

# Deeply Virtual Compton Scattering on the neutron

**Jefferson Lab**  
Thomas Jefferson National Accelerator Facility  
**Hall A**

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*For JLab Hall A & DVCS collaborations*

- Physics case
- n-DVCS experimental setup
- Analysis method
- Results and conclusions

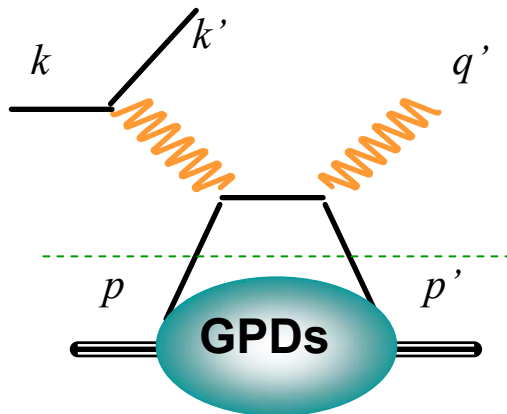
# Deeply Virtual Compton Scattering

GPDs give an access to **quark angular momentum** (Ji's sum rule)

$$J_q = \frac{1}{2} \Delta \Sigma_q + L_q = \frac{1}{2} \int_{-1}^1 x dx \left[ H^q(x, \xi, 0) + E^q(x, \xi, 0) \right]$$

**less constrained GPD** ← No link to DIS

➡ DVCS is the **simplest hard exclusive process** involving GPDs



Factorization theorem  
in the Bjorken regime

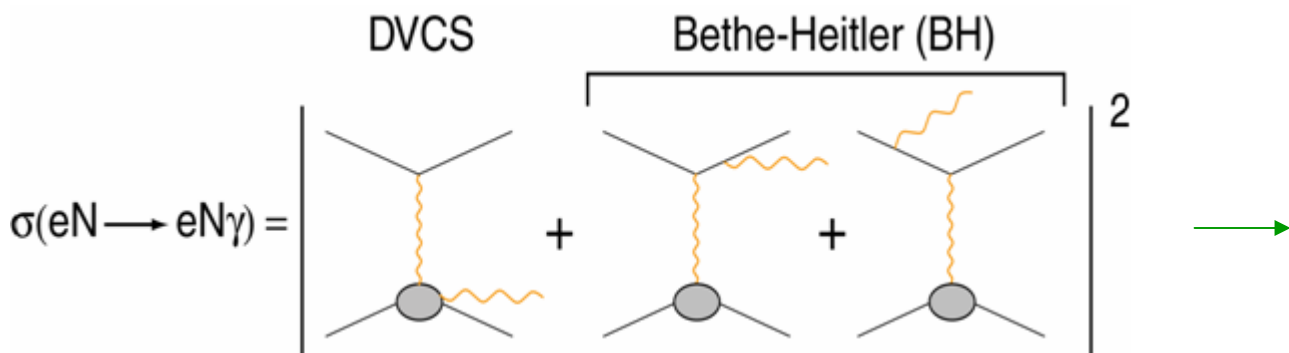
$$Q^2 = -q^2 = -(k - k')^2 \gg M^2$$

$$t = (p - p')^2 = \Delta^2 \ll Q^2$$



Non perturbative  
description by **GPDs**

# DVCS and Bethe-Heitler



The **total cross-section** accesses the **real** part of **DVCS** and therefore an **integral of GPDs over x**

The **polarized cross-section difference** accesses the **Imaginary** part of **DVCS** and therefore **GPDs at  $x=\pm\xi$**

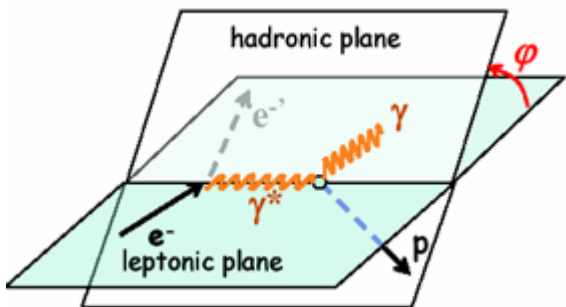
$$d^5 \vec{\sigma} - d^5 \bar{\sigma} = 2 \Im m(T^{BH} T^{DVCS}) + \left[ |\vec{T}^{DVCS}|^2 - |\vec{T}^{DVCS}|^2 \right]$$

Purely real and fully calculable

If handbag dominance **Small** at  $\perp$  lab energies (twist-3 term)

$$d^5 \vec{\sigma} - d^5 \bar{\sigma} = \Gamma(x_B, Q^2, t, \varphi) \Im m(C^I) \sin \varphi$$

$$C^I(F) = F_1(t) \mathcal{H} + \xi (F_1(t) + F_2(t)) \tilde{\mathcal{H}} - \frac{t}{4M^2} F_2(t) \mathcal{E}$$



# Neutron Target

## Model:

(Goeke, Polyakov  
and Vanderhaeghen)

Target	$\mathcal{H}$	$\tilde{\mathcal{H}}$	$\mathcal{E}$
neutron	0.81	-0.07	1.73

$$Q^2 = 2 \text{ GeV}^2$$

$$x_B = 0.3$$

$$-t = 0.3 \text{ GeV}^2$$

$$\Im(C^I) = \underbrace{F_1(t) \cdot \mathcal{H}}_{\text{Term 1}} + \underbrace{\frac{x_B}{2-x_B} \cdot (F_1(t) + F_2(t)) \cdot \tilde{\mathcal{H}}}_{\text{Term 2}} - \underbrace{\frac{t}{4M^2} F_2(t) \cdot \mathcal{E}}_{\text{Term 3}}$$

$-t$	$F_2^n(t)$	$F_1^n(t)$	$(F_1^n(t) + F_2^n(t)) \cdot x_B / (2 - x_B)$	$(-t / 4M^2) \cdot F_2^n(t)$
0.3	-0.91	-0.04	-0.17	-0.07

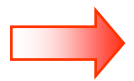
→

$$\Im(C^I) = \cancel{F_1(t) \cdot \mathcal{H}} + \frac{x_B}{2-x_B} \cdot \cancel{(F_1(t) + F_2(t)) \cdot \tilde{\mathcal{H}}} - \frac{t}{4M^2} F_2(t) \cdot \mathcal{E}$$

$$\Im(C^I) = \cancel{-0.03} + \cancel{0.01} - 0.13$$

# n-DVCS experiment

An **exploratory** experiment was performed at JLab Hall A on **hydrogen** target and **deuterium** target with **high luminosity** ( $4 \cdot 10^{37} \text{ cm}^{-2} \text{ s}^{-1}$ ) and **exclusivity**.



**Goal** : Measure the n-DVCS polarized cross-section difference which is mostly sensitive to **GPD E** (less constrained!)



E03-106 (n-DVCS) followed directly E00-110 (p-DVCS) which shows strong indications of handbag dominance at  $Q^2$  about  $2 \text{ GeV}^2$ .

(C. Muñoz-Camacho et al., PRL 97 (2006) 262002.)

$x_{Bj}=0.364$

s (GeV <sup>2</sup> )	Q <sup>2</sup> (GeV <sup>2</sup> )	P <sub>e</sub> (Gev/c)	Θ <sub>e</sub> (deg)	-Θ <sub>γ*</sub> (deg)	∫ Ldt (fb <sup>-1</sup> )
4.22	1.91	2.95	19.32	18.25	4365
4.22	1.91	2.95	19.32	18.25	24000

Hydrogen

Deuterium

# Experimental apparatus

Beam energy = 5.75 GeV

Beam polarization = 75%

Beam current =  $\sim 4 \mu\text{A}$

Luminosity =  $4 \cdot 10^{37} \text{ cm}^{-2} \cdot \text{s}^{-1} \text{ nucleon}^{-1}$



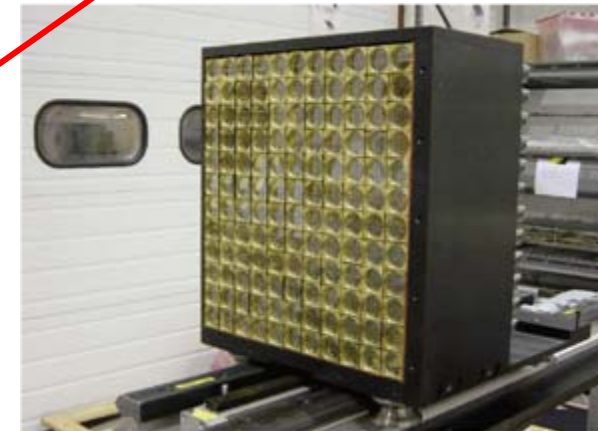
Scattered Electron

Left HRS

DVCS events are identified with  $M_x^2$

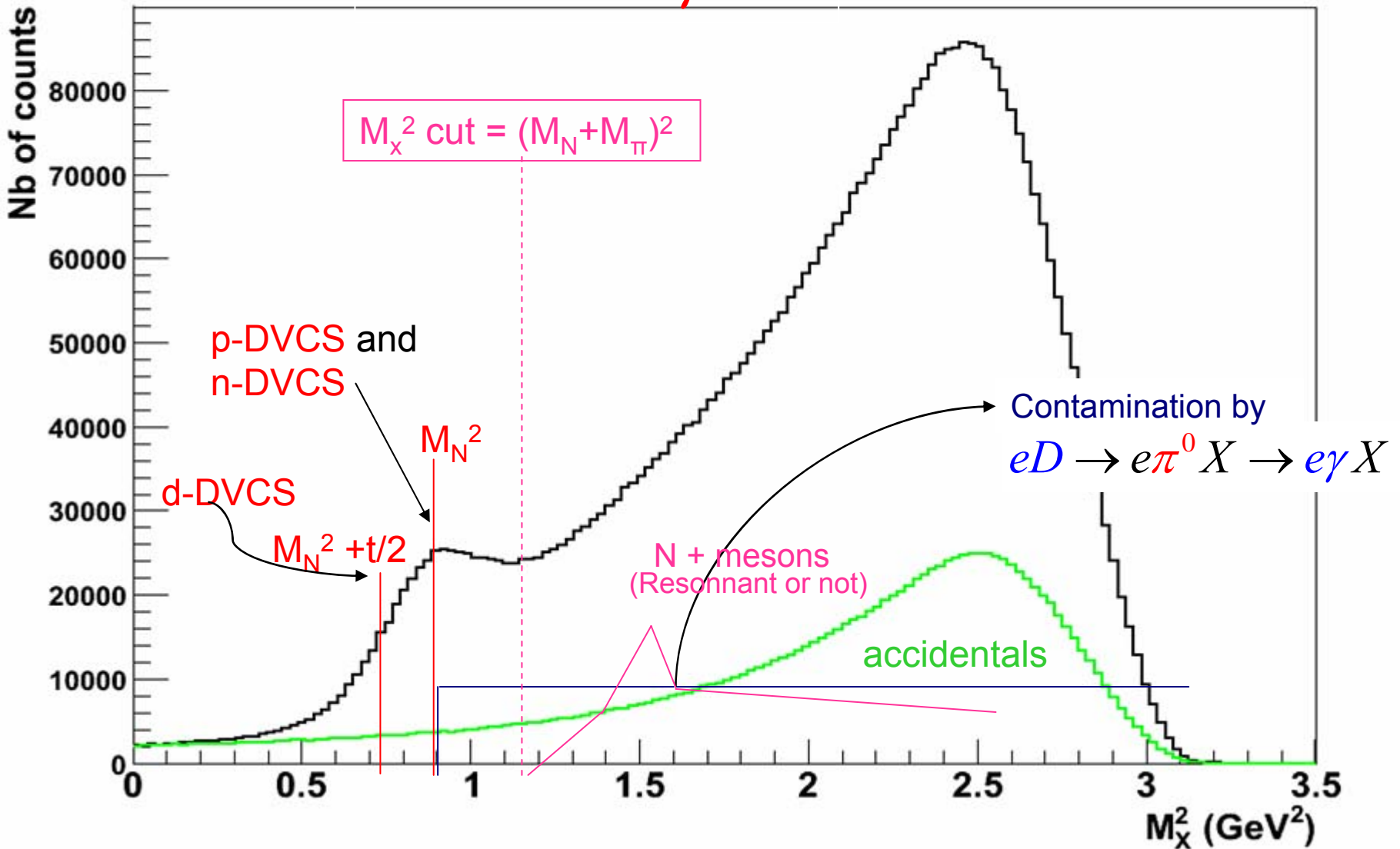
An experimental challenge

- Specific Scattering Chamber
- Čerenkov based Electromagnetic Calorimeter
- Customized Electronics & Data Acquisition

