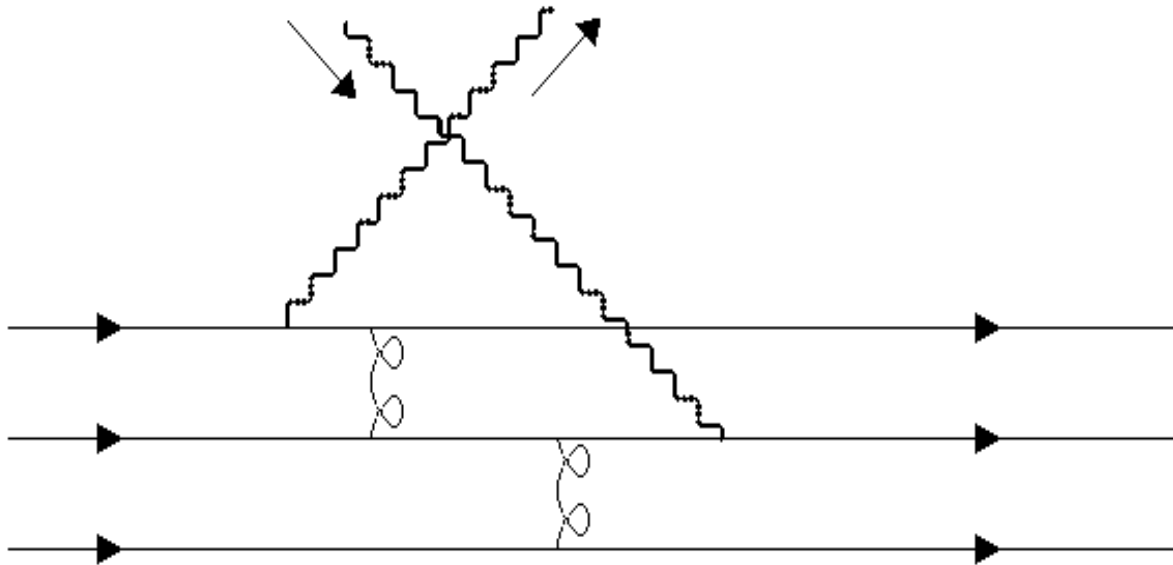


Real and Virtual Compton Scattering in Perturbative QCD



JLab, May 23, 2007

Motivation

- Forthcoming 12 GeV Upgrade and New JLab Data
 - Hall A Collab, PRL98, 152001(2007): $s=5-11\text{GeV}^2, -t=2-7\text{GeV}^2$
 - Cornell Data, PRD19, 1921(1979): $s=4.6-12.1\text{GeV}^2, -t=0.7-4.3\text{GeV}^2$
- Several Analyses of RCS but Different Results in the Past
 - E.Maina and G.R.Farra, PLB206, 120(1988)
 - G.R.Farrar and H.Zhang, PRD41, 3348(1990);42, 2413(E)(1990)
 - A.S.Kronfeld and B.Nizic, PRD44, 3445(1991)
 - M.Vanderhaeghen, P.Guichon and J.Van de Wiele, NPA622, 144(1997)
 - T.Brooks and L.Dixon, PRD62, 114021(2000)

Recent Agreement with Brooks and Dixon's Result

R.Thomson, A.Pang and C.Ji, PRD73,054023(2006)

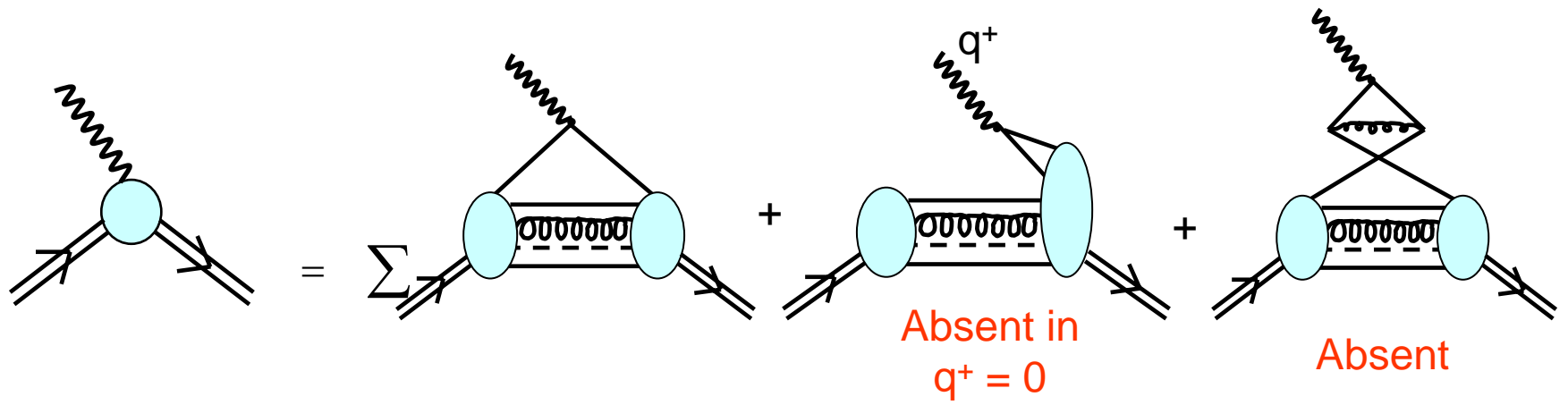
- Extension to Virtual Compton Scattering
 - Currently much interest in DVCS where GPD predictions are applicable
 - Only one previous PQCD calculation (Farrar and Zhang) and their RCS result is in disagreement with Brooks and Dixon's and ours.
 - GPD predictions can be compared with the PQCD results for DVCS.
This can shed light on the applicability of both GPD and PQCD methods.

