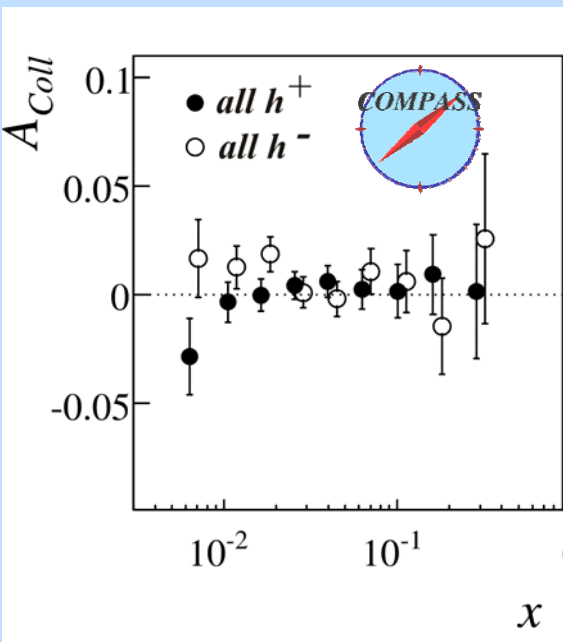
- from proton data (HERMES) non zero Collins asymmetry

→ unfavored Collins FF  $\approx$  - favored Collins FF

$$\Delta_T^0 D_2 \approx -\Delta_T^0 D_1 \quad \text{at variance with unpol case}$$

from proton: u quark dominance (d quark DF ~ unconstrained)

- deuteron data (COMPASS) small asymmetries compatible with 0



$$A_{Coll}^{d,\pi^+} \simeq \frac{\Delta_T u_v + \Delta_T d_v}{u_v + d_v} \frac{4\Delta_T^0 D_1 + \Delta_T^0 D_2}{4D_1 + D_2}$$

$$A_{Coll}^{d,\pi^-} \simeq \frac{\Delta_T u_v + \Delta_T d_v}{u_v + d_v} \frac{\Delta_T^0 D_1 + 4\Delta_T^0 D_2}{D_1 + 4D_2}$$

→ Cancellation between  $\Delta_T u(x)$  and  $\Delta_T d(x)$

→ Deuteron data give access to  $\Delta_T d(x)$

for a global analysis of Hermes, Belle and Compass data see the works of Vogelsang and Yuan, Efremov et al., Anselmino et al.

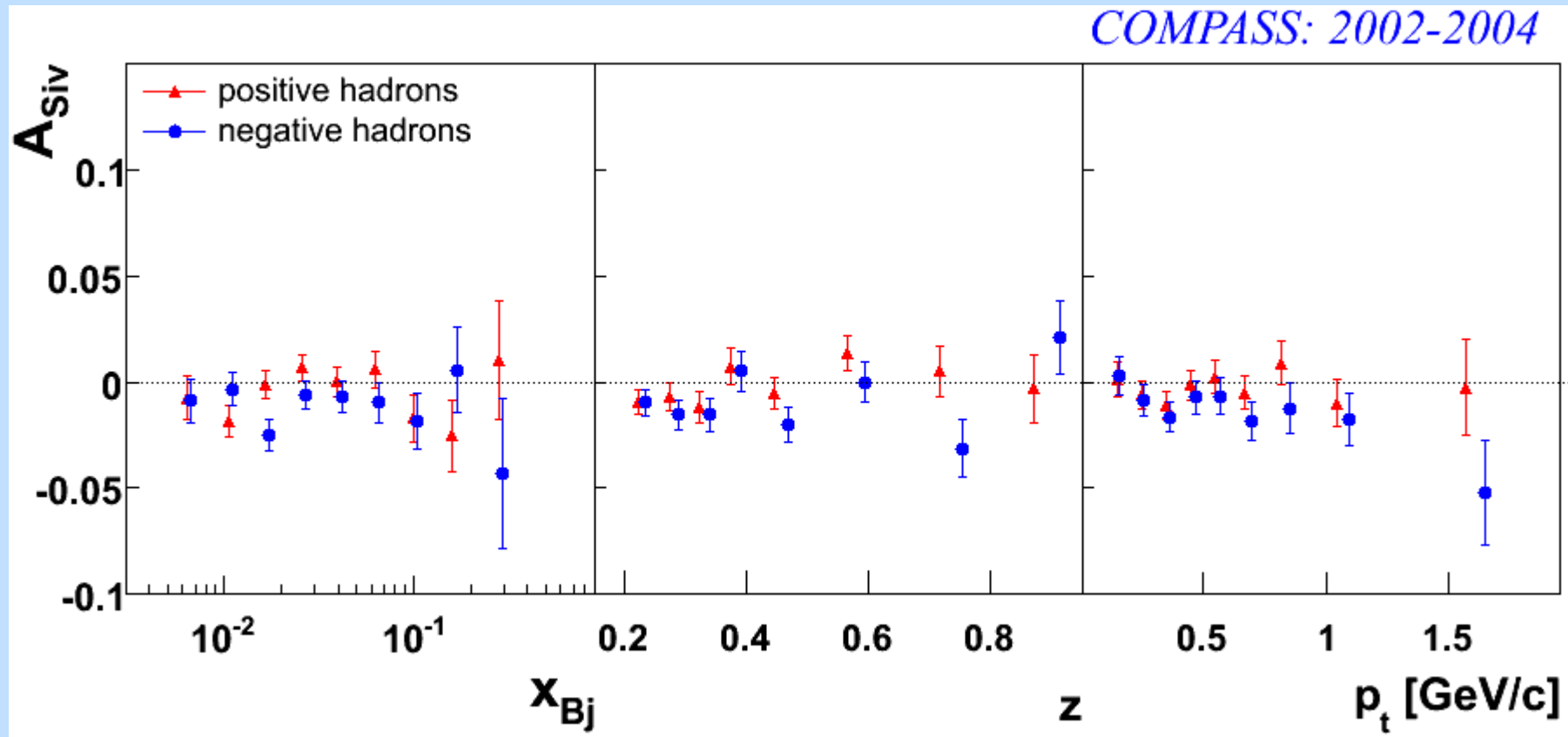
# SIVERS mechanism

- The Sivers DF  $\Delta_0^T q$  is probably the most famous between TMDs...
- gives a measure of the **correlation between the transverse momentum and the transverse spin**
- requires final/initial state interactions  
quark rescattering via soft gluon exchange
- it is related to the **parton orbital angular momentum** in a transversely polarized nucleon
- should change sign from SIDIS to DY  $\Delta_0^T q(x, k_T^2)_{SIDIS} = -\Delta_0^T q(x, k_T^2)_{DY}$

In SIDIS:

$$\mathbf{N}_h^\pm(\Phi_s) = \mathbf{N}_h^0 \cdot \left\{ \mathbf{1} \pm \mathbf{A}_S^h \cdot \sin\Phi_s \right\}$$

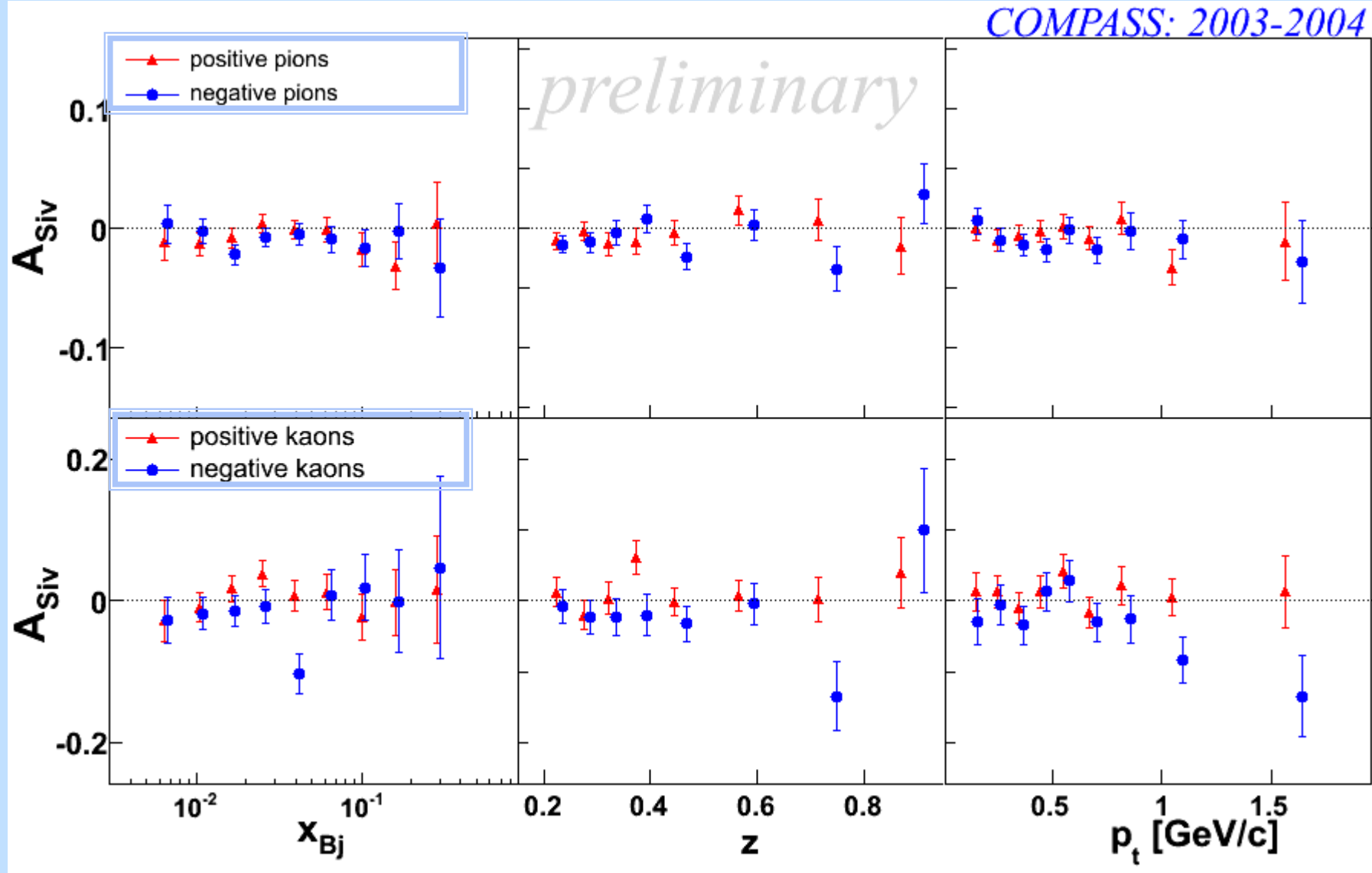
$$\mathbf{A}_{Siv} = \frac{\mathbf{A}_S^h}{\mathbf{f} \cdot \mathbf{P}_T} = \frac{\sum_q e_q^2 \Delta_0^T q \cdot \mathbf{D}_q^h}{\sum_q e_q^2 \cdot \mathbf{q} \cdot \mathbf{D}_q^h}$$



- only statistical errors shown (~1%), systematic errors considerably smaller
- **small asymmetries compatible with 0 for both + and - hadrons**

[[NP B765 \(2007\) 31-70](#)]

# SIVERS asymmetries 2003-2004 data





# naïve interpretation (parton model, valence region)

## ◆ proton data (HERMES)

asymmetry for  $\pi^+ > 0$ ,

asymmetry for  $\pi^- \approx 0$

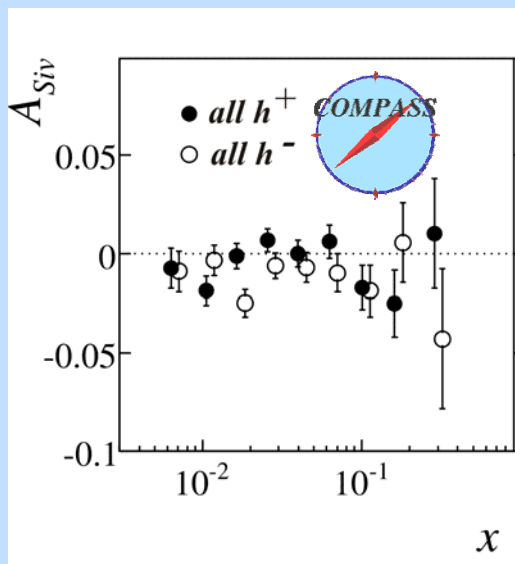
$$A_{Siv}^{p,\pi^+} \simeq \frac{4\Delta_0^T u_v D_1 + \Delta_0^T d_v D_2}{4u_v D_1 + d_v D_2}$$

$$A_{Siv}^{p,\pi^-} \simeq \frac{4\Delta_0^T u_v D_2 + \Delta_0^T d_v D_1}{4u_v D_2 + d_v D_1}$$

→ **Sivers DF for d-quark  $\approx -2$  Sivers DF for u-quark**

$$\Delta_0^T d_v \simeq -2 \cdot \Delta_0^T u_v$$

## ◆ deuteron data (COMPASS)



$$A_{Siv}^{d,\pi^+} \simeq A_{Siv}^{d,\pi^-} \simeq \frac{\Delta_0^T u_v + \Delta_0^T d_v}{u_v + d_v}$$

the measured asymmetries  
**compatible with zero** suggest

$$\Delta_0^T d_v \simeq -\Delta_0^T u_v$$



**the measured asymmetry on deuteron compatible with zero has been interpreted as**

## **Evidence for the Absence of Gluon Orbital Angular Momentum in the Nucleon**

**S.J. Brodsky and S. Gardner, PLB643 (2006) 22**

The approximate cancellation of the SSA measured on a deuterium target suggests that the gluon mechanism, and thus the orbital angular momentums carried by gluons in the nucleon, is small.