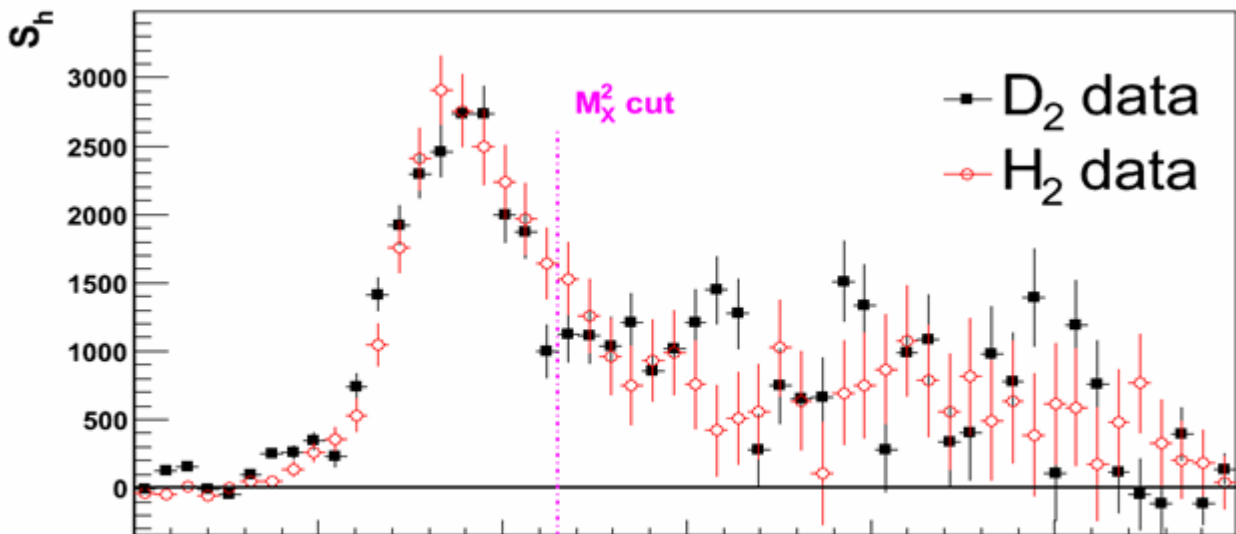


Electromagnetic Calorimeter

# Analysis method

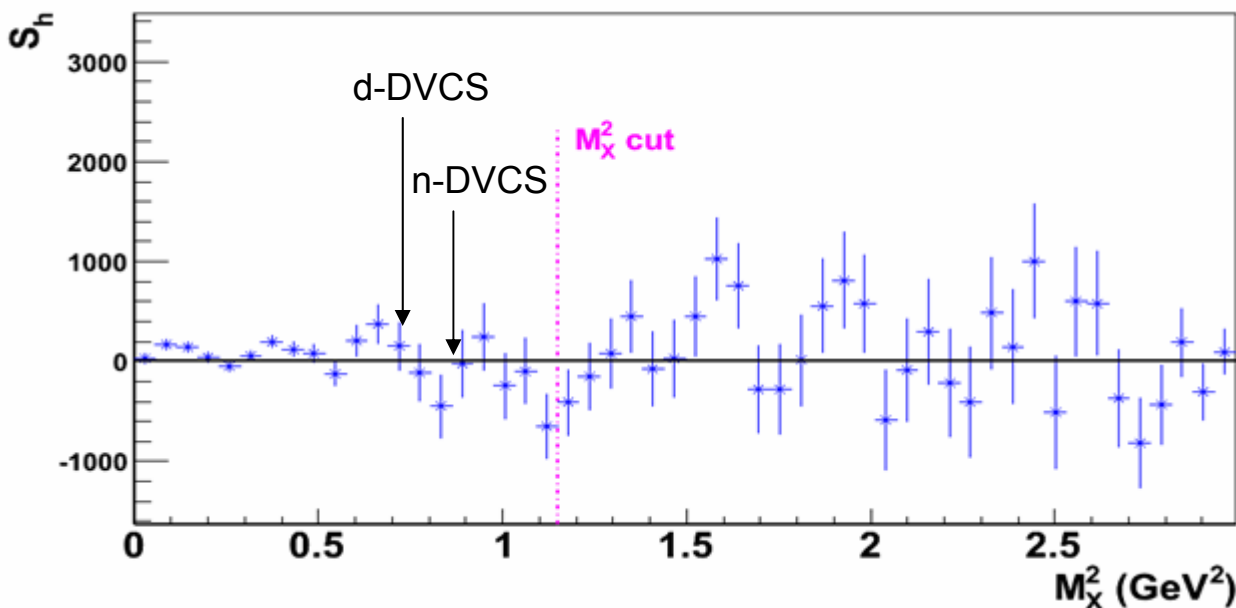


# Helicity signal and exclusivity



After :

- Normalizing  $H_2$  and  $D_2$  data to the same luminosity
- Adding Fermi momentum to  $H_2$  data



2 principle sources of systematic errors :

- The contamination of  $\pi^0$  electroproduction on the neutron (and deuteron).
- The uncertainty on the relative calibration between  $H_2$  and  $D_2$  data



# Extraction of observables

$$\frac{1}{2} \left[ \frac{d\bar{\sigma}}{dQ^2 dx_B d\Delta^2 d\varphi_e d\varphi_{\gamma\gamma}} - \frac{d\bar{\sigma}}{dQ^2 dx_B d\Delta^2 d\varphi_e d\varphi_{\gamma\gamma}} \right] =$$

$$\Gamma_n(x_B, \varphi_e, \Delta^2, \varphi) \cdot \Im(C_n^{I-\text{exp}}) \sin \varphi + \Gamma_d(x_B, \varphi_e, \Delta^2, \varphi) \cdot \Im(C_d^{I-\text{exp}}) \sin \varphi$$

A. V. Belitsky, D. Muller, A. Kirchner, Nucl. Phys. B629, 323 (2002).

$$\Delta N^{\text{Exp}}(i_e) = N_{i_e}^+ - N_{i_e}^- \quad \text{with } i_e = 20 \otimes 12 \otimes 7 \text{ bins in } (M_X^2, \varphi, t)$$

$$\Delta N^{\text{MC}}(i_e) = L \left[ \underbrace{\Im(C_n^{I-\text{exp}}) \int_{x \in i_e} \Gamma_n \cdot \sin \varphi \otimes \text{Acc}}_{\text{MC sampling}} + \underbrace{\Im(C_d^{I-\text{exp}}) \int_{x \in i_e} \Gamma_d \cdot \sin \varphi \otimes \text{Acc}}_{\text{MC sampling}} \right]$$

Luminosity
MC sampling
MC sampling

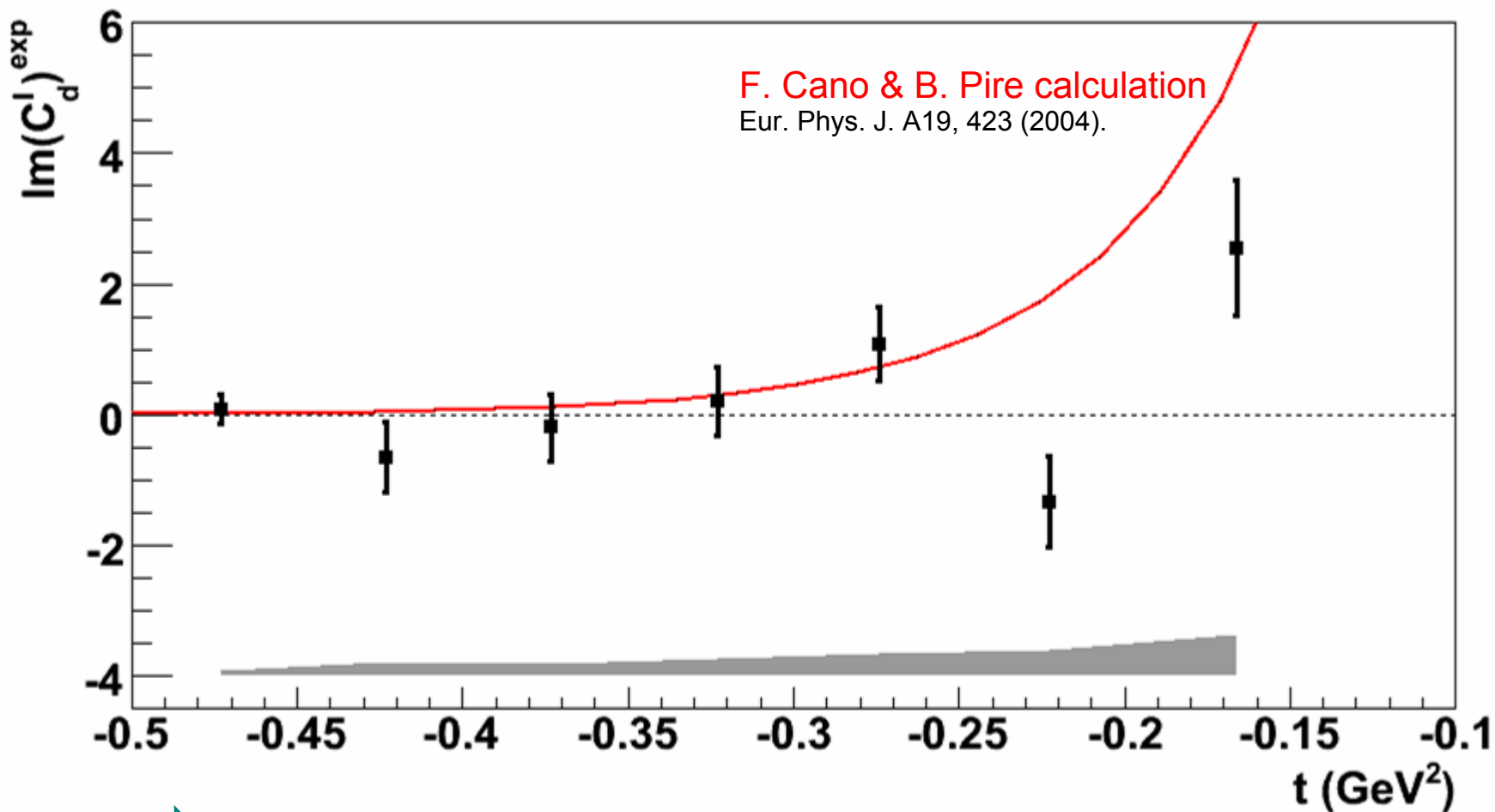
MC includes real radiative corrections (external+internal)

$$\chi^2 = \sum_{i_e} \frac{[\Delta N^{\text{Exp}}(i_e) - \Delta N^{\text{MC}}(i_e)]^2}{[\sigma^{\text{Exp}}(i_e)]^2} \quad \longrightarrow \quad \begin{cases} \Im(C_n^{I-\text{exp}}) \\ \Im(C_d^{I-\text{exp}}) \end{cases}$$

