

A Possible $\vec{H}\vec{D}$ Target for Electro-production Experiments

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(BNL \rightarrow JLab)

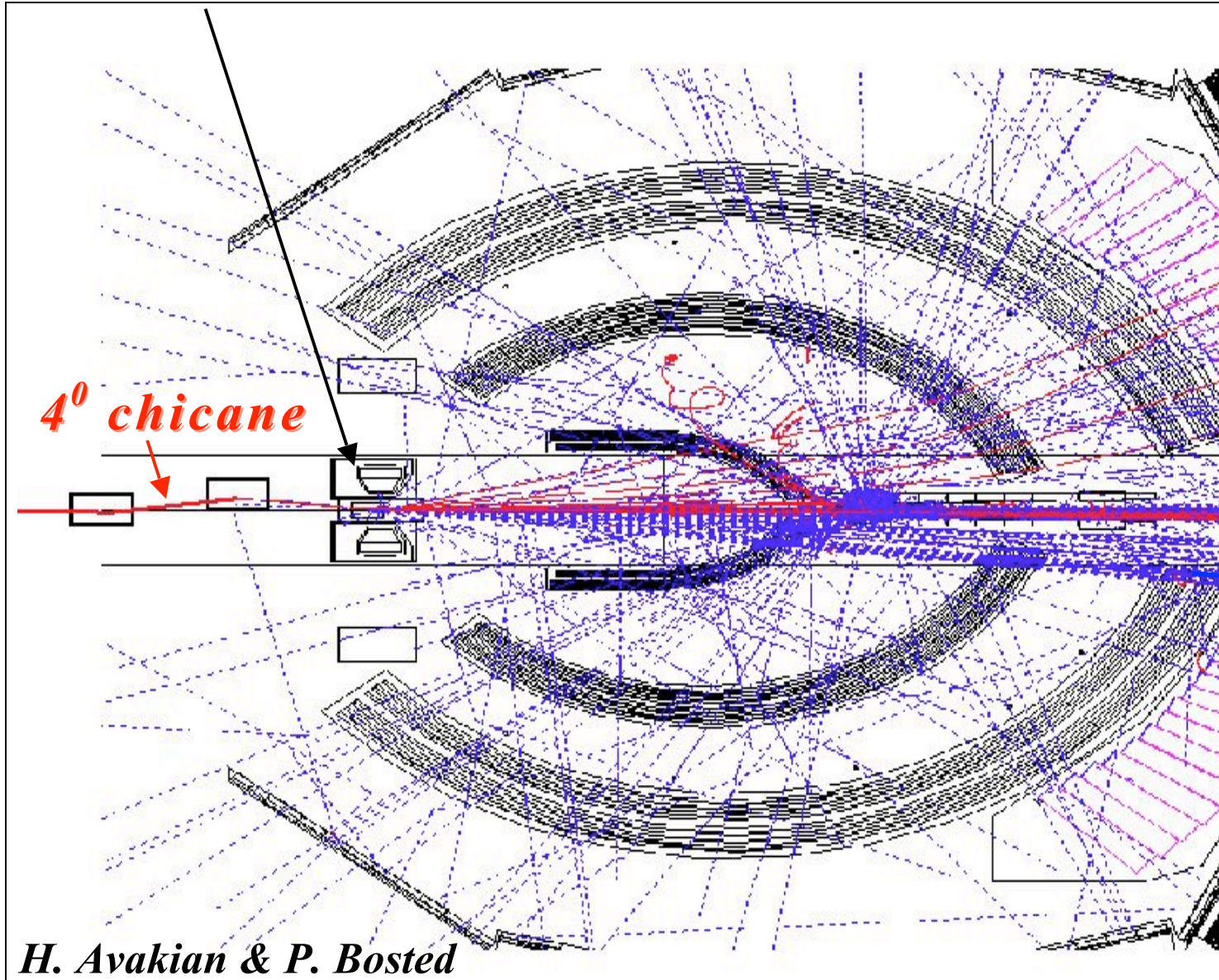
- *motivating factors for transversely polarized targets*
- *frozen-spin $\vec{H}\vec{D}$ and performance with photon beams*
- *factors limiting depolarization with electrons*
- *advantages for transversely polarized \vec{H} and \vec{D}*

Electron experiments with transversely polarized \vec{H} and \vec{D}

- *DVCS, DVMP :*
 $\Rightarrow \tilde{E}$ GPD \rightarrow 2+1 dimensional tomography
 \rightarrow quark orbital angular momentum
- *Semi-inclusive-DIS :*
 \Rightarrow Collins function \rightarrow transverse $\vec{q}q \rightarrow$ Asy in hadron fragmentation
 \rightarrow transverse quark orbital angular momentum
 \Rightarrow Sivers function \rightarrow u-d separation in $\vec{N}^\perp \rightarrow$ single-spin Asy
- *Inclusive-DIS :*
 $\Rightarrow g_2, A_2$ PDF \rightarrow color-polarizability of the gluon field
- *N^* transition form factors :*
 \rightarrow constraining structure of baryons