# two hadrons - the coordinate system

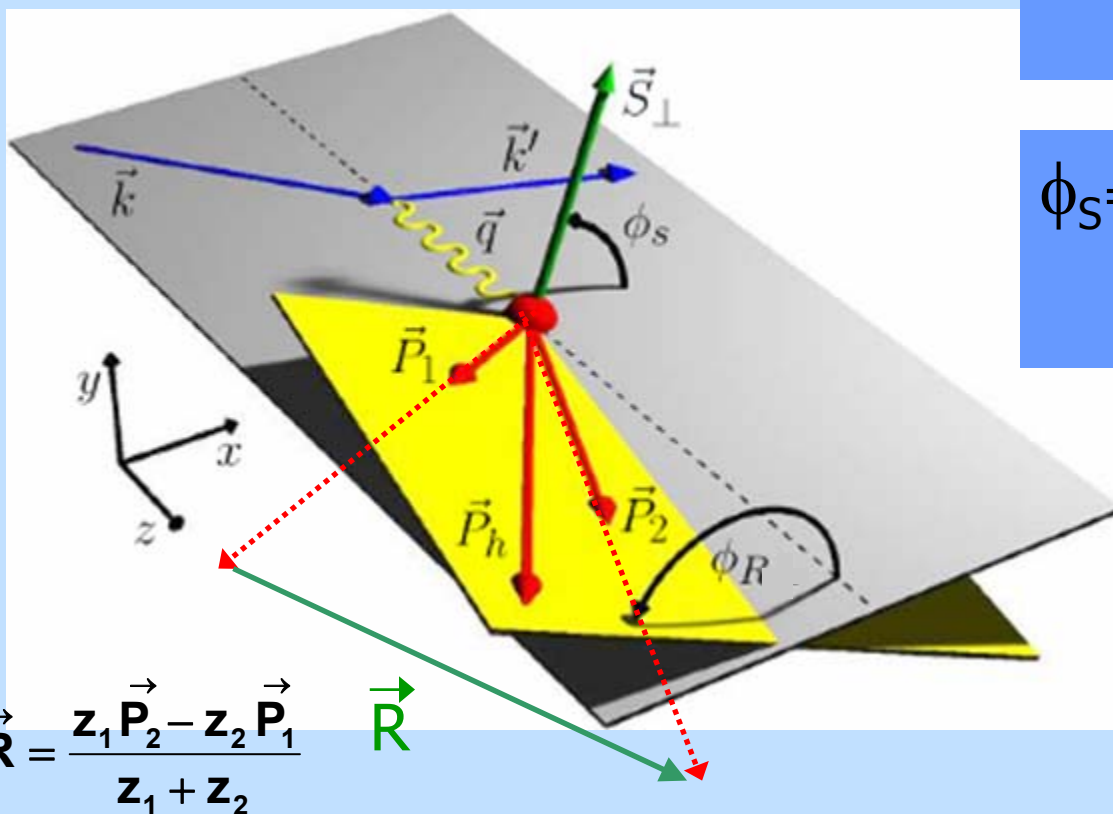
z-axis = virtual photon direction  
 x-z plane = lepton scattering plane

$\phi_R$  = angle between lepton scattering plane and two-hadron plane

$\phi_S$  = azimuthal angle of initial quark versus lepton scattering plane

$\phi_{S'} = \pi - \phi_S$   
 (fragmenting quark)

$$\begin{aligned} \phi_{RS} &= \phi_R - \phi_{S'} \\ &= \phi_R + \phi_S - \pi \end{aligned}$$



(A. Bacchetta, M. Radici, hep-ph/0407345)

(X. Artru, hep-ph/0207309)

# azimuthal asymmetry for two-hadron production



Target single spin asymmetry  $A_{RS}(x, z, M_h^2)$ :

$$z = z_1 + z_2$$

$$N^\pm(\Phi_{RS}) = N_0 \left\{ 1 \pm A_{UT}^{\sin\Phi_{RS}} \sin\Phi_{RS} \right\} \quad \text{and}$$

$$A_{RS} = \frac{1}{f P_T D} A_{UT}^{\sin\Phi_{RS}}$$

$N^\pm(\Phi_{RS})$ : Number of events for target spin up (+) and down (-)

f: Dilution factor  $\approx 0.38$

D: Depolarisation factor  
 $D = (1-y)/(1-y+y^2/2)$

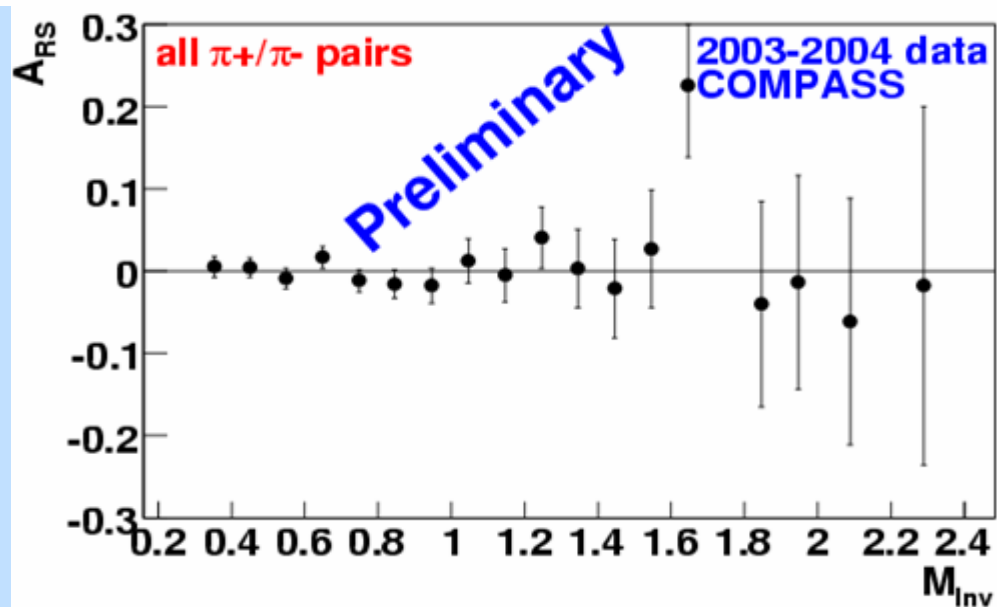
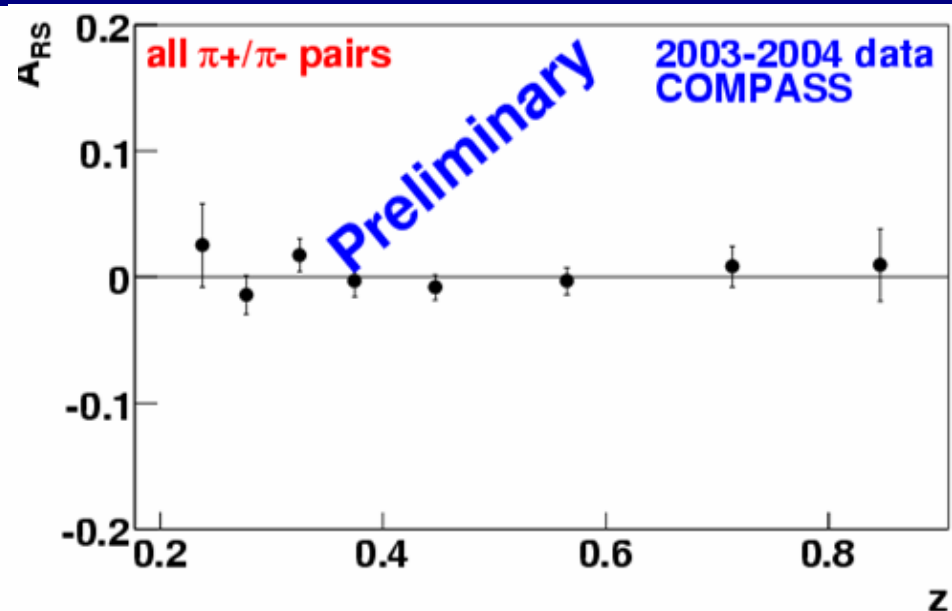
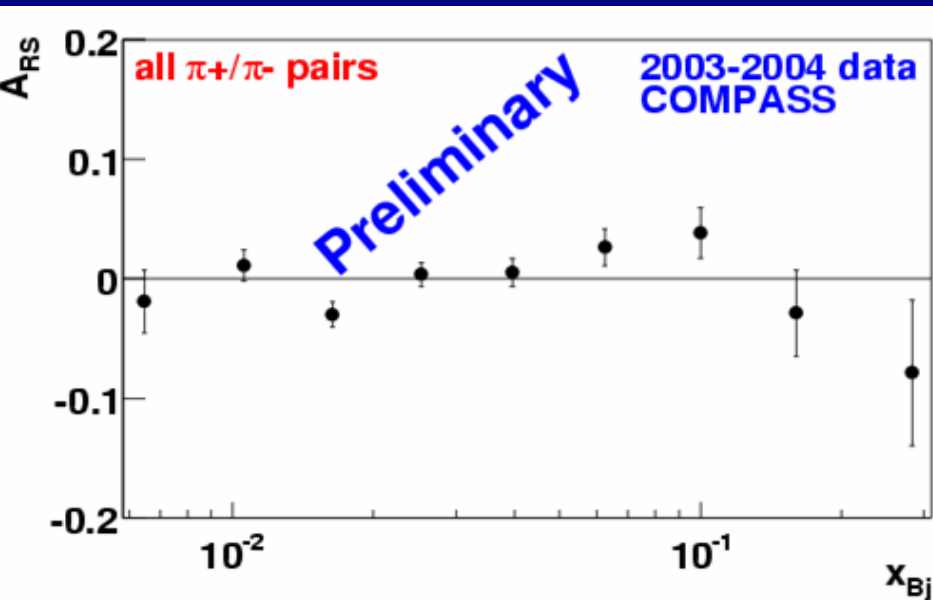
$P_T$ : Target polarisation  $\approx 0.5$

$$A_{RS}(x, z, M_h^2) = \frac{\sum_q e_q^2 \Delta_T q(x) H_q^{\perp\angle h}(z, M_h^2)}{\sum_q e_q^2 q(x) D_q^h(z, M_h^2)}$$

$D_q^h$ ,  $H_q^{\perp\angle h}$  presently unknown can be measured in  $e^+e^-$  (BELLE)  
 expected to depend on the hadron pair invariant mass