# Extraction results

d-DVCS extraction results

**PRELIMINARY**



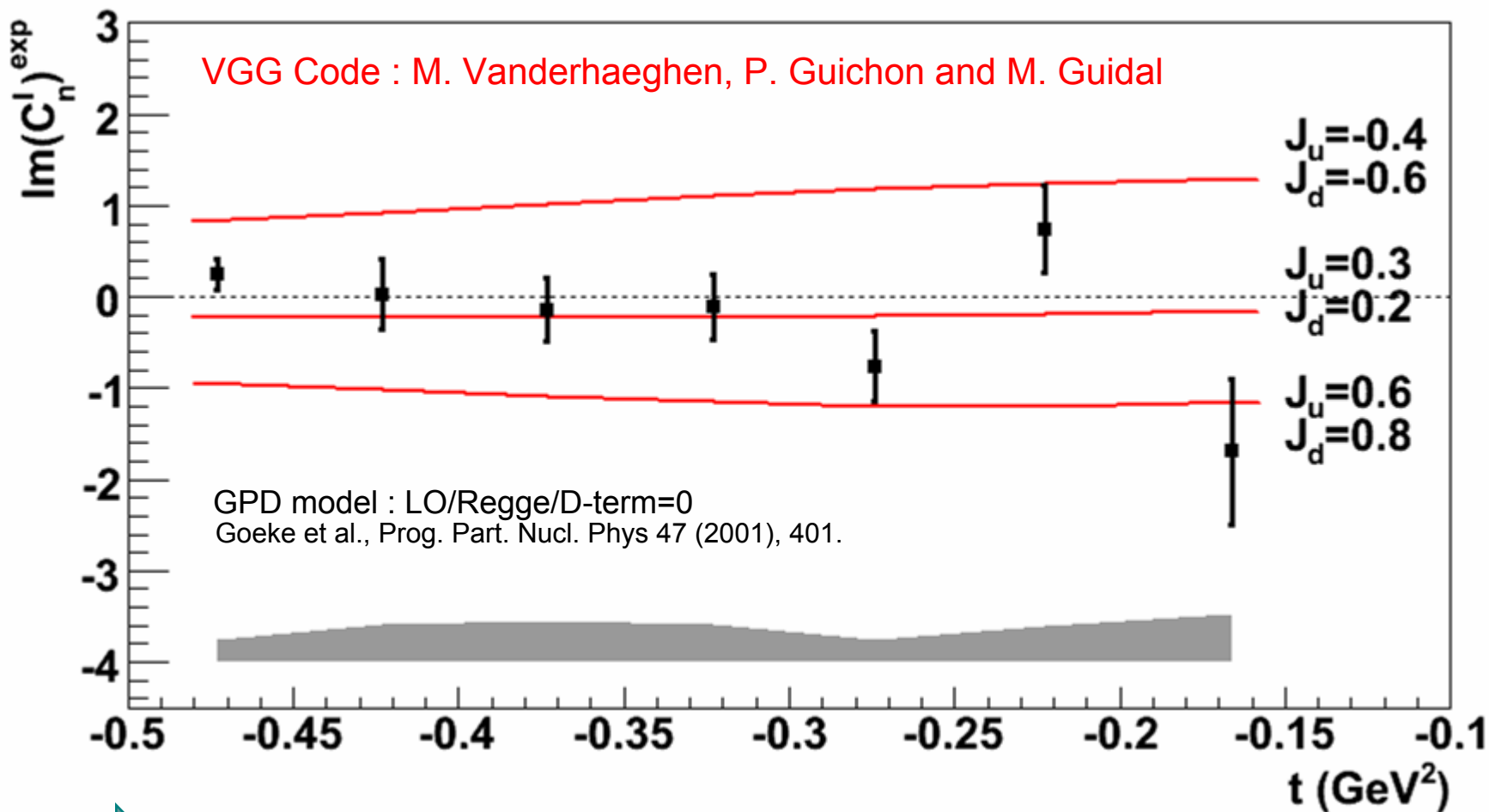
➡ Deuteron moments compatible with zero at large -t

➡ Exploration of small -t regions in future experiments is interesting

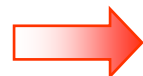
# Extraction results

n-DVCS extraction results

**PRELIMINARY**



Neutron contribution is small and compatible with zero



Results can constrain GPD models (and therefore GPD E)

# n-DVCS experiment results

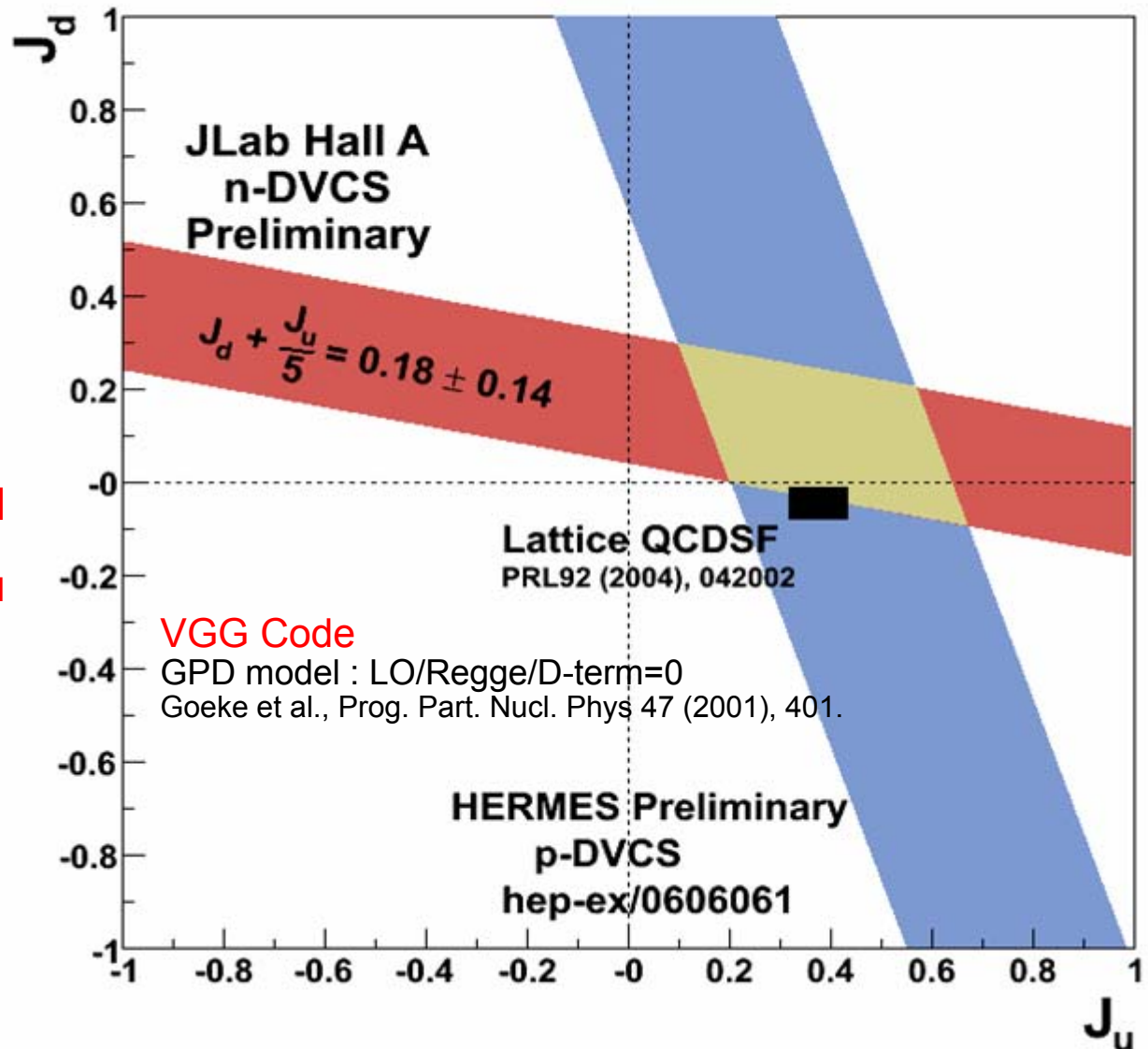
Systematic errors  
of models are not  
shown

n-DVCS is sensitive to  $J_d$

p-DVCS is sensitive to  $J_u$



**Complementarity**  
between **neutron** and  
**transversally polarized**  
**proton** measurements



# Summary and conclusion

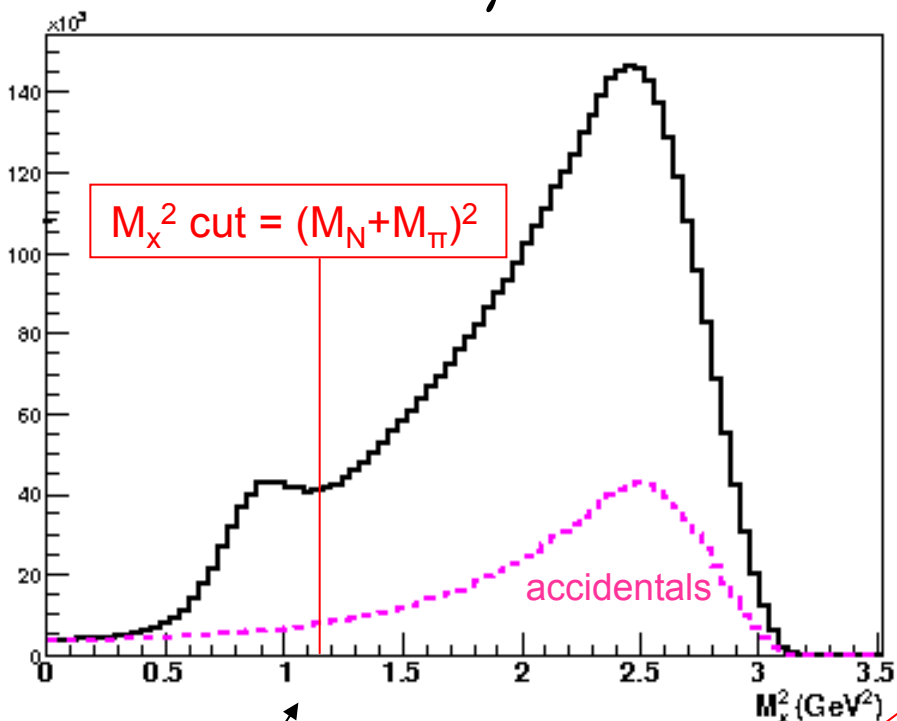
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- ➡ Our experiment is **exploratory** and is dedicated to n-DVCS. **n-DVCS and d-DVCS** contributions are obtained after a subtraction of Hydrogen data from Deuterium data (no recoil detectors needed).
- ➡ n-DVCS and d-DVCS polarized cross-sections difference are compatible with zero.
- ➡ Neutron results can **constrain GPD models** (GPD E parametrization)
- ➡ Neutron has a **different flavor sensitivity** to GPD E than transversally polarized proton.
- ➡ Neutron experiments are mandatory complements to proton ones.
- ➡ **Re(DVCS)** from unpolarized cross-section should be measured.

