

# Regge in exclusive electro-production (e.g. DVCS, meson production)

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Indiana University



Adam Szczepaniak, J.T. Londergan, Phys.Lett. B643, 17 (2006)

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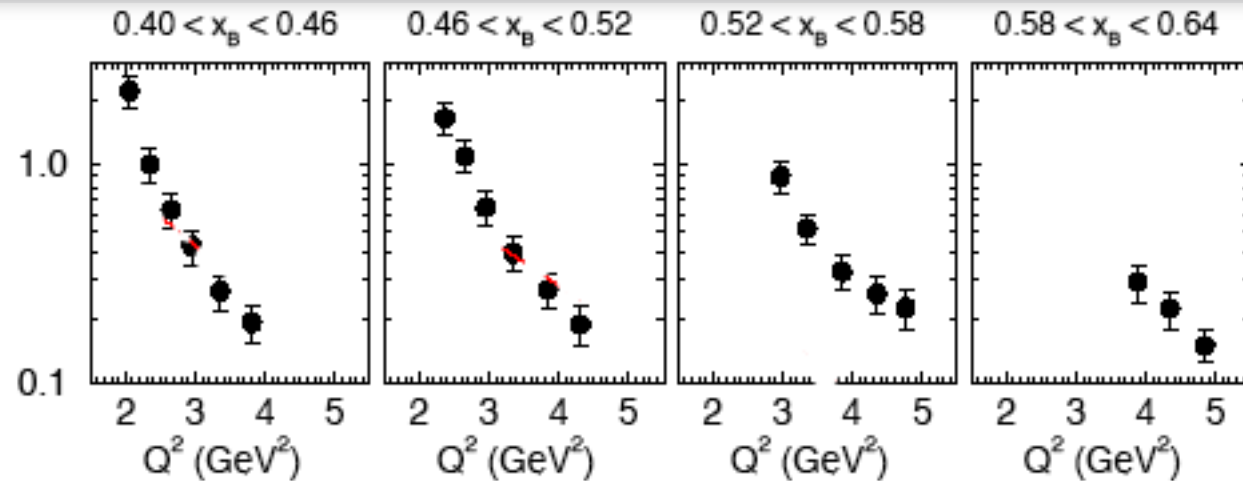
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- Why ?

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- Why ?
- there exist “regge motivated” parametrizations of gpd’s, but are they properly “motivated”

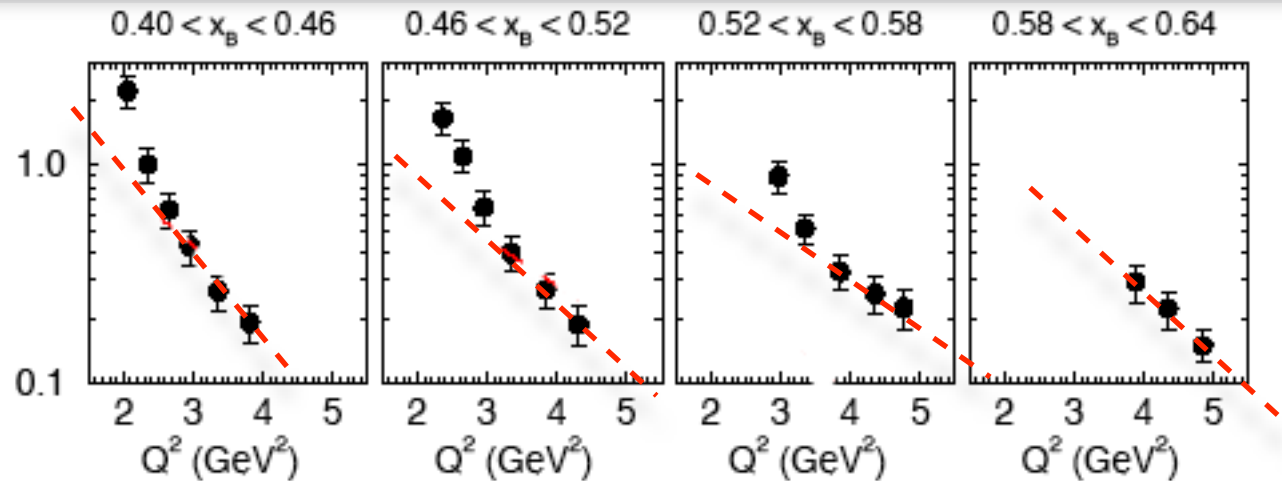


**Fig. 10.** (Color online) Reduced cross sections  $\gamma^* p \rightarrow \omega p$  as a function of  $Q^2$  for different bins in  $x_B$ , in units of  $\mu b$ . Full circles: this work; open circles: Ref. [17]. The red cross and curves correspond to the JML model [29] discussed in sect. 5.

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scaling (collinear factorization)