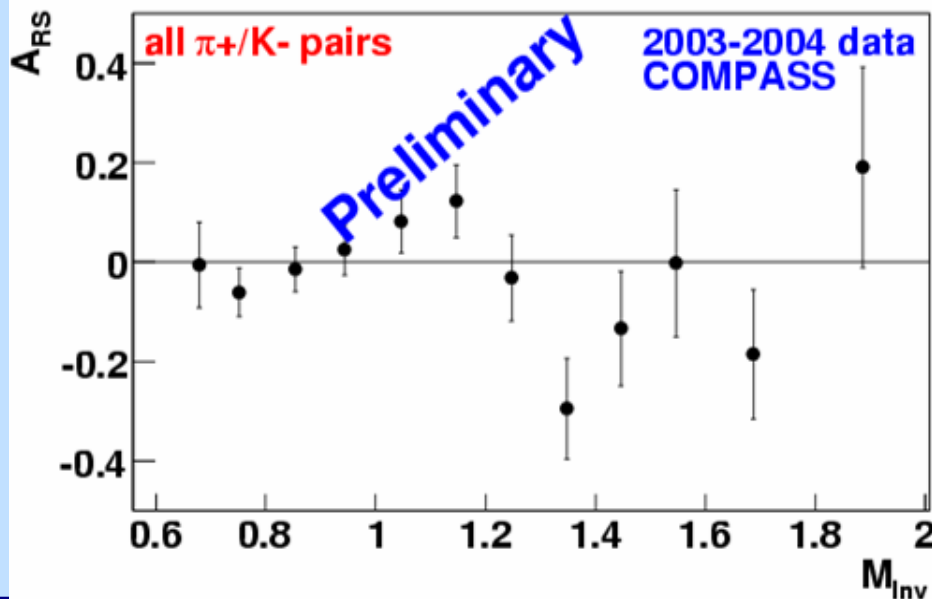
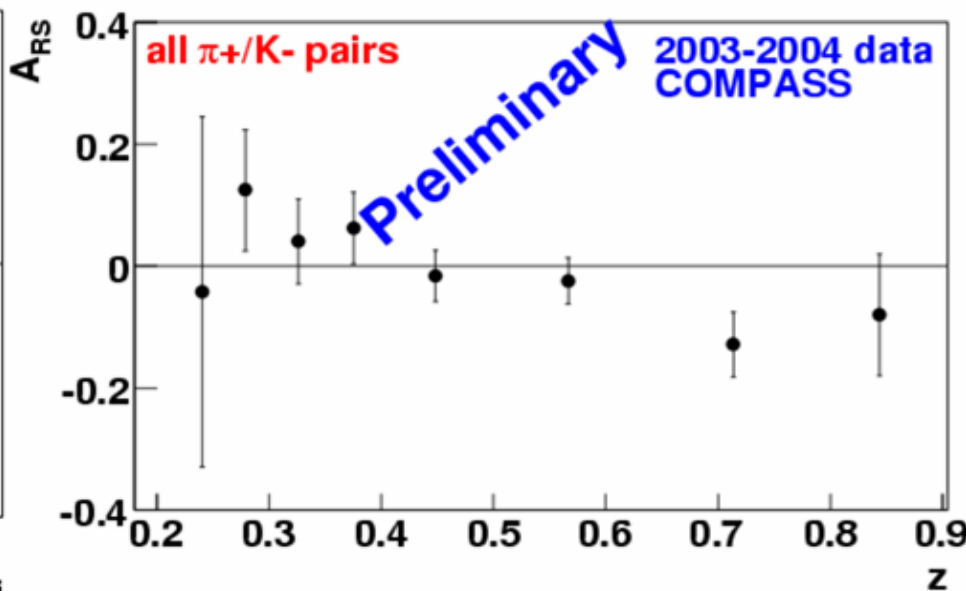
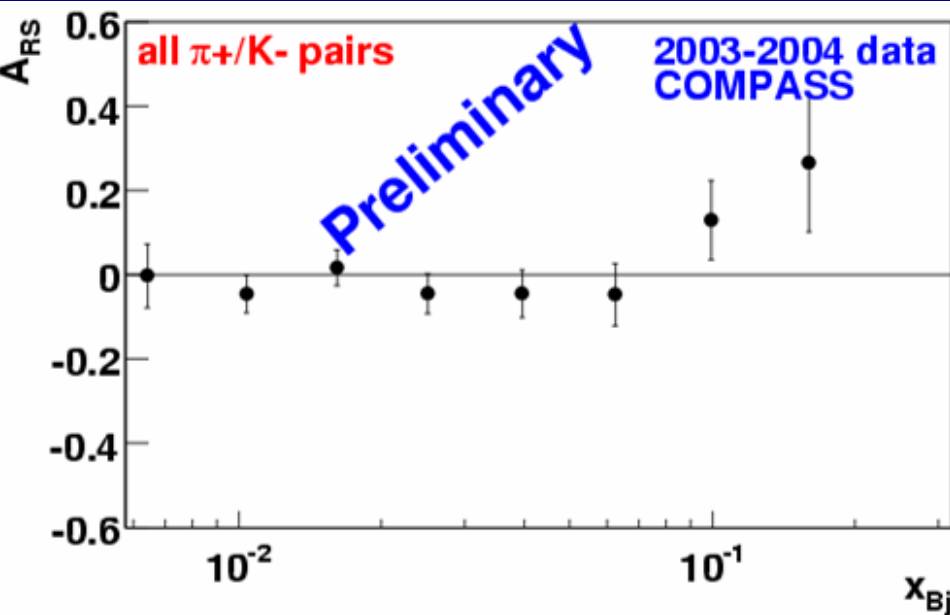
(X. Artru, hep-ph/0207309)

# results for $\pi^+ \pi^-$ pairs (RICH identification)

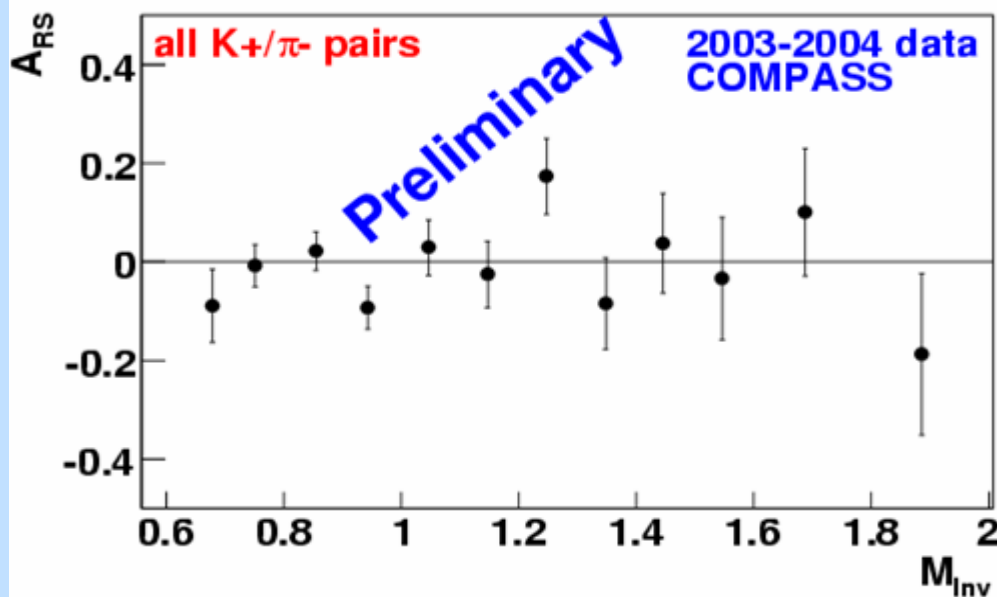
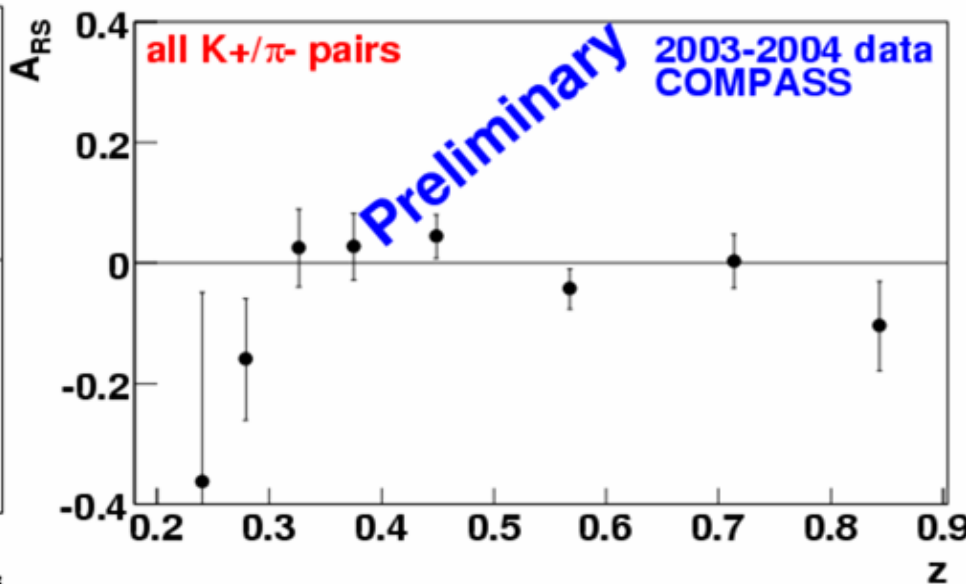
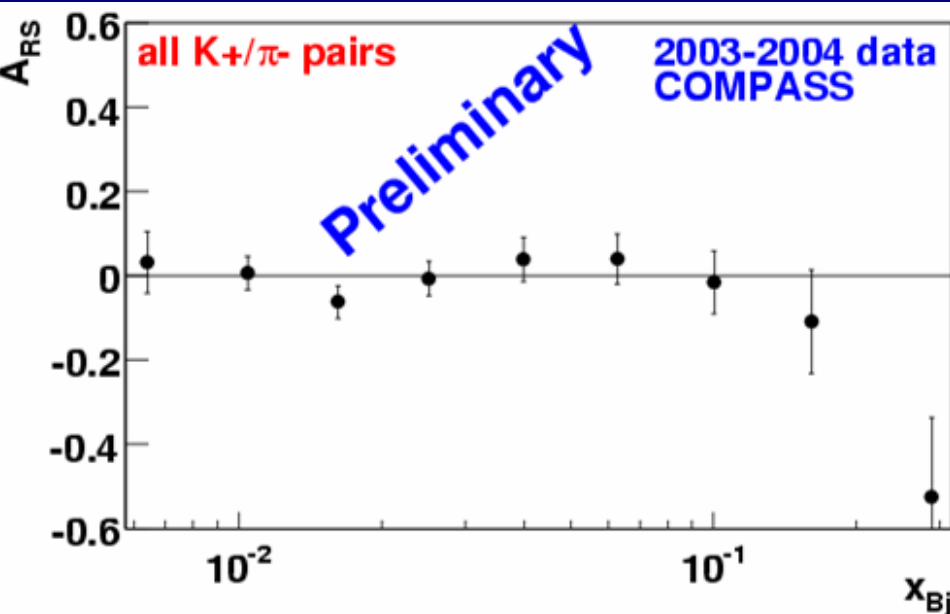