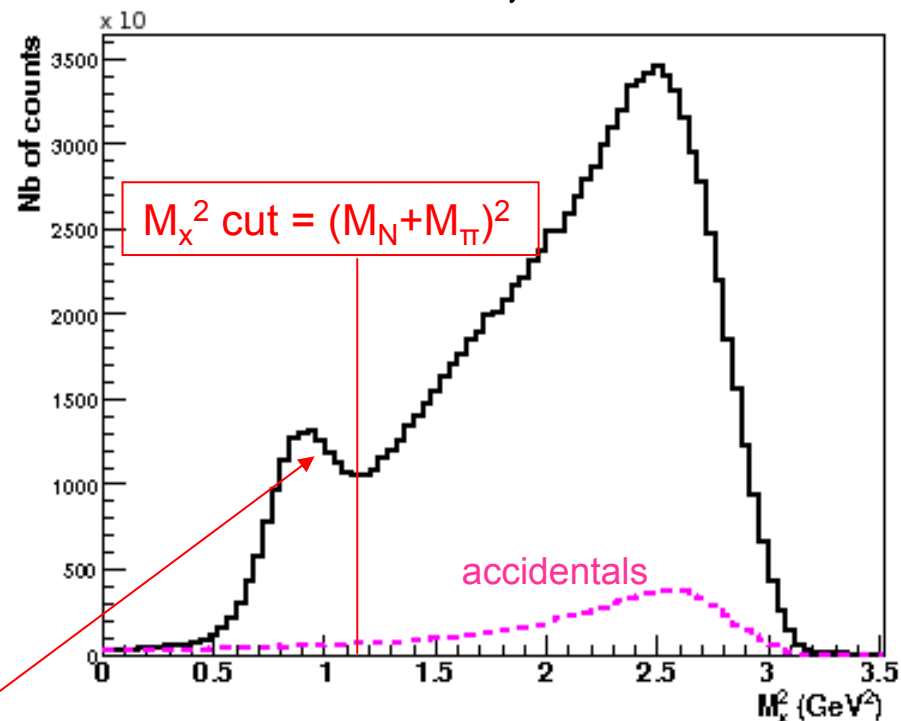


# Analysis method

$eD \rightarrow e\gamma X$



$eH \rightarrow e\gamma X$



$$D(e, e' \gamma) X = p(e, e' \gamma) p + n(e, e' \gamma) n + d(e, e' \gamma) d + \dots$$

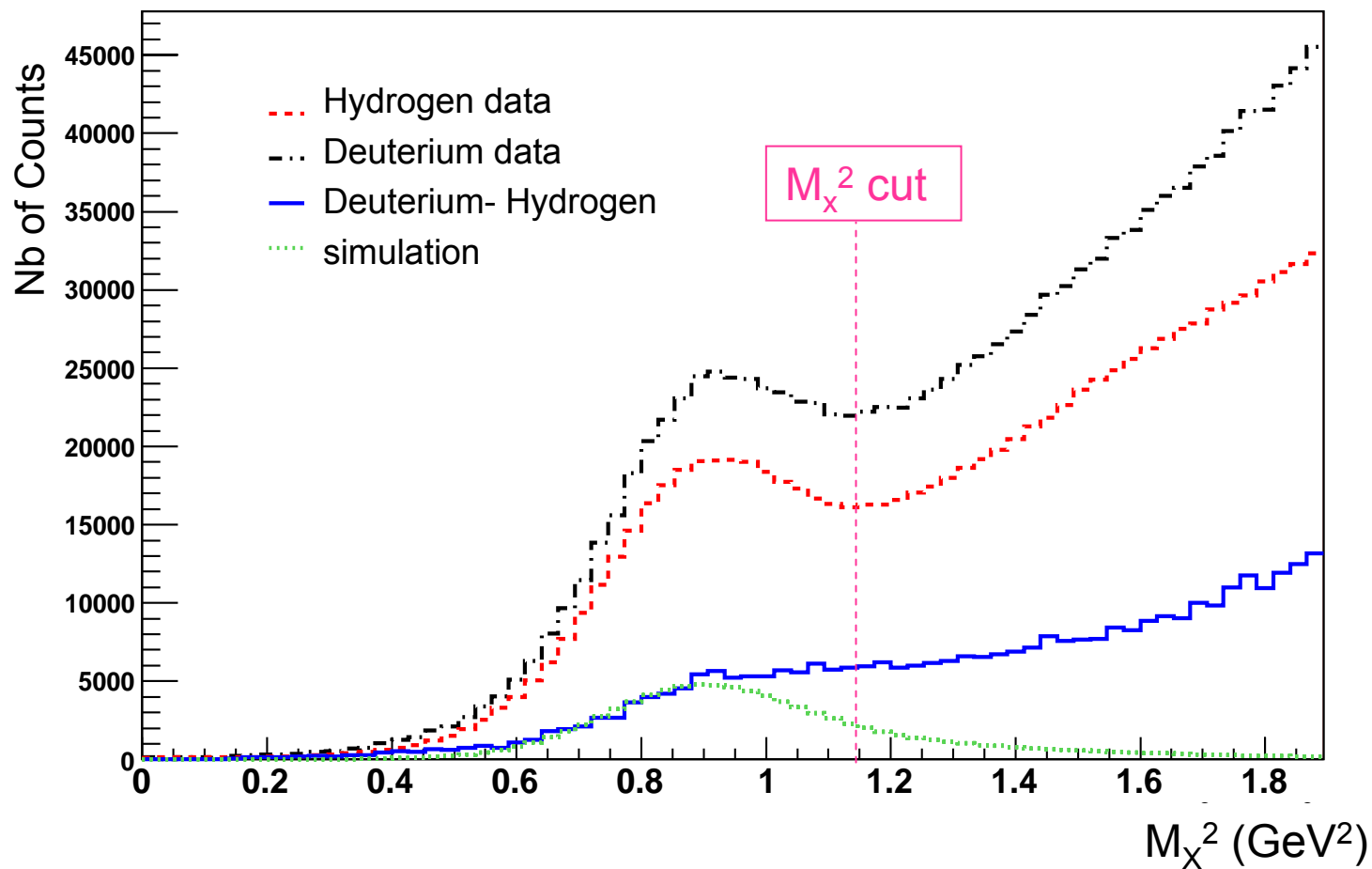
↑  
p-DVCS  
events

↑  
n-DVCS  
events

↑  
d-DVCS  
events

↑  
Mesons  
production

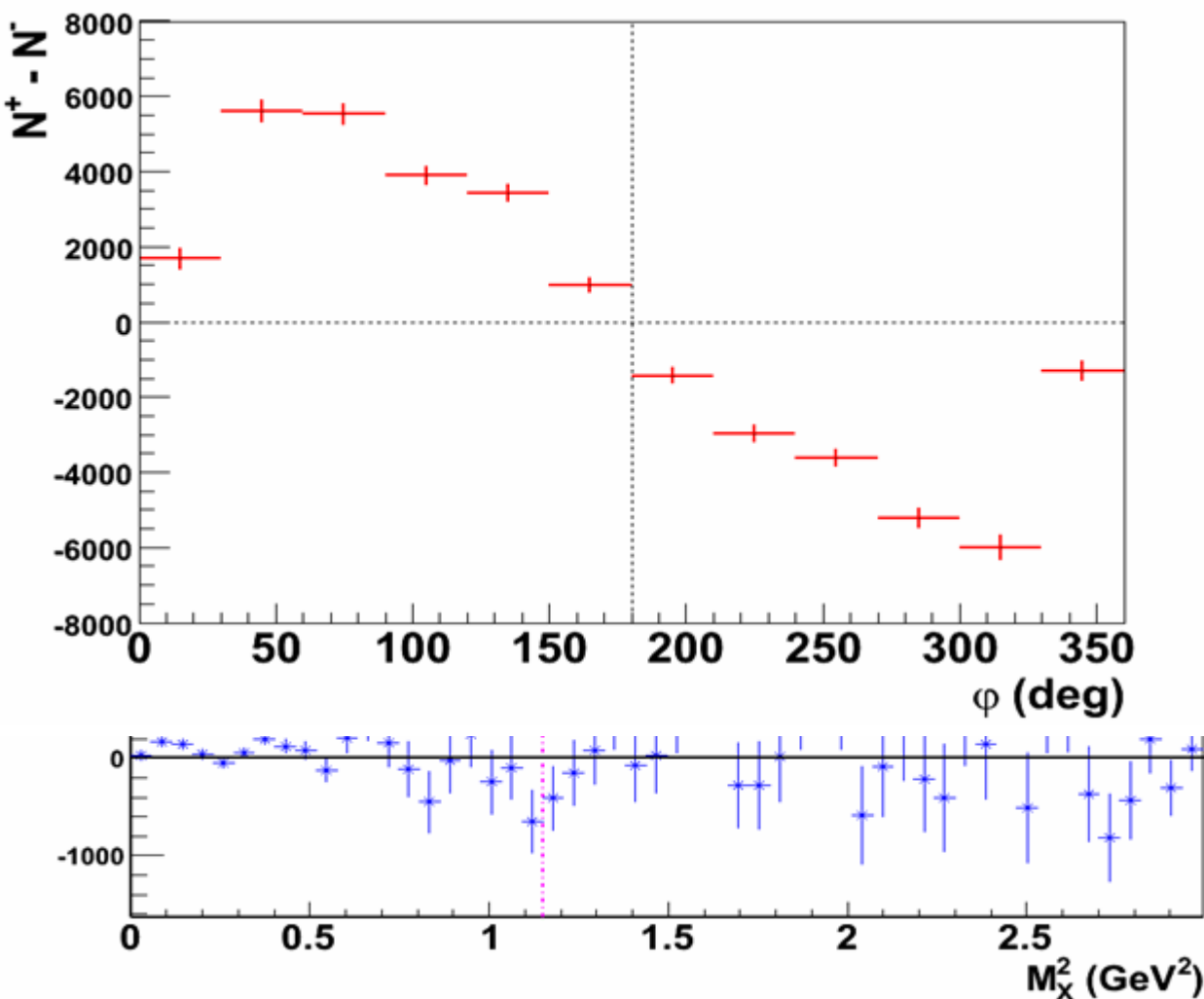
# Double coincidence analysis





# Helicity signal and exclusivity

$$S_h = \int_0^\pi (N^+ - N^-) d\phi - \int_\pi^{2\pi} (N^+ - N^-) d\phi$$



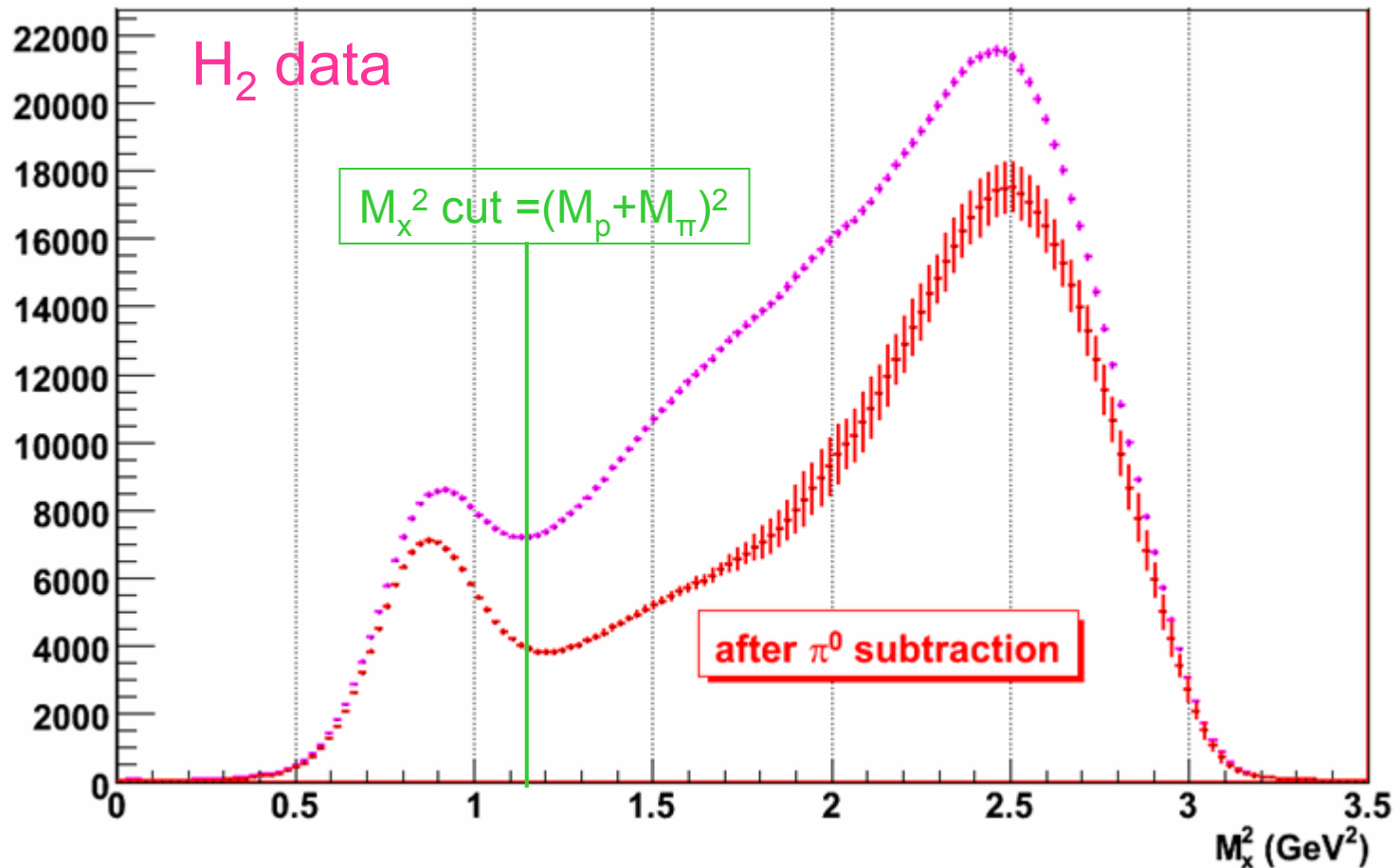
After :