$$n = 3$$



$$n \sim 1.7$$

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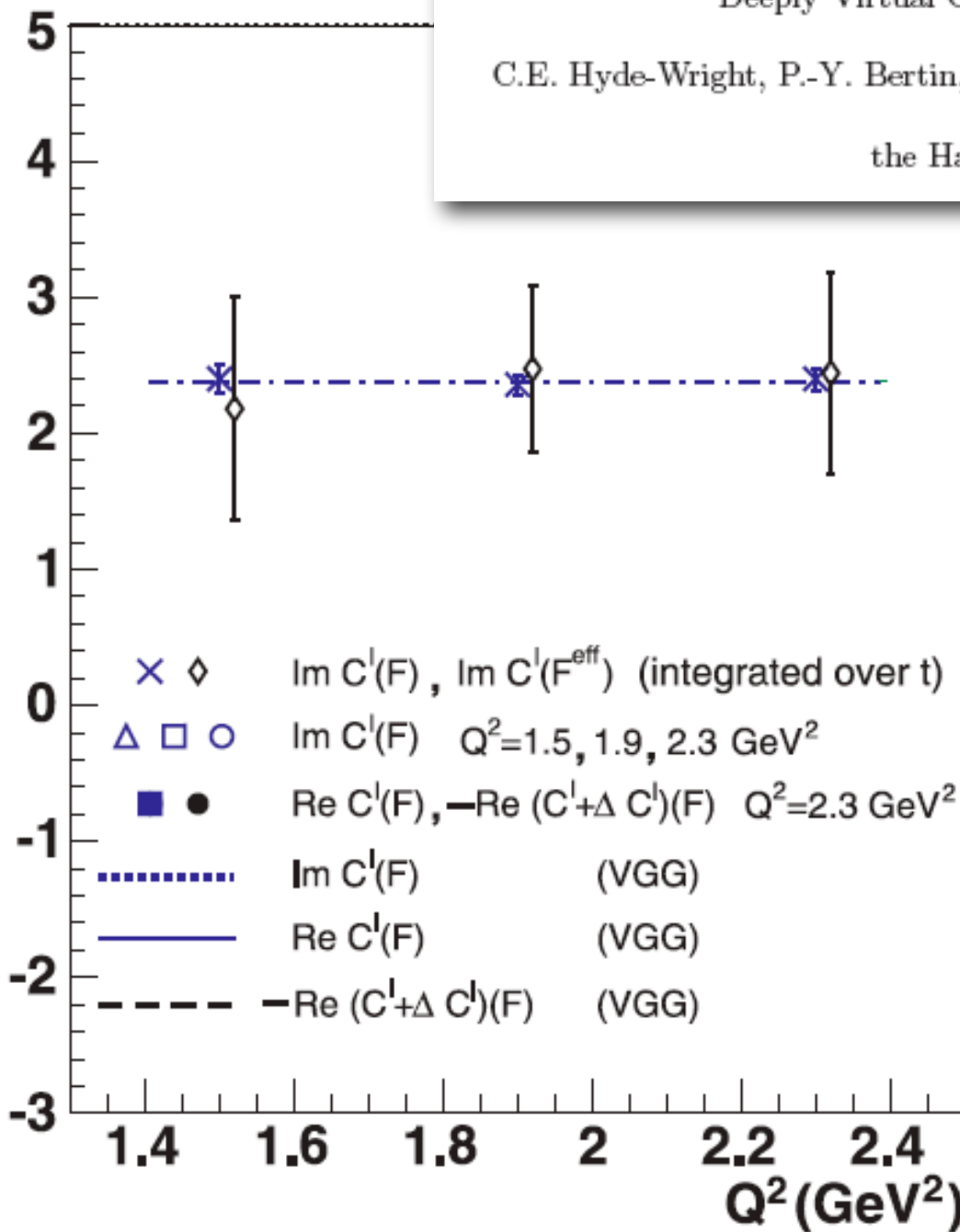
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### 3.7 E00-110

#### Deeply Virtual Compton Scattering at 6 GeV

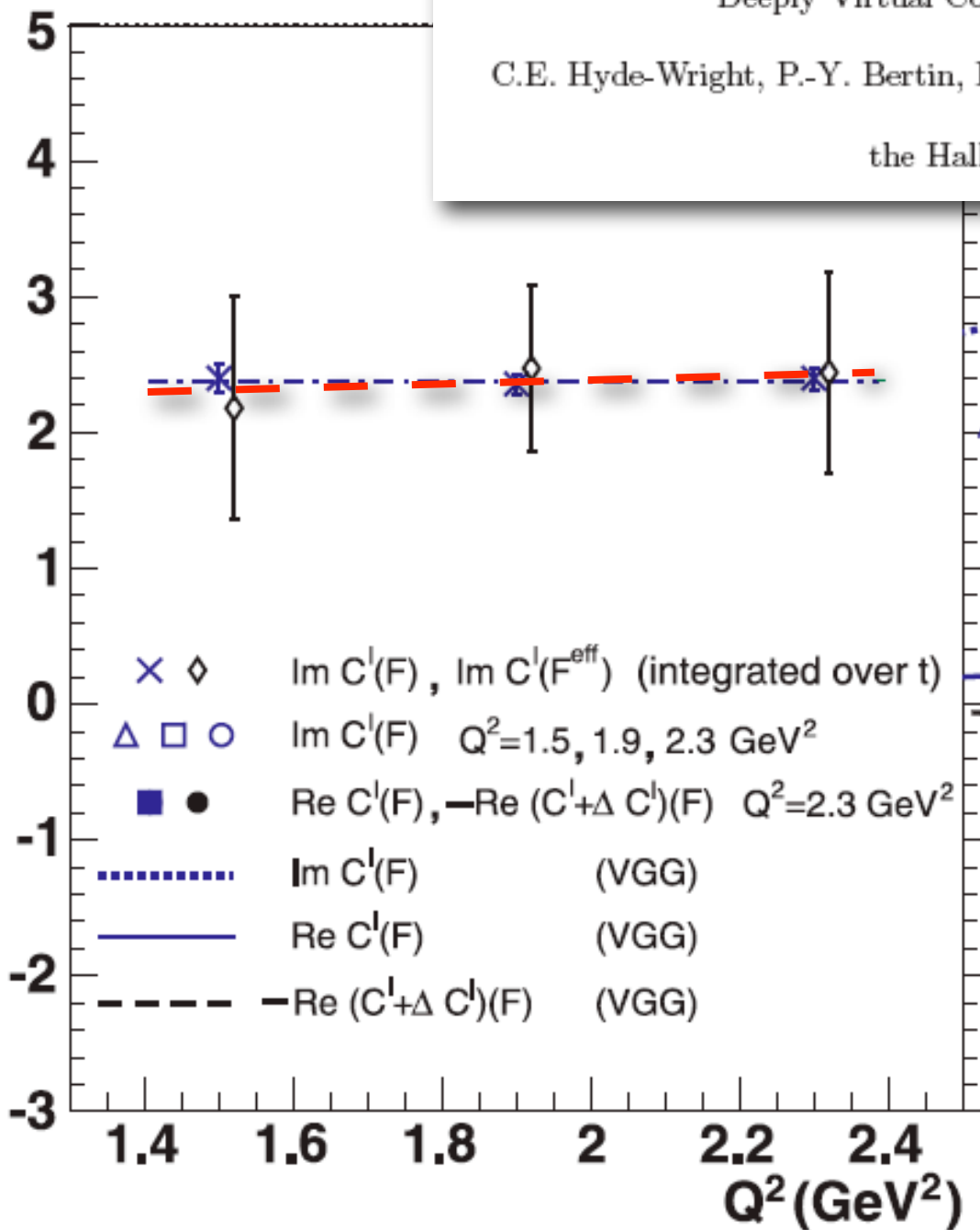
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$\cdots$