$$z = z_1 + z_2$$

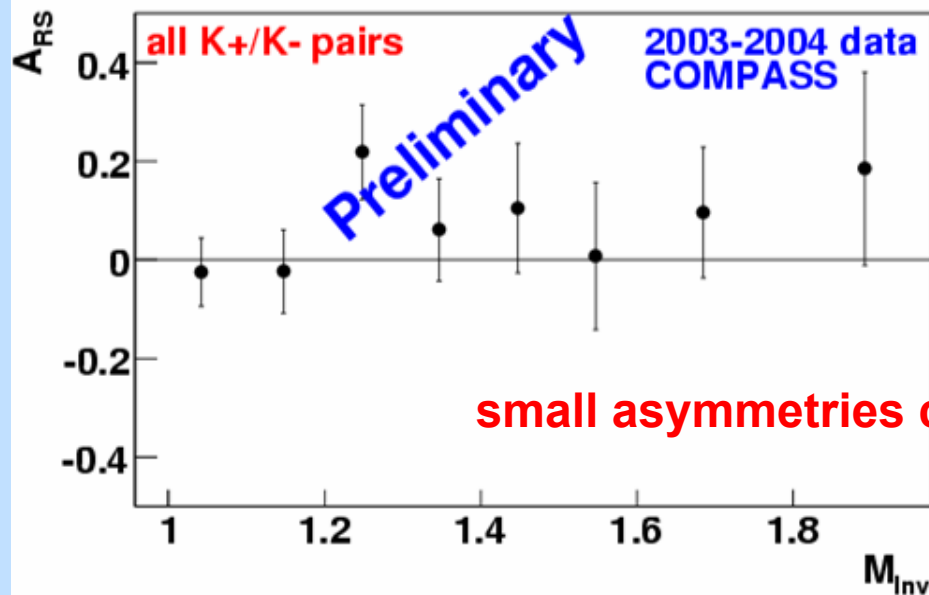
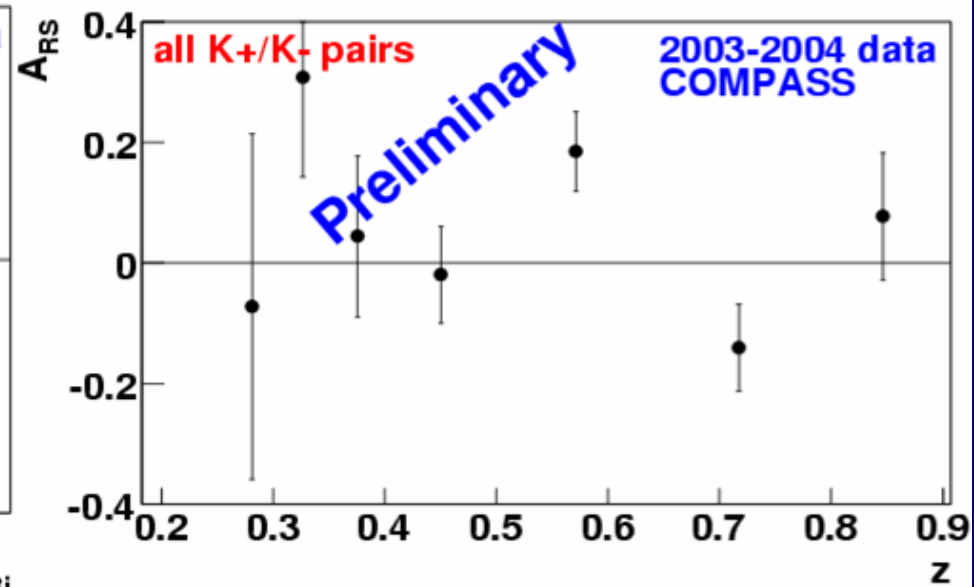
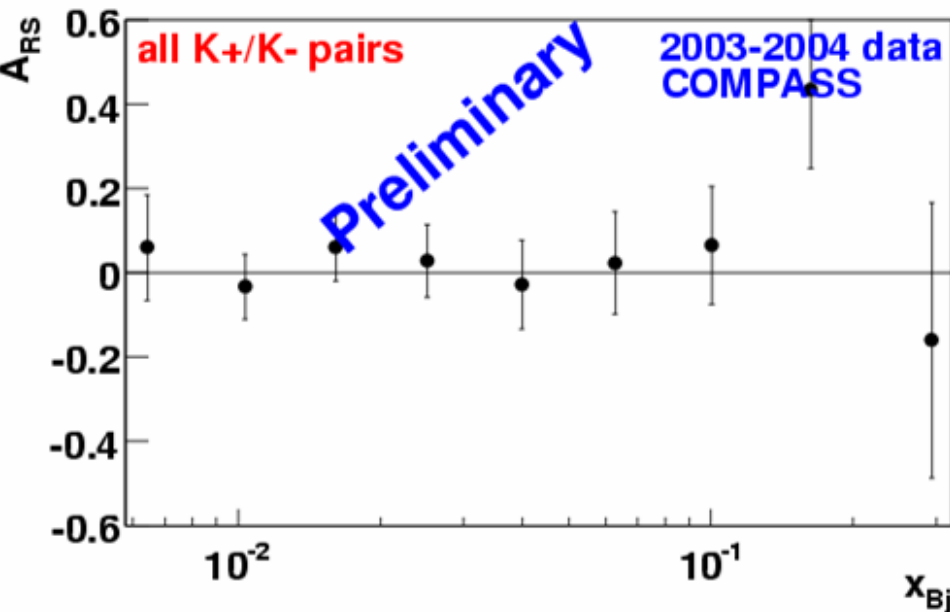
# results for $\pi^+ K^-$ pairs (RICH identification)



# results for $K^+ \pi^-$ pairs (RICH identification)



# results for $K^+ K^-$ pairs (RICH identification)





# other TMDs: general expression of polarized SIDIS xSections



$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \\
 \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) & \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + P_{beam} \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} & \\
 + P_L \left[ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] & \\
 + P_L P_{beam} \left[ \sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] & \\
 + |P_T| \left[ \sin(\phi_h - \phi_s) \left( F_{UT,T}^{\sin(\phi_h - \phi_s)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_s)} \right) \right. & \\
 \text{Sivers} & \\
 + \varepsilon \sin(\phi_h + \phi_s) F_{UT}^{\sin(\phi_h + \phi_s)} + \varepsilon \sin(3\phi_h - \phi_s) F_{UT}^{\sin(3\phi_h - \phi_s)} & \\
 \text{'Collins'} & \\
 + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_s F_{UT}^{\sin\phi_s} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_s) F_{UT}^{\sin(2\phi_h - \phi_s)} & \\
 + |P_T| P_{beam} \left[ \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_s) F_{LT}^{\cos(\phi_h - \phi_s)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_s F_{LT}^{\cos\phi_s} \right. & \\
 + \left. \left. \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_s) F_{LT}^{\cos(2\phi_h - \phi_s)} \right] \right\}, &
 \end{aligned}$$

## Azimuthal modulations:

2 polarization independent

1 single beam polarization dependent

2 single target longitudinal polarization dependent

1 double beam + target longitudinal polarization dependent

+

5 single target transverse polarization dependent

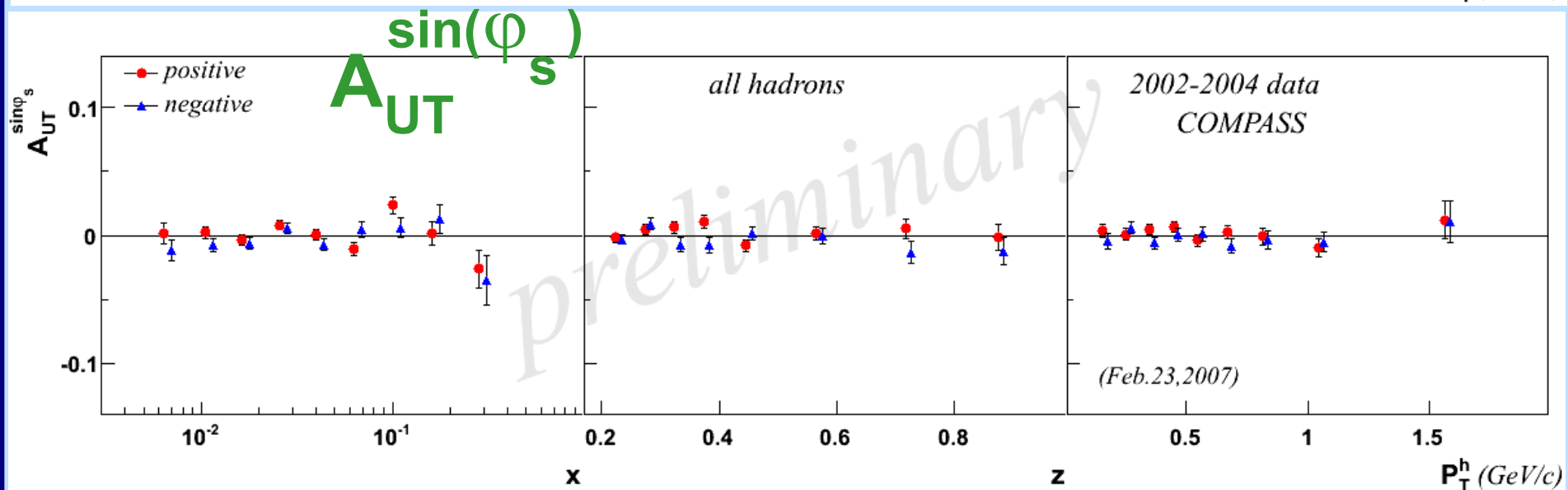
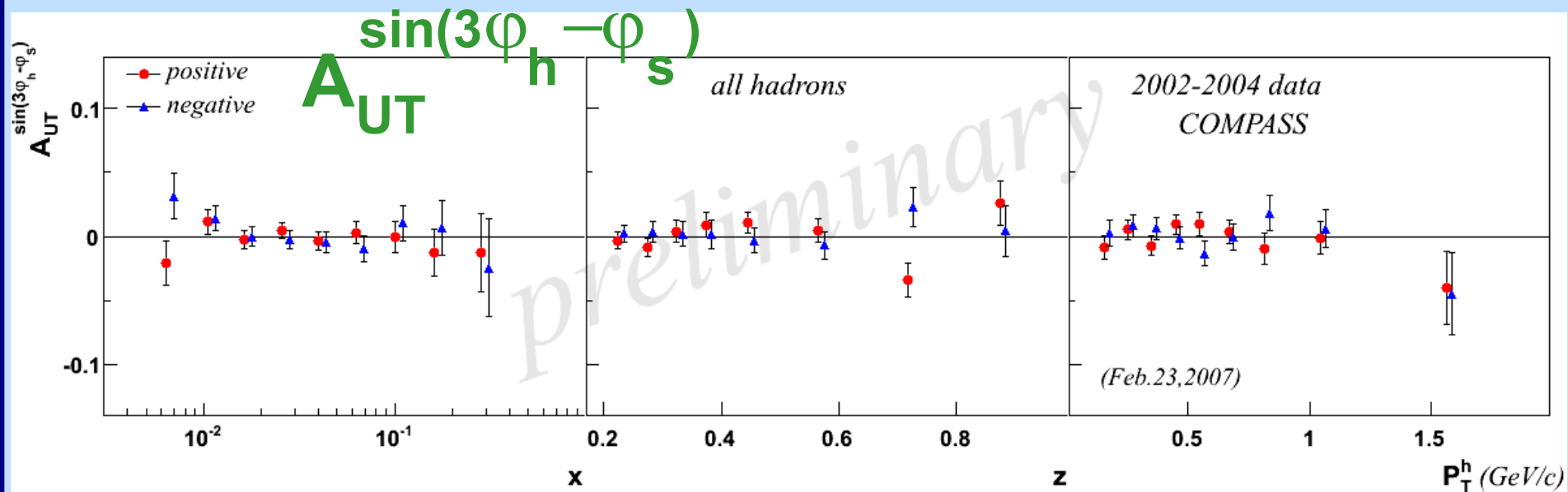
3 double beam + target transverse polarization dependent

