

Deeply Virtual Compton Scattering @ JLab

Franck Sabatié

CEA Saclay

For the Hall A and Hall B collaborations

Exclusive'07 - JLab

May 21st 2007

Introduction

Non-dedicated measurements

E00-110 experiment in Hall A

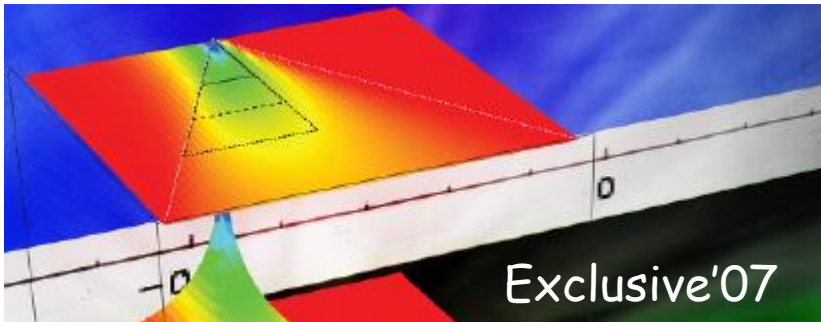
E03-106 experiment in Hall A

[Prelim. deep- π^0 E00-110 result]

E1-DVCS experiment in Hall B

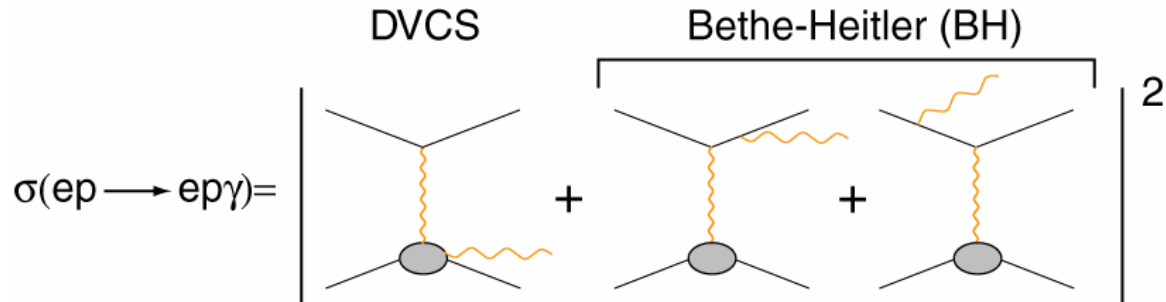
What I didn't cover

Summary



Experimental observables linked to GPDs

Experimentally, DVCS is undistinguishable with Bethe-Heitler



However, we know FF at low t and **BH is fully calculable**

Using a polarized beam on an unpolarized target, 2 observables can be measured:

$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} \approx |T^{BH}|^2 + 2T^{BH} \cdot \text{Re}(T^{DVCS}) + |T^{DVCS}|^2$$

$$\frac{d^4\vec{\sigma} - d^4\overleftarrow{\sigma}}{dx_B dQ^2 dt d\varphi} \approx 2T^{BH} \cdot \text{Im}(T^{DVCS}) + \left[|T^{DVCS\rightarrow}|^2 - |T^{DVCS\leftarrow}|^2 \right]$$

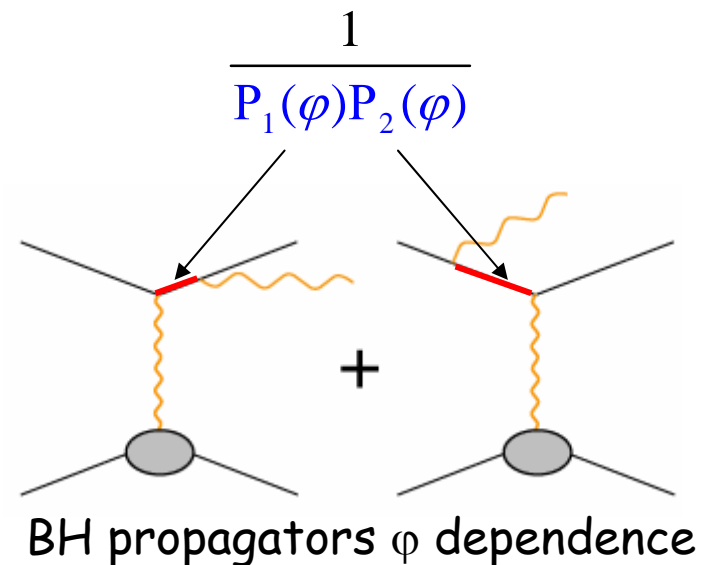
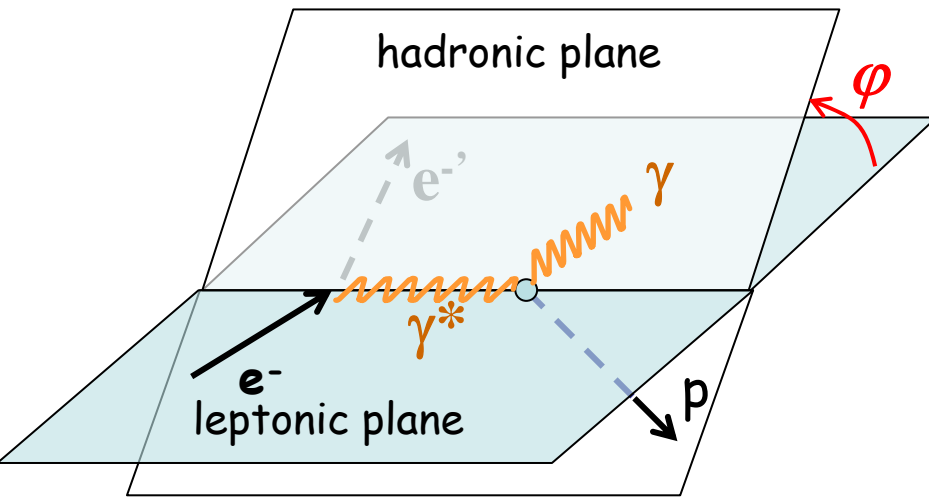
At JLab energies,
 $|T^{DVCS}|^2$ supposed small

Into the harmonic structure of DVCS

$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} = \frac{1}{P_1(\varphi)P_2(\varphi)} \Gamma_1(x_B, Q^2, t) \left\{ c_0^{BH} + c_1^{BH} \cos \varphi + c_2^{BH} \cos 2\varphi \right\} \leftarrow |T^{BH}|^2$$

$$+ \frac{1}{P_1(\varphi)P_2(\varphi)} \Gamma_2(x_B, Q^2, t) \left\{ c_0^I + c_1^I \cos \varphi + c_2^I \cos 2\varphi + c_3^I \cos 3\varphi \right\}$$

$$\frac{d^4 \vec{\sigma} - d^4 \overleftarrow{\sigma}}{dx_B dQ^2 dt d\varphi} = \frac{\Gamma(x_B, Q^2, t)}{P_1(\varphi)P_2(\varphi)} \left\{ s_1^I \sin \varphi + s_2^I \sin 2\varphi \right\} \leftarrow \text{Interference term}$$



Tests of scaling

$$\frac{d^4 \sigma}{dx_B dQ^2 dt d\varphi} = \frac{1}{P_1(\varphi)P_2(\varphi)} \Gamma_1(x_B, Q^2, t) \{c_0^{BH} + c_1^{BH} \cos \varphi + c_2^{BH} \cos 2\varphi\} \\ + \frac{1}{P_1(\varphi)P_2(\varphi)} \Gamma_2(x_B, Q^2, t) \{c_0^I + c_1^I \cos \varphi + c_2^I \cos 2\varphi + c_3^I \cos 3\varphi\}$$

$$\frac{d^4 \vec{\sigma} - d^4 \overleftarrow{\sigma}}{dx_B dQ^2 dt d\varphi} = \frac{\Gamma(x_B, Q^2, t)}{P_1(\varphi)P_2(\varphi)} \{s_1^I \sin \varphi + s_2^I \sin 2\varphi\}$$

1. Twist-2 terms should dominate σ and $\Delta\sigma$
2. All coefficients have Q^2 dependence which can be tested!

Special case of the asymmetry

The asymmetry can be written as:

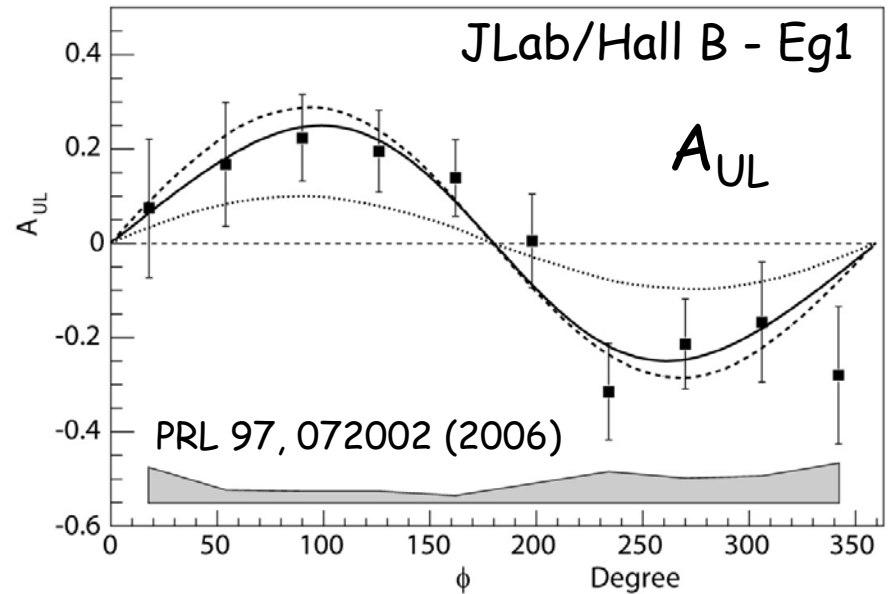
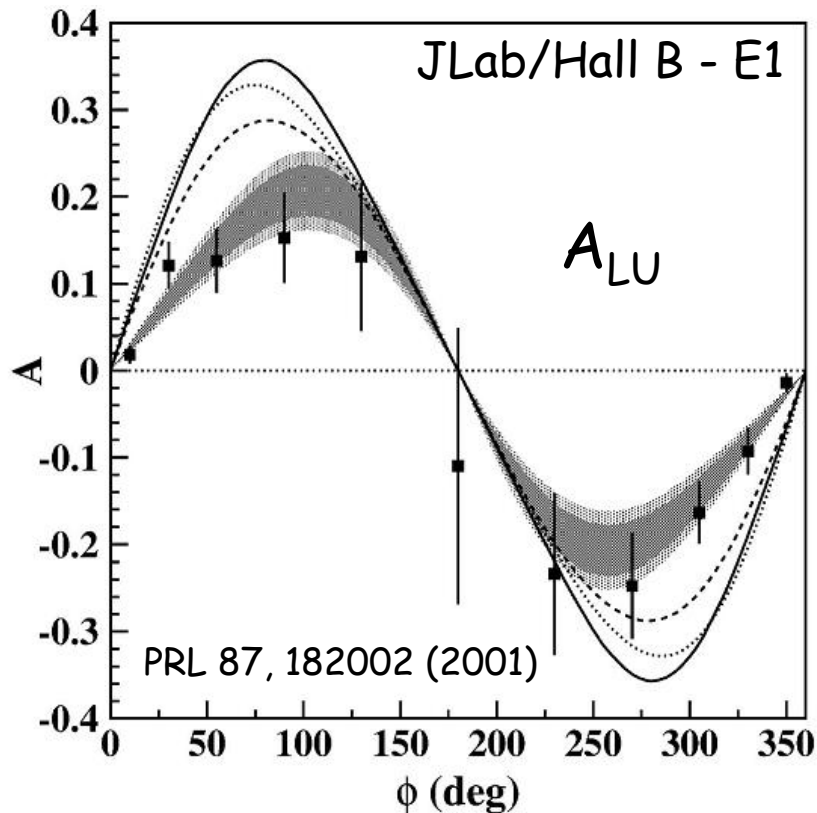
$$\frac{d^4 \overset{\rightarrow}{\sigma} - d^4 \overset{\leftarrow}{\sigma}}{d^4 \overset{\rightarrow}{\sigma} + d^4 \overset{\leftarrow}{\sigma}} = \Gamma_A(x_B, Q^2, t) \frac{s_1^I \sin \varphi + s_2^I \sin 2\varphi}{c_0^I + c_0^{BH} + (c_1^I + c_1^{BH}) \cos \varphi + \dots}$$

Pro: easier experimentally, smaller RC, smaller systematics

Con: direct extraction of GPDs is model- (or hypothesis-) dependent
(denominator complicated and unknown)

It was naturally the first observable extracted from non-dedicated experiments...

Published non-dedicated JLab/Hall B results on A_{LU} and A_{UL}

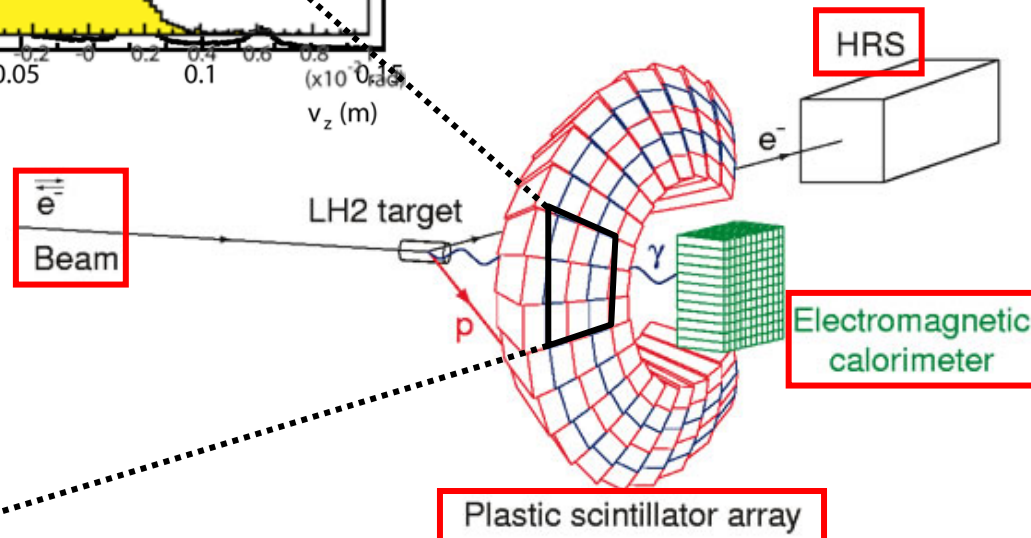
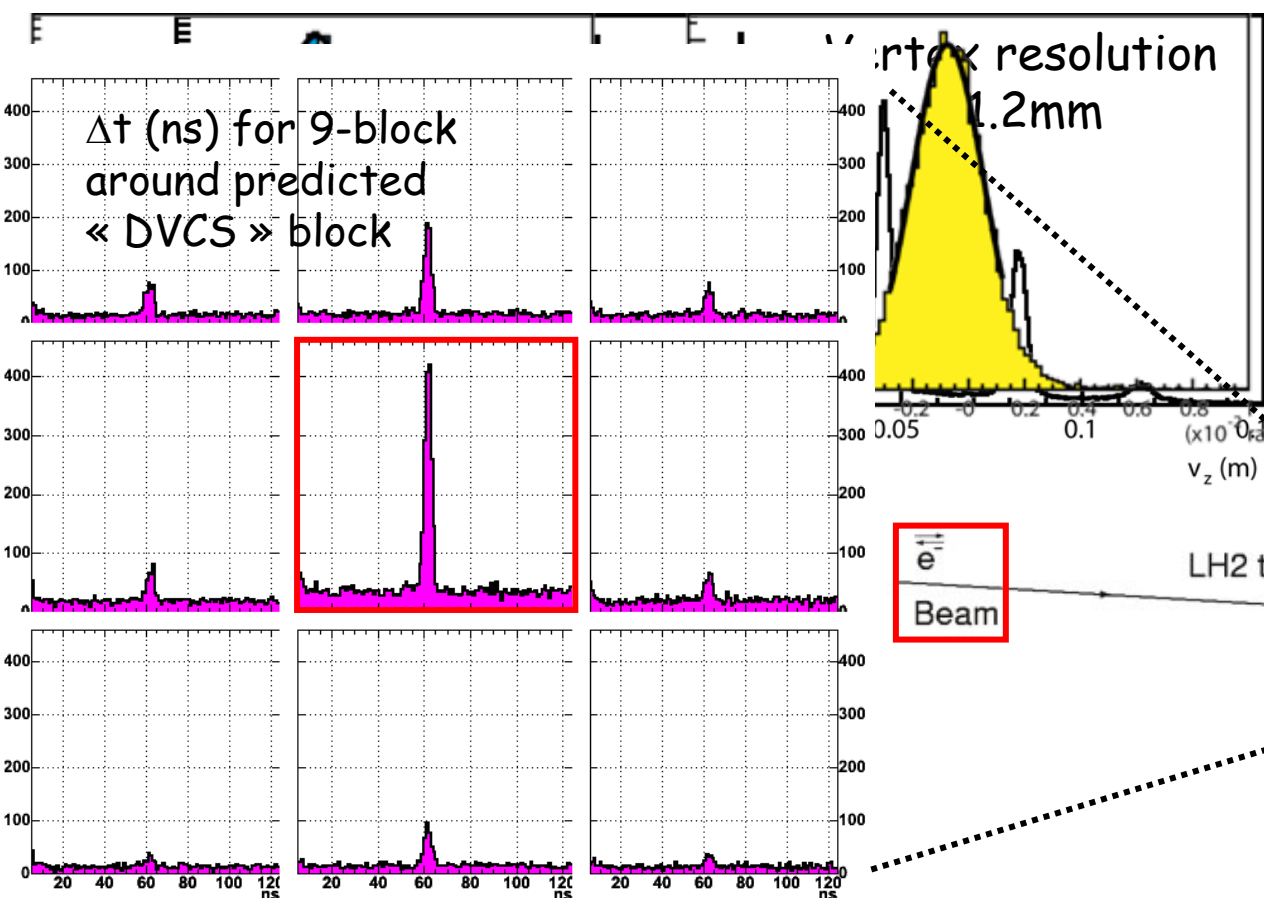
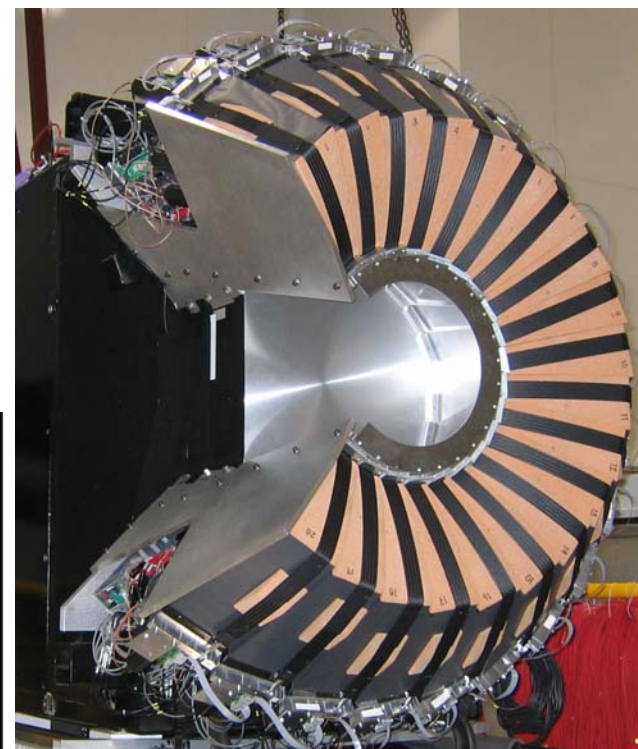


Both results show, with a limited statistics, a $\sin \phi$ behavior (necessary condition for handbag dominance)