$$\propto (Q^2)^\alpha$$

$$\alpha = 0.25$$

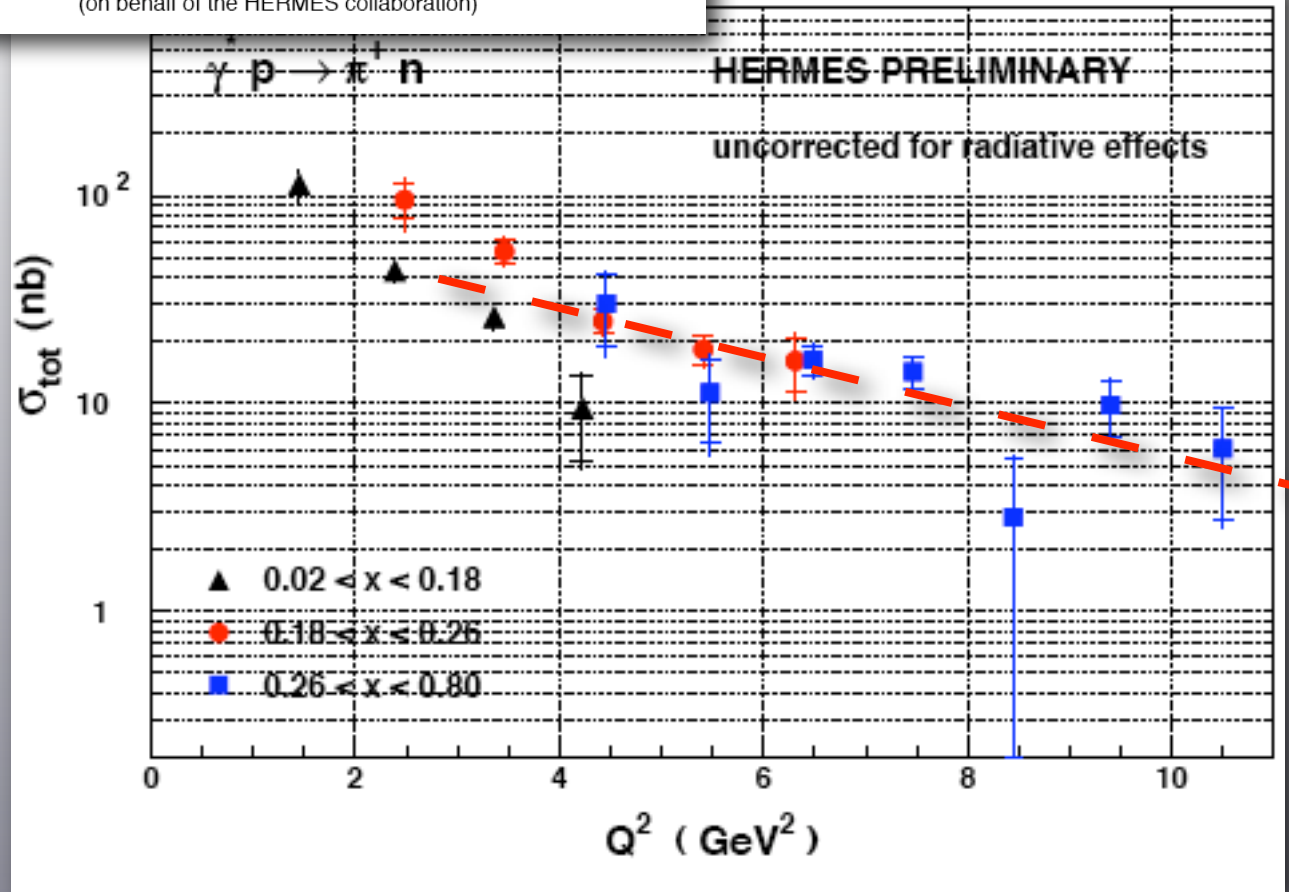


# Exclusive Meson Production at HERMES

Jefferson Lab, VA, USA, May 2007

Armine Rostomyan

(on behalf of the HERMES collaboration)



$$\sigma_{x=\text{fixed}} \sim \left( \frac{1}{Q^2} \right)^n \quad n = 3 \quad \longrightarrow \quad n = 1.6$$

- regge parametrizations motivated by what DIS structure functions

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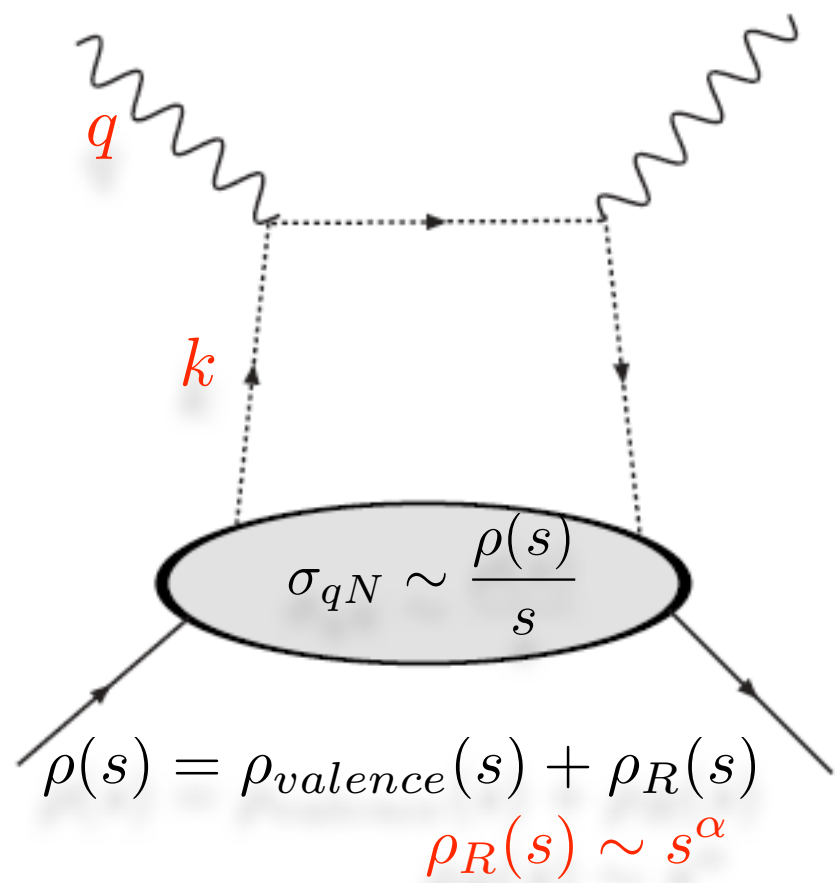
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- large masses are “problematic” so put regge at the level of parton-nucleon amplitude and see what happens to the handbag diagram