UVa (Oxford) Transverse $N\vec{H}_3 / N\vec{D}_3$ target with CLAS

BdL $\approx 4.2 T \times 0.3 m$



- *large transverse field compensated by chicane*

- *brem γ 's peaked along incoming e at $\sim 4^\circ$*

\Rightarrow "Sheet of flame" 4°

\Rightarrow large background

- *limited acceptance in θ and Q^2*

H. Avakian & P. Bosted

Polarizing HD: the rotational levels of the solid hydrogens

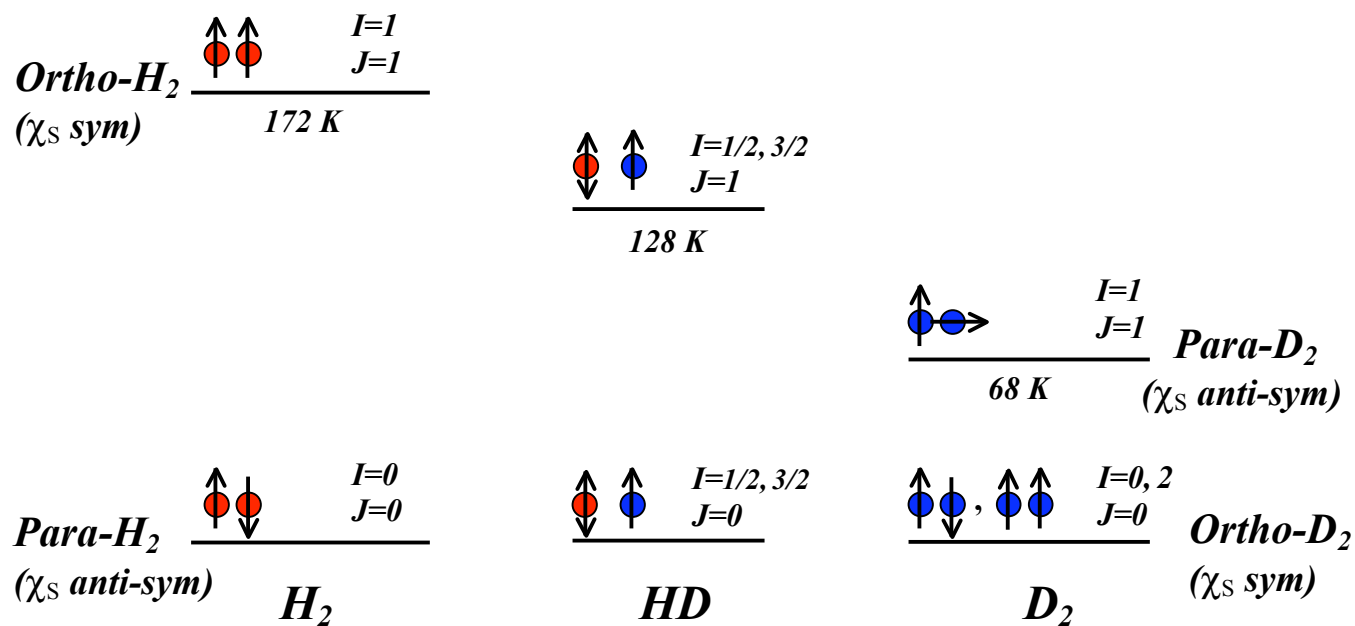


Figure. The relative energy spacing of the low-lying nuclear spin (I) and molecular orbital angular momentum (J) levels in H_2 , HD and D_2 . The symmetries of the nuclear spin wavefunction (χ_s) are indicated.

