



Exclusive Electroproduction of π^0 Mesons

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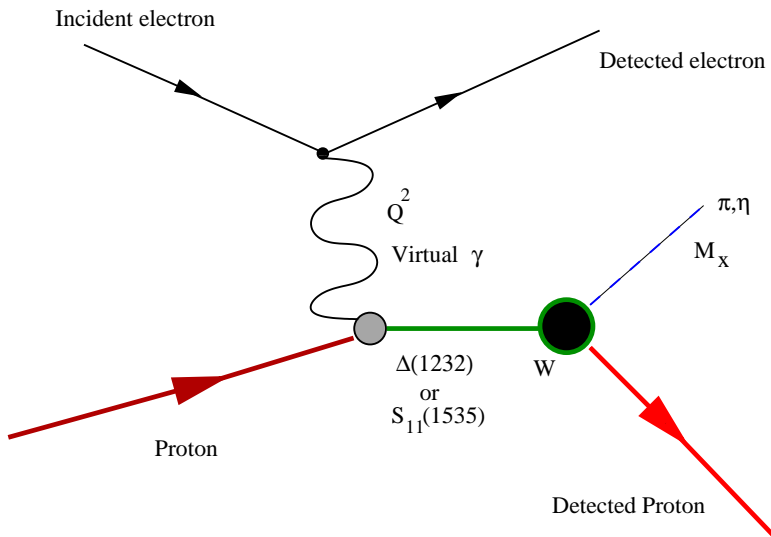
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Overview

- Physical motivations
- Brief experimental overview
- Backgrounds, corrections, and data exclusions
- Differential cross sections

Exciting the Δ in E01-002



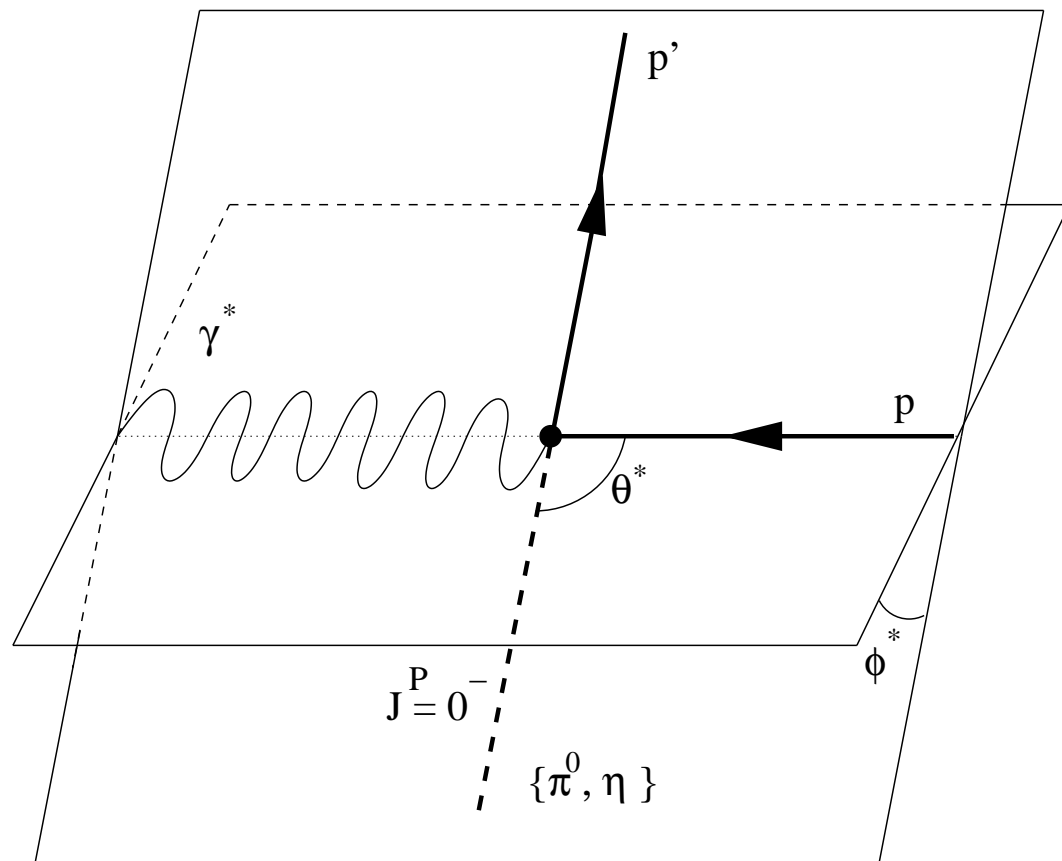
- Kinematic variables

$$Q^2 = 6.3 \text{ GeV}^2; 7.4 \text{ GeV}^2$$

$$W = \text{Elastic} \rightarrow 1.8 \text{ GeV}$$

- Single photon exchange
- Full angular coverage in COM

Center of Mass Kinematics



Center of Mass Cross Section

- The cross section can be written in the center of mass :

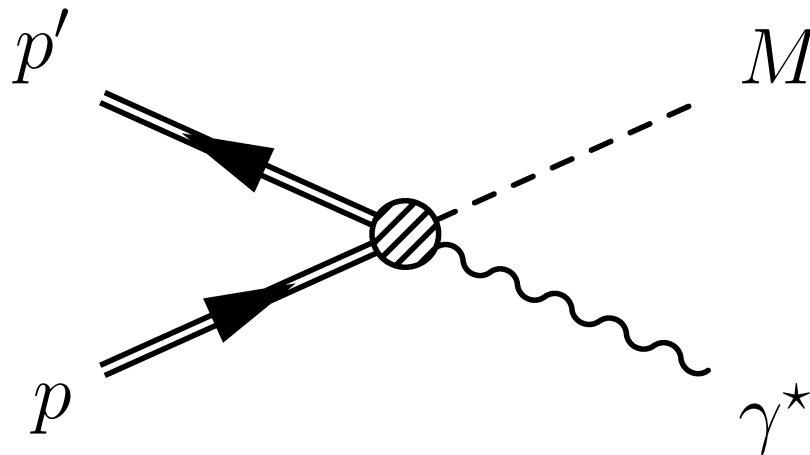
$$\frac{d\sigma}{dQ^2 dW d\Omega_{cm}} = \Gamma \frac{d\sigma_\nu}{d\Omega_{cm}}$$

- The virtual photon flux factor Γ relates to kinematic quantities in center of mass

$$\Gamma = \frac{\alpha}{2\pi} \frac{E'}{E} \frac{W^2 - m_p^2}{2m_p Q^2} \frac{1}{1 - \epsilon}$$
$$\epsilon = \frac{1}{1 + 2 \tan^2 \left(\frac{\theta_e}{2} \right) \frac{|\vec{q}|^2}{Q^2}}$$

- In *virtual* photoproduction one may have
 - virtual photons having "mass"
 - virtual photons having longitudinal polarization states

Virtual Photoproduction



- Virtual photoproduction amplitude can be written

$$\mathcal{M}_{ph} = \epsilon_{\mu}^{\lambda} \langle h_f | J_{had}^{\mu} | h_i \rangle$$

$$\frac{d\sigma}{d\Omega_{cm}} = \sigma_T + \epsilon\sigma_L + \epsilon\sigma_{TT} \cos 2\phi^* + \sqrt{\frac{\epsilon(1+\epsilon)}{2}} \sigma_{LT} \cos \phi^*$$