In the A_{LU} result, models (VGG) tend to over-estimate the data

E00-110 experimental setup and performances

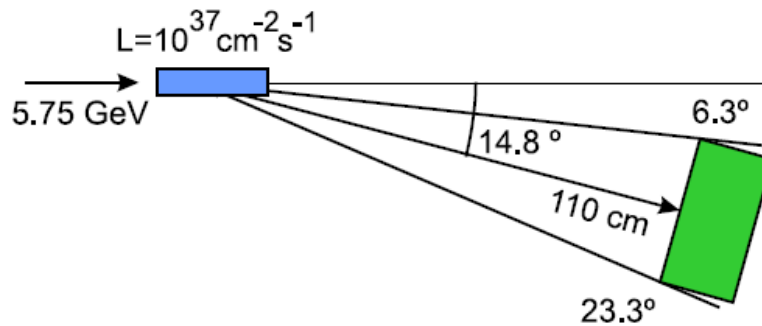
- 75% polarized 2.5uA electron beam
- 15cm LH2 target
- Left Hall A HRS with electron package
- 11x12 block PbF2 electromagnetic calorimeter
- 5x20 block plastic scintillator array



E00-110 kinematics

Kin	Q^2 (GeV ²)	x_B	θ_{γ^*} (deg.)	W (GeV)
1	1.5	0.36	22.3	1.9
2	1.9	0.36	18.3	2.0
3	2.3	0.36	14.8	2.2

The calorimeter is centered on the virtual photon direction



50 days of beam time in the fall 2004, at 2.5 μA intensity

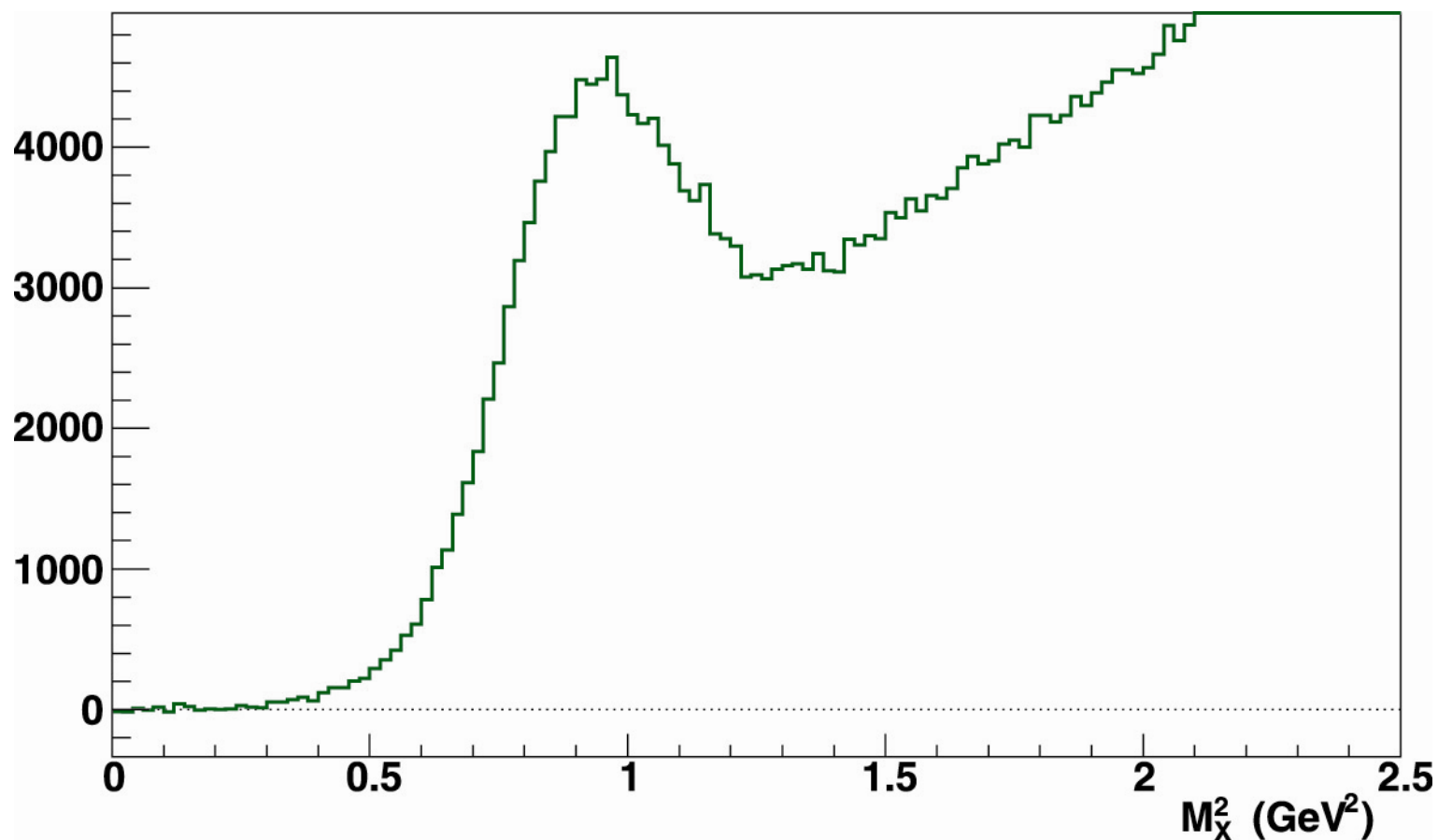
$$\int Lu \cdot dt = 13294 \text{ fb}^{-1}$$

Analysis - Looking for DVCS events

HRS: Cerenkov, vertex, flat-acceptance cut with R-functions

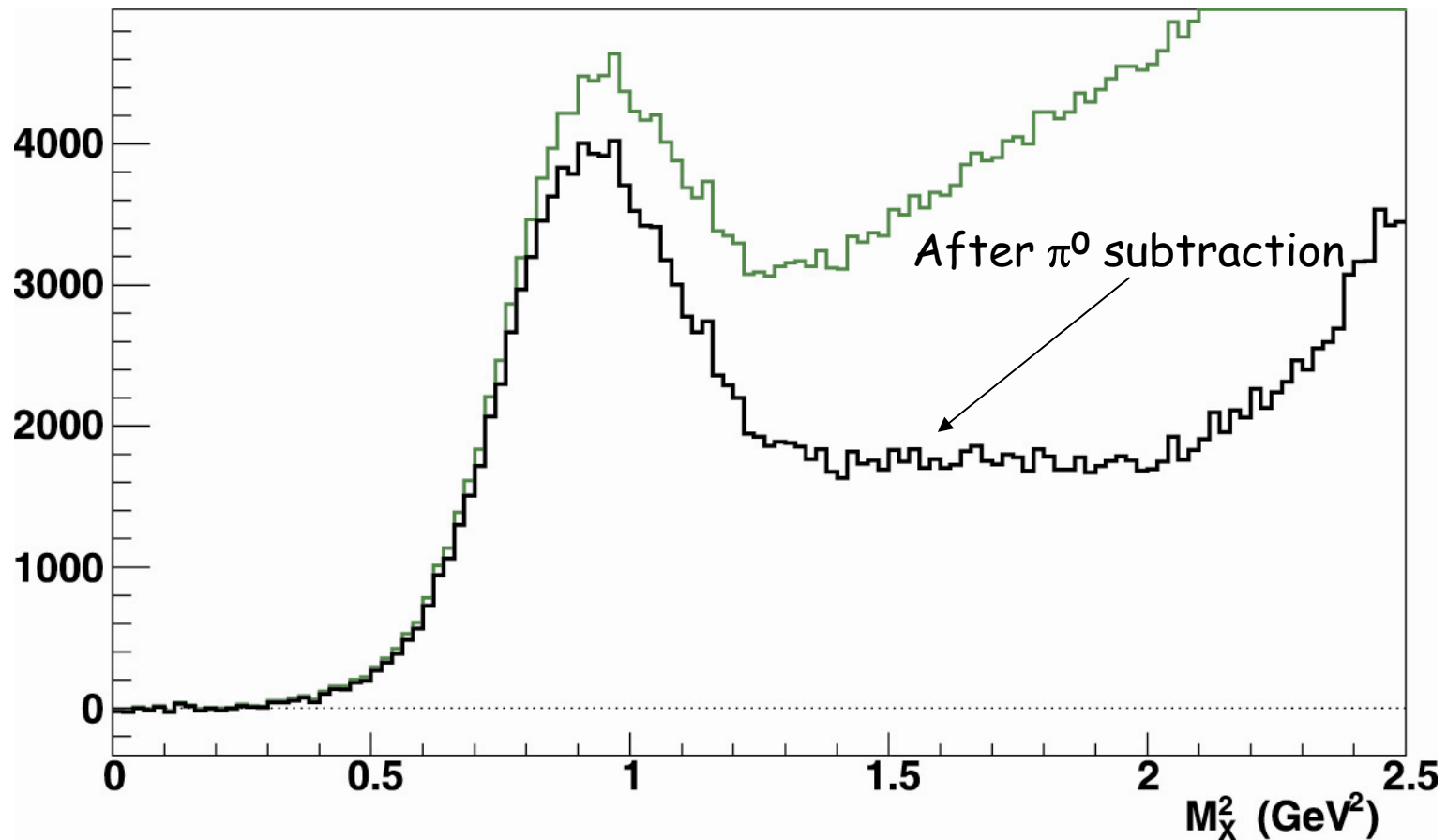
Calo: 1 cluster in coincidence in the calorimeter above 1 GeV