Polarizing HD: the rotational levels of the solid hydrogens

- **Rapidly polarizable levels: nuclear spin $I \neq 0$ AND orbital $J \neq 0$**

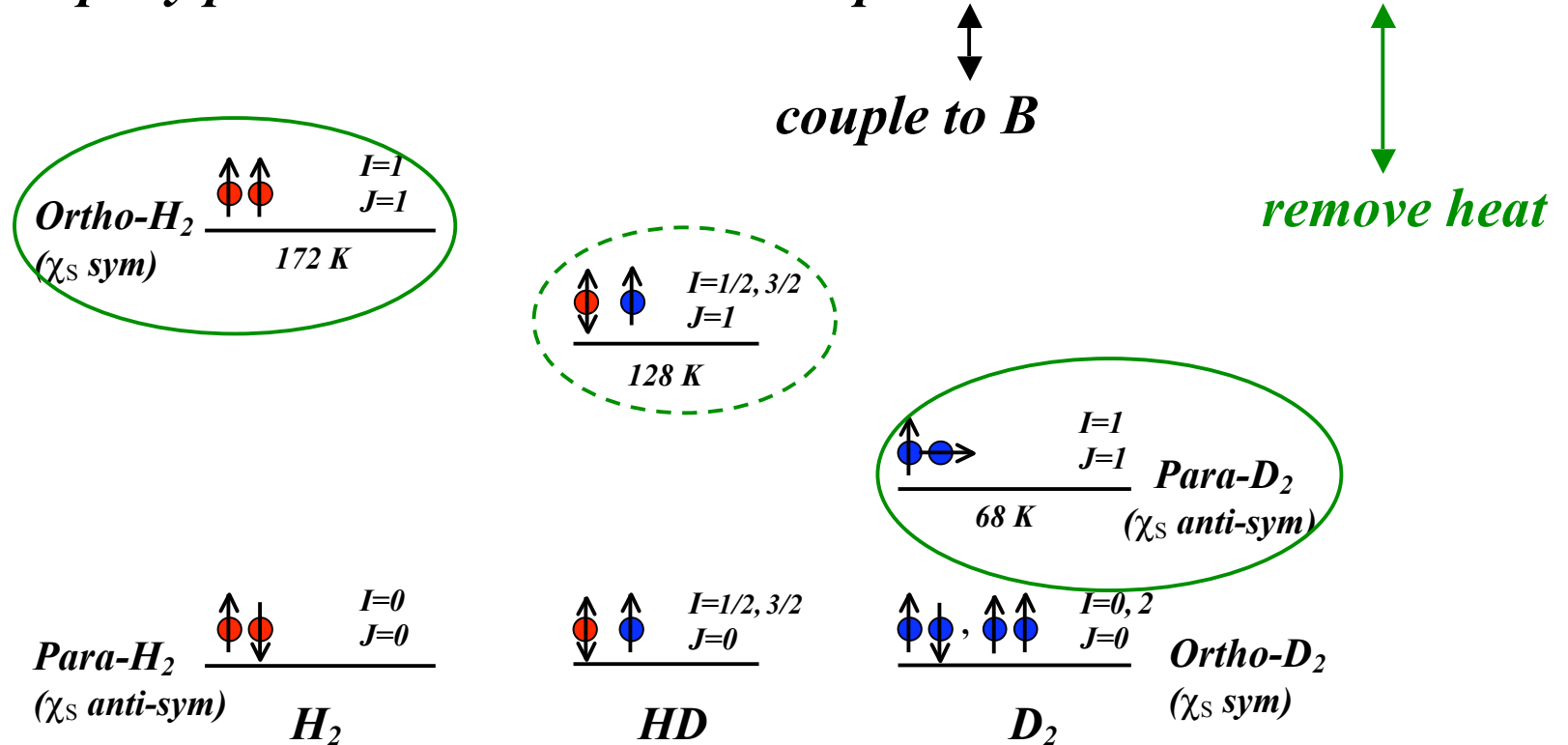
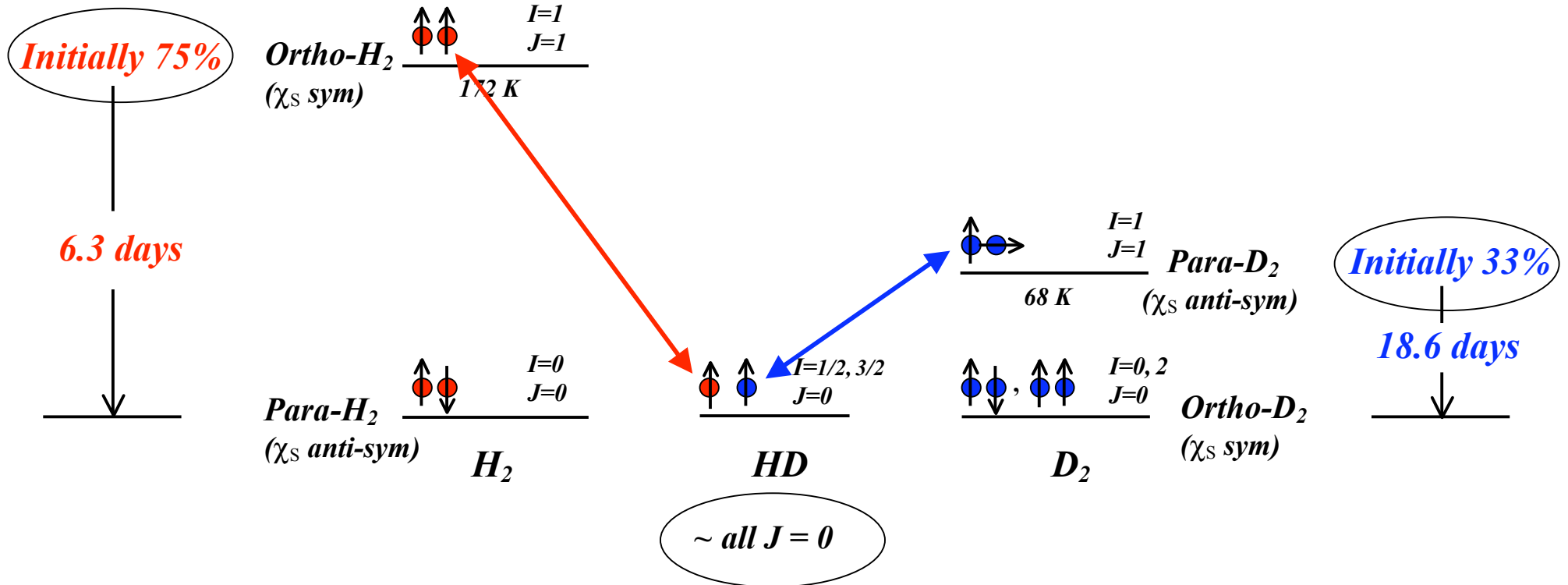


Figure. The relative energy spacing of the low-lying nuclear spin (I) and molecular orbital angular momentum (J) levels in H_2 , HD and D_2 . The symmetries of the nuclear spin wavefunction (χ_s) are indicated.

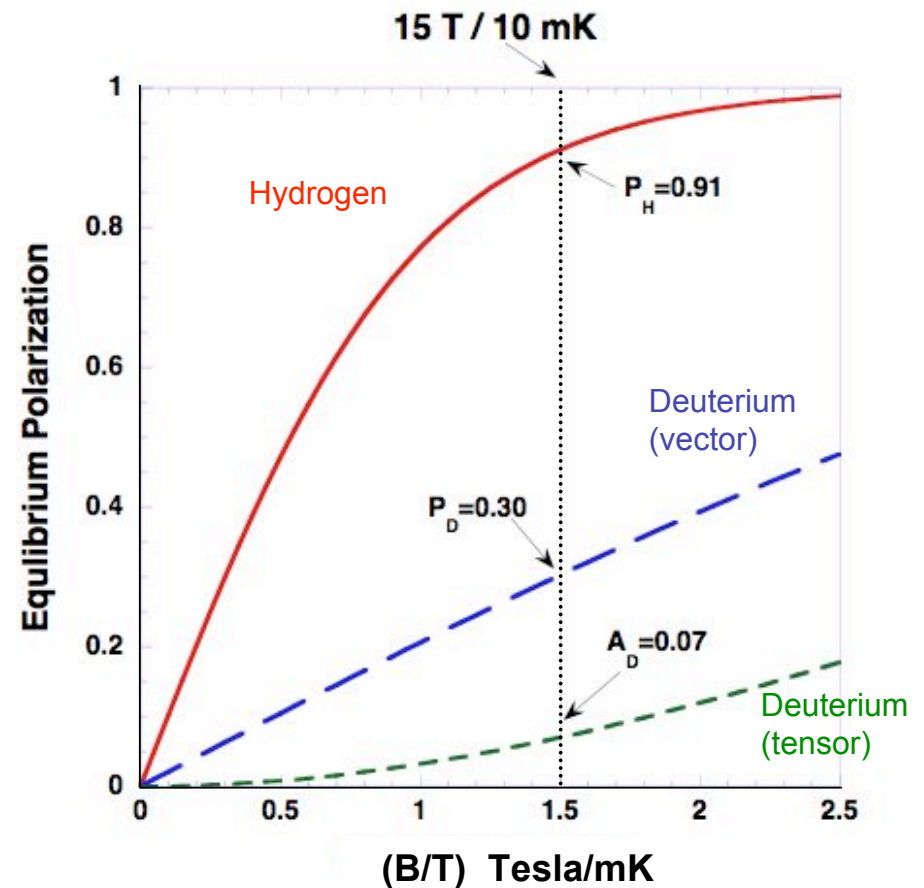
*External Magnetic field rapidly aligns **Ortho-H₂** and **Para-D₂**
then spin-exchanges with **H** and **D** in **HD***



- relaxation switch – A. Honig, *Phys. Rev. Lett.* **19** (1967).