Current Decomposition

- Current matrix element from virtual photoproduction amplitude can be decomposed

$$\langle p' M | J_{had}^\mu | p \rangle = i \bar{u}_f(p'_p) \gamma_5 \left[\gamma^\mu \not{q} B_1 + (p_p + p'_p)^\mu B_2 + p_p^\mu B_3 + q^\mu B_4 + \gamma^\mu B_5 \right. \\ \left. + (p_p + p'_p)^\mu \not{q} B_6 + q^\mu \not{q} B_7 + p_p^\mu \not{q} B_8 \right] u_i(p_p)$$

- Scalar functions B_i can be expressed in terms of multipole expansions

$$B_1 \propto \sum_{l \geq 0} \left[(l M_{l+} + E_{l+}) P'_{l+1} + ((l+1) M_{l-} + E_{l-}) P'_{l-1} \right]$$

- Quantum numbers l and \pm specify $J = |l \pm \frac{1}{2}|$ in final state

Multipoles from Atomic Physics

- Wave function of incident photon can be decomposed as vector spherical harmonics

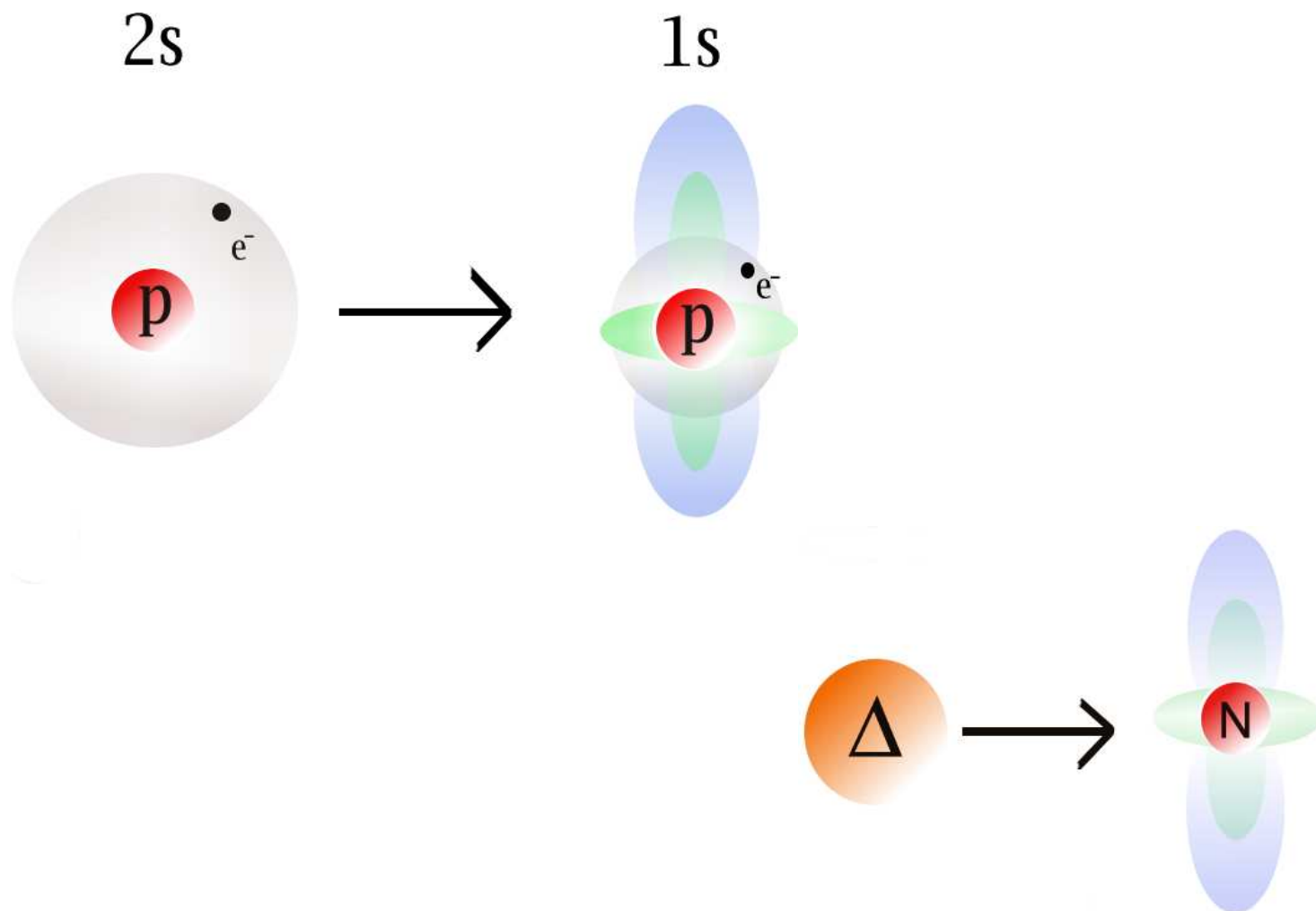
$$\mathbf{Y}_{\bar{l}LM} = \sum_{\nu} C(1\lambda, \bar{l}\nu | LM) \hat{e}_{\lambda} Y_{\bar{l}\nu}$$

- EL (electric) and CL (coulomb) type radiation are made of parity even combinations and ML (magnetic) is parity odd
- Since angular momentum is conserved one has $J = |l \pm \frac{1}{2}| = |L \pm \frac{1}{2}|$
- Parity arguments can then give the following relations

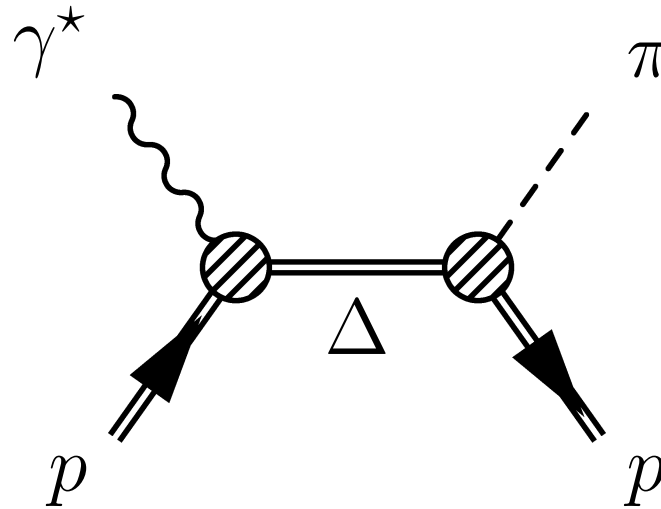
$$EL, CL : |L - l| = 1$$

$$ML : L = l$$

The Atomic Analog



Resonance Production (Δ)



- Restriction to Δ decreases the number of independent functions to three
- Functions can be represented:

$$G^{\pm,0} = \frac{1}{2M} \langle (\Delta), \lambda_{res} | \epsilon_{\mu}^{\pm,0} J_{had}^{\mu} | P, \lambda_p = \pm \frac{1}{2} \rangle$$