With both: subtract accidentals, build missing mass of (e, γ) system



Analysis - π^0 subtraction effect on missing mass spectrum

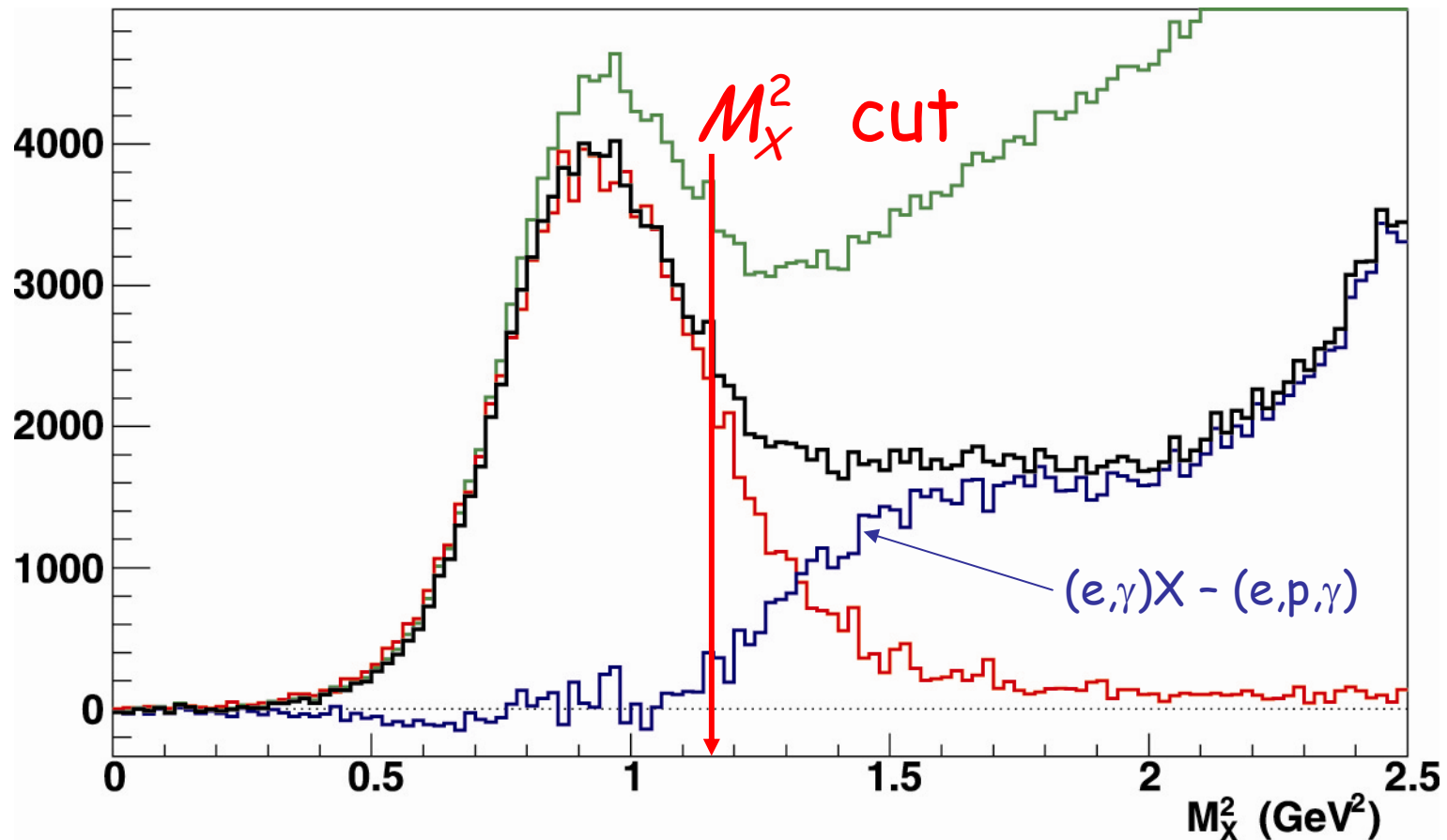
Using $\pi^0 \rightarrow 2\gamma$ events in the calorimeter, the π^0 contribution is subtracted bin by bin



Analysis - Exclusivity check using Proton Array and MC

Using Proton-Array, we compare the missing mass spectrum of the triple and double-coincidence events.

The missing mass spectrum using the Monte-Carlo gives the same position and width. **Using the cut shown on the Fig., the contamination from inelastic channels is estimated to be under 3%.**



Difference of cross-sections

PRL97, 262002 (2006)

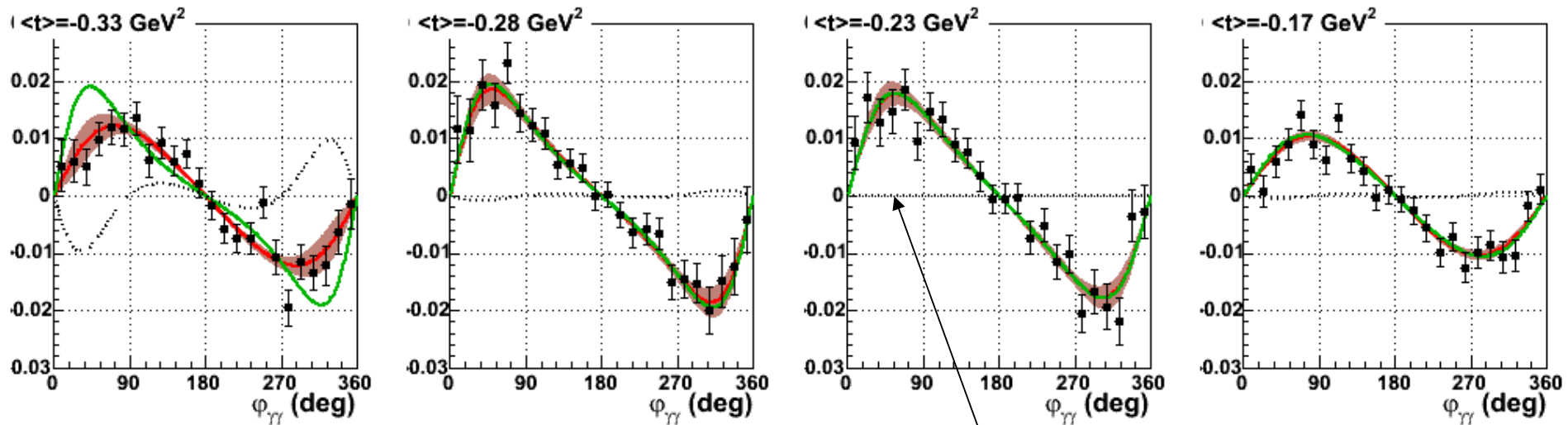
$$\langle Q^2 \rangle = 2.3 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.36$$

$$\frac{1}{2} \left(\frac{d^4\sigma^+}{dQ^2 dx_B dt d\phi_{\gamma\gamma}} - \frac{d^4\sigma^-}{dQ^2 dx_B dt d\phi_{\gamma\gamma}} \right) \text{ (nb/GeV}^4\text{)}$$

* E00-110
 — Fit
 ■ 1- σ

— Im(C¹) Twist-2
 Im(C_{eff}) Twist-3



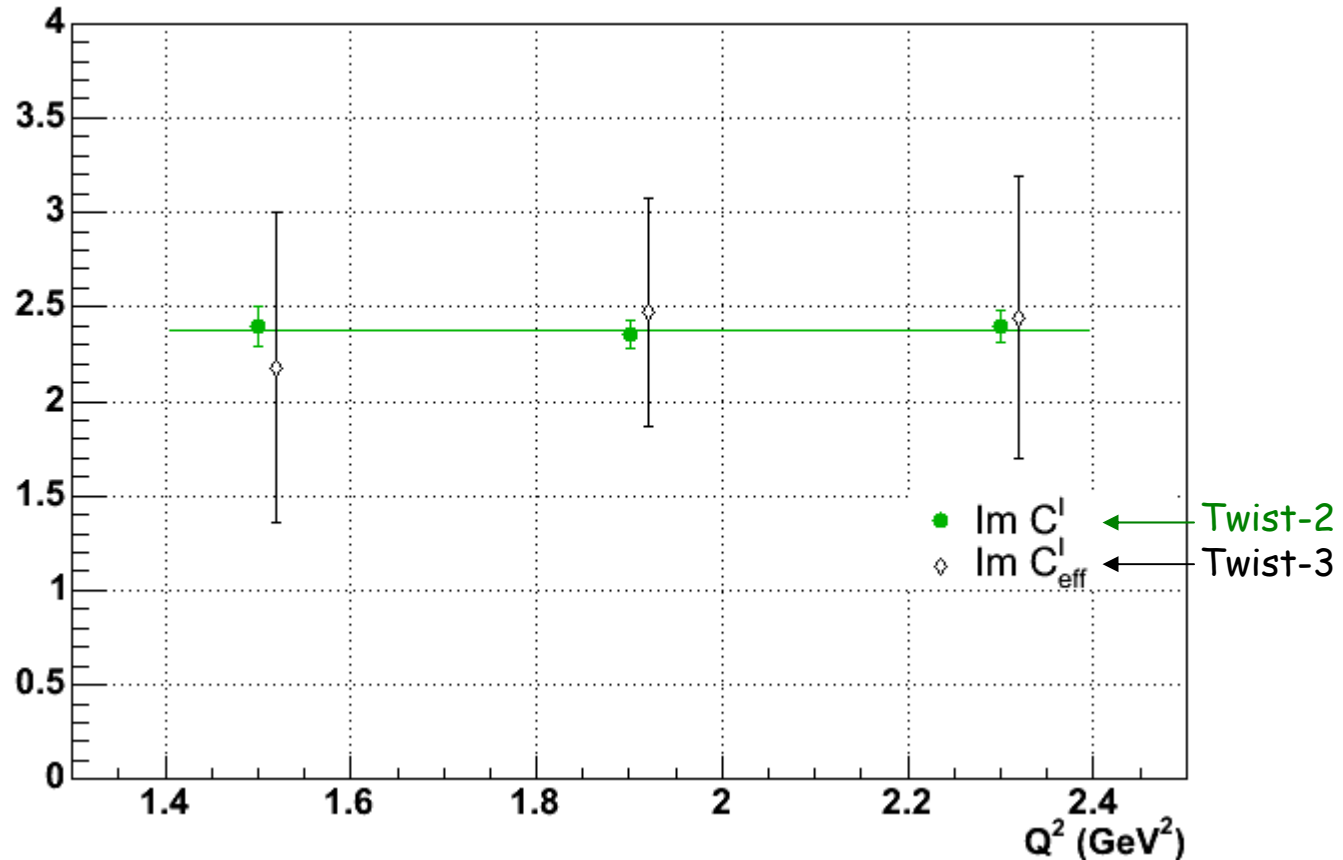
Corrected for real+virtual RC
 Corrected for efficiency
 Corrected for acceptance
 Corrected for resolution effects
 Checked elastic cross-section @ ~1%

Extracted Twist-3
 contribution small!