- Normalizing H<sub>2</sub> and D<sub>2</sub> data to the **same luminosity**
- Adding **Fermi momentum** to H<sub>2</sub> data

**2 principle sources of systematic errors :**

- The contamination of  $\pi^0$  electroproduction on the neutron (and deuteron).
- The uncertainty on the relative calibration between H<sub>2</sub> and D<sub>2</sub> data

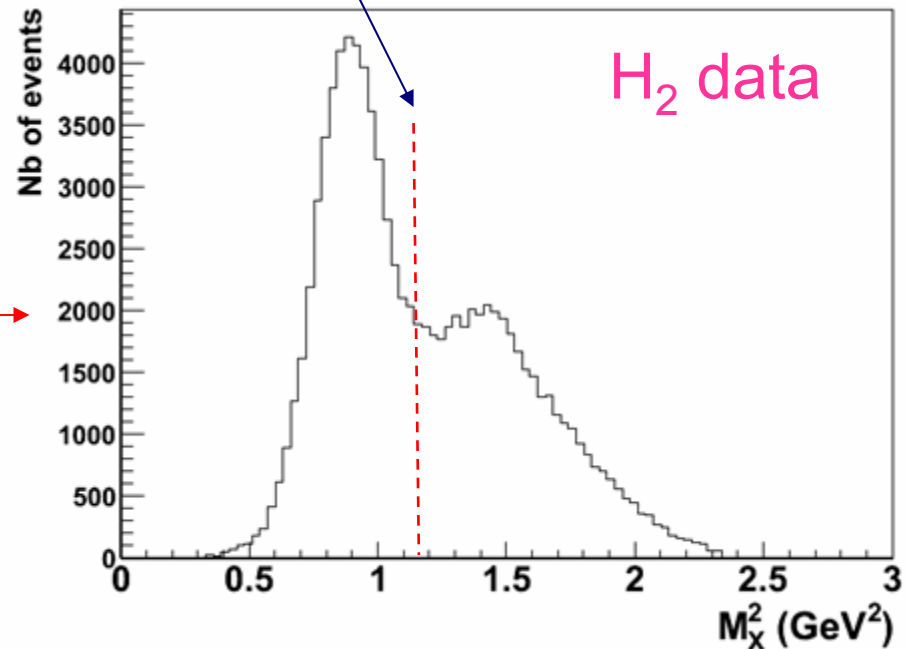
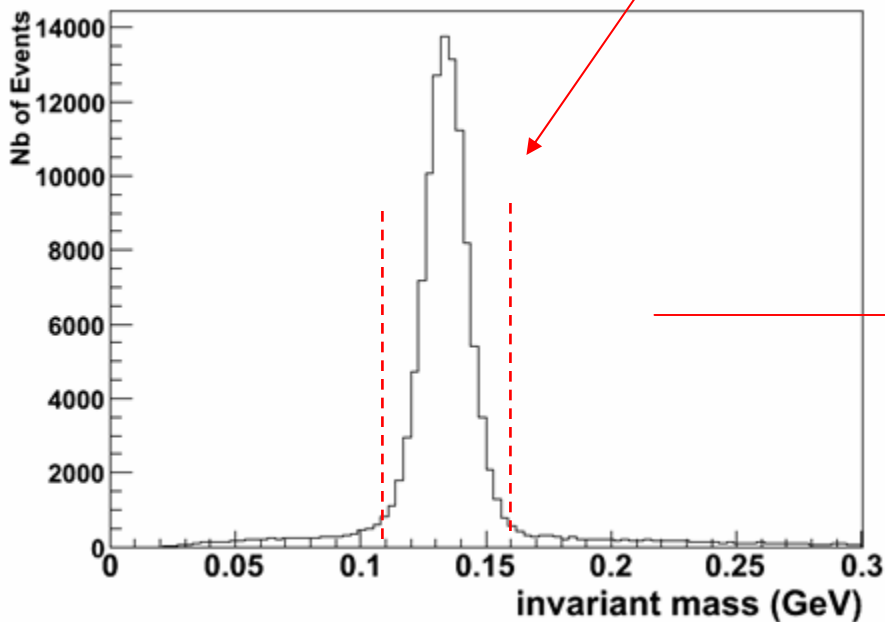
# $\pi^0$ contamination subtraction



Subtraction of  $\pi^0$  contamination ( $1\gamma$  in the calorimeter) is obtained from a phase space simulation which weight is adjusted to the experimental  $\pi^0$  cross section ( $2\gamma$  in the calorimeter).

# $\pi^0$ contamination subtraction

Unfortunately, the high trigger threshold during **Deuterium** runs did not allow to record **all exclusive  $\pi^0$  events** ( $M_X^2 < 1.15 \text{ GeV}^2$ )



**But:** In order to find the procedure of  $\pi^0$  contamination subtraction, **we must have** :

$$\frac{\sigma(e d \rightarrow e p \pi^0 X) \sigma(e p \rightarrow e p \pi^0 n)}{\sigma(e p \rightarrow e p \pi^0 X) \sigma(e p \rightarrow e p \pi^0 p)} = 0.95 \pm 0.06$$

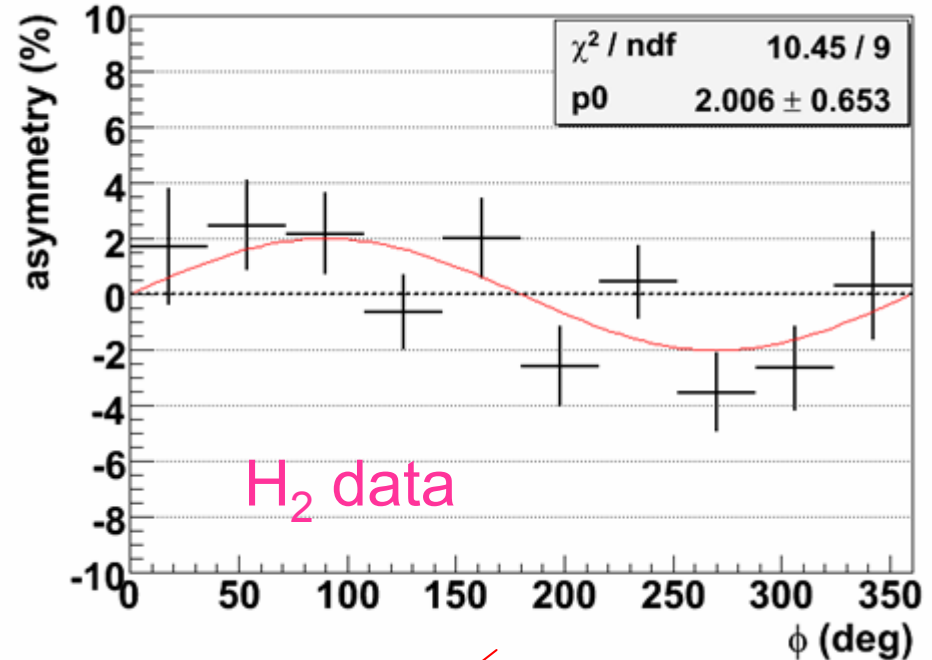
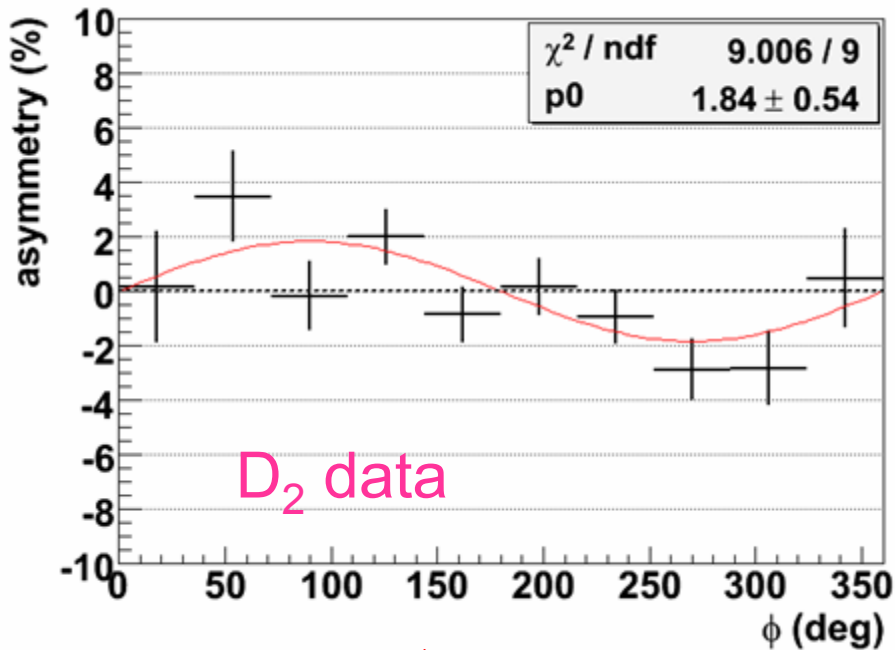
with  $M_X^2 < 1.15 \text{ GeV}^2$