Multipole Definition

- Considering only Δ production reduces the number of multipoles

$$A_{\frac{1}{2}} \propto G^+$$

$$A_{\frac{3}{2}} \propto G^-$$

$$S_{\frac{1}{2}} \propto G^0$$

- Also These can be related (through $E_{l\pm}$ and $M_{l\pm}$) to E2 and M1:

$$A_{\frac{1}{2}} = -\frac{1}{2}(M1 + 3E2)$$

$$A_{\frac{3}{2}} = \frac{\sqrt{3}}{2}(E2 - M1)$$

$$S_{\frac{1}{2}} = -C2$$

Measurement and Prediction

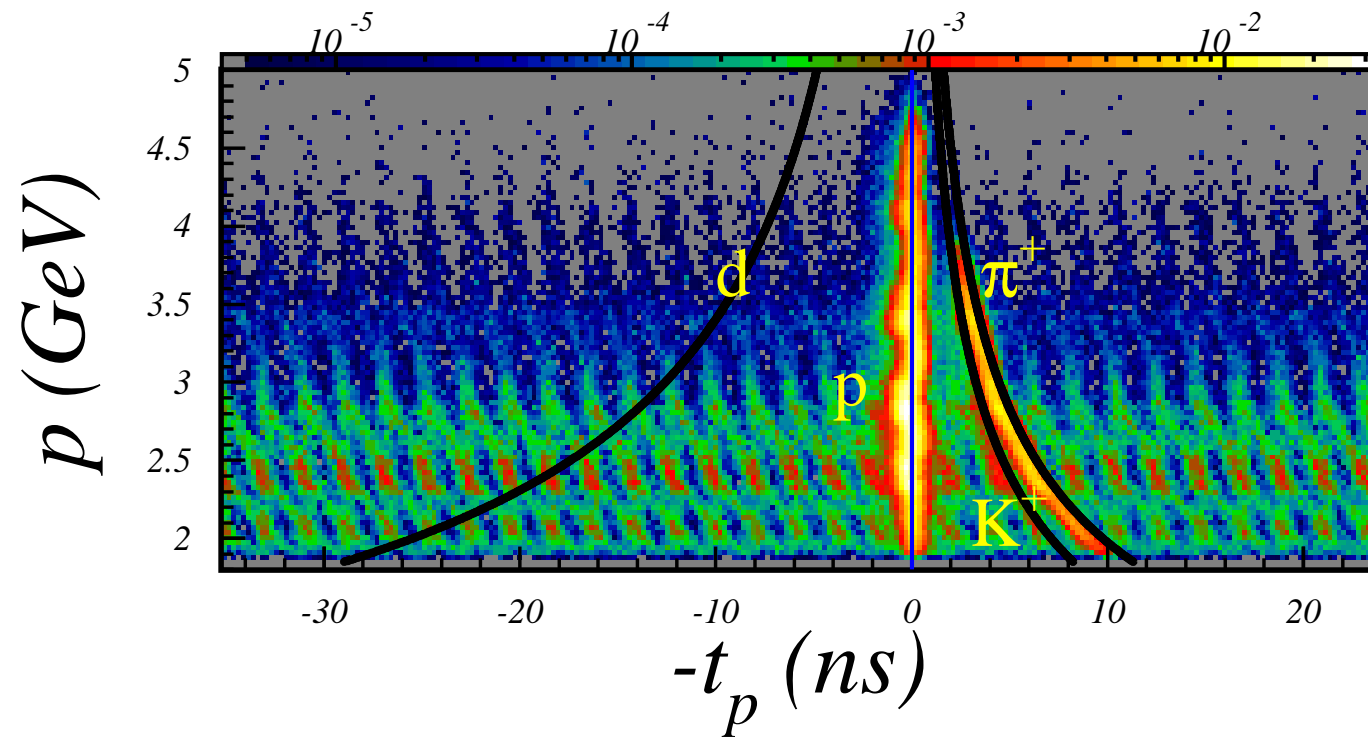
Measured Quantities

- E_{1+} , M_{1+} and S_{1+} extracted for Δ
- Above used to infer information about $\frac{E2}{M1}$

Predictions

- Perturbative QCD predicts that $\frac{E2}{M1} \rightarrow 1$ as Q^2 becomes large
- Constituent quark model predicts $M1$ dominance because Δ transition is viewed as a simple spin flip excitation

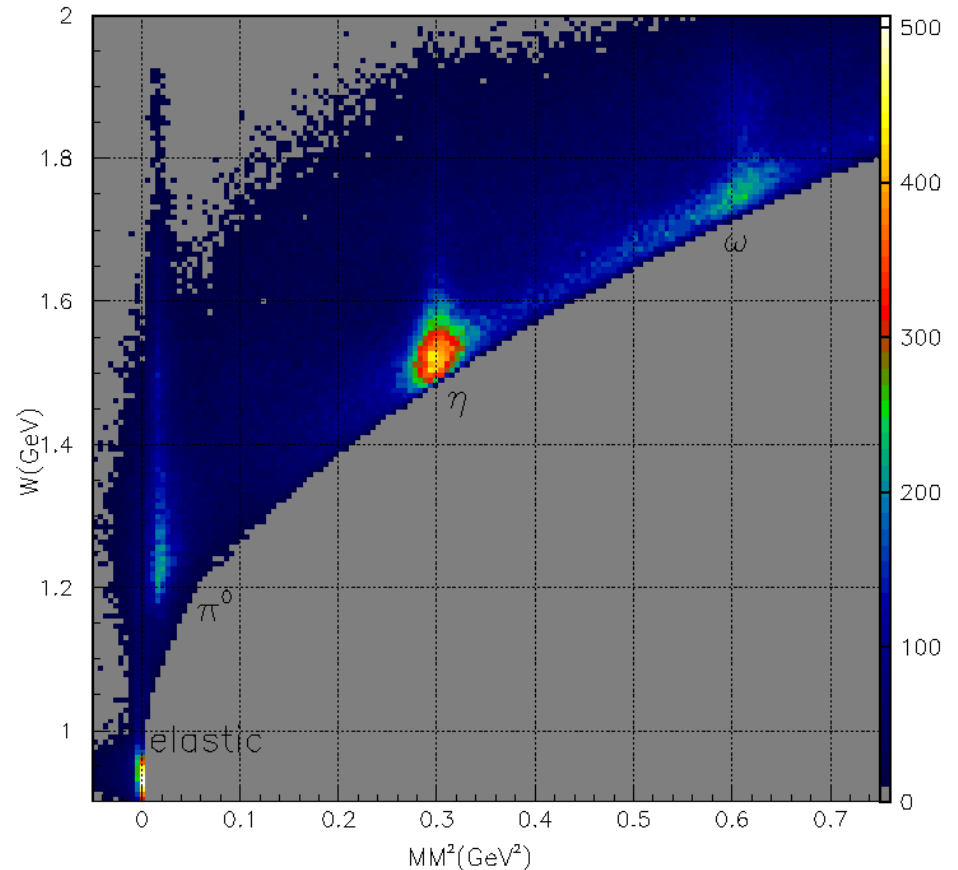
Particle Selection



- Clean separation of p and π^+ events
- Enough π^+ events for $\gamma^*p \rightarrow n\pi^+$