5

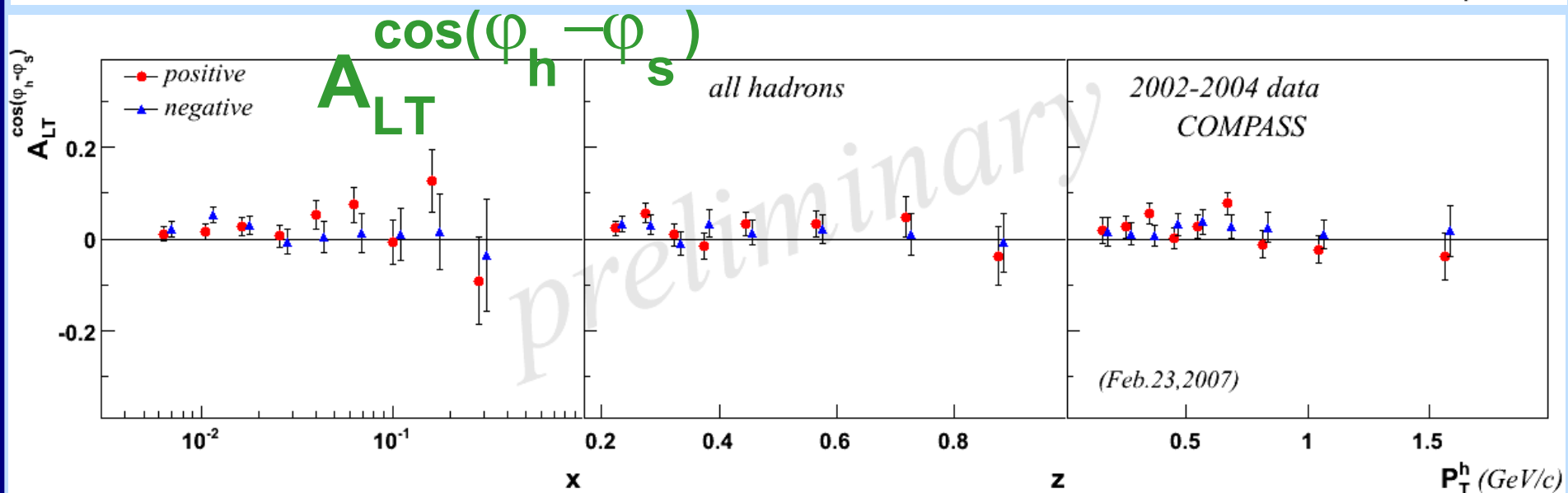
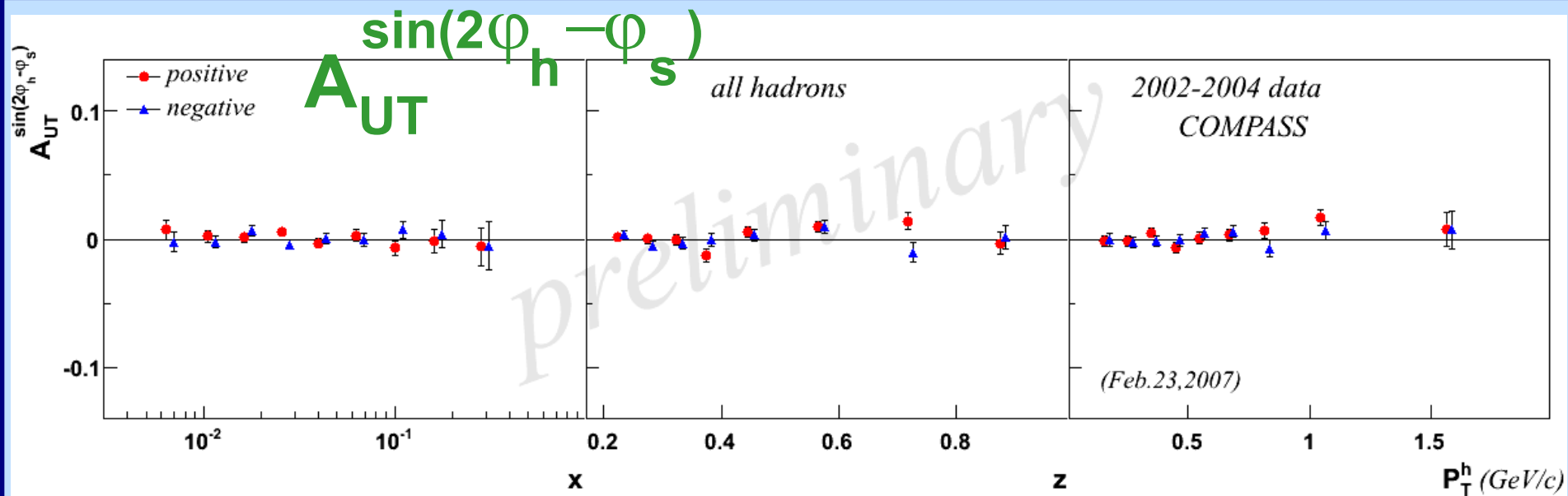
independent angles

$$\begin{aligned}
 \Phi_1 &= \phi_h - \phi_s \\
 \Phi_2 &= \phi_h + \phi_s \\
 \Phi_3 &= 3\phi_h - \phi_s \\
 \Phi_4 &= \phi_s \\
 \Phi_5 &= 2\phi_h - \phi_s
 \end{aligned}$$

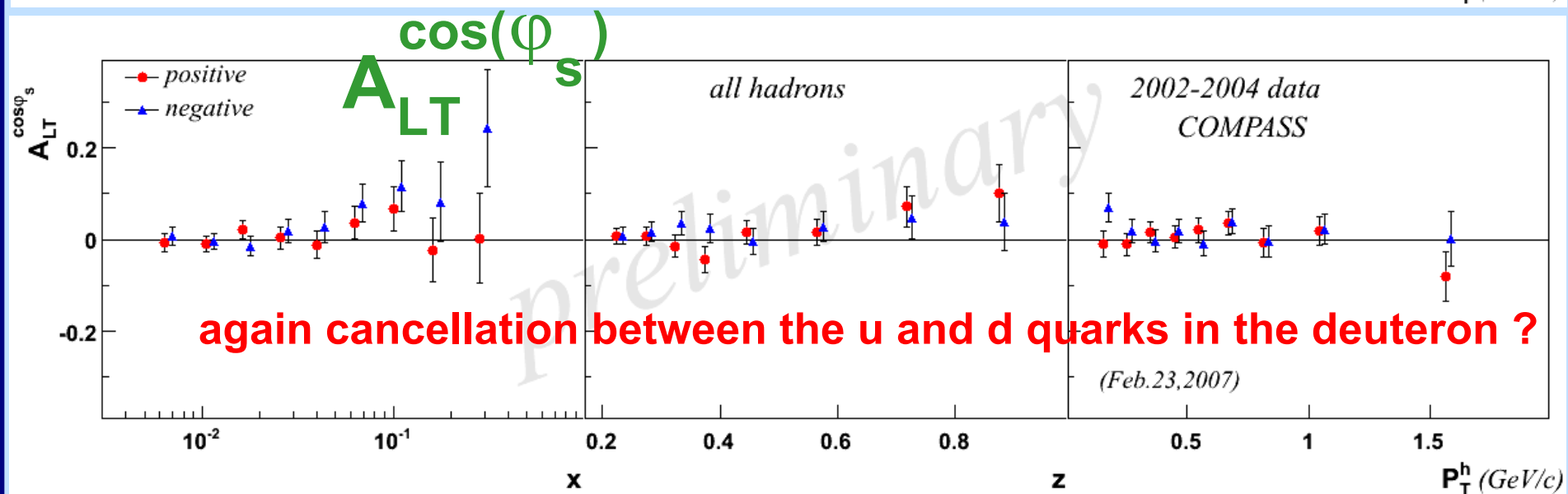
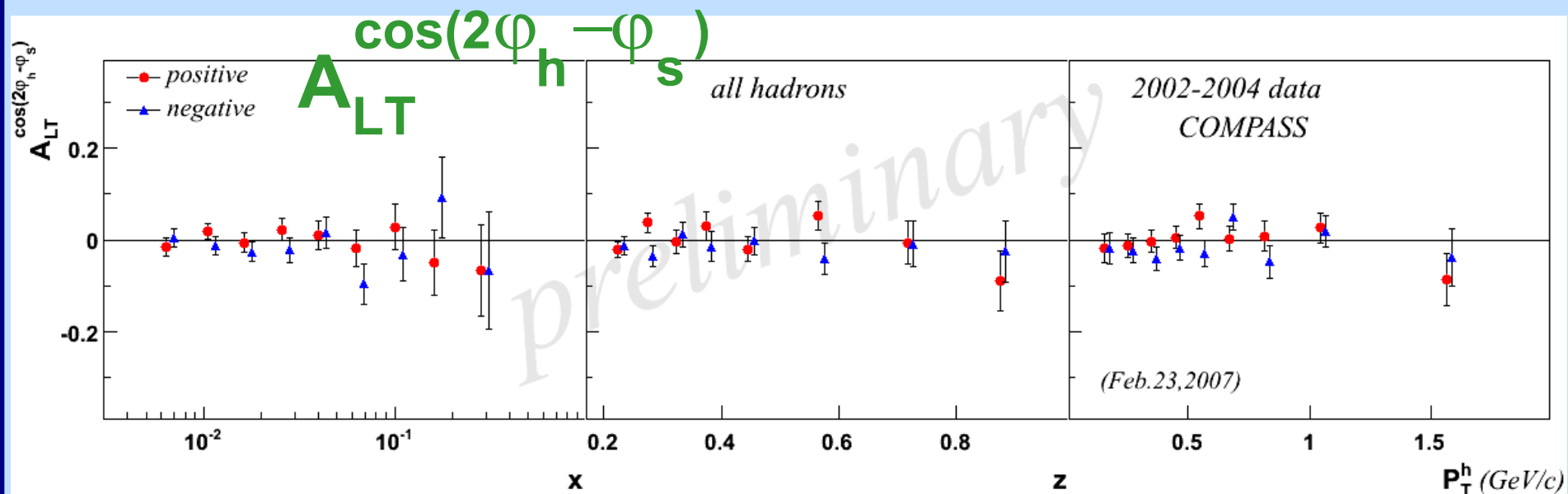
# asymmetry extraction - I



# asymmetry extraction - II



# asymmetry extraction - III





**now: precise deuteron data from COMPASS are available**

*Collins and Sivers asymmetries  $h^\pm$ ,  $\pi^\pm$ ,  $K^\pm$*

*Two hadron asymmetries*

*all other TMD SSA azimuthal asymmetries*

**→ deuteron asymmetries very small, compatible with zero**

**→ first extractions of the u and d quark DFs**

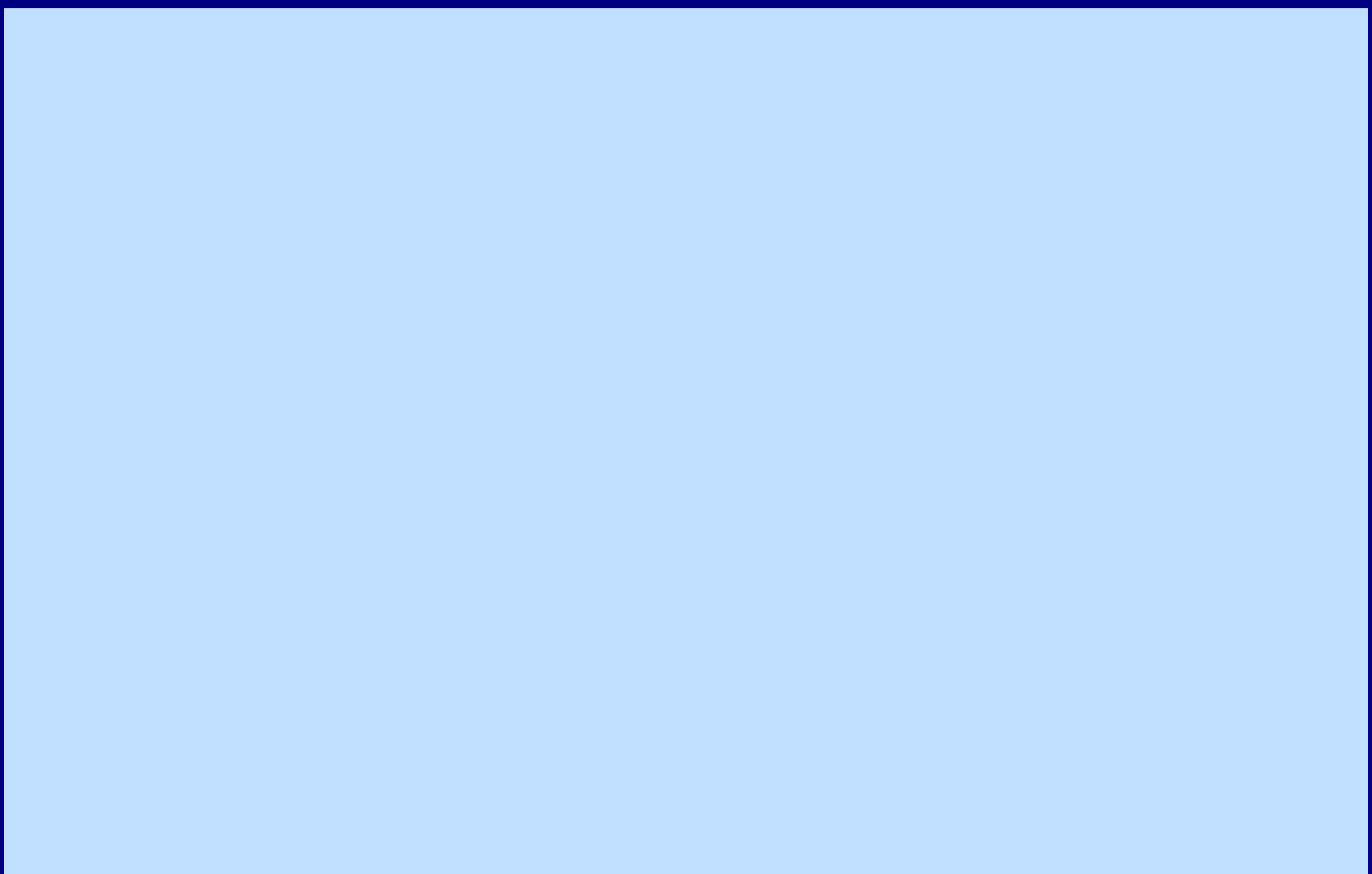
**near future: COMPASS will complete the analysis of the recorded deuteron data**

- *$K^0$  asymmetries*
- *exclusive  $\rho$  asymmetries on transversely polarised target → GPD E*
- *transverse effects from longitudinal data*
- *Cahn and Boer-Mulders effect*