HD field/low-temp Polarization

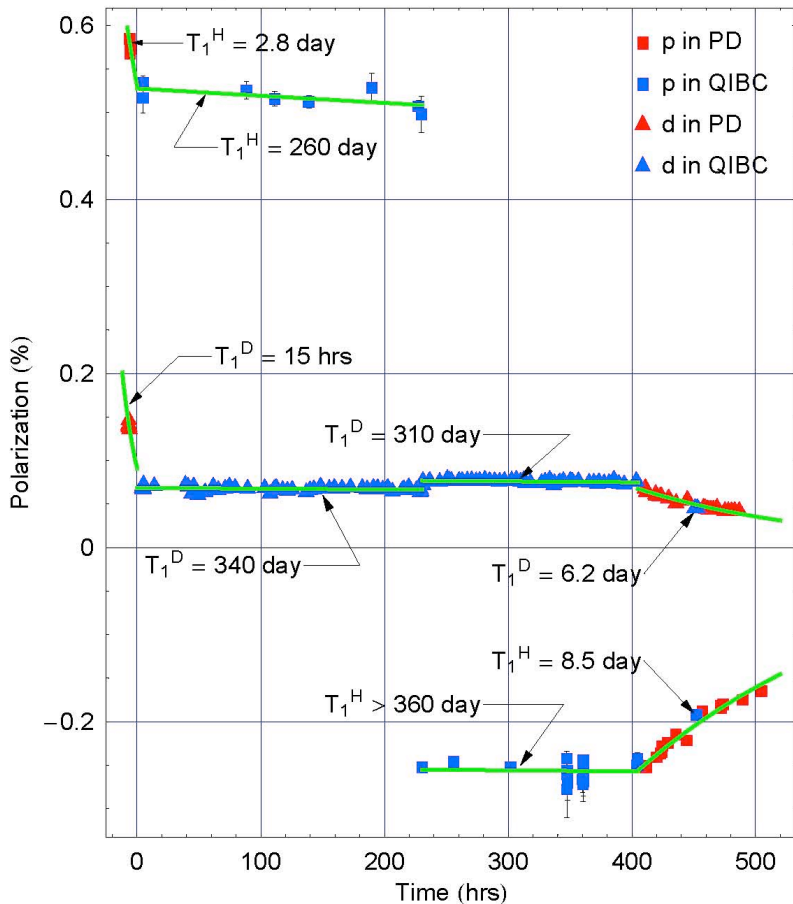
- align spins with high B (15 Tesla) and low T (~12 mK)
- polarize small concentrations of J=1 H₂ and D₂
- o-H₂ and p-D₂ spin-exchanges and polarizes HD
- wait for J=1 H₂ and D₂ to decay



HD polarize/run sequence:

- condense HD gas \rightarrow liquid \rightarrow solid in 2-4 °K dewar ; *calibrate NMR*
- transfer to dilution refrigerator
 - *polarize at 15 tesla and 12-16 mK*
 - *hold for 2-6 months, waiting for ortho- H_2 and para- D_2 to decay away*
- transfer to 2-4 K dewar for polarization measurement
- transfer to In-Beam-Cryostat
 - *hold target for experiment at 0.2 – 0.7 °K and \sim 0.1 to 0.9 tesla*
 - \Rightarrow *Spin-relaxation (T_1) decay times \sim a year*

Target #3 Polarization during Fall 2004 data run



HD polarize/run sequence:

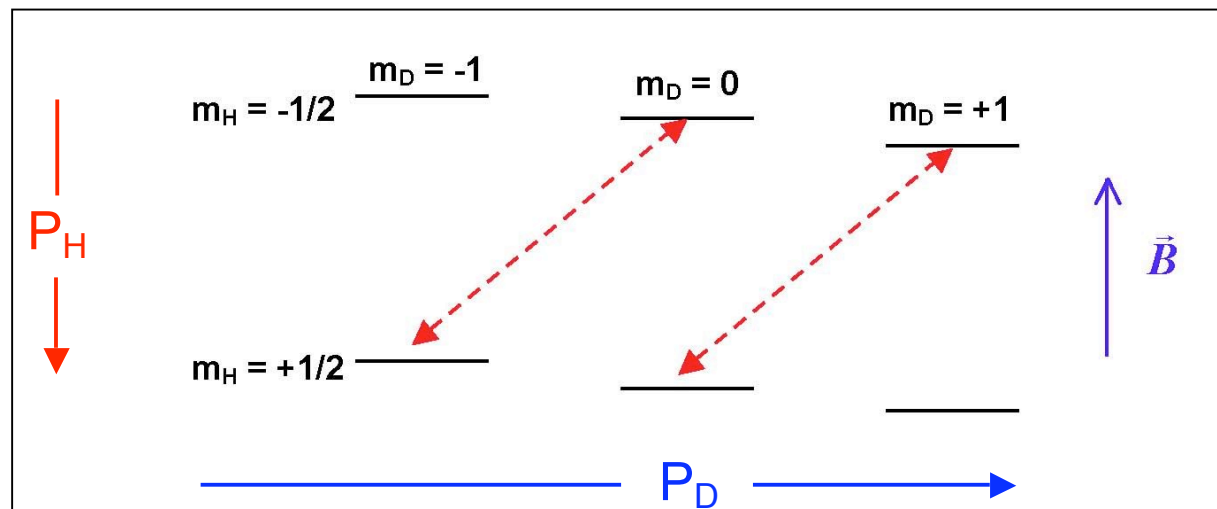
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 - *RF transfer $P_H \rightarrow P_D$ at 0.2 °K and 0.05 tesla*
 - *RF flip spins as needed*