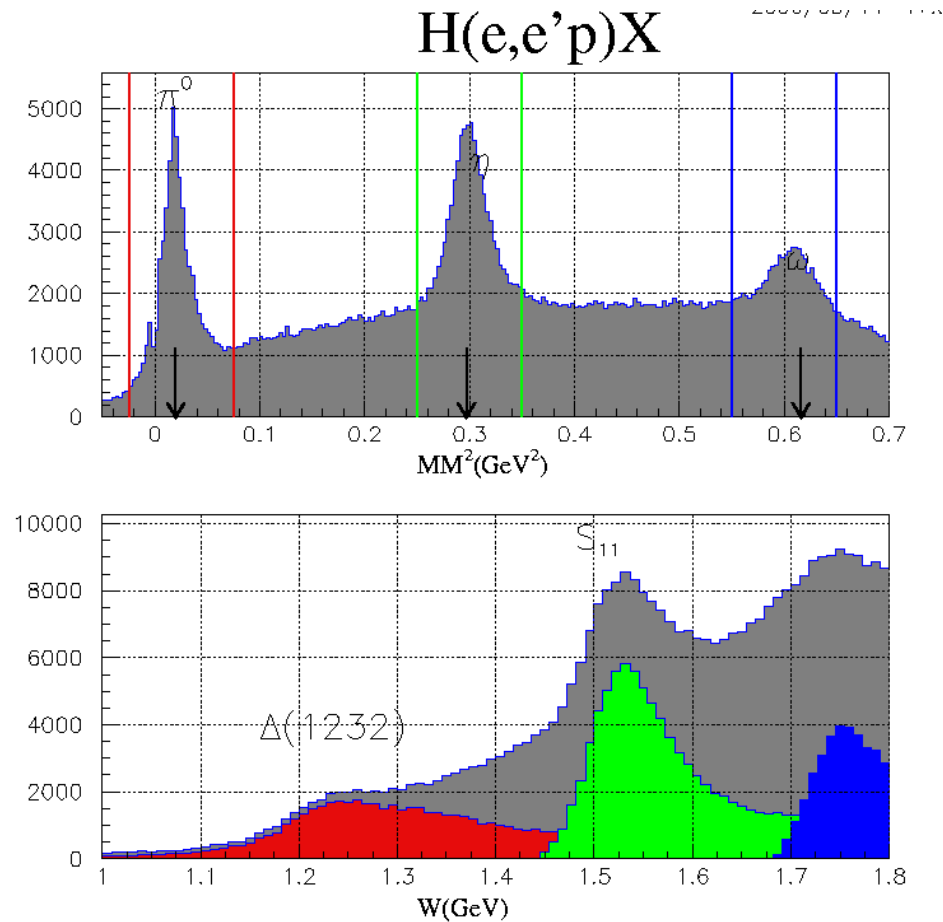
Baryon Resonances

- The $\Delta(1232)$ and S_{11} resonances are clearly correlated with the π^0 and η M_x^2 peaks
- The elastic events clearly come from lower W with some overlap into a higher W region due to pre or post radiation
- The ω meson comes from the largest W region for the experiment

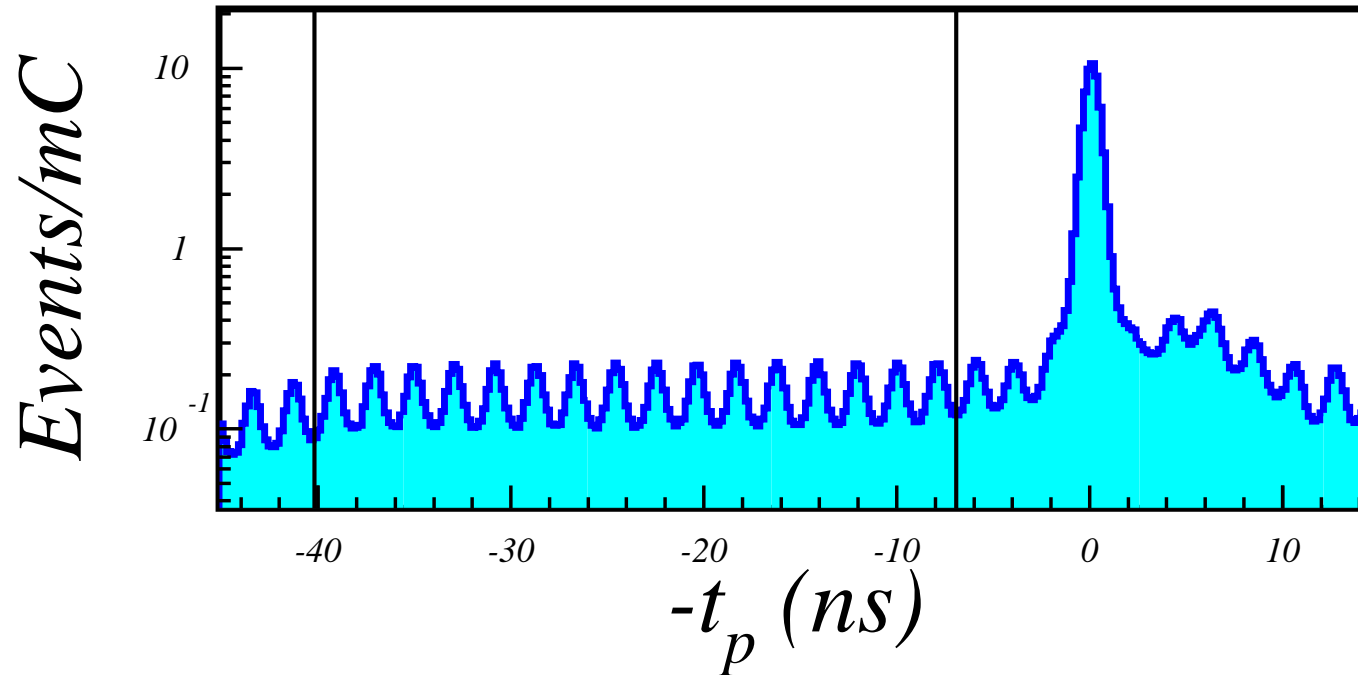


Exclusive Studies

- The M_x^2 peaks can be used to constrain the reaction and/or baryon resonance
- The M_x^2 resolution for the π^0 allows detailed study of the reaction ${}^1H(e, e'p)\pi^0$
- Exclusive cross sections and amplitudes will be compared to models and previous data