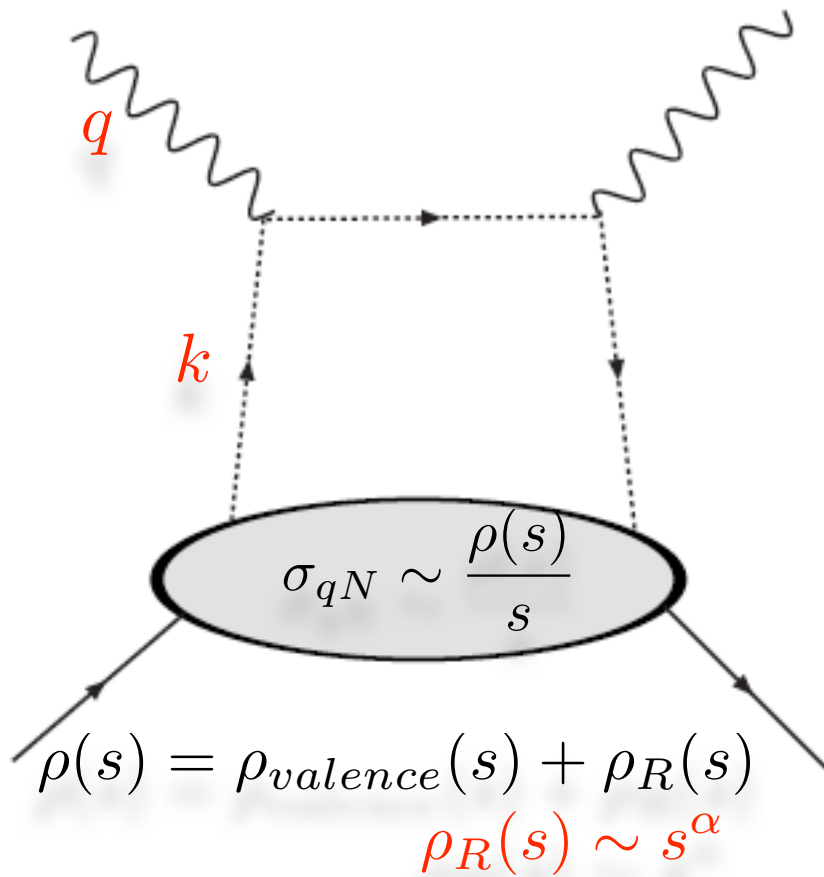




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S.J. Brodsky, F. Close, J.F. Gunion,  
Phys. Rev. 8, 3678 (1973)

$$f(x_{BJ}) = f_v(x_{BJ}) + \sum_{0 < \alpha < 1} x_{BJ}^{-\alpha}$$



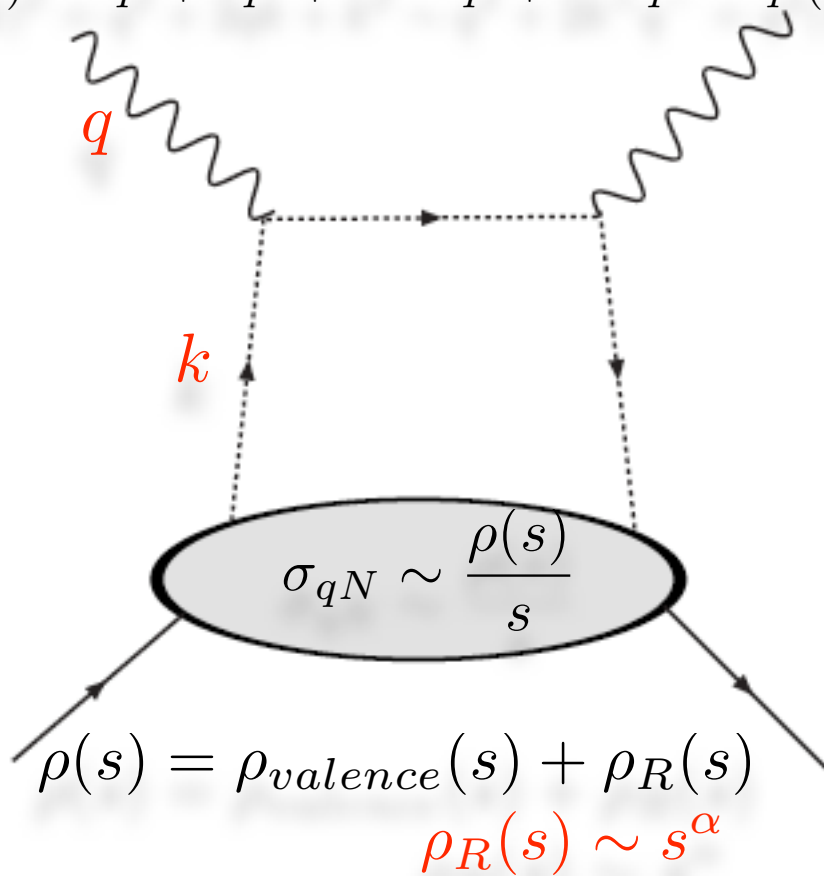
one can start from collinear factorization and get the same result