Increasing D polarization:

- *Brute force* (high B/low T) $\Rightarrow P_D \sim 15\% - 25\%$ ($\mu_D / \mu_H \sim 1/3$)
- *1st forbidden adiabatic fast passage* (FAFP) to invert state populations;

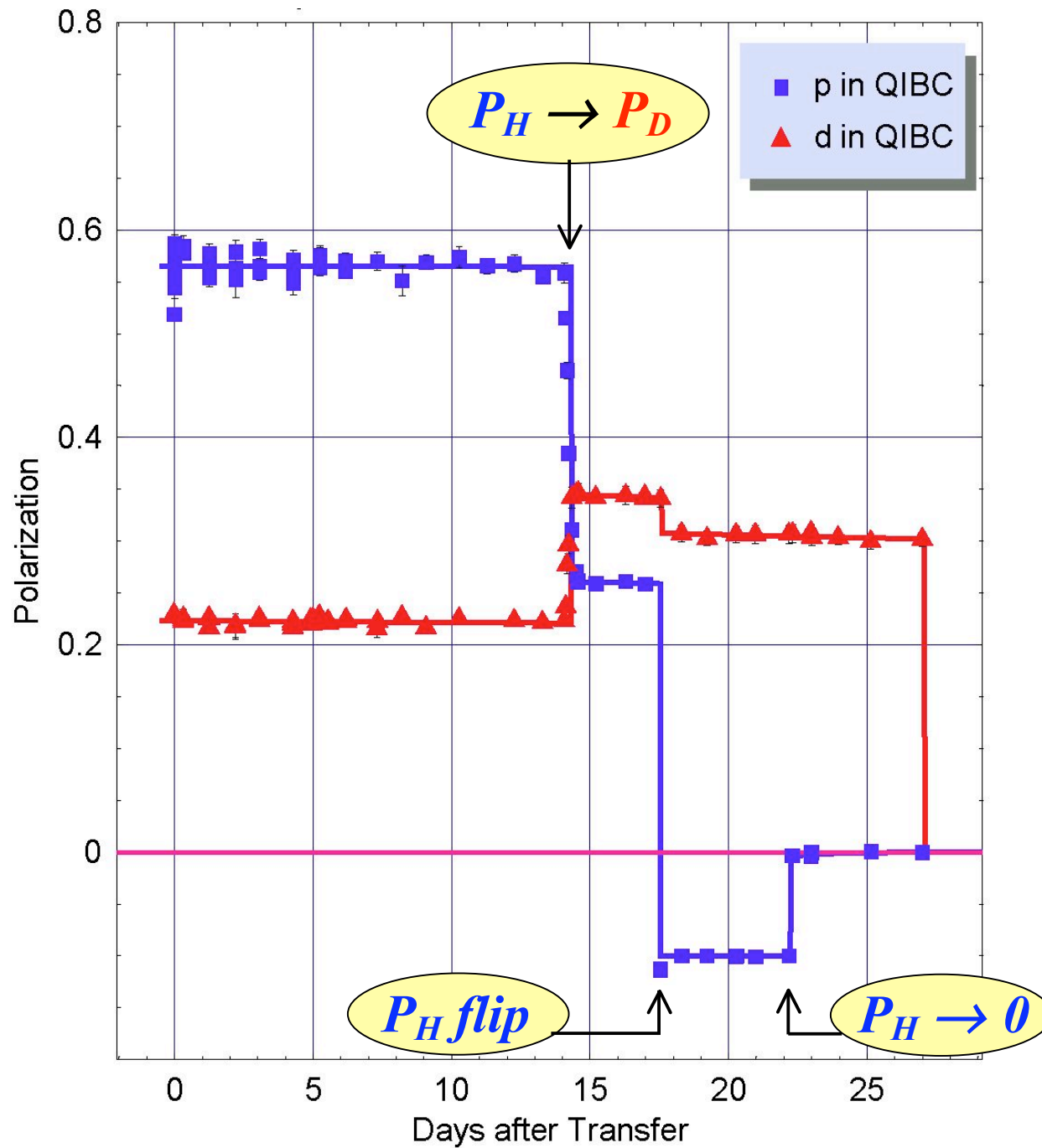
Zeeman levels of HD

- polarize H
- RF transfer $P(H) \rightarrow P(D)$



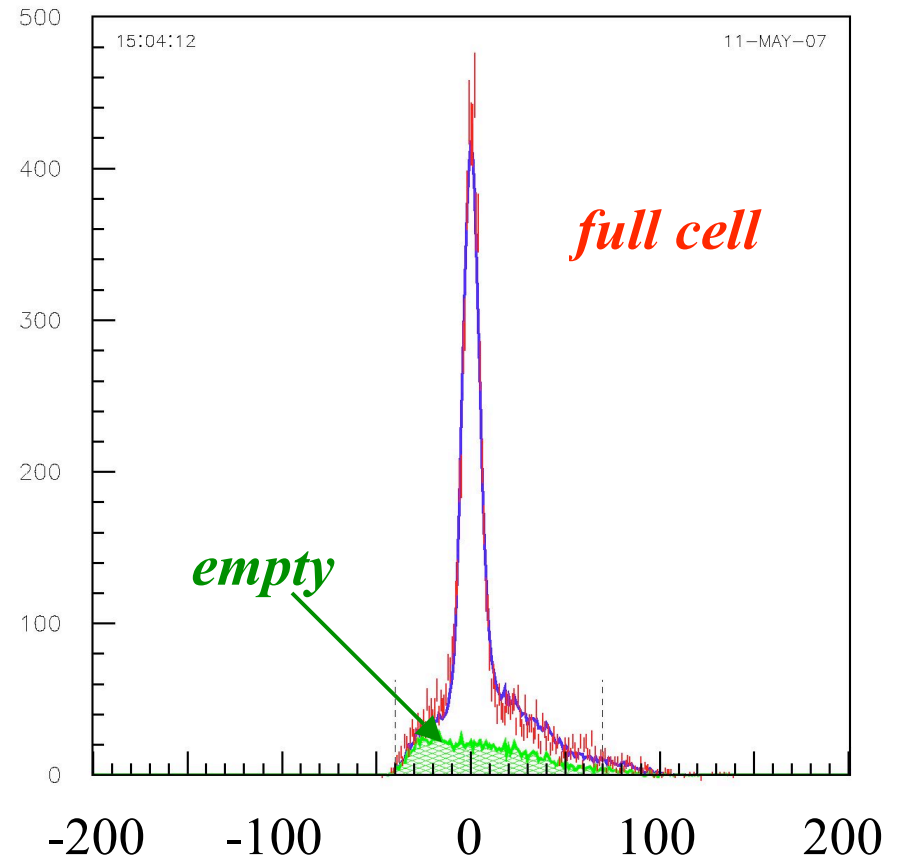
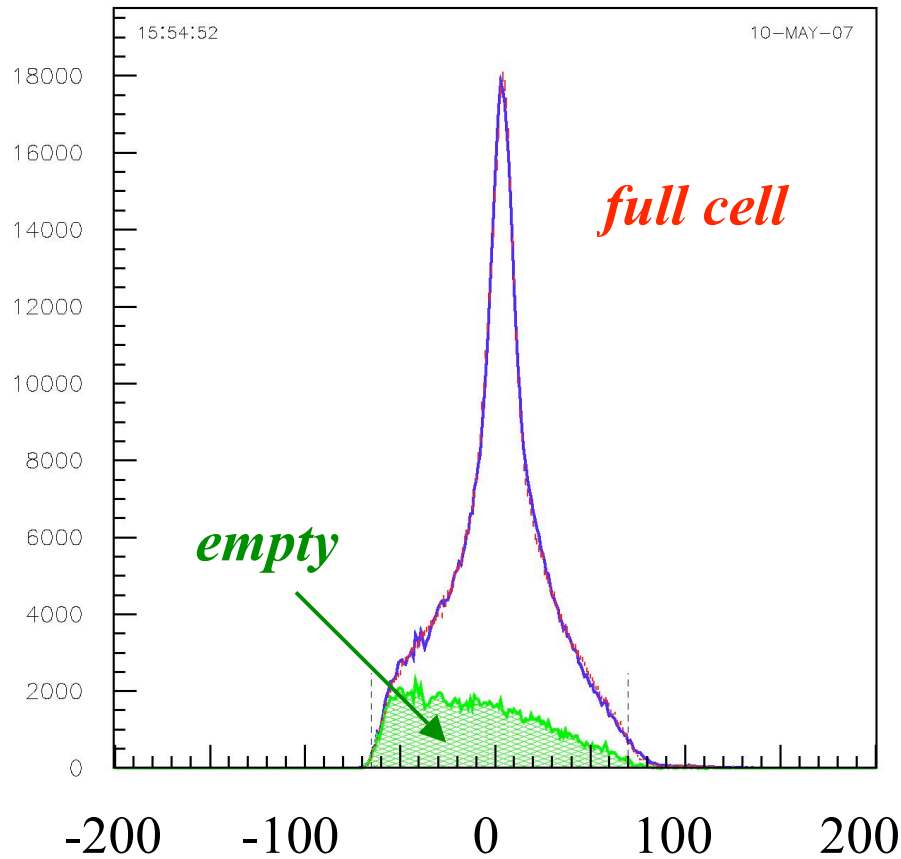
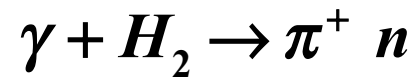
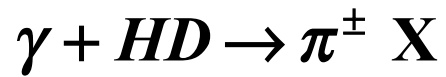
- P_D should reach 50% (limited by NMR field uniformity) \leftarrow requires R&D
- **saturating FAFP transition**
 \rightarrow equalize $\{ m_H = +1/2; m_D = -1, 0 \} \Leftrightarrow \{ m_H = -1/2; m_D = 0, +1 \}$
 $\Rightarrow P_D = 37\% \leftarrow$ today

BNL – Fall'2006



- target cell contribution can be measured and subtracted

$$E_\gamma = 300 \text{ MeV}$$



missing 2 - body energy (MeV)

Expected spin-relaxation times for appropriately prepared targets

measured (γ)

projected

<i>B</i>	<i>0.89 tesla</i>	<i>0.01 tesla</i>	<i>0.40 tesla</i>	<i>0.04 tesla</i>
<i>B × dL (for L=0.12m)</i>	<i>0.108 tesla-m</i>	<i>0.001 tesla-m</i>	<i>0.048 tesla-m</i>	<i>0.005 tesla-m</i>
<i>orientation</i>	<i>solenoid</i>	<i>solenoid</i>	<i>saddle</i>	<i>saddle</i>
<i>T₁(H)</i>	<i>> 300 d</i>	<i>8 d</i>	<i>>200 d</i>	<i>~ 30 d</i>
<i>T₁(D)</i>	<i>> 500 d</i>	<i>55 d</i>	<i>>300d</i>	<i>~200 d</i>

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*compare to 1.4 T-m
with NH₃/ND₃*