by comparing two samples of high energy  $\pi^0$  in each case

# Exclusive $\pi^0$ asymmetry

$ed \rightarrow ep(n)\pi^0 + ed \rightarrow en(p)\pi^0 + ed \rightarrow ed\pi^0$

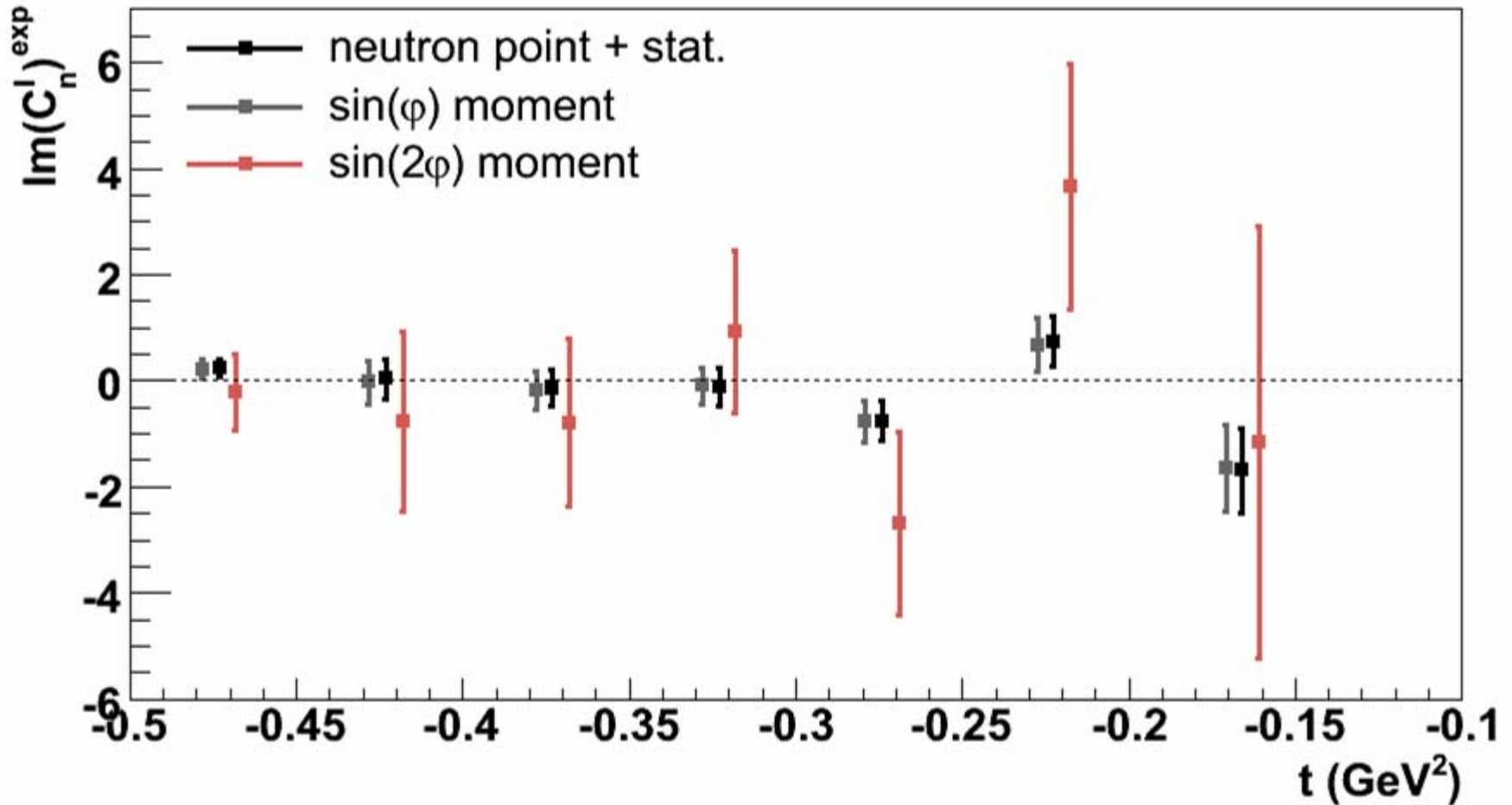
$ep \rightarrow ep\pi^0$



$$-0.5 < \frac{S_h(en \rightarrow en\pi^0 + ed \rightarrow ed\pi^0)}{S_h(ep \rightarrow ep\pi^0)} < 1$$

Well known from H<sub>2</sub> data

# $\sin(\varphi)$ and $\sin(2\varphi)$ moments



Results are coherent with the fit of a single  $\sin(\varphi)$  contribution

# Test of the handbag dominance : E00-110

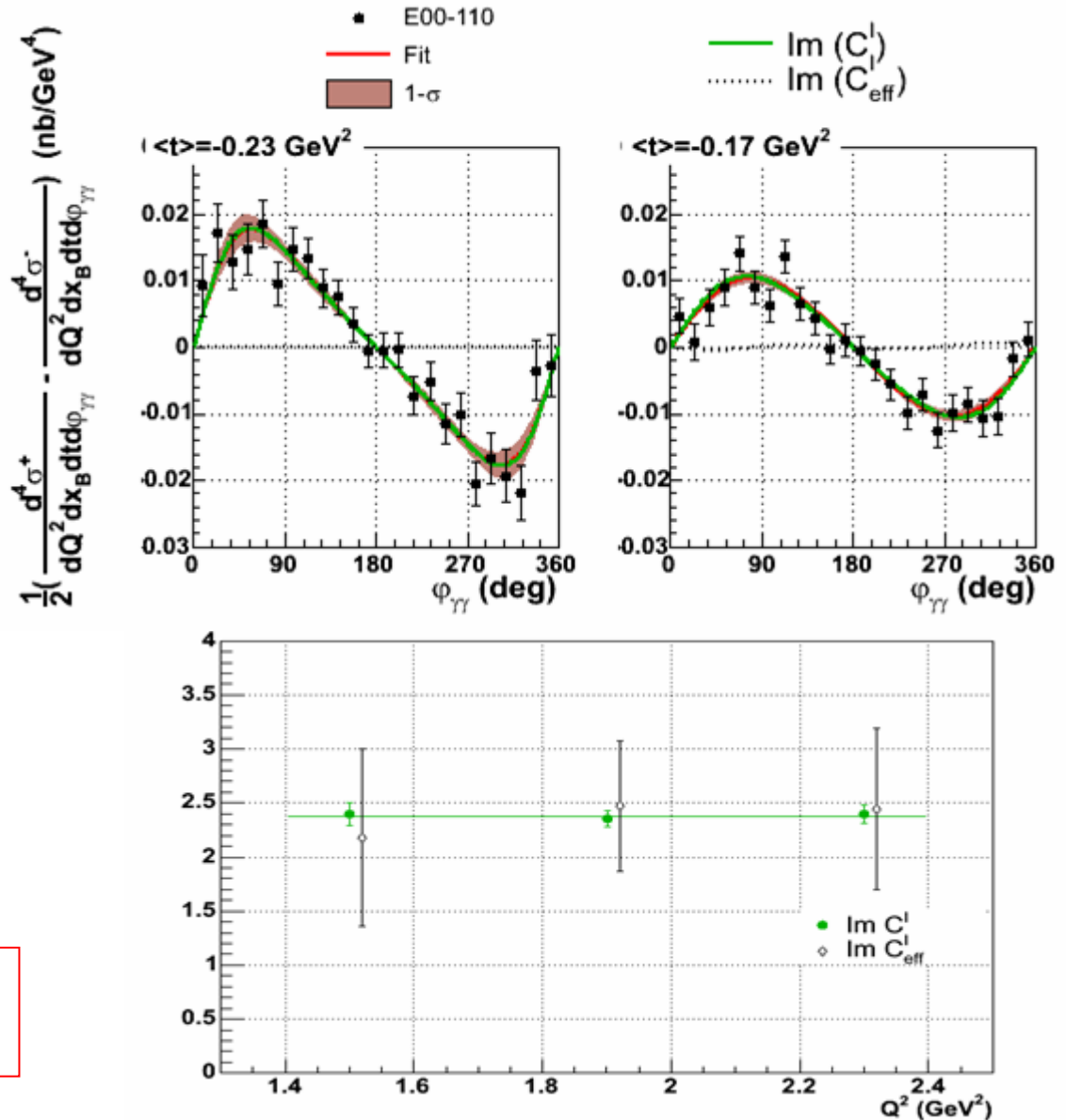
## p-DVCS experiment results

C. Muñoz-Camacho *et al.*,  
to appear in PRL (2007)

Twist-2 contribution dominates  
the total cross-section and the  
cross-section difference.

No  $Q^2$  dependence of twist-2  
and twist-3 terms

Strong indications for  
handbag dominance



# VGG parametrisation of GPDs

Vanderhaeghen, Guichon, Guidal,  
Goeke, Polyakov, Radyushkin, Weiss ...

## Non-factorized $t$ dependence

$$H^q(x, \xi, t) = \int_{-1}^1 d\beta \int_{-1+|\beta|}^{1-|\beta|} d\alpha \delta(x - \beta - \alpha\xi) F^q(\beta, \alpha, t) + \theta(\xi - |x|) \underbrace{D^q\left(\frac{x}{\xi}\right)}_{\text{D-term}}$$

Double distribution :  
 $\alpha' = 0.8 \text{ GeV}^{-2}$  for quarks

$$F^q(\beta, \alpha, t) = \frac{1}{|\beta|^{\alpha't}} h(\beta, \alpha) q(\beta)$$

→ Parton distribution

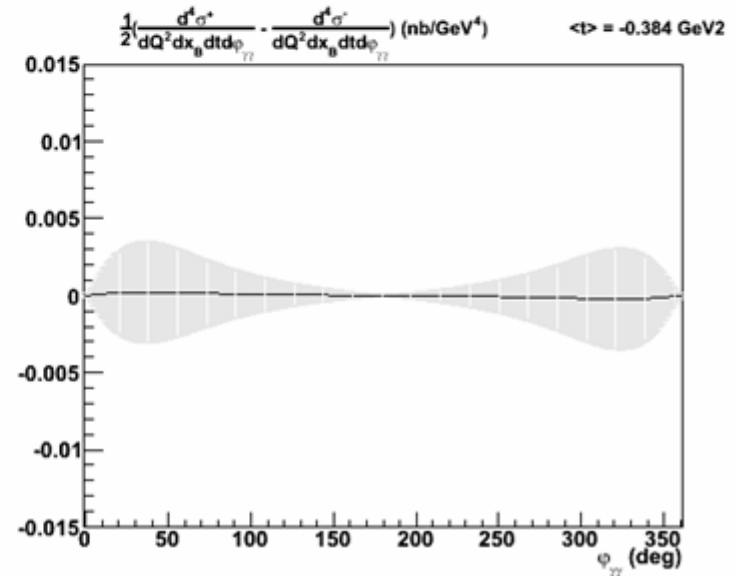
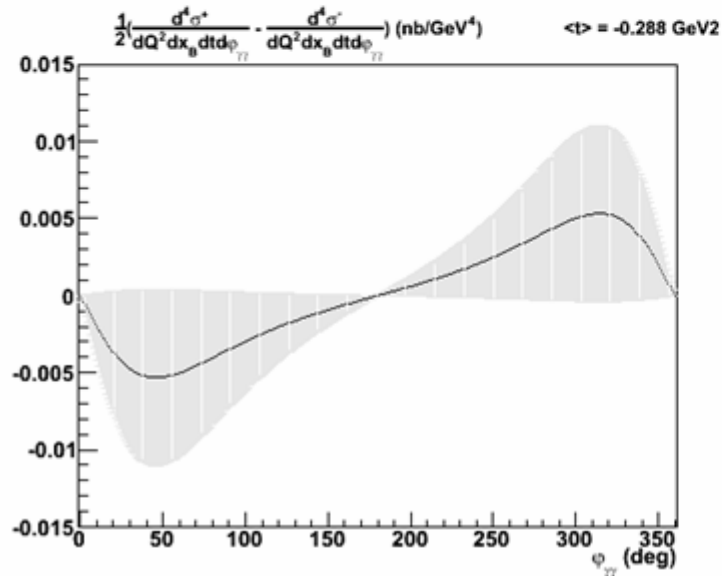
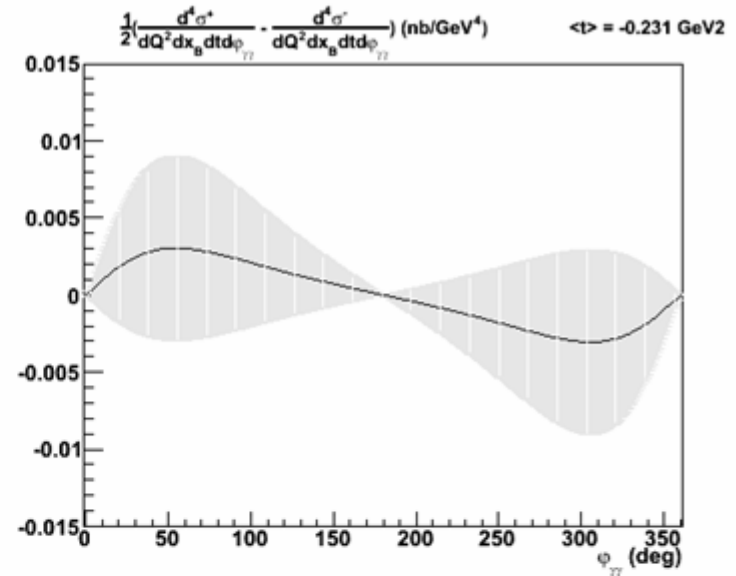
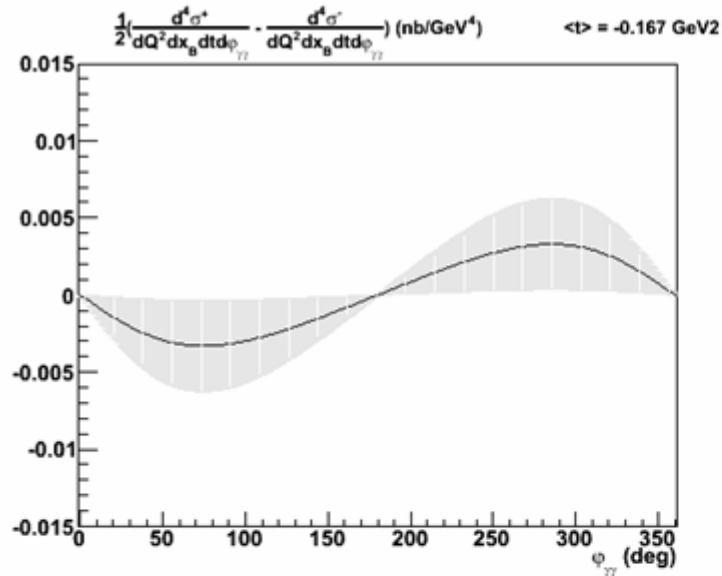
Profile function :