**2007: COMPASS data will be with a transversely polarised proton target ( $\text{NH}_3$ )**

**with 50 days, same precision at small x as for deuterium, better at “large” x for the new PT magnet**

**on a longer time scale: good perspective for a measurement of Drell-Yan pairs in COMPASS  
and further measurements of transverse spin effects in SIDIS**



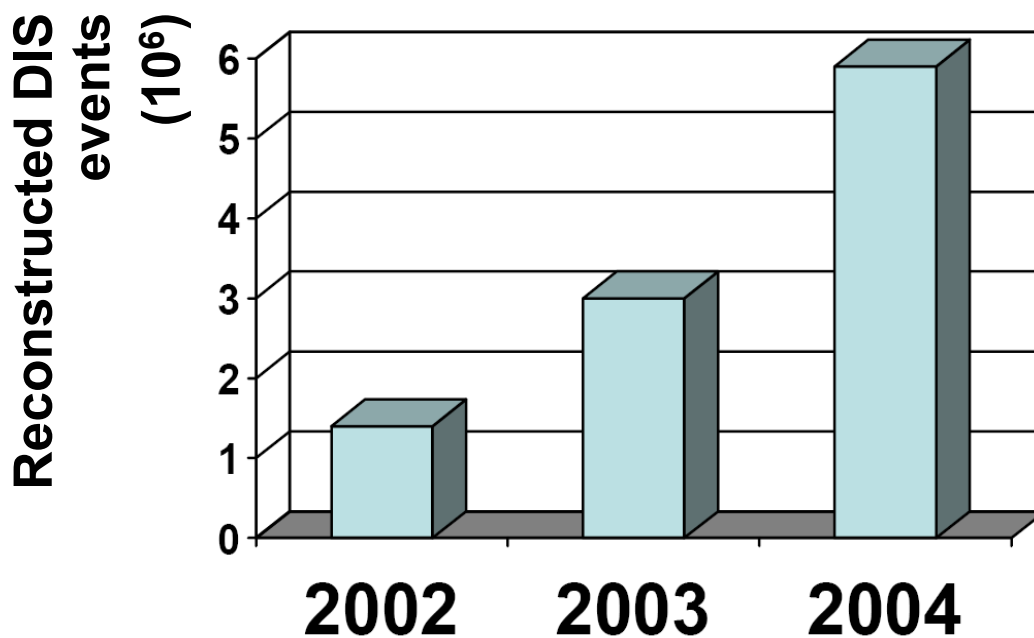


**transversely polarised deuteron target ( $^6\text{LiD}$ )  
~ 20% of the running time**

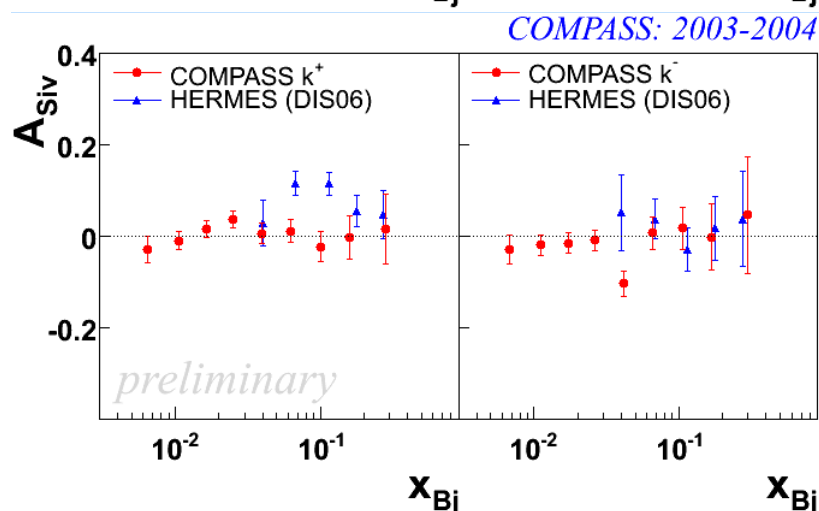
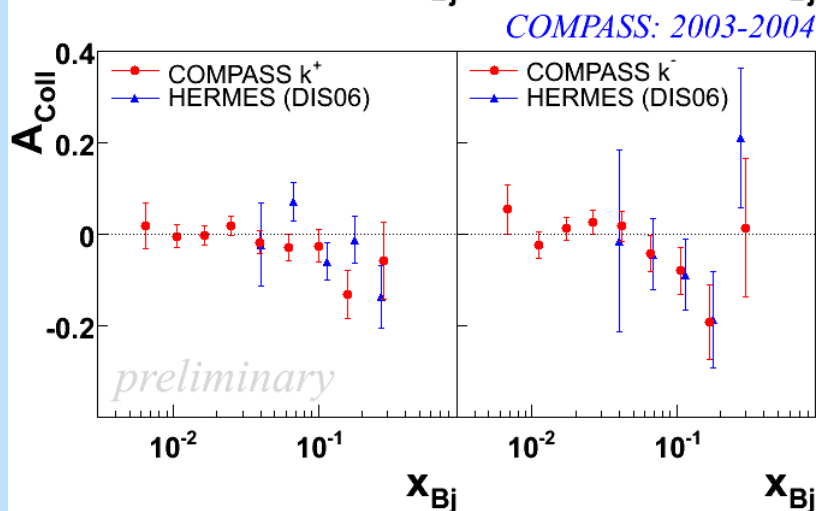
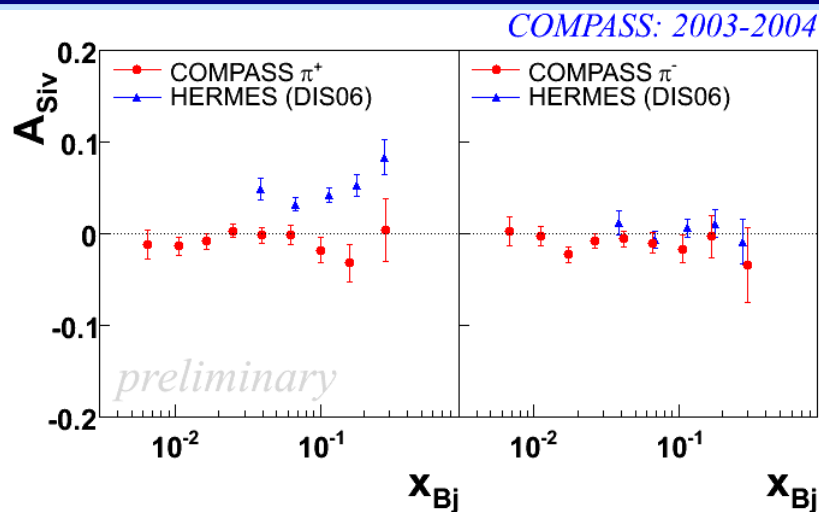
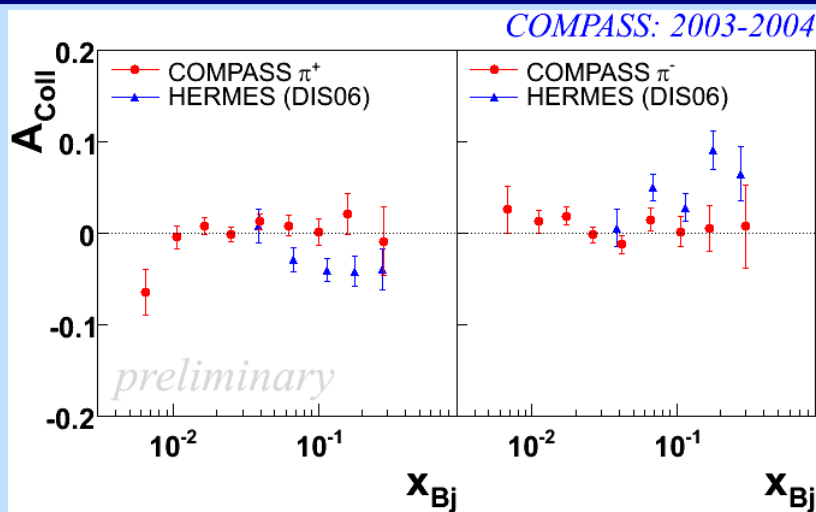
<b>2002</b>	<b>11 days of data taking,</b>	<b>2 periods</b>
<b>2003</b>	<b>9 days of data taking,</b>	<b>1 period</b>
<b>2004</b>	<b>14 days of data taking,</b>	<b>2 periods</b>

trigger (large  $x$ ,  $Q^2$ )

DAQ, on line filter



# comparison with HERMES for Collins and Sivers asym



HERMES data from 'Transversity results from HERMES', L.Pappalardo et al., to appear in the proceedings of the XIV International Workshop on Deep Inelastic Scattering, Tsukuba city, Japan, April 20-24, 2006., courtesy of the HERMES Collaboration



# SIVERS mechanism

- The Sivers DF  $\Delta_0^T q$  is probably the most famous between TMDs...
- gives a measure of the correlation between the transverse momentum and the transverse spin
- requires final/initial state interactions of the struck quark with the spectator system and the interference between different helicity Fock states to survive time-reversal invariance
- Time-reversal invariance implies:

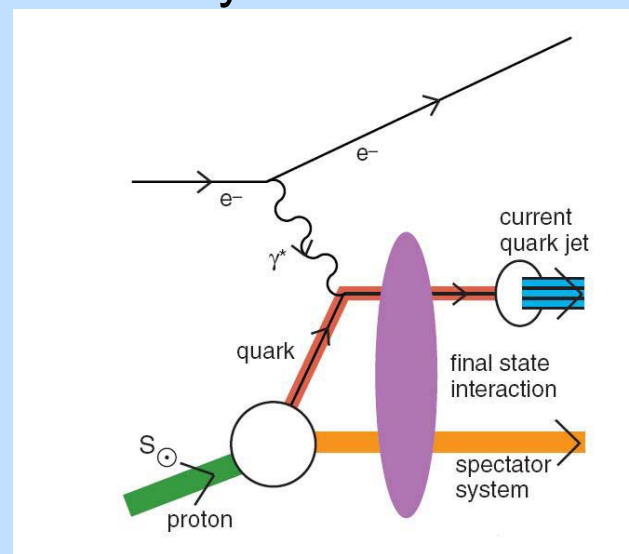
$$\Delta_0^T q(x, k_T^2)_{SIDIS} = -\Delta_0^T q(x, k_T^2)_{DY}$$

...to be checked

- In SIDIS:

$$\mathbf{N}_h^\pm(\Phi_s) = \mathbf{N}_h^0 \cdot \left\{ \mathbf{1} \pm \mathbf{A}_S^h \cdot \sin\Phi_s \right\}$$

$$\mathbf{A}_{Siv} = \frac{\mathbf{A}_S^h}{\mathbf{f} \cdot \mathbf{P}_T} = \frac{\sum_q e_q^2 \Delta_0^T q \cdot \mathbf{D}_q^h}{\sum_q e_q^2 \cdot \mathbf{q} \cdot \mathbf{D}_q^h}$$