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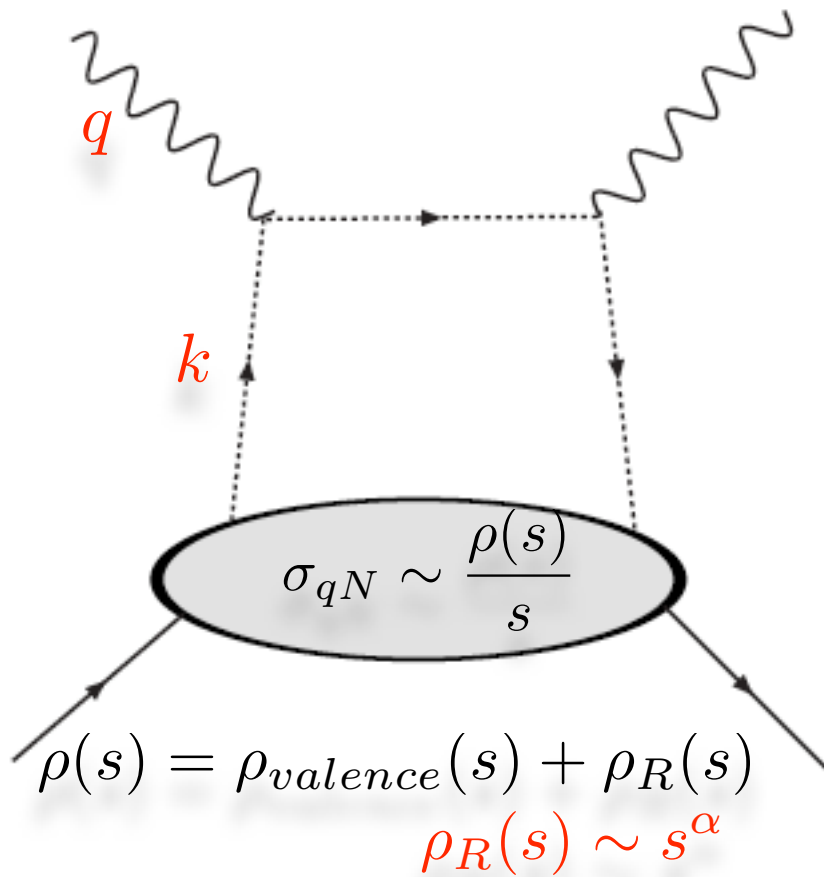
$$(q+k)^2 = q^2 + 2qk + k^2 \sim q^2 + 2k^+q^- = q^2(1 - x/x_B)$$



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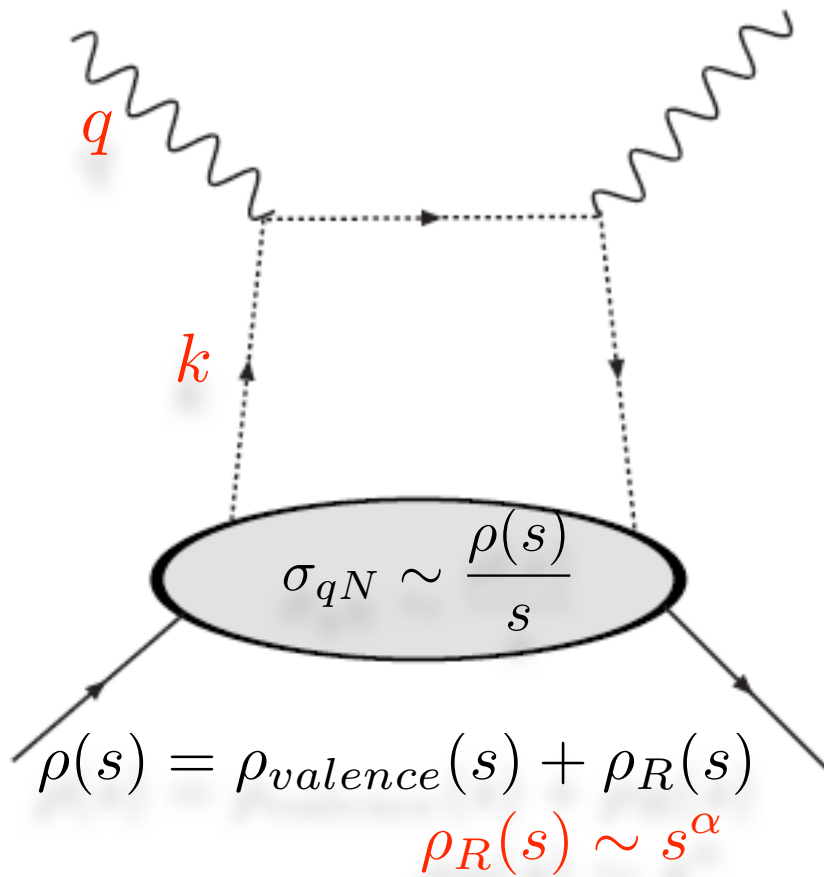
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$$f\left(\frac{\nu}{Q^2}\right) \rightarrow f\left(\frac{\nu}{\Lambda^2}\right)$$