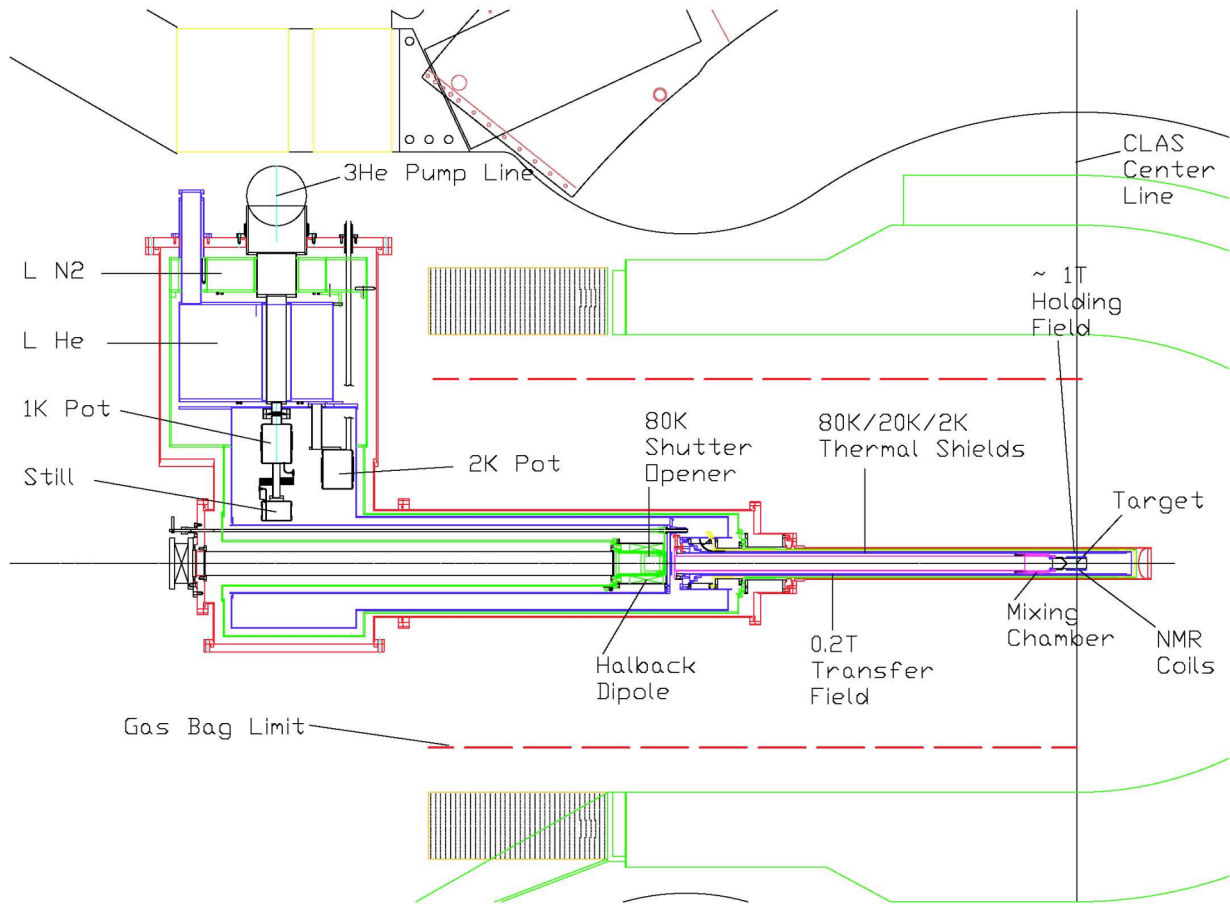


Figure A8. Conceptual design of modified version of the BNL IBC for use in the CLAS.

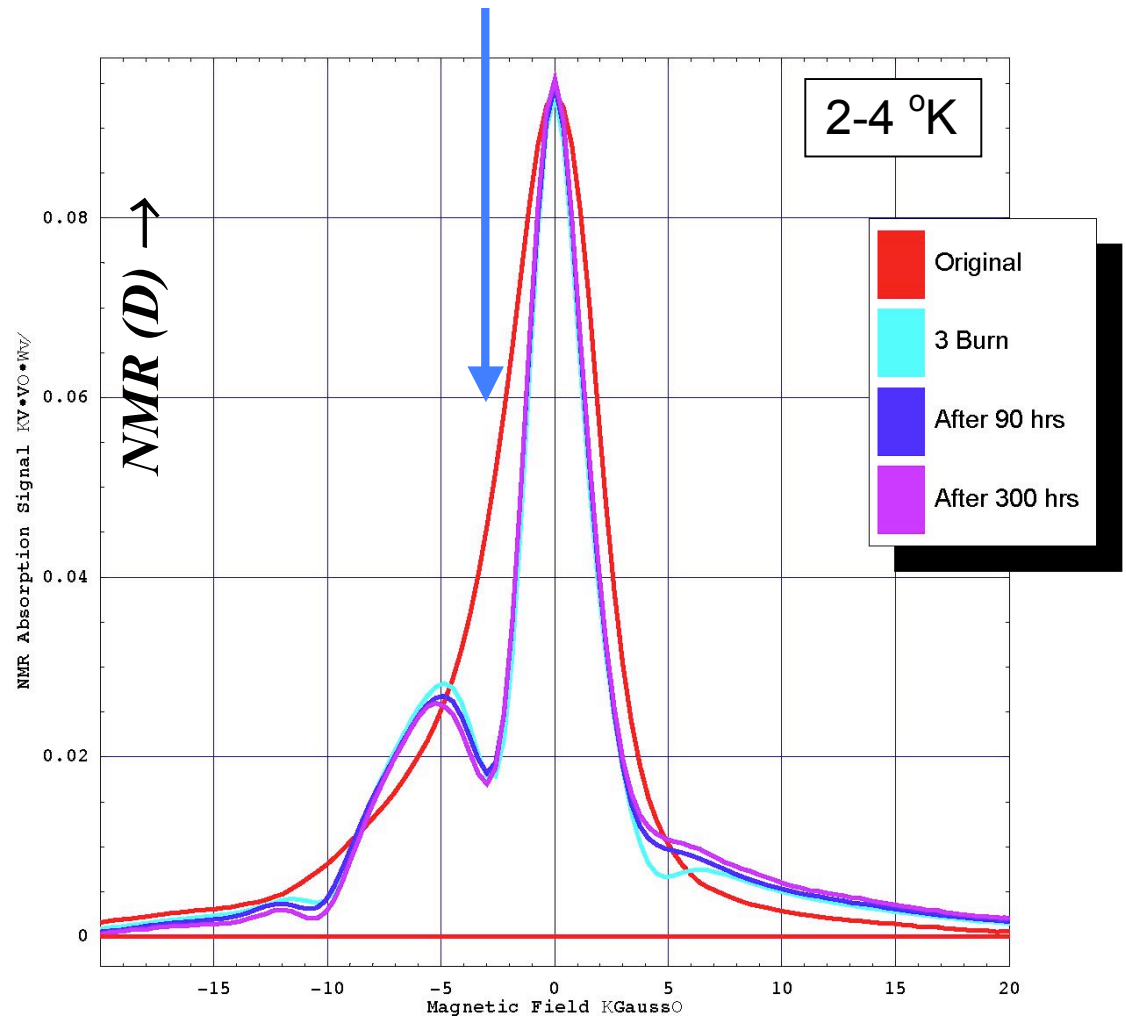
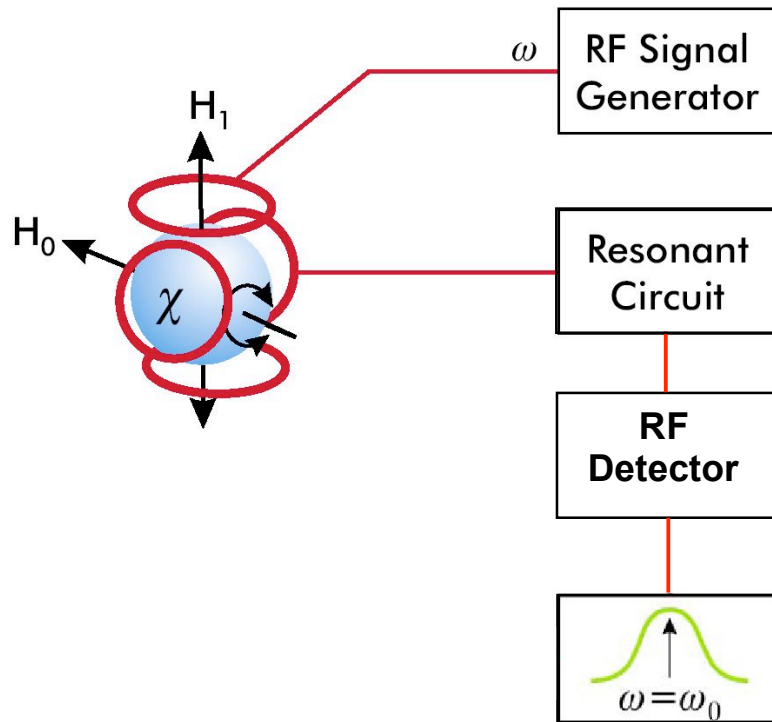
Depolarization of frozen-spin \vec{H} \vec{D} with electrons