($Q^2 \gg -t \gg \Lambda^2$)

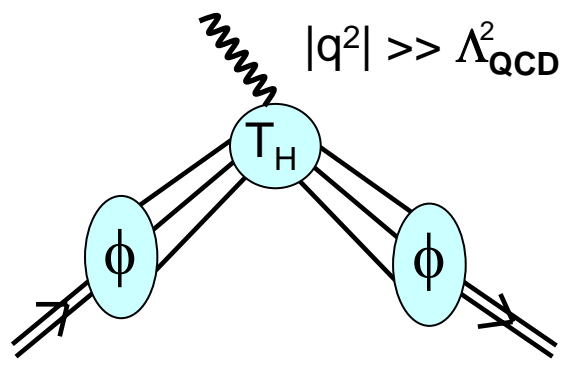
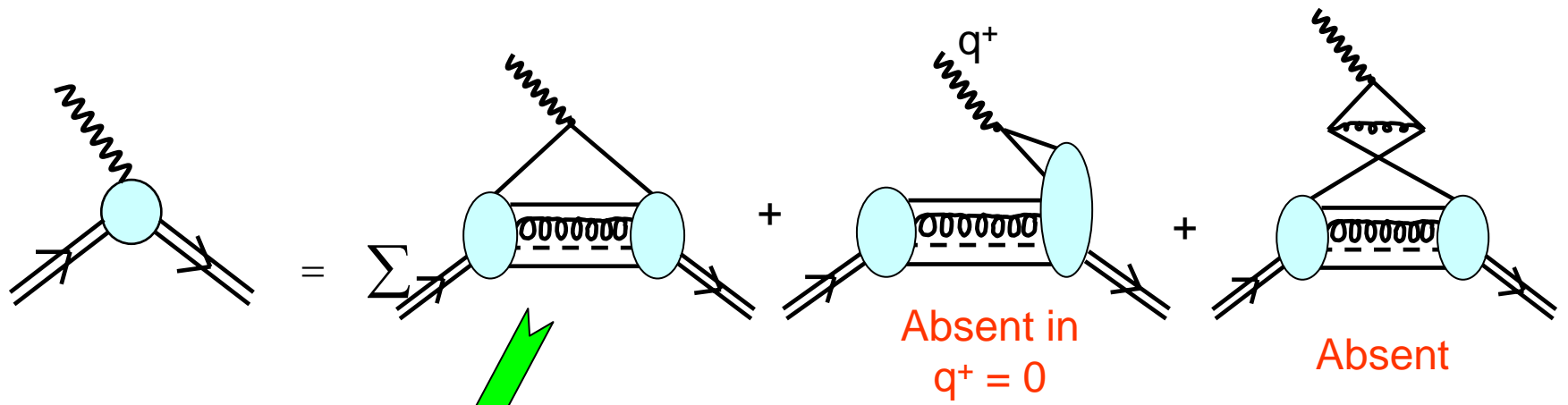
Outline

- PQCD in Light-Front Dynamics(LFD)
 - Hard Scattering Amplitude
 - Distribution Amplitude
 - Remarks on Computational Methods
- Real Compton Scattering Results
 - Comparison with Previous Computations
 - Comparison with JLab Data
- Extension to Virtual Compton Scattering
 - Comparison with Previous Computations
 - Link to GPD and Handbag Dominance
- Conclusions

LFD in Exclusive Processes



LFD in Exclusive Processes



$$T_H = \sum \left[\begin{array}{c} x_1 \text{---} \text{---} y_1 \\ x_2 \text{---} \text{---} y_2 \\ x_3 \text{---} \text{---} y_3 \end{array} \right] + \dots$$

$$= \frac{\alpha_s^2}{Q^4} f(x_i, y_i)$$

Classification of Diagrams

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

A, C, E \longleftrightarrow B,D,F under 1 \longleftrightarrow 3;

Color factor for all type: $C^{(d)} = 4/9$;

Triple gluon contribution is absent due to the color factor.