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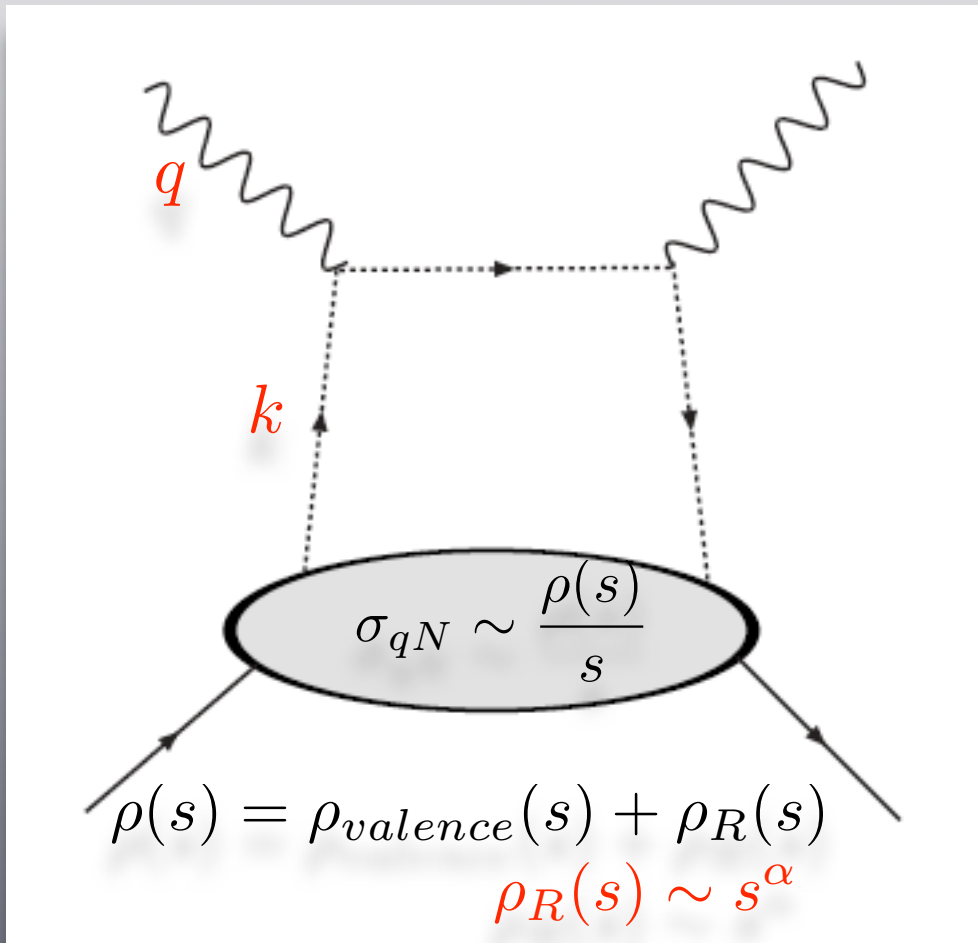
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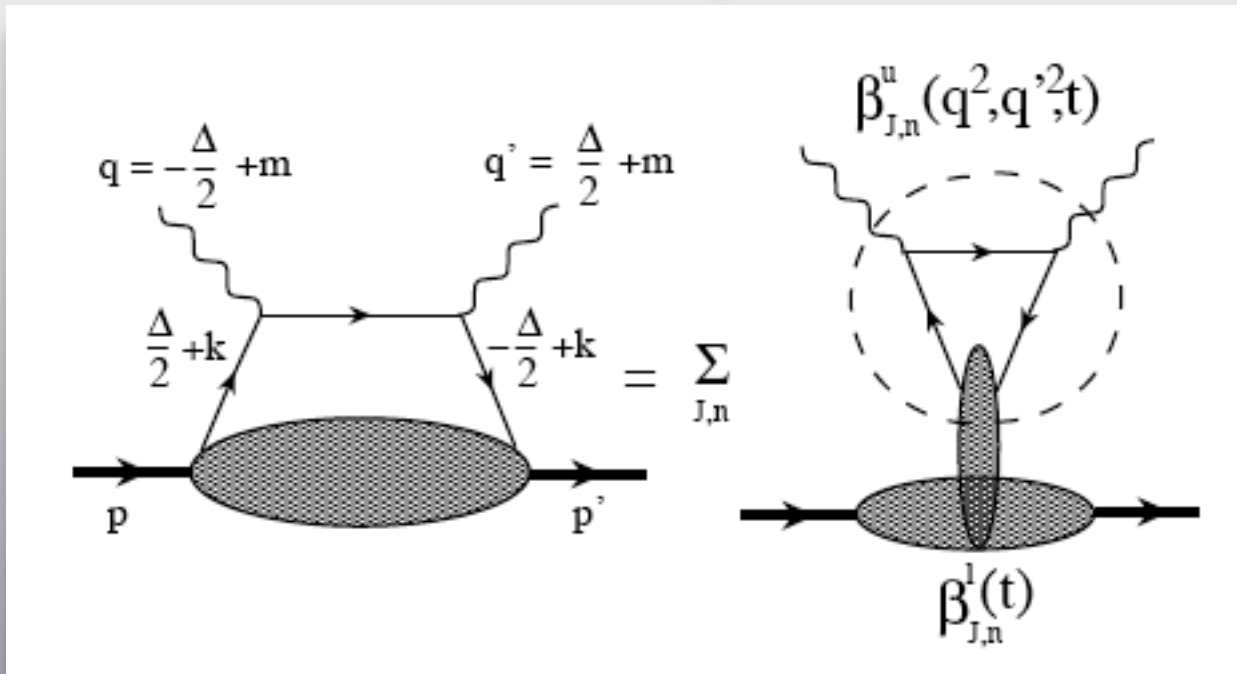
singular contribution to  
collinear factorization  
from subtraction terms



$$T_{partonN \rightarrow partonN}(s) = \int dm^2 \frac{\rho_R(m^2)}{s - m^2} - \int dm^2 \frac{\rho_R(m^2)}{-m^2}$$



# what about t-channel description ?



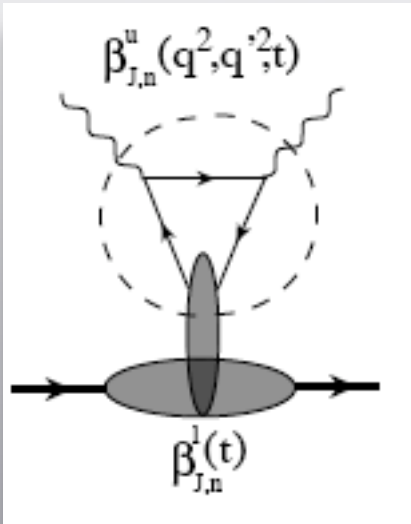
$$\int d^4z e^{-ikz} \langle p' | T \left[ \phi^\dagger \left( \frac{z}{2} \right) \phi \left( -\frac{z}{2} \right) \right] | p \rangle = \frac{\beta_{Jn}^u(t)}{t - M_{Jn}^2} \Phi_{Jn}(p - p', k)$$