$$h(\beta, \alpha) = \frac{\Gamma(2b+2)}{2^{2b+1} \Gamma^2(b+1)} \frac{\left[ (1-|\beta|)^2 - \alpha^2 \right]^b}{(1-|\beta|)^{2b+1}}$$

for GPD  $E$ , the spin-flip parton densities is used :  $e_q(\beta)$

Modelled using  $J_u$  and  $J_d$  as free parameters

# n-DVCS polarized cross-section difference



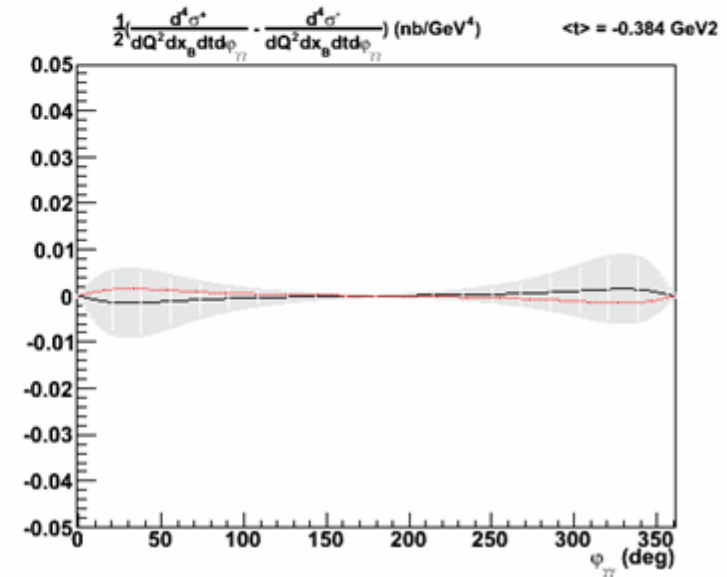
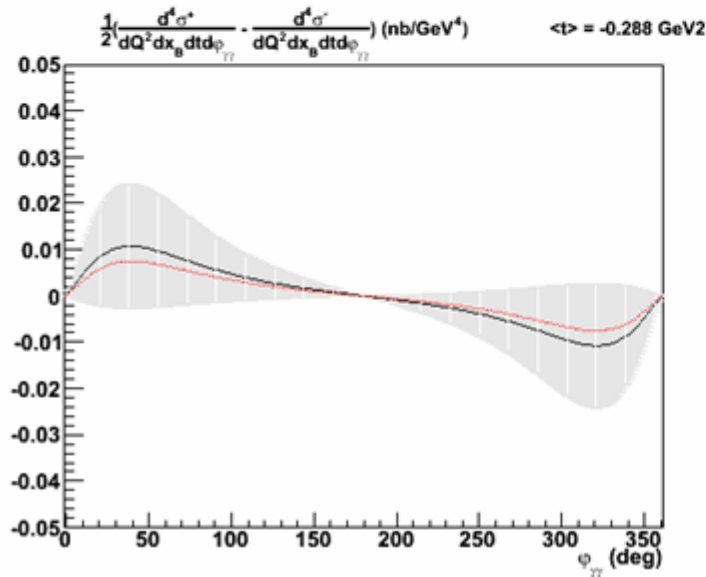
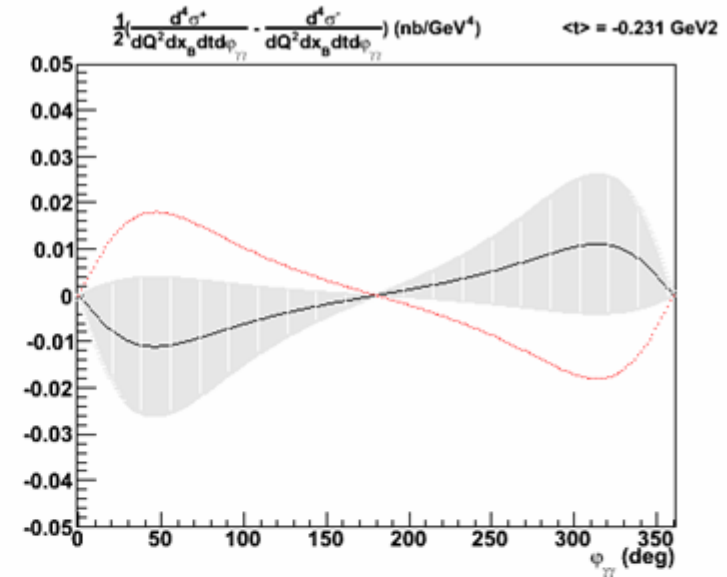
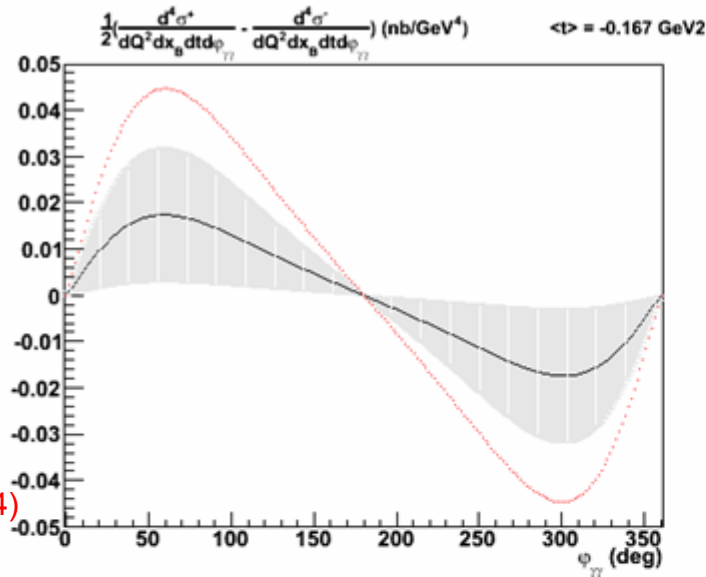


# d-DVCS polarized cross-section difference

Experimental  
results

+

Prediction from F.  
Cano and B. Pire.  
Eur. Phys. J. A19, 423 (2004)



# $\pi^0$ electroproduction on the neutron

Pierre Guichon, private communication (2006)

Amplitude of pion electroproduction :

$$T(N, \alpha) = \delta(\alpha, 3) T^+ + \tau_N^\alpha T^0 + i \varepsilon_{3\alpha\beta} \tau^\beta T^- \quad \alpha \text{ is the pion isospin}$$

  
 nucleon isospin matrix

➡  $\pi^0$  electroproduction amplitude ( $\alpha=3$ ) is given by :

$$\begin{aligned}
 T(p, 3) &= T^+ + T^0 \propto \frac{2}{3} \Delta u + \frac{1}{3} \Delta d \\
 T(n, 3) &= T^+ - T^0 \propto \frac{1}{3} \Delta u + \frac{2}{3} \Delta d
 \end{aligned}
 \left. \vphantom{\begin{aligned} T(p, 3) \\ T(n, 3) \end{aligned}} \right\} \boxed{\frac{T(p, 3) + T(n, 3)}{T(p, 3)} \approx \frac{3 + 3\Delta d / \Delta u}{2 + \Delta d / \Delta u} \approx 1.15}$$

  
 Polarized parton distributions in the proton

# Triple coincidence analysis

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Proton Array and Tagger (**hardware**) work properly during the experiment, but :

Identification of n-DVCS events with the recoil detectors is **impossible** because of the **high background rate**.

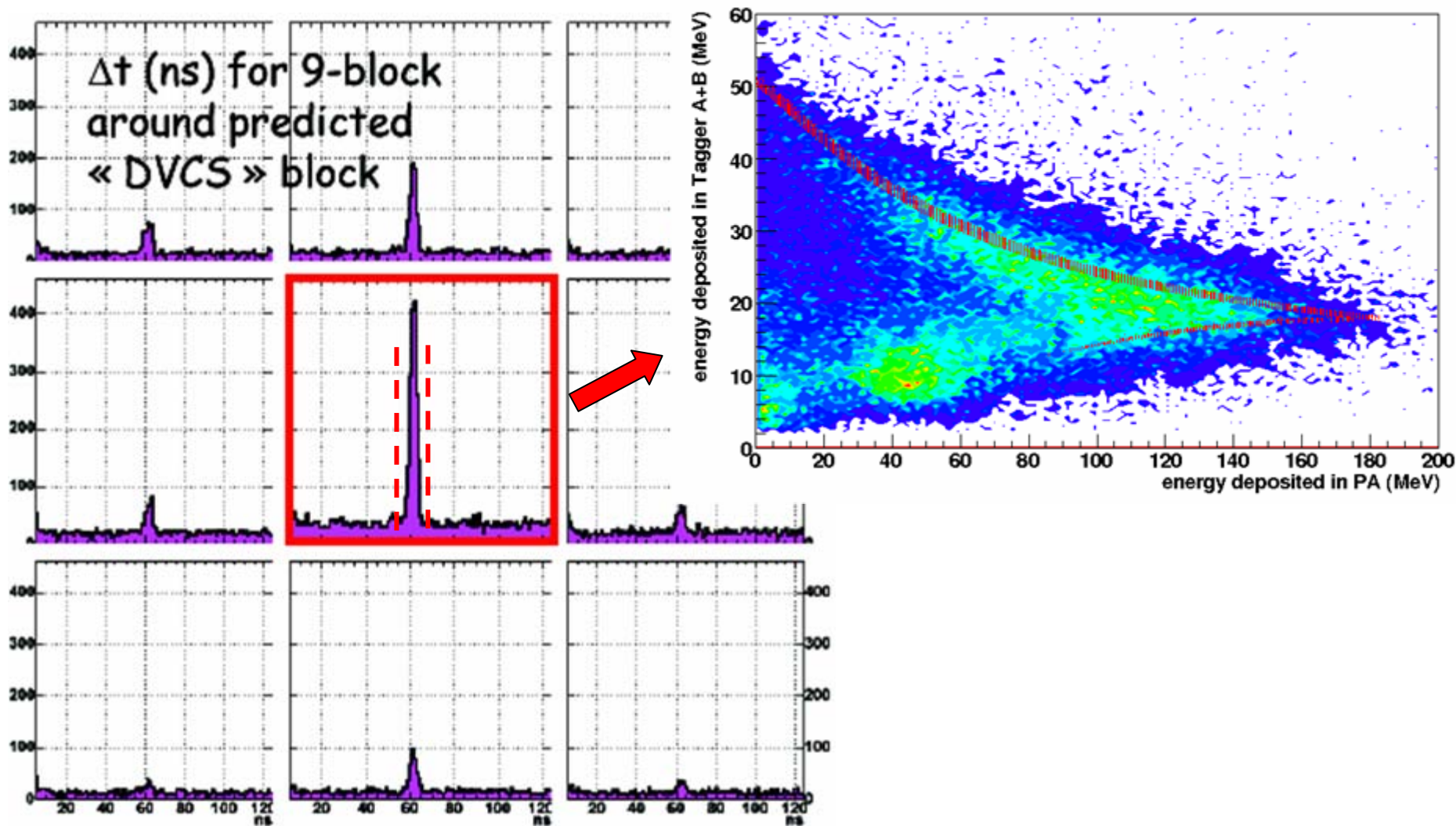
- ➡ Many Proton Array blocks contain signals on time for each event .
- ➡ **Accidental subtraction** is made for **p-DVCS** events and gives **stable** beam spin asymmetry results. The same subtraction method gives **incoherent results** for **neutrons**.