New work by P. Guichon !

Q^2 dependence and test of scaling

$$\langle -t \rangle = 0.26 \text{ GeV}^2, \langle x_B \rangle = 0.36$$



No Q^2 dependence: strong indication for scaling behavior and handbag dominance

Twist 4+ contributions are smaller than 10%

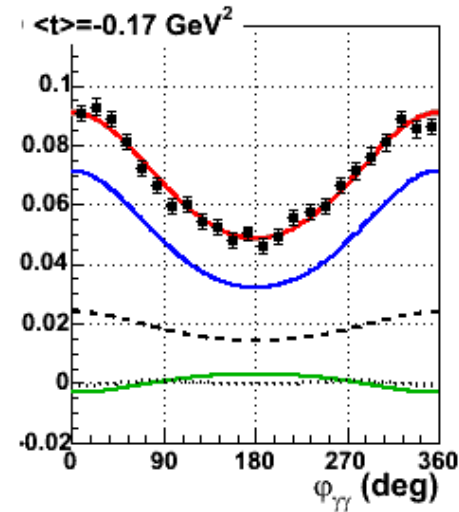
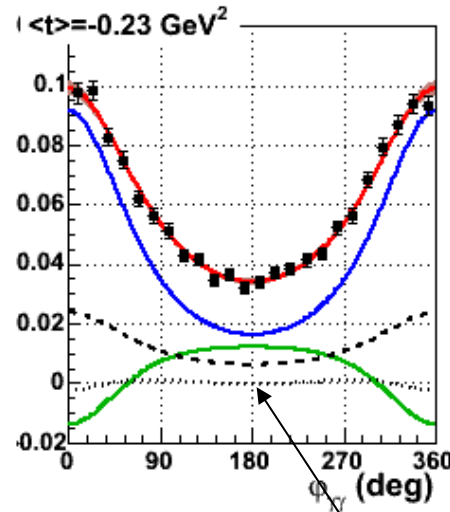
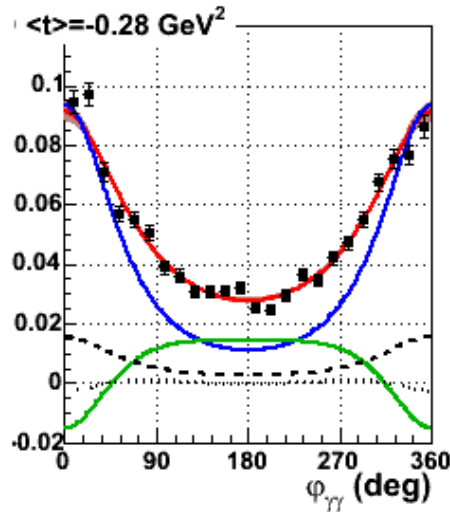
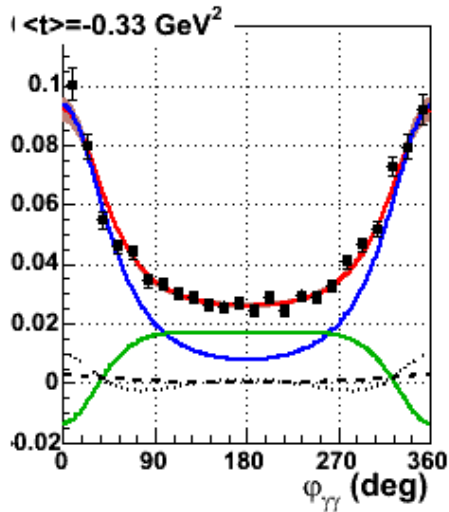
Total cross-section

PRL97, 262002 (2006)

$$\langle Q^2 \rangle = 2.3 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.36$$

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi_{\gamma\gamma}} \text{ (nb/GeV}^4\text{)}$$



Corrected for real+virtual RC
 Corrected for efficiency
 Corrected for acceptance
 Corrected for resolution effects

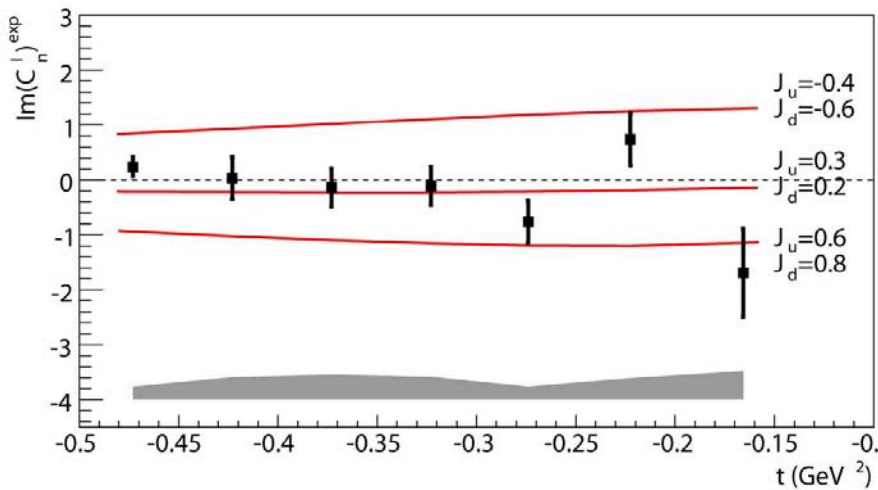
Extracted Twist-3
 contribution small !

but impossible to disentangle DVCS²
 from the interference term
 (more on this in J. Roche talk)

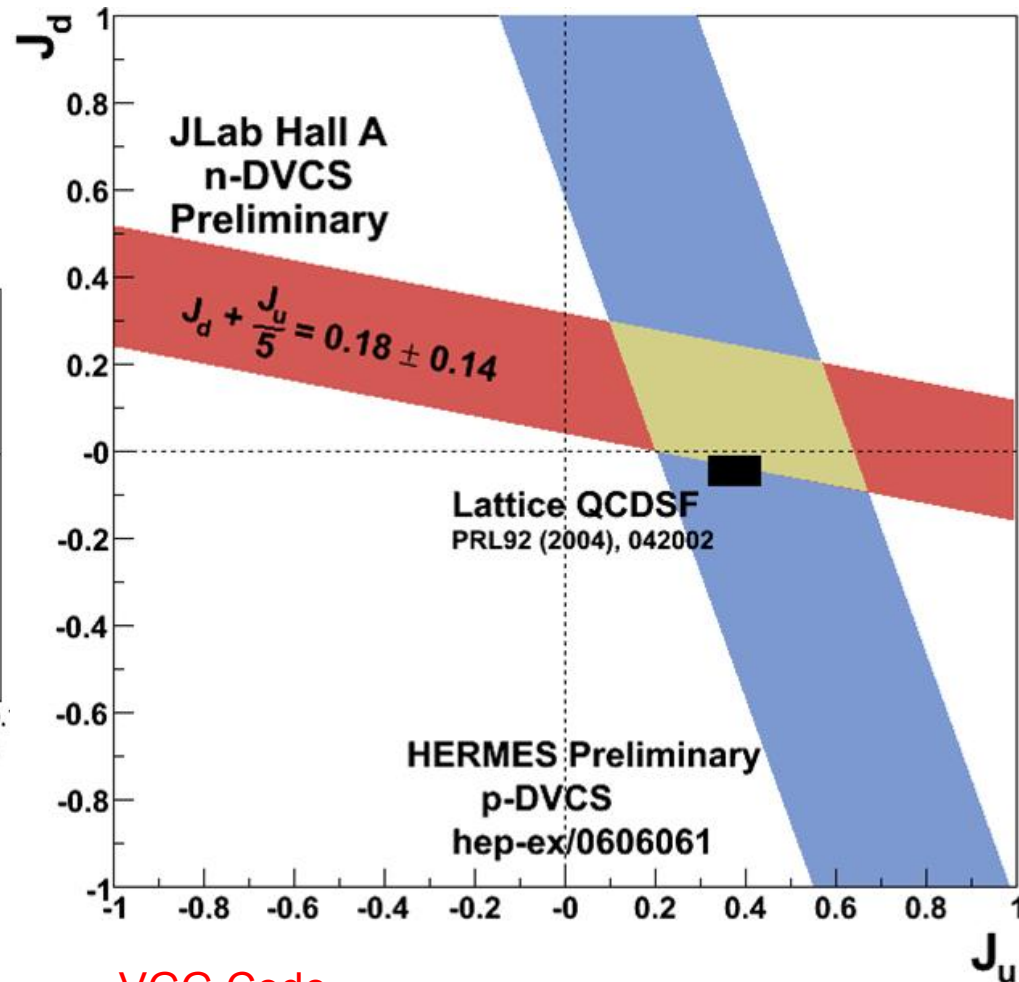
DVCS on the neutron in JLab/Hall A: E03-106