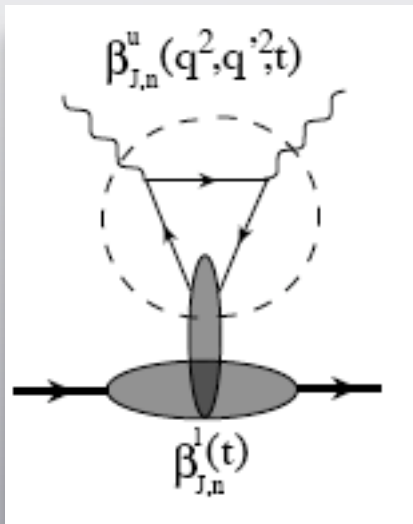
$$\sum_{\lambda=-J}^J [k^{\nu_1} \dots k^{\nu_J} \epsilon_{\nu_1 \dots \nu_J}^\lambda(p' - p)]^* \left[ \frac{(p' + p)^{\mu_1}}{2} \dots \frac{(p' + p)^{\mu_J}}{2} \epsilon_{\mu_1 \dots \mu_J}^\lambda(p' - p) \right]$$

$$a(J, t) \sim \frac{\beta(t)}{J - \alpha(t)} \left( \frac{1}{Q^2} \right)^J \left( \frac{Q^2}{x} \right)^J$$

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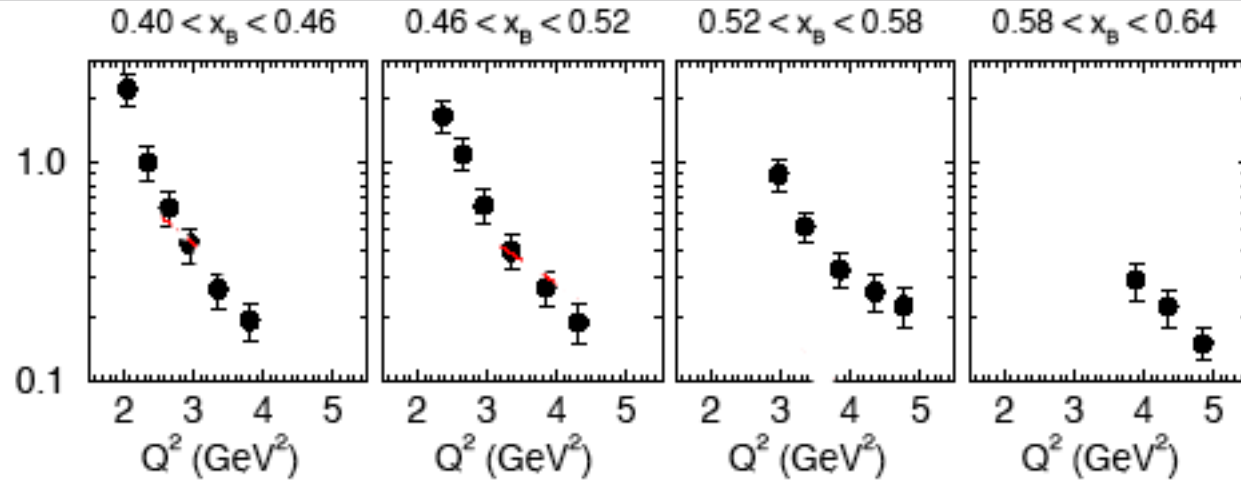
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DVCS

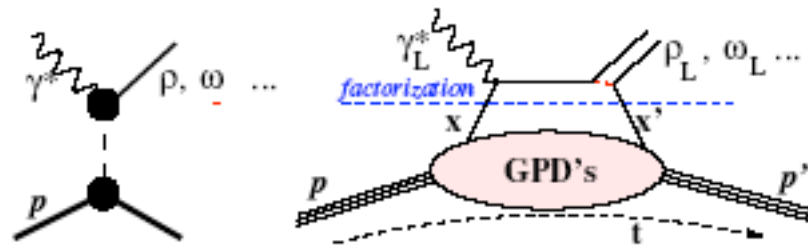
$$A(s, t) \sim \sum_J a(J, t) \sim \left( \frac{Q^2}{x} \right)^{\alpha(t)}$$

Leading Regge exchange dominates not only at small-x but for all-x



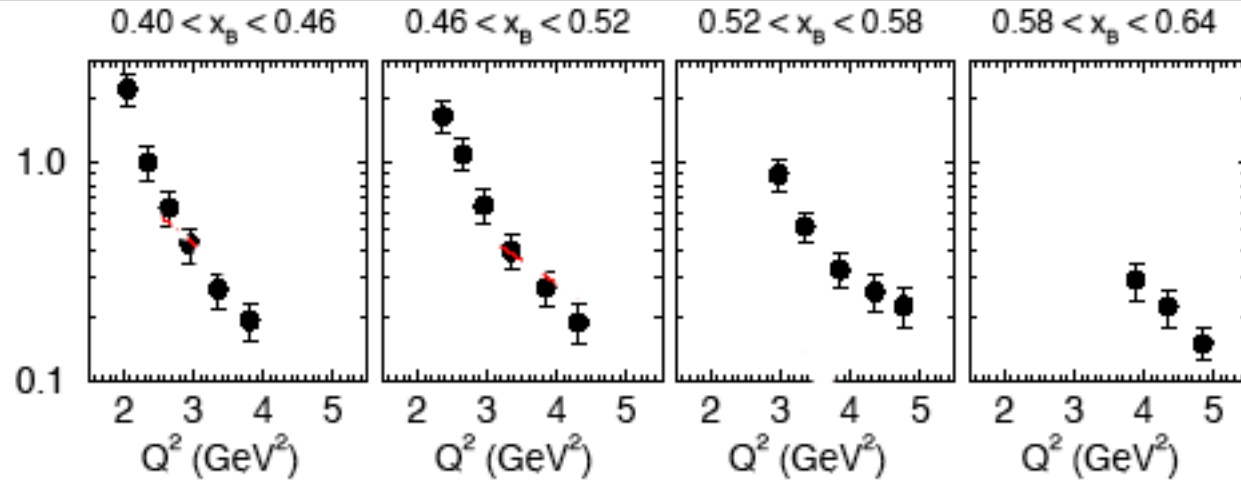


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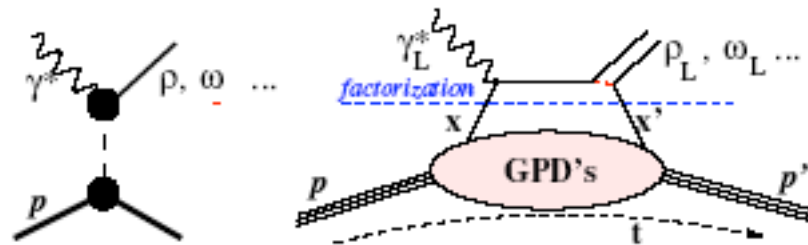