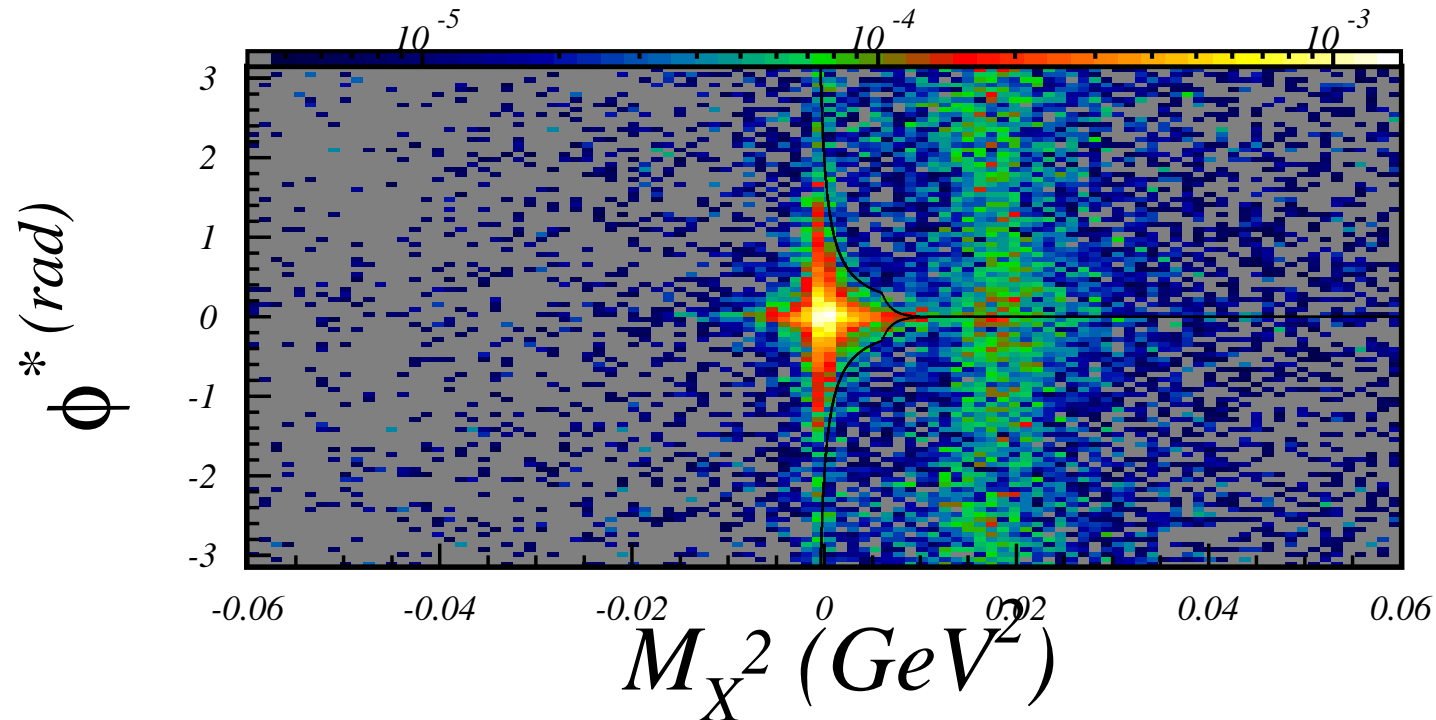


Accidental Corrections



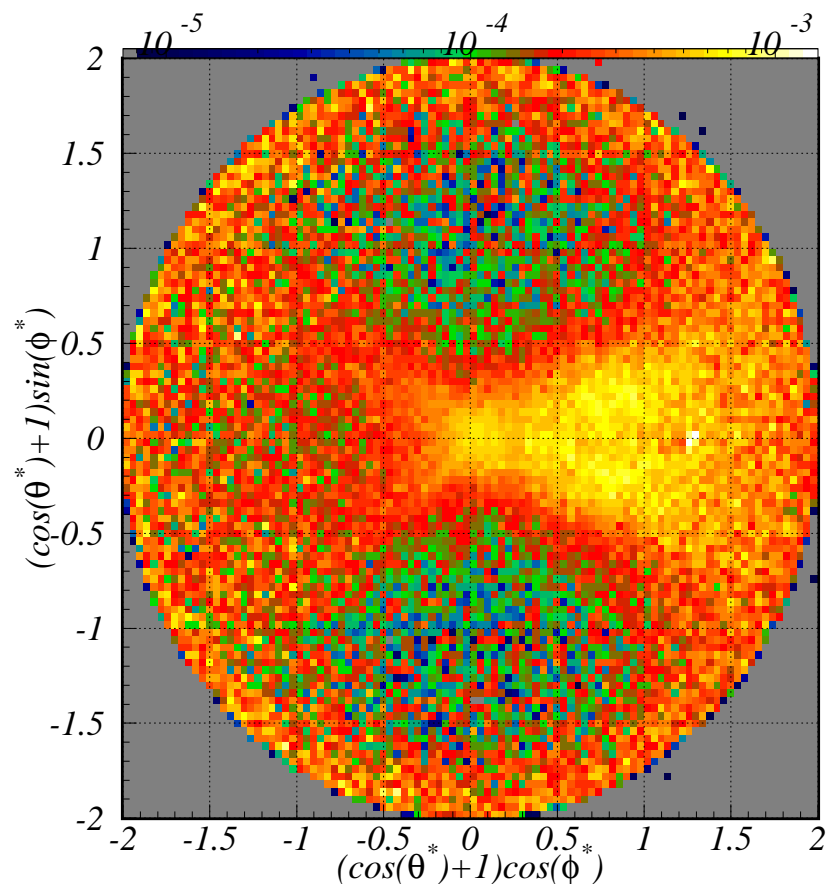
- Use the beam structure away from coincidence proton peak
- Extract angular distribution by above cut and subtract
- Low momentum hadrons dominate beam structure

Elastic Radiative



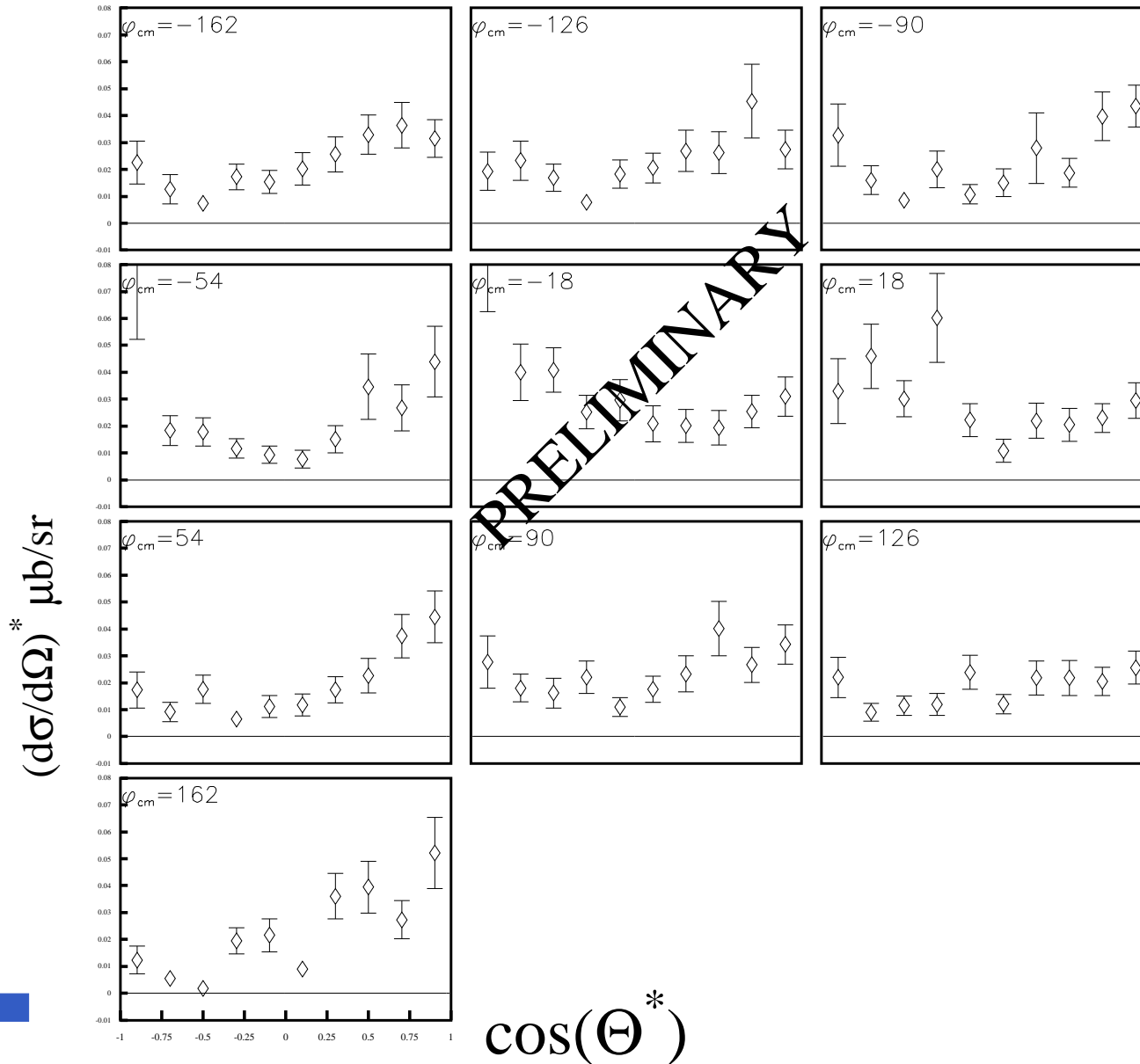
- Above example for $-1.0 \leq \cos \theta^* < -0.6$
- For lowest $\cos \theta^*$ elastic radiative concentrated
- Use same form for other $\cos \theta^*$ bins