See talk from M. Mazouz

LD₂ target
24000 fb⁻¹
 $x_B=0.36, Q^2=1.9 \text{ GeV}^2$



MODEL-DEPENDENT
J_u-J_d extraction



VGG Code

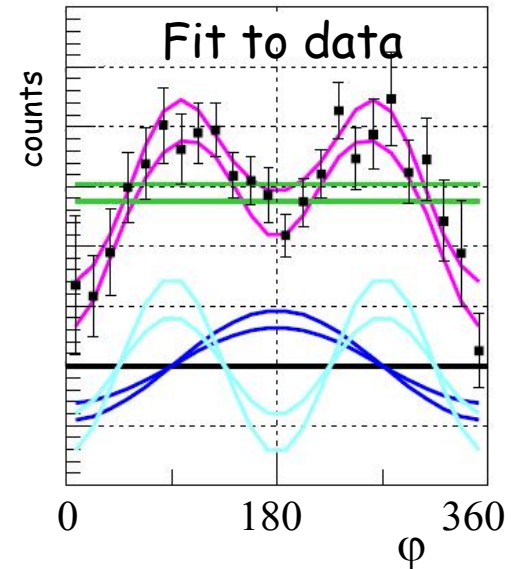
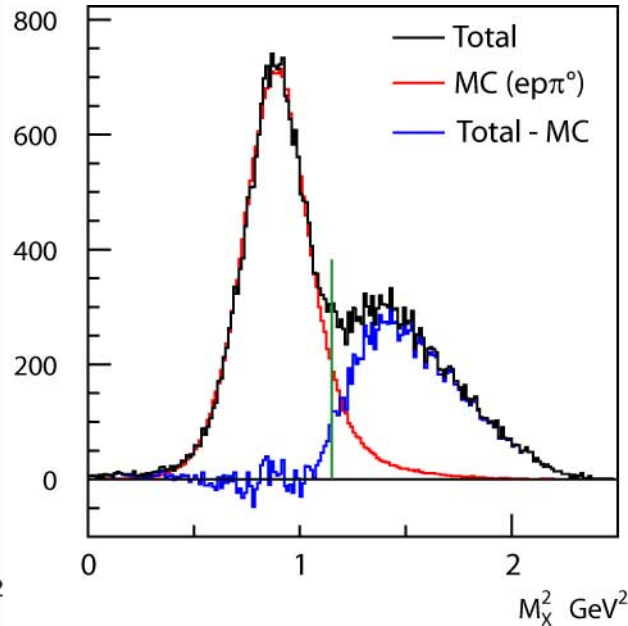
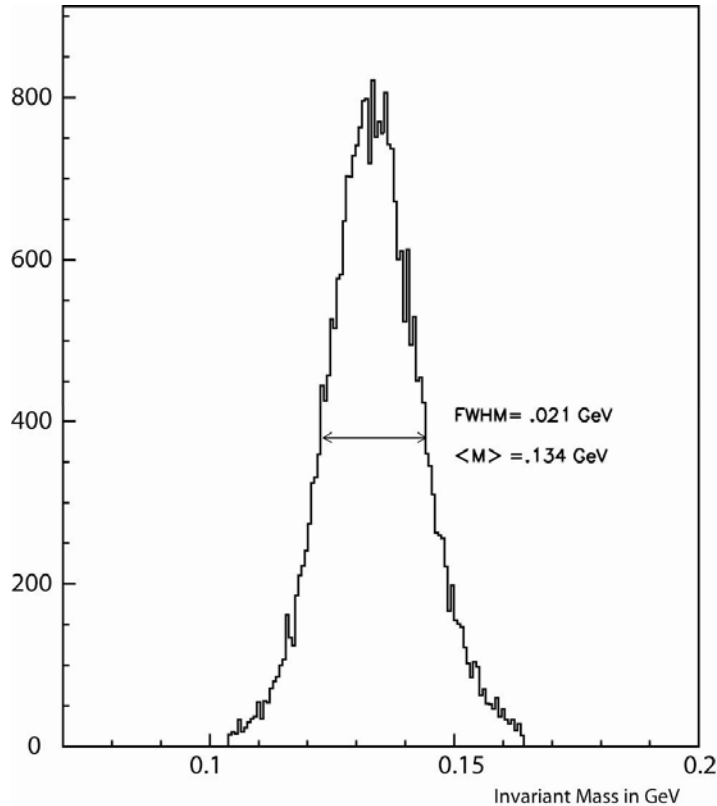
GPD model : LO/Regge/D-term=0

Goeke et al., Prog. Part. Nucl. Phys 47 (2001), 401.

Deep- π^0 electroproduction in JLab/Hall A

Same data as E00-110, but:

- 2γ requirement in the calorimeter at π^0 mass,
- $e\pi^0X$ at proton mass.



$$\frac{d\sigma}{dt} = \frac{d\sigma_T}{dt} + \varepsilon \frac{d\sigma_L}{dt} + \sqrt{2\varepsilon(1+\varepsilon)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

Cross-sections for deep- π^0 production in JLab/Hall A

$$\langle Q^2 \rangle = 2.3 \text{ GeV}^2$$

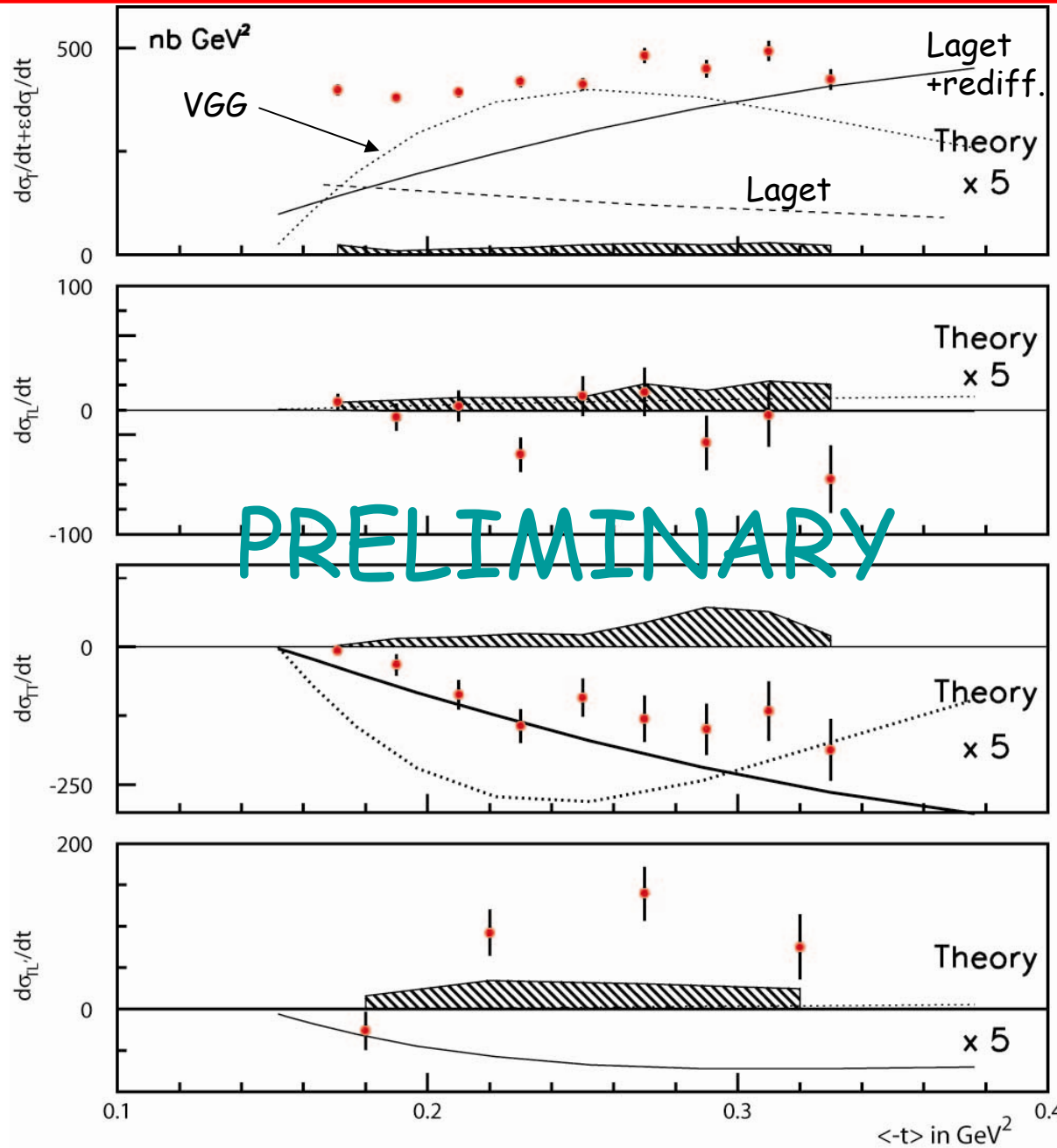
$$\langle x_B \rangle = 0.36$$

-Low- t cross-section largely overshoots GPD & JML models.

-TT term is large, suggesting a large transverse component.

-But π^0 cross-section ratio for proton and deuteron targets is found 0.95...

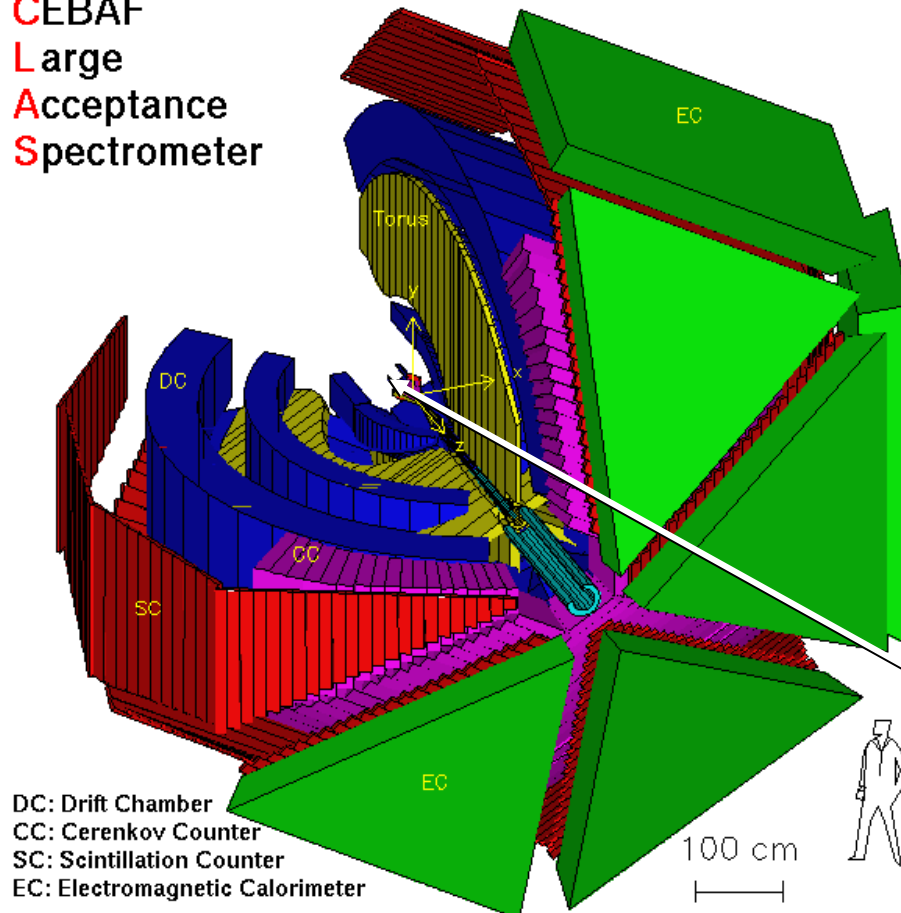
More data needed!
(J. Roche talk)