**Fig. 1.** Schematic representations of the  $t$ -channel exchange (left) and of the handbag diagram (right) for exclusive vector meson electroproduction.

$$\sigma_{x=fixed} \sim \frac{|A_{\gamma^* \rightarrow V}|^2}{Q^4} \sim \frac{1}{Q^6}$$



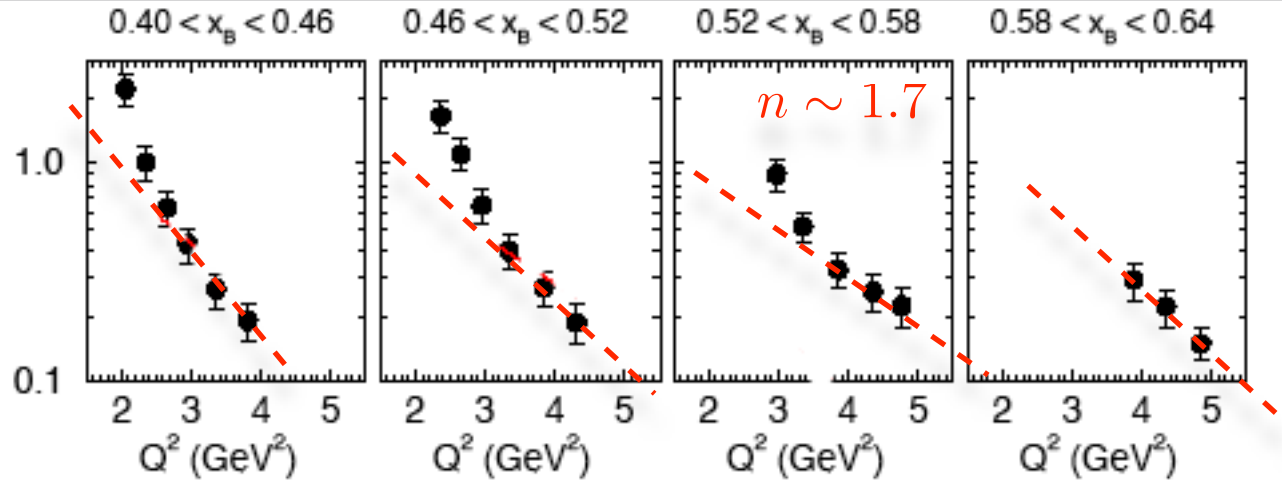
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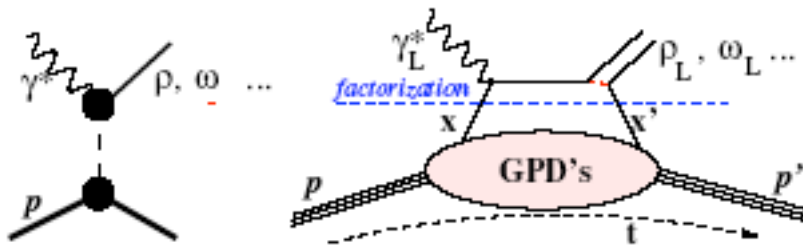
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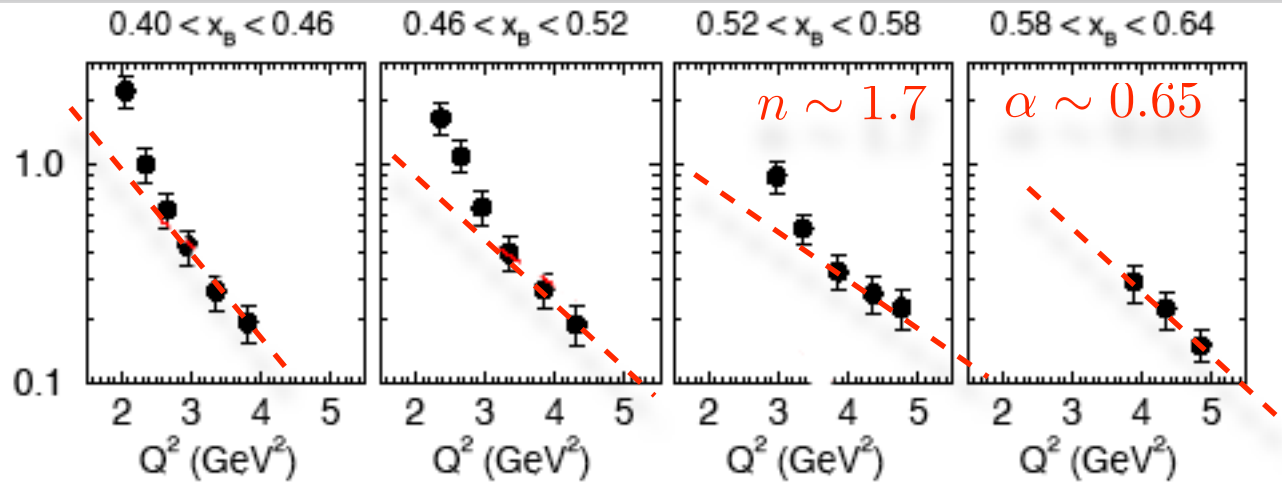
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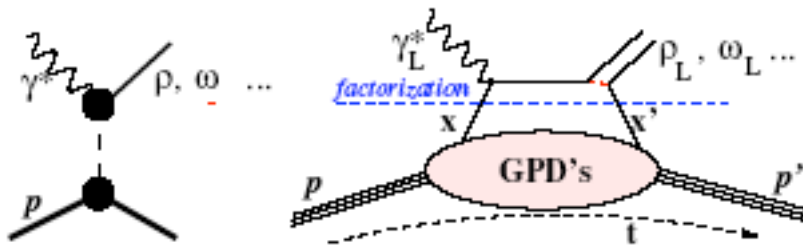
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$$n = 3 \rightarrow 3 - 2\alpha(t)$$

## Summary:

- in DIS regge dominates only when  $x_{BJ} \ll 1$
- in DVCS leading regge dominates even in the valence region