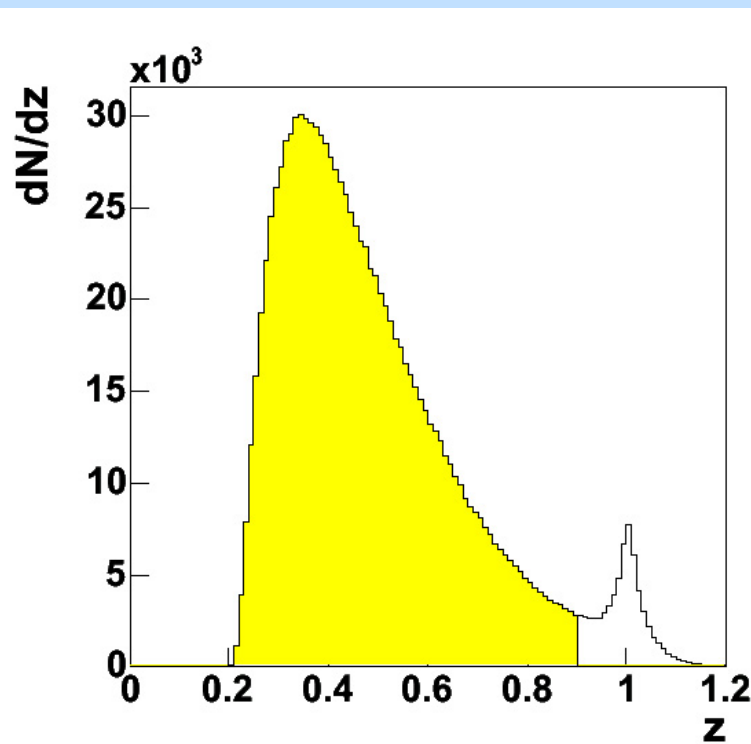
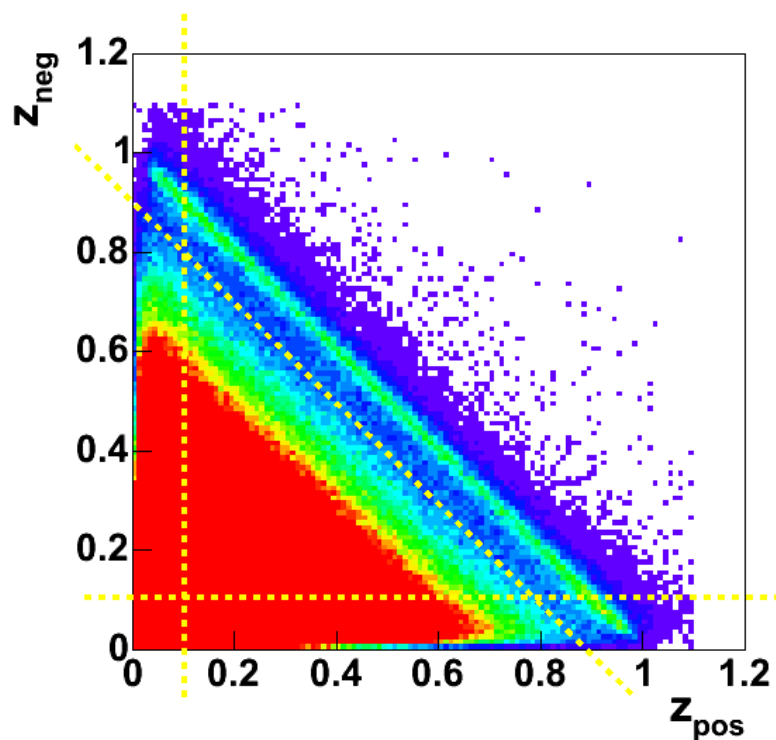
# event selection for 2 hadrons

## DIS cuts:

- $Q^2 > 1 \text{ GeV}^2/c^2$
- $0.1 < y < 0.9$
- $W > 5 \text{ GeV}/c^2$

## Hadron selection:

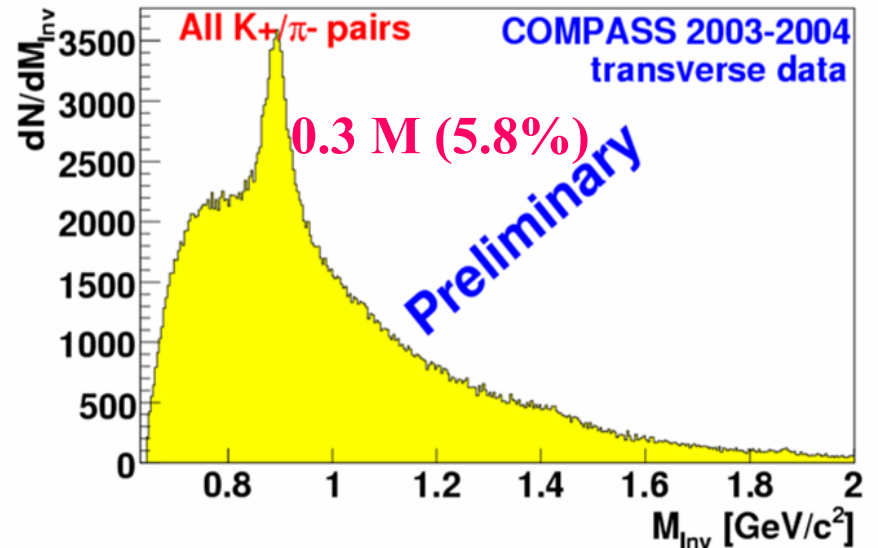
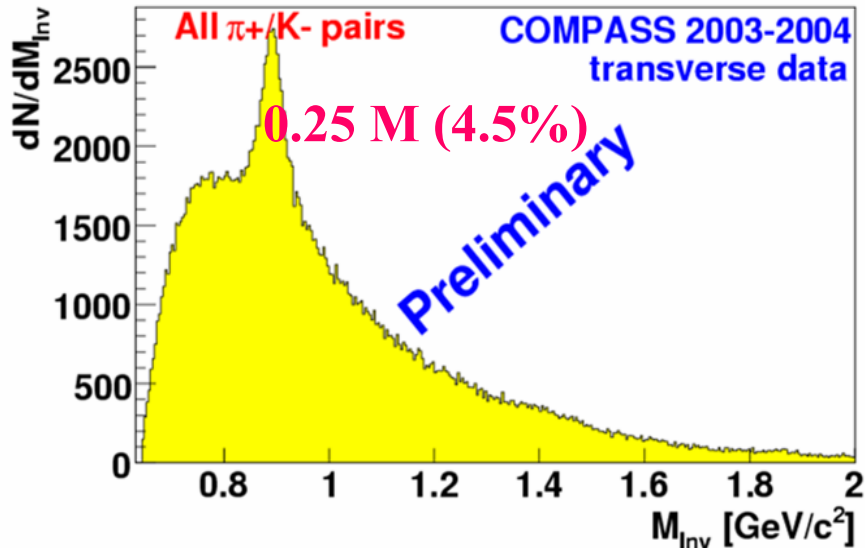
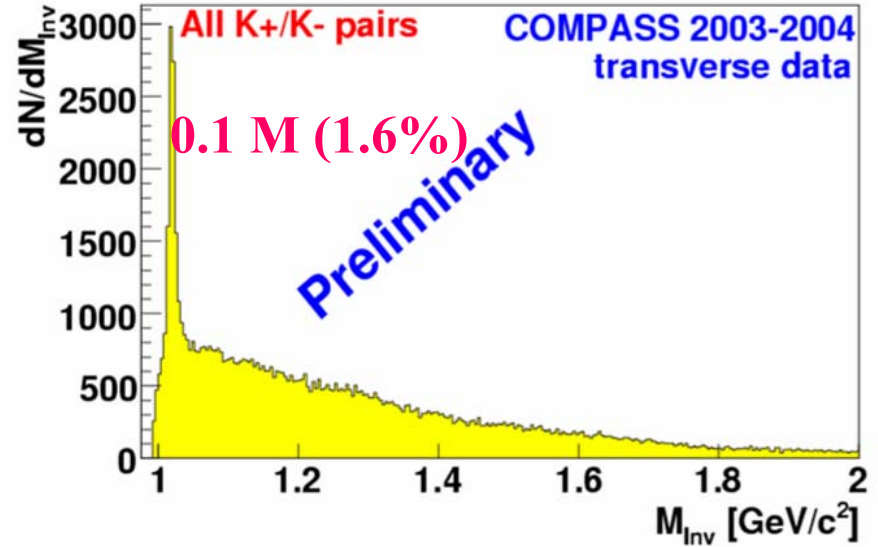
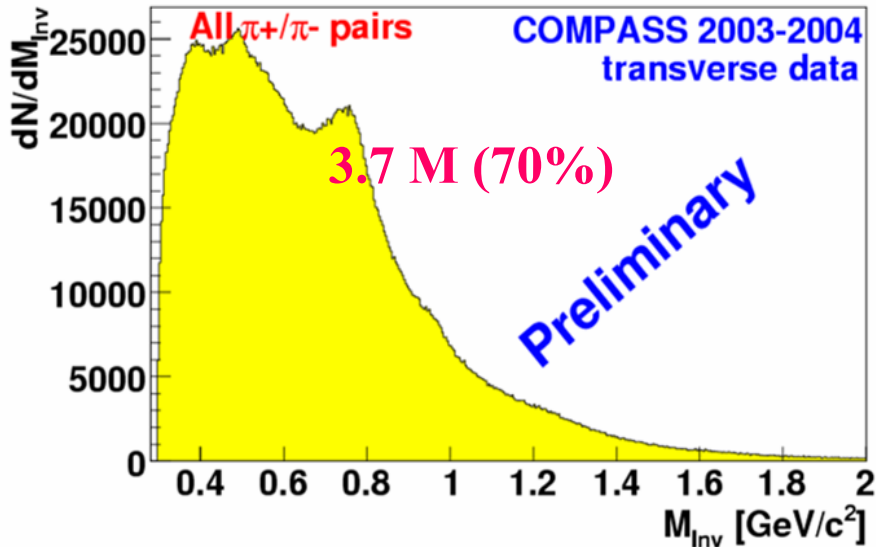
- $z_{1,2} > 0.1$  (current fragmentation)
- $x_{F1,2} > 0.1$
- $z = z_1 + z_2 < 0.9$  (exclusive rho)
- RICH identification of  $\pi, K$



# RICH hadron identification



all hadron pairs: 5.3 M

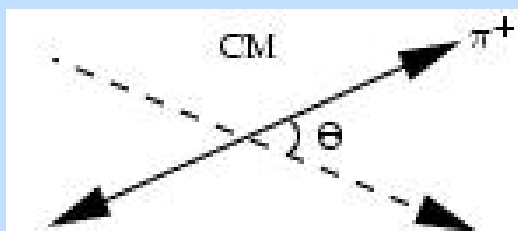


# sinθ dependance

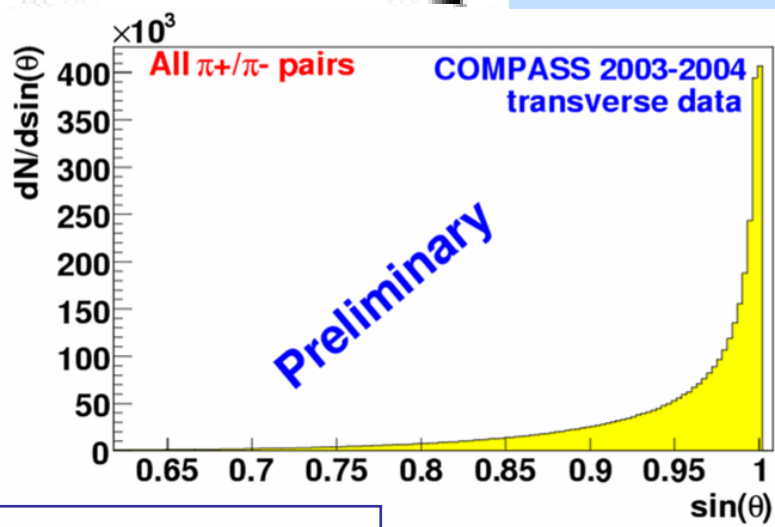
Cross section  $\sigma_{UT}$  for two- $\pi$  fragmentation depends on  $\sin\theta$ :  
 (Interference of s- and p-wave of the  $2\pi$ -state)

$$\sigma_{UT} \propto \sum_q e_q^2 |S_T| \sin\theta \sin\phi_{RS} \Delta_T q(x) H_q^{\perp \angle h}(z, M_h^2)$$

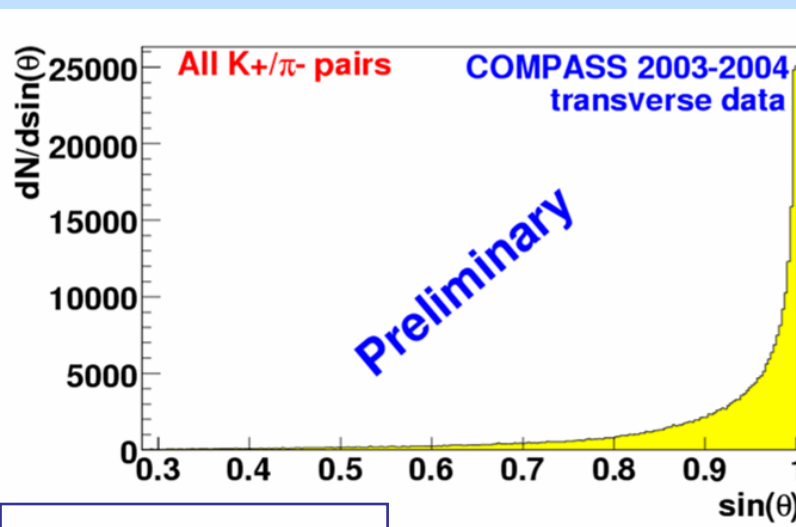
(A. Bacchetta and M. Radici, hep-ph/0212300)