E1-DVCS with CLAS : a dedicated DVCS experiment in Hall B

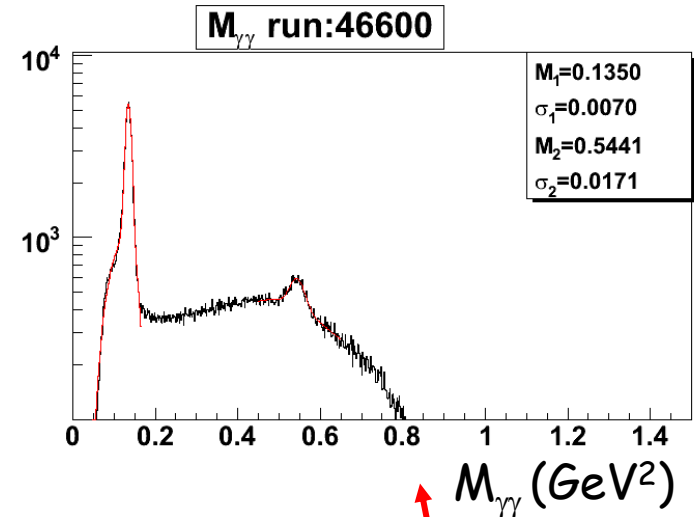
Beam energy: ~ 5.8 GeV
Beam Polarization: 75-85%
Integ. Luminosity: 45 fb^{-1}

CEBAF
Large
Acceptance
Spectrometer

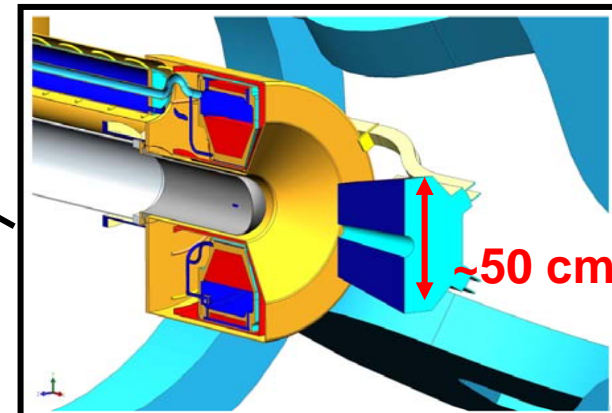


DC: Drift Chamber
CC: Cerenkov Counter
SC: Scintillation Counter
EC: Electromagnetic Calorimeter