## Other major difficulties of this analysis:

- ➡ **proton-neutron conversion in the tagger shielding.**  
Not enough statistics to subtract this contamination correctly
- ➡ **The triple coincidence statistics** of n-DVCS is at least a **factor 20 lower** than the available statistics in the double coincidence analysis.

# Triple coincidence analysis

One can **predict** for each (e, $\gamma$ ) event the **Proton Array block** where the missing nucleon is supposed to be (assuming DVCS event)

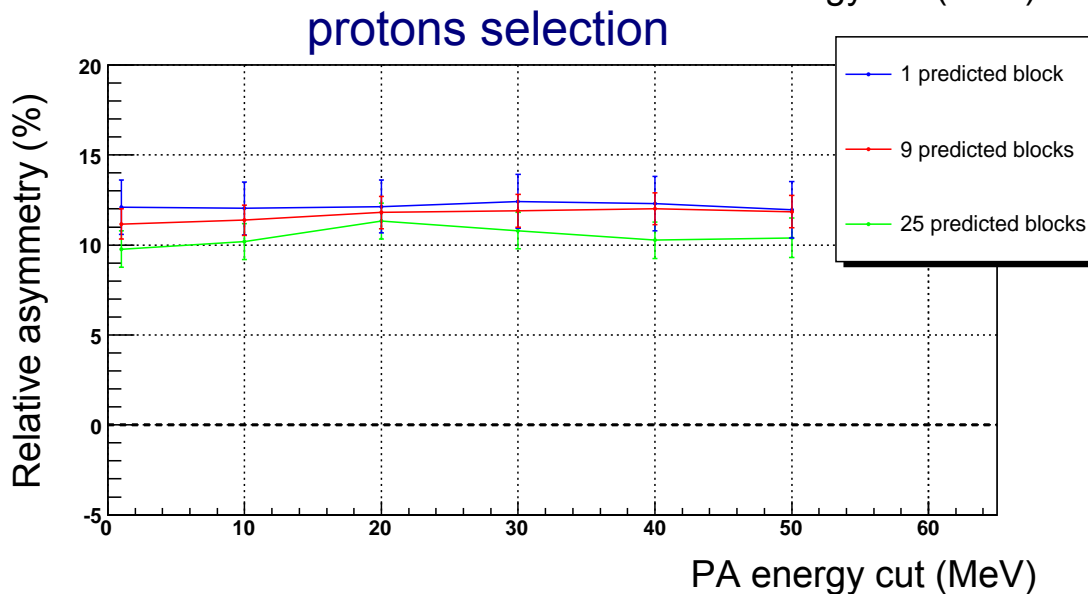
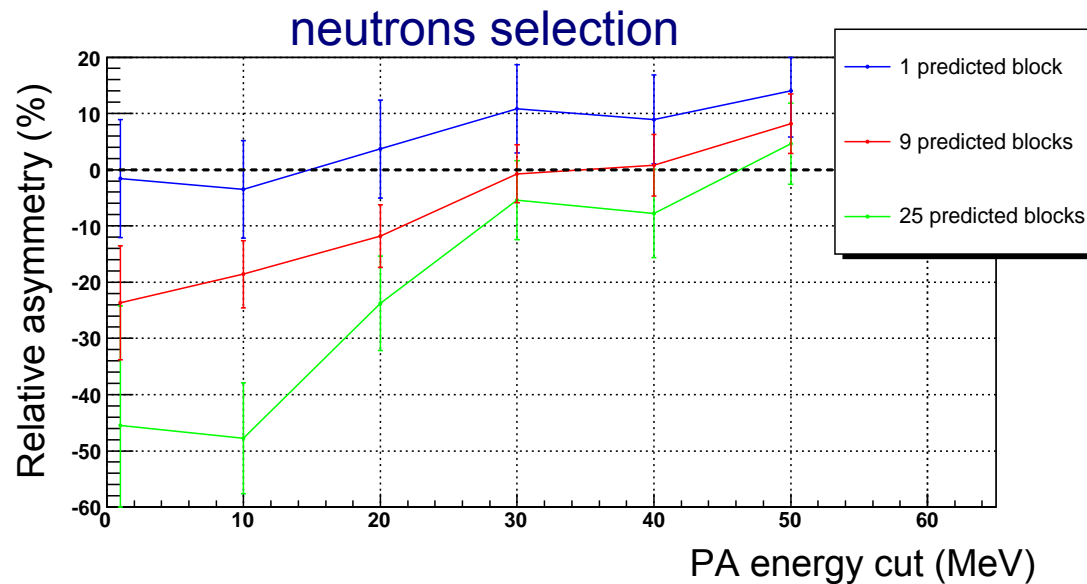


# Triple coincidence analysis

After accidentals subtraction

- proton-neutron conversion in the tagger shielding
- accidentals subtraction problem for neutrons

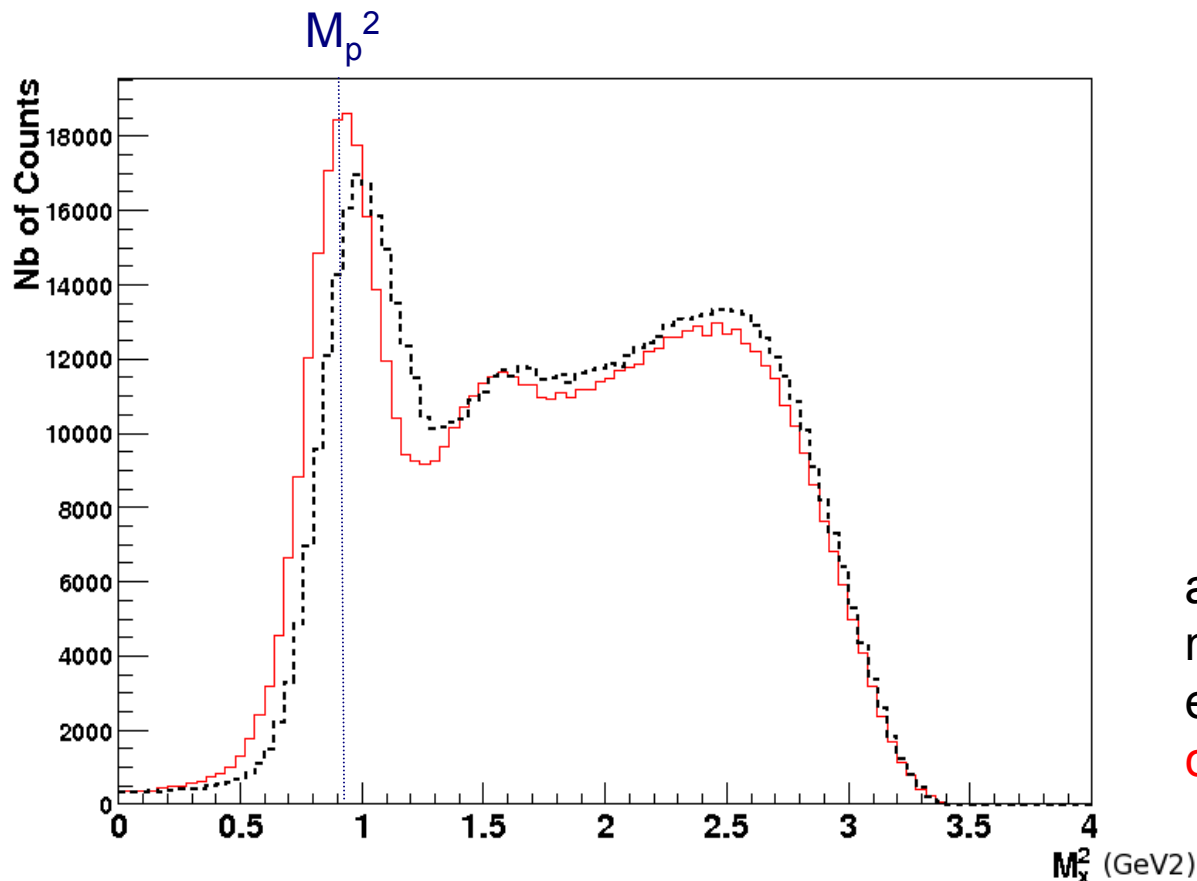
p-DVCS events (from LD2 target) asymmetry is stable



# Calorimeter energy calibration

We have **2 independent methods** to check and correct the calorimeter calibration

➔ **1<sup>st</sup> method** : missing mass of  $D(e, e' \pi^-) X$  reaction



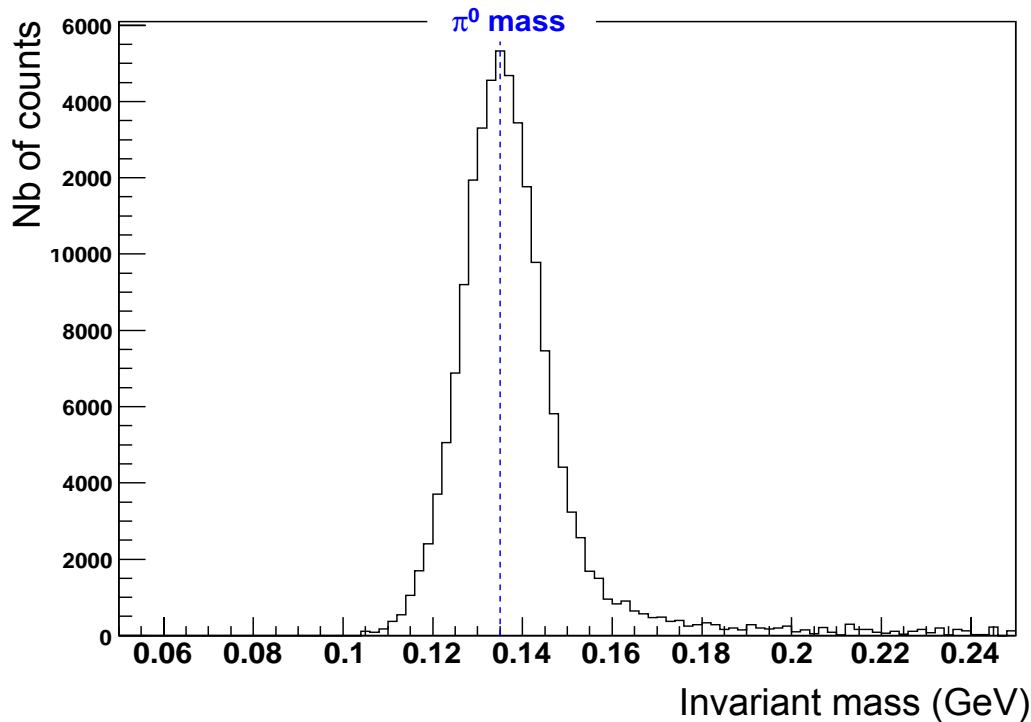
By selecting  $n(e, e' \pi^-) p$  events, one can predict the energy deposit in the calorimeter using only the cluster position.



a  $\chi^2$  minimisation between the measured and the predicted energy gives a better calibration.

# Calorimeter energy calibration

➔ **2<sup>nd</sup> method** : Invariant mass of 2 detected photons in the calorimeter ( $\pi^0$ )



$\pi^0$  invariant mass position  
check the quality of the  
previous calibration for  
each calorimeter region.



Corrections of the previous  
calibration are possible.



Differences between the results of the 2 methods introduce a systematic error of **1%** on the calorimeter calibration.