- *beam heating*
 - 5 nA of 10 GeV electrons \Rightarrow 5 mW heat in 2 cm of HD (GEANT)
 \approx cooling power of BNL In-Beam Cryostat at 0.5 K (can be increased)
 - 4 times lower heating than FROST(Butanol), due to lower Z
 - spin-relaxation time (T_1) for HD \sim a year at these temperatures
- *spin-diffusion of paramagnetic centers*
 - *e brem* creates free radicals with randomly oriented nuclear spin; absolute number are small, but these can be *sinks* for polarization
 - *spin-diffusion* time measured at 2 K: \sim 1 day for \vec{H}
 $\sim \infty$ for \vec{D} (unmeasurable in 2 weeks)

(*spin-diffusion* times could increase at lower T ?)

Burning an RF polarization hole

- *cross-coil NMR*
- *field scan at fixed frequency*



- H_0 inhomogeneity \Rightarrow *D*-line width
- *field and position are correlated*
- *no change in the D-polarization hole after 2 weeks \Rightarrow D spin diffusion extremely slow*

Potential advantages with frozen-spin transverse $\vec{H} \vec{D}$

- *very low BdL (almost none for \vec{D})
⇒ no “sheet of flame”*
- *better figure of merit
→ almost no dilution
→ small nuclear background (sampled with empty cell)
→ long Radiation Length (625 cm) ⇒ few brem γ 's*
- *wide acceptance in θ and Q^2
→ open geometry cryostat centered in CLAS(6/12)
(but, will have to deal with low-momentum Møller electrons)*

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caveat:

$e^+ \vec{H} \vec{D}$ test is necessary to verify polarization retention with electrons

Extras

Ortho \Leftrightarrow para decays generate heat, which must be removed to polarize

- *HD condensed into target cell with ~ 2000 $50\ \mu\text{m}$ Al cooling wires soldered into 60 holes in copper cooling ring*



- *Composition of a standard target cell with 4 cm of HD (0.9 moles):*

Material	gm/cm ²	mass fraction
HD	0.735	77%
Al	0.155	16%
CTFE (C ₂ ClF ₃)	0.065	7 %

Frozen-Spin \vec{H} \vec{D} - summary

- *pure target, high nucleon polarizations*
- *very low-background - cell contains only unpolarizable nucleons (20%)
 \Leftrightarrow conventional empty-cell subtractions*
 \Rightarrow *E06-101 HD figure of merit $> 20 \times \text{FROST}(\text{C}_4\text{H}_9\text{OH})$*
- *in- γ -beam life-times $>$ year*
- *RF moves spins $\vec{H} \Leftrightarrow \vec{D}$ as needed*
- *In-Beam Cryostat centered in CLAS; open acceptance at back angles*
- *developed at BNL/LEGS; migrating to JLab*

Table A2. Factors contributing to the systematic error on target polarization.

Source	$\delta P(H)$	$\delta P(D)$
thermal equilibrium calibration - noise, temperature, bkg, ...	0.9%	1.0%
frozen-spin measurement - white noise	0.4%	2.0%
- holding field noise	0.5%	0.5%
- non-linearities, homogeneity, ...	1.0%	1.0%
calibration transfer - circuit drift, differential ramp	1.6%	1.6%
- Lock-in gain differential error	2.8%	2.8%
- cold-transfer loss	1.0%	1.0%
Total fractional error:	3.7%	4.2%