100 cm

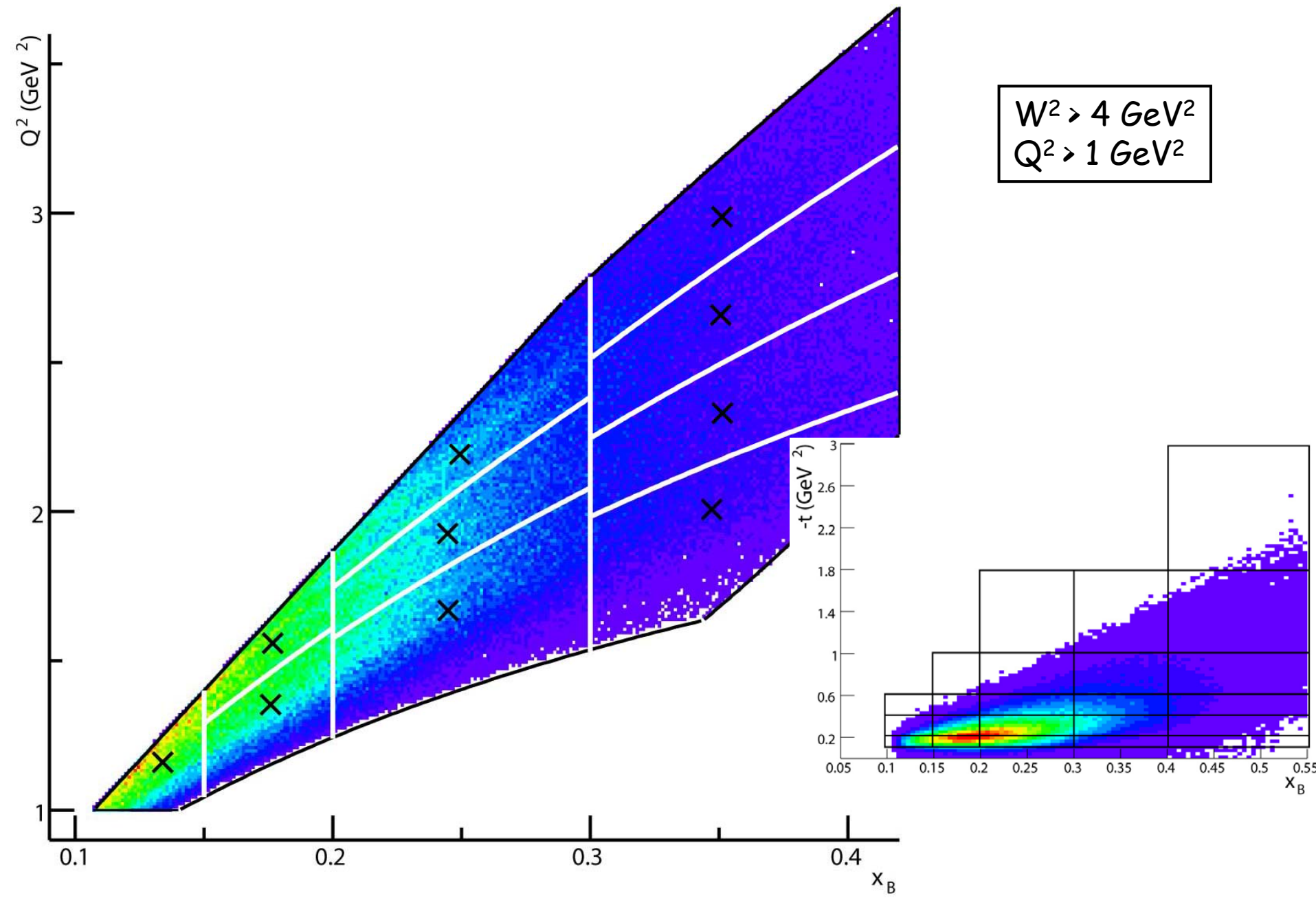


Inner Calorimeter
+ Moller shielding solenoid

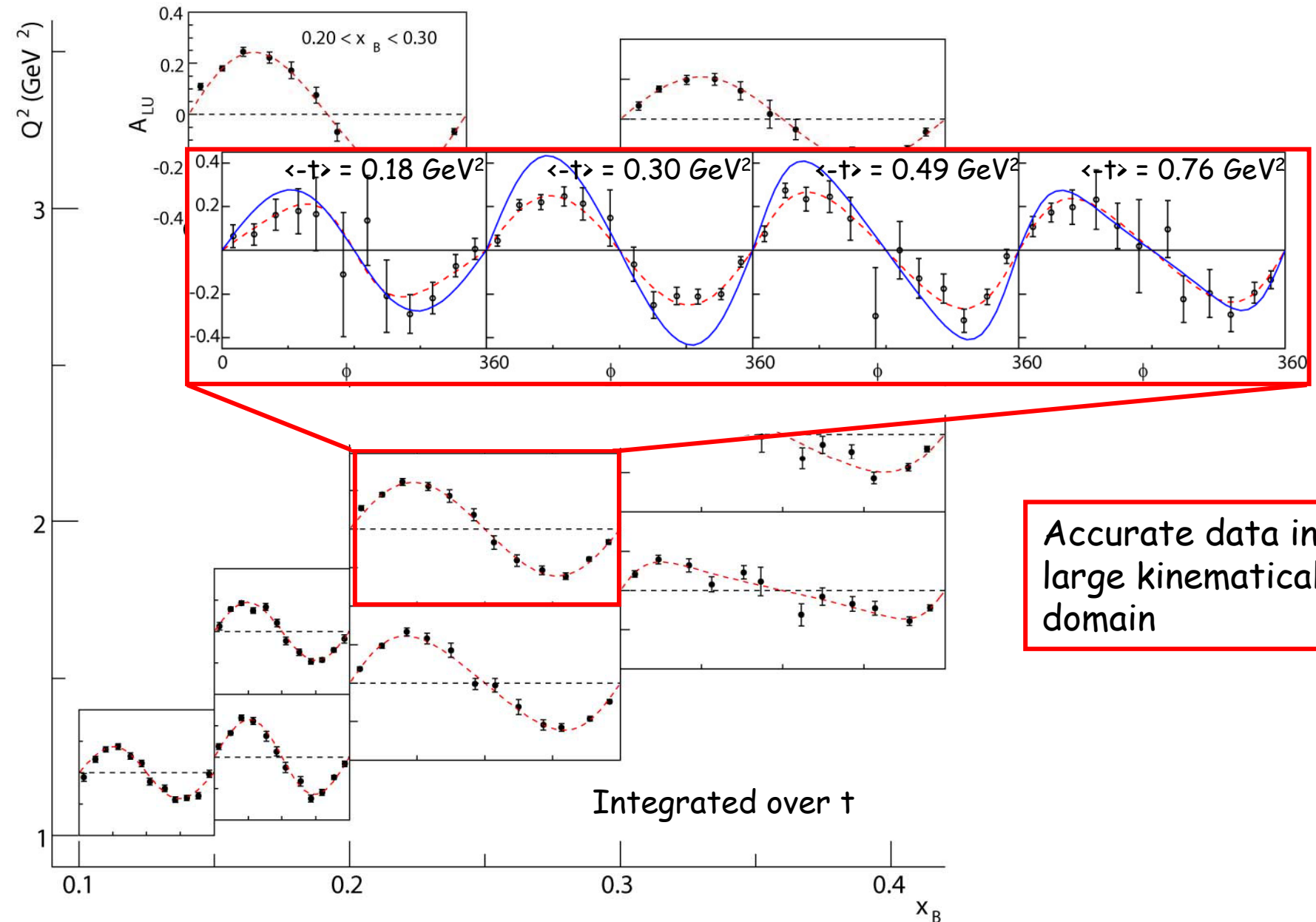


More details in FX Girod's talk

E1-DVCS kinematical coverage and binning

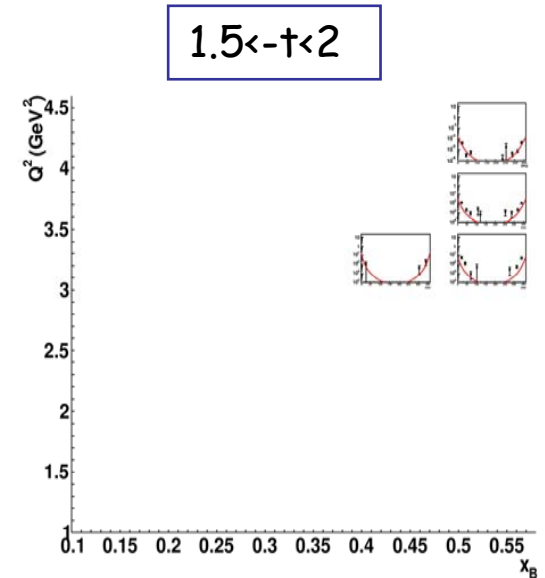
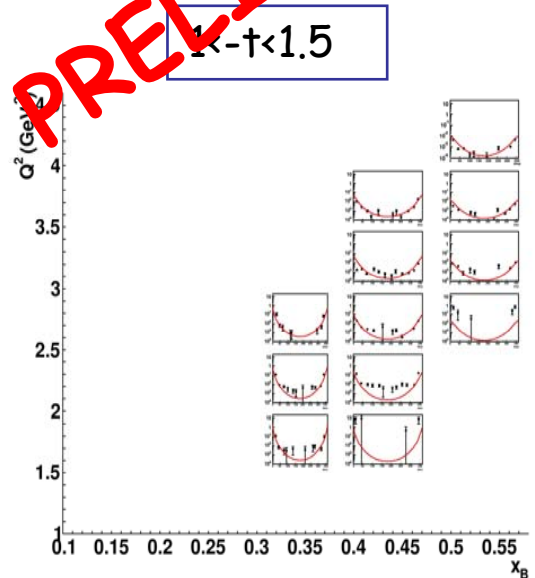
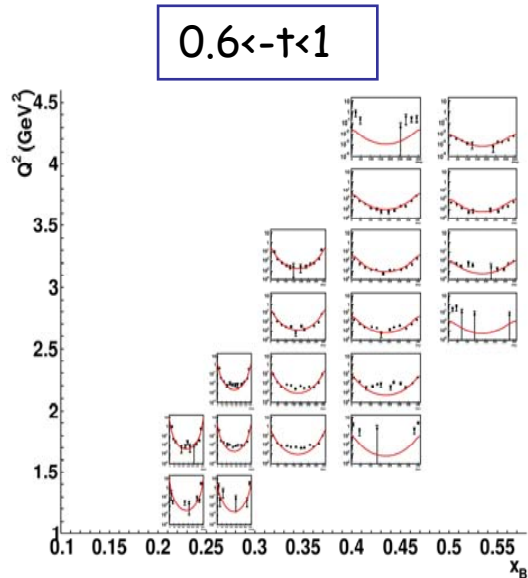
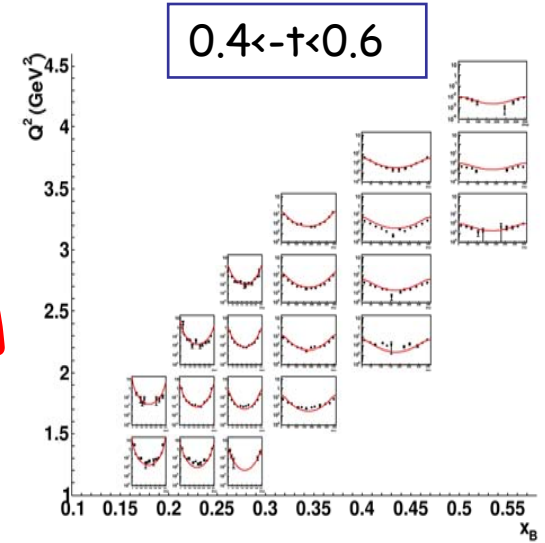
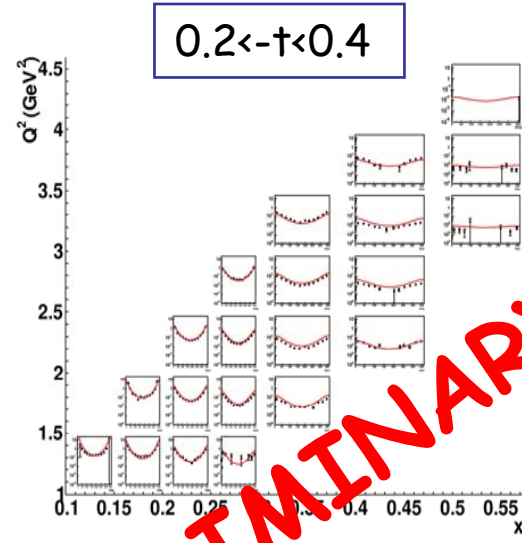
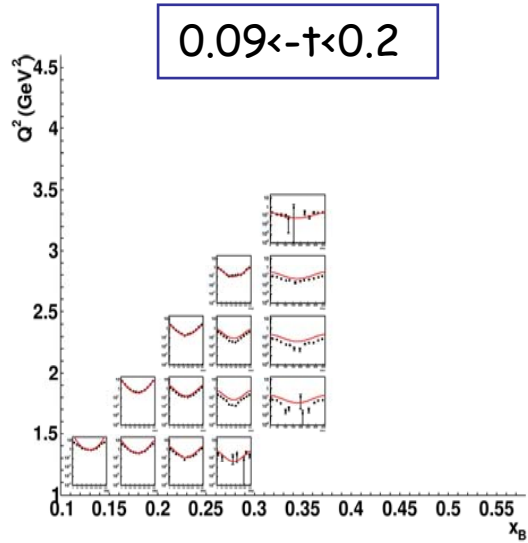


E1-DVCS : Asymmetry as a function of x_B and Q^2



E1-DVCS : Cross-sections over a wide kinematical range

PhD Thesis H.S. Jo (IPNO)



PRELIMINARY

What I did not have time to talk about

- ❑ Unpublished non-dedicated data from CLAS (G. Gavalian, H. Avakian, ...)
- ❑ More details about E1-dvcs in Hall B (FX Girod talk)
- ❑ More details about E03-106 in Hall A (M. Mazouz talk)
- ❑ Plans for the short and long term studies of DVCS at JLab in both Hall A (J. Roche talk) and B (L. Elouadrhiri talk)

Summary

DVCS BSA (Hall B/CLAS):

- ❑ Data in a large kinematical range and good statistics. It will give a very hard time to models, to fit the whole range in Q^2 , x_B and t .
- ❑ Data seem to favor JM Laget model at low- t and VGG model at high t .

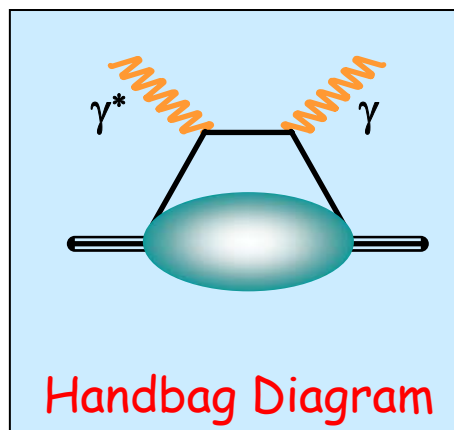
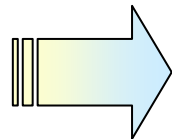
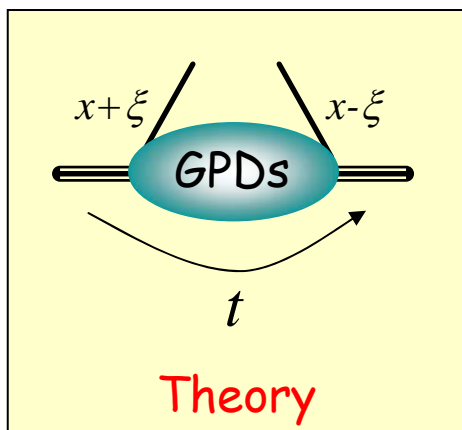
DVCS Cross-section difference (Hall A):

- ❑ High statistics test of scaling: **Strong support for twist-2 dominance**
- ❑ First model-independent extraction of GPD linear combination from DVCS data in the twist-3 approximation
- ❑ Upper limit set on twist-4+ effects in the cross-section difference: **twist \geq 3 contribution is smaller than 10%**

DVCS Total cross-section (Hall A):

- ❑ Bethe-Heitler is not dominant everywhere
- ❑ $|DVCS|^2$ terms might be sizeable (more on this in J. Roche's talk)

GPDs from Theory to Experiment



2. The GPDs enter the DVCS amplitude as an integral over x :

- GPDs appear in the **real part** through a PP integral over x
- GPDs appear in the **imaginary part** but at the line $x=\xi$

$$T^{DVCS} = \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi + i\epsilon} dx + \dots$$

$$= P \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi} dx - i\pi GPD(x = \xi, \xi, t) + \dots$$

Observables and their relationship to GPDs

$$T^{DVCS} = \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi + i\varepsilon} dx + \dots$$

The **cross-section difference** accesses the imaginary part of DVCS and therefore GPDs at $x = \xi$

$$= P \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi} dx - i\pi GPD(x = \xi, \xi, t) + \dots$$

The **total cross-section** accesses the real part of DVCS and therefore an integral of GPDs over x