$\theta$  : Angle of  $h_1$  in the two-hadron CMS  
 to the direction of  $P_h = P_{h1} + P_{h2}$



$\langle \sin\theta \rangle = 0.95$



$\langle \sin\theta \rangle = 0.90$

→ small contribution in the kinematic region of COMPASS

The number-of-events  
asymmetries

$$A_{UT, raw}^{w_i(\phi_h, \phi_s)} = D^{w_i(\phi_h, \phi_s)}(y) f_{S_T} A_{UT}^{w(\phi_h, \phi_s)}, \quad (i = 1, 5),$$

$$A_{LT, raw}^{w(\phi_h, \phi_s)} = D^{w(\phi_h, \phi_s)}(y) f_{P_{beam}|S_T} A_{LT}^{w(\phi_h, \phi_s)}, \quad (i = 6, 8)$$

Independent angles

$$\begin{aligned} \Phi_1 &= \phi_h - \phi_s \\ \Phi_2 &= \phi_h + \phi_s \\ \Phi_3 &= 3\phi_h - \phi_s \\ \Phi_4 &= \phi_s \\ \Phi_5 &= 2\phi_h - \phi_s \end{aligned}$$

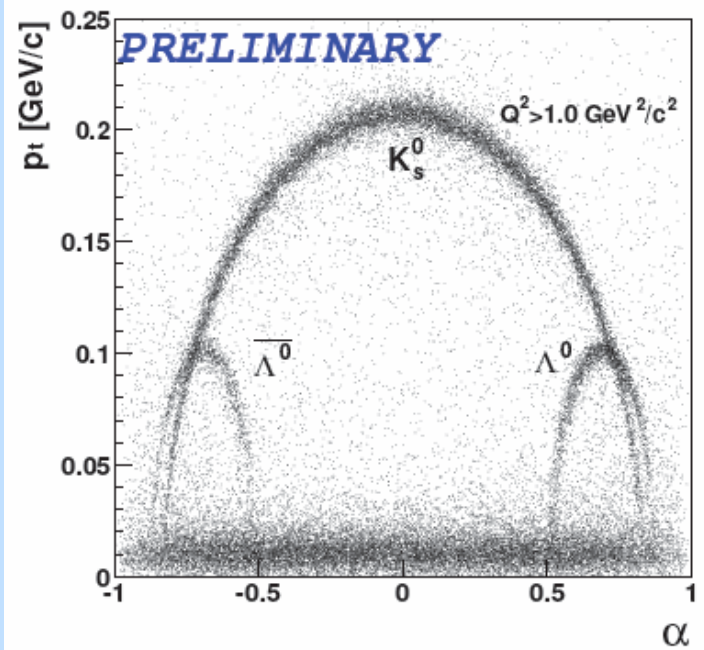
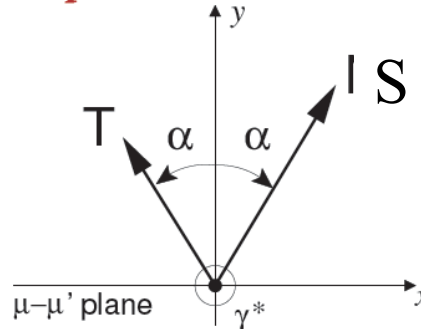
Azimuthal  
modulations

$$\begin{aligned} W_1(\Phi_1) &= A_{raw}^{w_1(\phi_h, \phi_s)} \sin(\Phi_1) + A_{raw}^{w_6(\phi_h, \phi_s)} \cos(\Phi_1) \\ W_2(\Phi_2) &= A_{raw}^{w_2(\phi_h, \phi_s)} \sin(\Phi_2) \\ W_3(\Phi_3) &= A_{raw}^{w_3(\phi_h, \phi_s)} \sin(\Phi_3) \\ W_4(\Phi_4) &= A_{raw}^{w_4(\phi_h, \phi_s)} \sin(\Phi_4) + A_{raw}^{w_7(\phi_h, \phi_s)} \cos(\Phi_4) \\ W_5(\Phi_5) &= A_{raw}^{w_5(\phi_h, \phi_s)} \sin(\Phi_5) + A_{raw}^{w_8(\phi_h, \phi_s)} \cos(\Phi_5) \end{aligned}$$

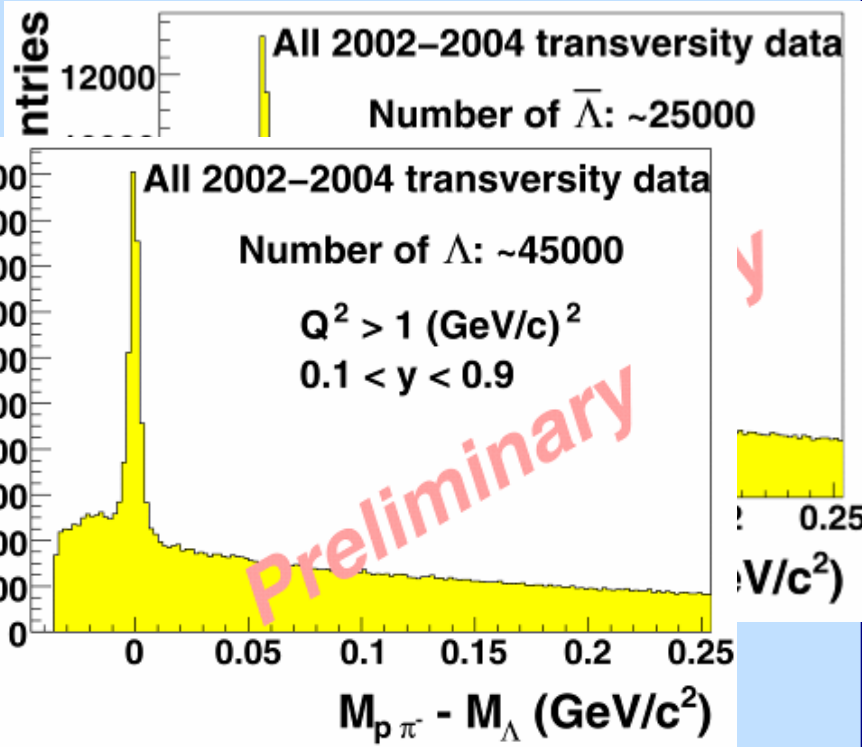
# transversity from $\Lambda$ polarimetry

$$\begin{aligned}
 P_{T,exp}^{\Lambda} &= \frac{d\sigma^{\mu N^{\uparrow} \rightarrow \mu' \Lambda^{\uparrow} X} - d\sigma^{\mu N^{\downarrow} \rightarrow \mu' \Lambda^{\uparrow} X}}{d\sigma^{\mu N^{\uparrow} \rightarrow \mu' \Lambda^{\uparrow} X} + d\sigma^{\mu N^{\downarrow} \rightarrow \mu' \Lambda^{\uparrow} X}} \\
 &= f P_N D(y) \frac{\sum_q e_q^2 \Delta_T q(x) \Delta_T D_{\Lambda/q}(z)}{\sum_q e_q^2 q(x) D_{\Lambda/q}(z)}
 \end{aligned}$$

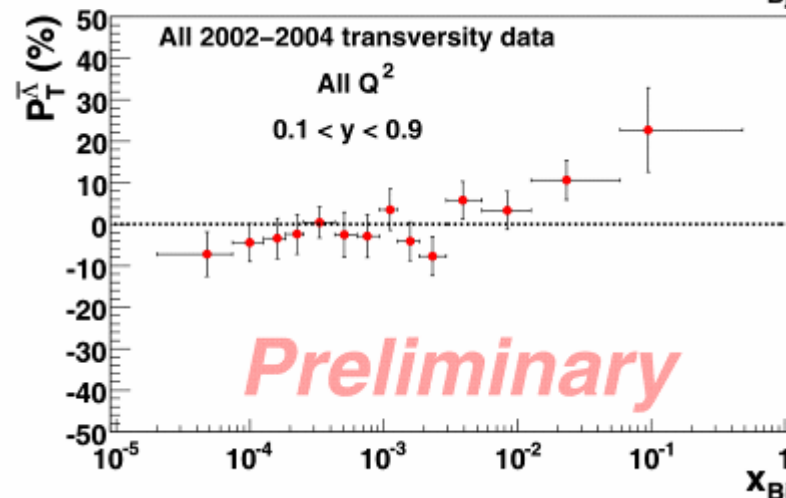
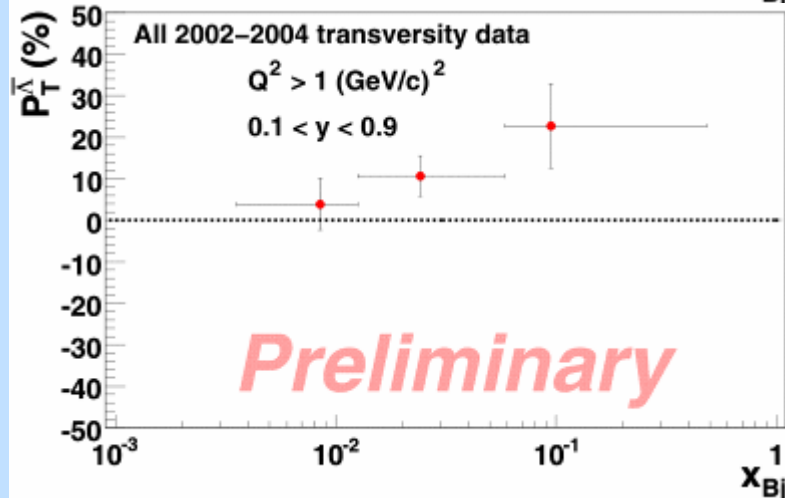
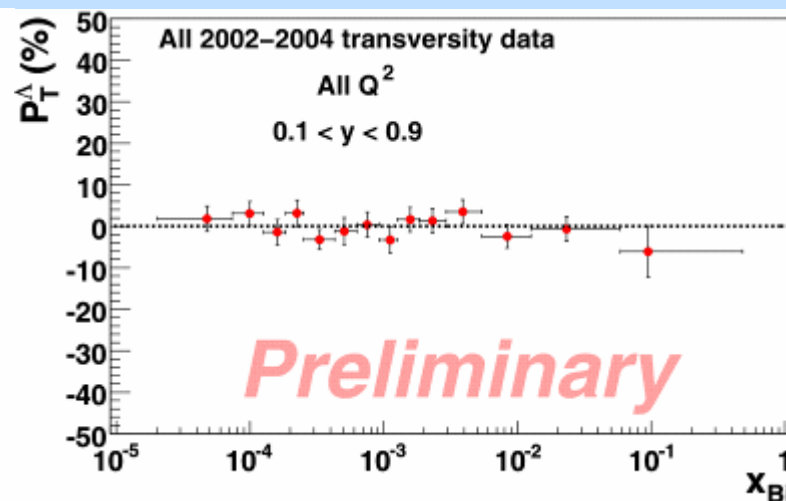
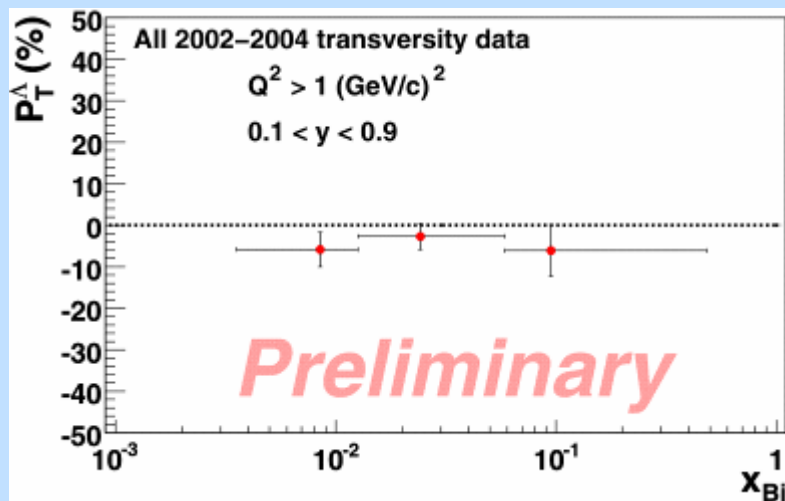
$\Lambda$  polarization axis



x 10 all  $Q^2$



# $\Lambda$ polarimetry



systematic errors not larger than statistical errors