Full Angular Coverage

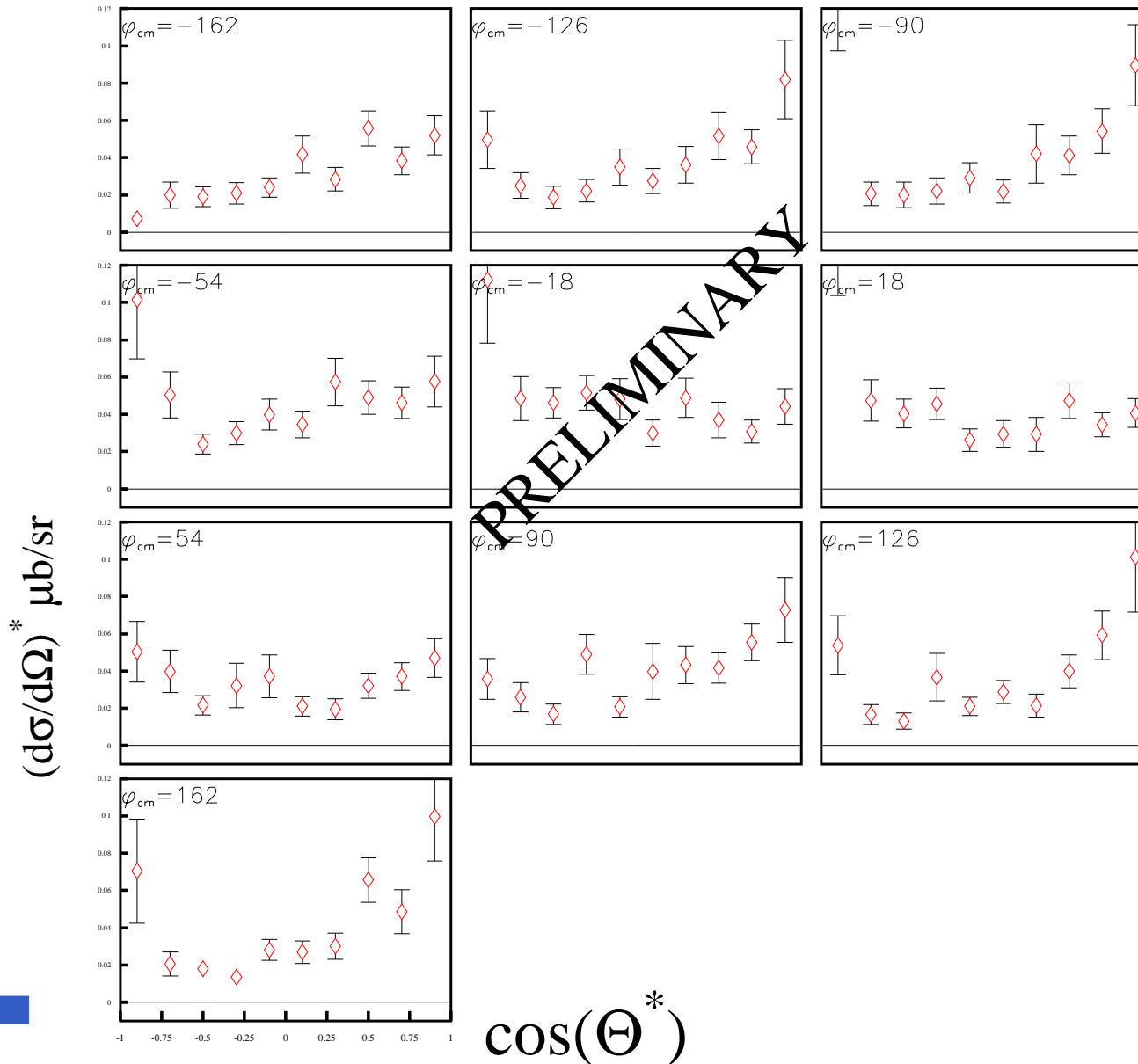


- Uniform coverage in the center of mass angles
- All W plotted
- Elastic radiative cut removes bad $\phi^* = 0$ data

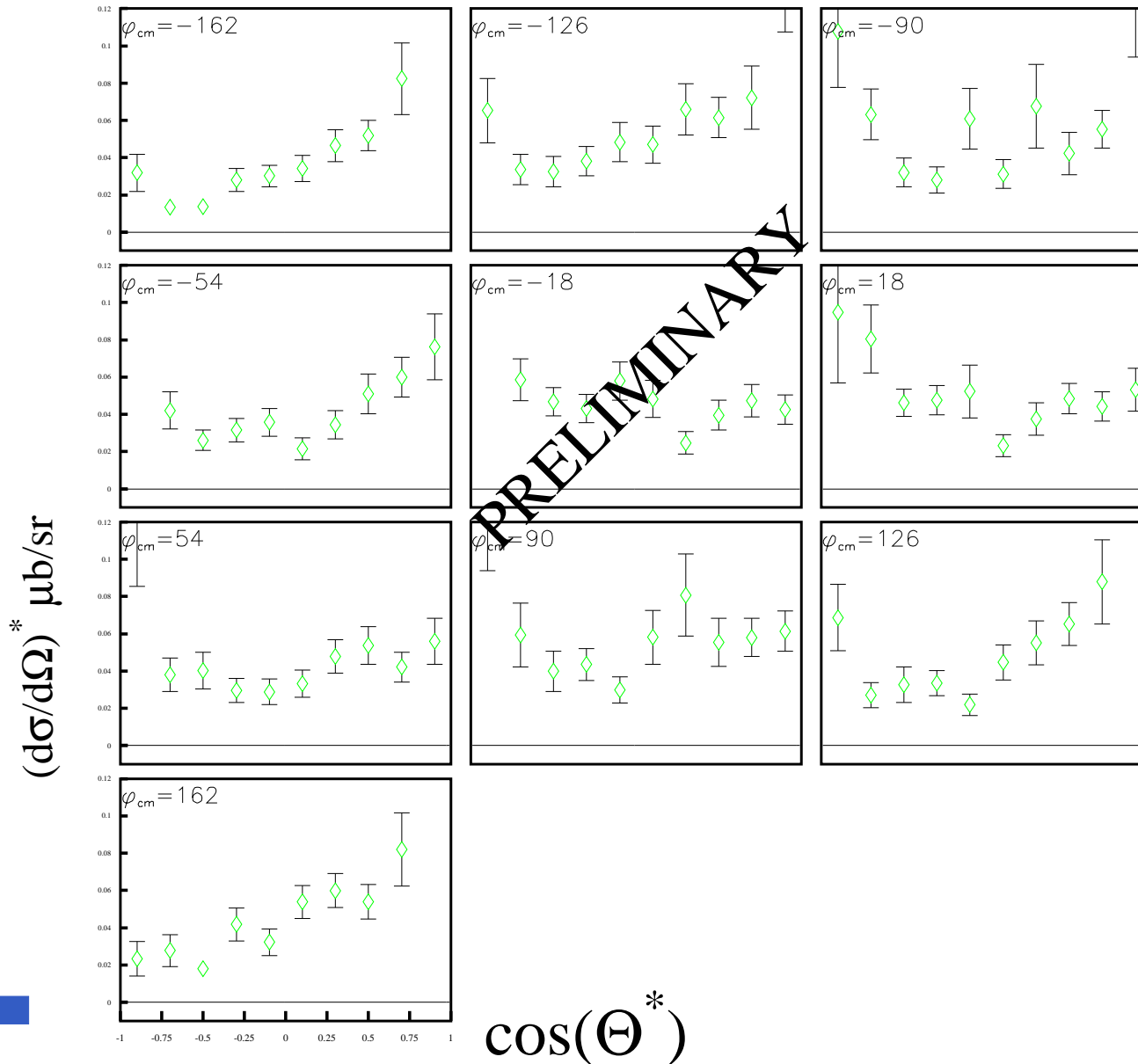
$W=1.172$ GeV ; 30 MeV Bin ; $Q^2=6.3$



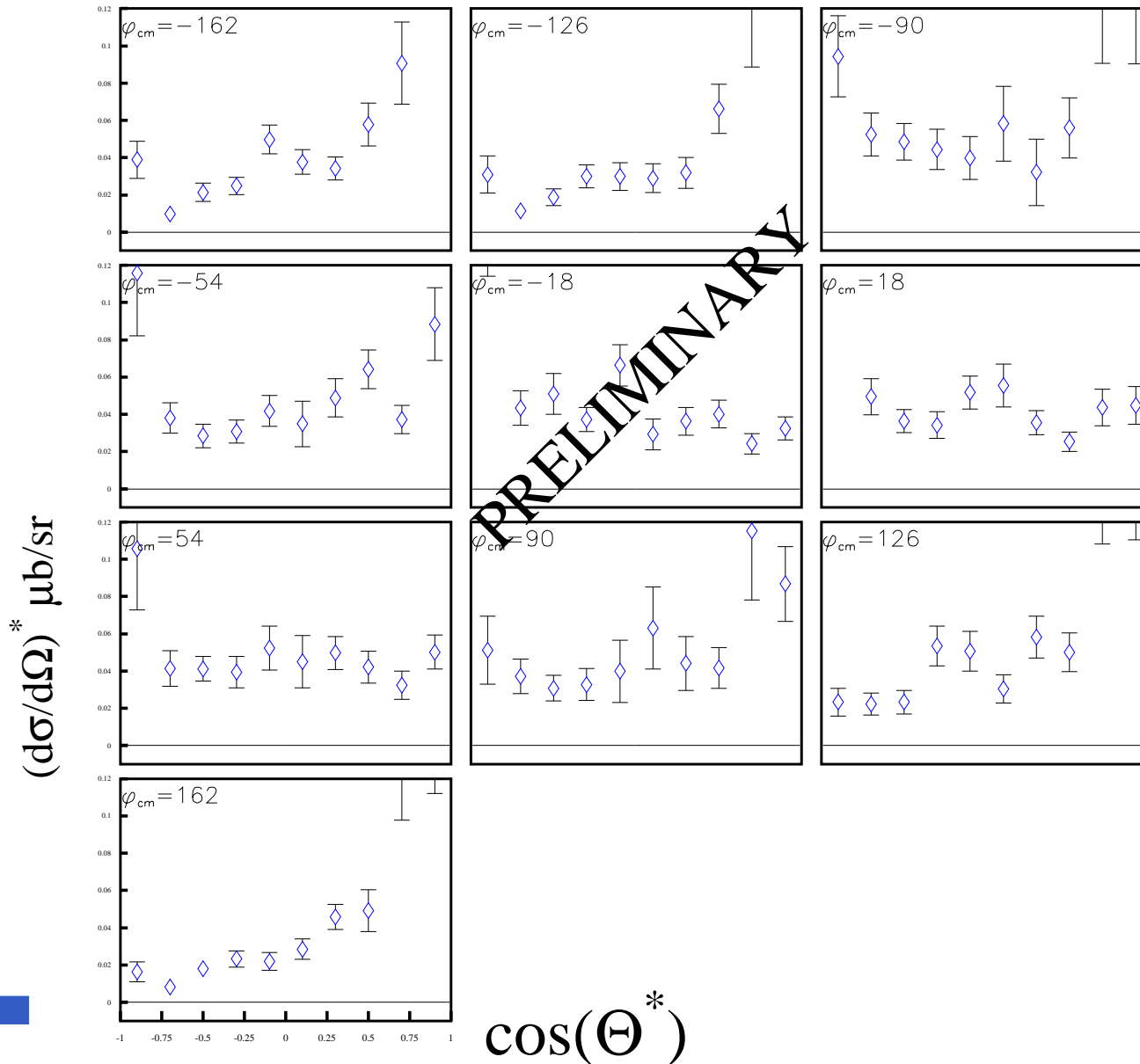
$W=1.202$ GeV ; 30 MeV Bin ; $Q^2=6.3$



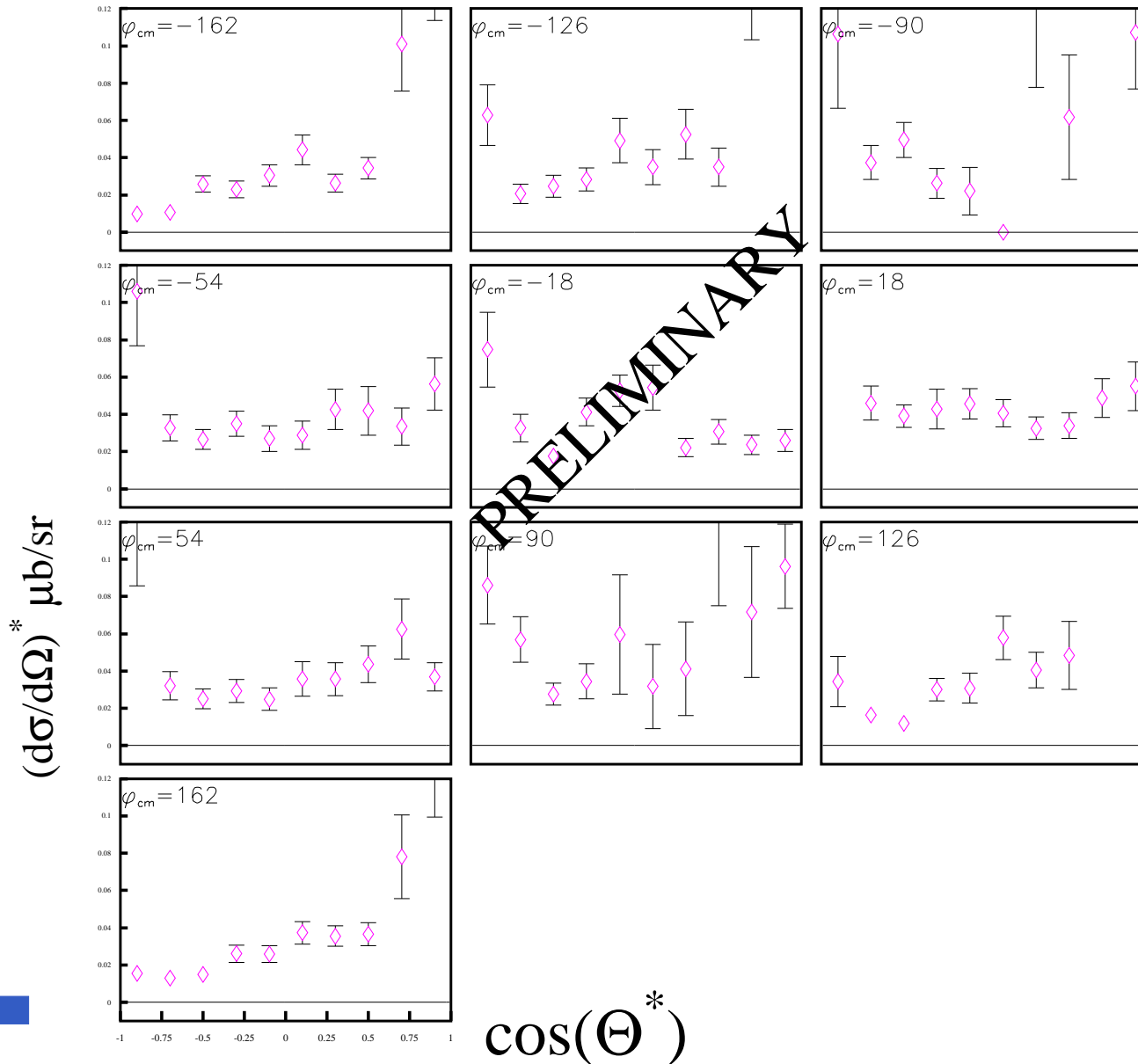
$W=1.232$ GeV ; 30 MeV Bin ; $Q^2=6.3$



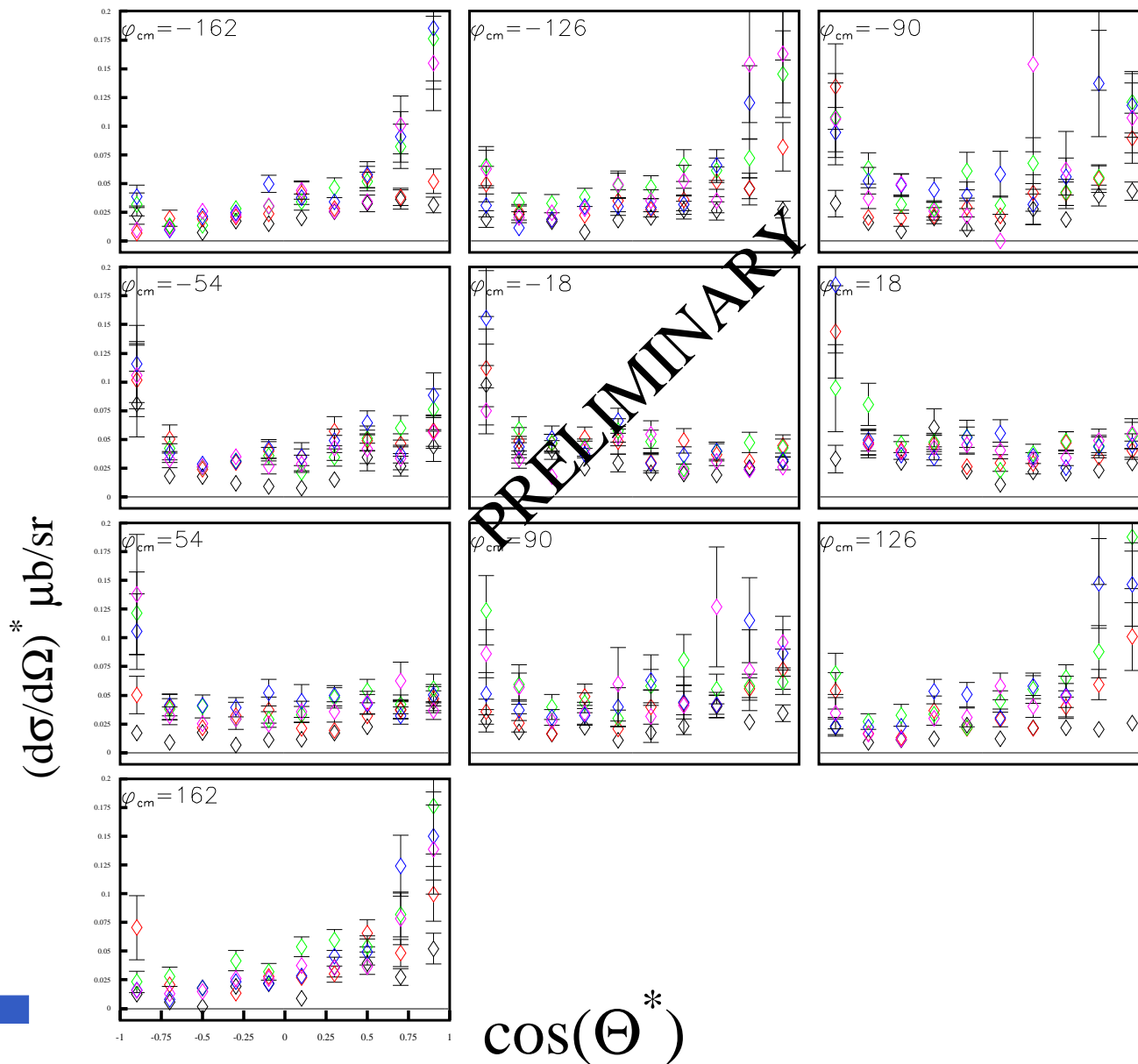
$W=1.262$ GeV ; 30 MeV Bin ; $Q^2=6.3$



$W=1.292$ GeV ; 30 MeV Bin ; $Q^2=6.3$



Exclusive Cross Section



Summary and Progress

- Beam energy of 5.5GeV with two Q^2 settings
 - Measure the cross sections for $^1H(e, e'p)X$, $X = \{\pi^0, \eta, \omega\}$
 - X identified by missing mass, M_x
 - Q^2 of 6.3 and 7.7GeV for Δ resonance
 - Varied proton arm angle and momentum to cover wide range of θ^* and ϕ^* bins for W up to 2GeV
- Physics
 - Showed angular distributions for low Q^2
 - Distributions will be fit though it is not clear that the pure Δ multipoles will show
 - Systematic uncertainties need to be brought under control