For RCS,

$$M_{hh'}^{\lambda\lambda'}(s,t) = M_{\bar{h}\bar{h}'}^{\bar{\lambda}\bar{\lambda}'}(s,t) \quad \text{due to Parity Inv.}$$
$$M_{hh'}^{\lambda\lambda'}(s,t) = M_{h'h}^{\lambda'\lambda}(s,t) \quad \text{Time-reversal Inv.}$$

Number of Contributing Diagrams

Single Photon (e.g. F_1^P)

Attachment of a photon: $(6/2) \times 7 = 21$

Nonzero Diagrams: A1, A4, A7, C2, C7, E1, E5

Two Photons: $(6/2) \times 7 \times 8 = 168$

RCS needs only A and C types due to T-inv.

52 Nonzero Diagrams: A11, A11, ..., C12, ..., C77.

-A.S.Kronfeld and B.Nizic, PRD44, 3445(1991)

Use $h=h'=1$ for proton helicities and

$|\gamma_{in}\rangle = \alpha|\uparrow\rangle + \beta|\downarrow\rangle$, $|\gamma_{out}\rangle = \gamma|\uparrow\rangle + \delta|\downarrow\rangle$.

Virtual Compton Scattering can't take advantage of T-inv.

96 Nonzero Diagrams: A11, ..., C77; A77, A77, ..., E12, ..., E55.

-Alex (Chiu-Yan) Pang's Thesis, NCSU (1995)

Virtual photon has also the longitudinal polarization.

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Summary of Theory

Leading twist PQCD approximation for proton Compton scattering gives helicity amplitude

$$M_{hh'}^{\lambda\lambda'} = \sum_{i,d} \int dx_1 dx_2 dx_3 dy_1 dy_2 dy_3 \phi_i(x) T_i^{(d)} \phi_i^*(y)$$

subject to constraints $x_1 + x_2 + x_3 = 1$ and $y_1 + y_2 + y_3 = 1$

$h, \lambda(h', \lambda')$ helicities of i/c (o/g) proton, photon

$x(y)$ longitudinal momentum fractions of i/c (o/g) quarks

i labels independent Fock states of the proton with distribution amplitudes $\phi_i(x)$

d labels the Feynman diagrams that contribute to hard scattering amplitude $T_i^{(d)}$

Distribution Amplitude

$$|p_\uparrow\rangle = \frac{f_N}{8\sqrt{6}} \int [dx] \sum_i \phi_i(x_1, x_2, x_3) |i, x_1, x_2, x_3\rangle;$$

$$|1; x_1, x_2, x_3\rangle = |u_\uparrow(x_1)u_\downarrow(x_2)d_\uparrow(x_3)\rangle,$$

$$|2; x_1, x_2, x_3\rangle = |u_\uparrow(x_1)d_\downarrow(x_2)u_\uparrow(x_3)\rangle,$$

$$|3; x_1, x_2, x_3\rangle = |d_\uparrow(x_1)u_\downarrow(x_2)u_\uparrow(x_3)\rangle;$$

$$\phi_2(x_1, x_2, x_3) = -[\phi_1(x_1, x_2, x_3) + \phi_1(x_3, x_2, x_1)],$$

$$\phi_3(x_1, x_2, x_3) = \phi_1(x_3, x_2, x_1);$$

$$[dx] = dx_1 dx_2 dx_3 \delta(1 - x_1 - x_2 - x_3);$$

$$f_N = (5.2 \pm 0.3) \times 10^{-3} \text{ (GeV/c)}^2.$$

V.Chernyak and I.Zhitnitsky, Nucl.Phys.B246,52(84)

$$\phi_{CZ} = 120x_1x_2x_3(1.69 - 9.26x_1 - 10.94x_3 + 22.70x_1^2 + 13.45x_3^2 + 9.26x_1x_3)$$

V.Chernyak,A.Ogloblin and I.Zhitnitsky,Sov.J.Nucl.Phys.48,536(88)

$$\phi_{COZ} = 120x_1x_2x_3(5.880 - 25.956x_1 - 20.076x_3 + 36.792x_1^2 + 19.152x_3^2 + 25.956x_1x_3)$$

I.King and C.Sachrajda,Nucl.Phys.B279,785(87)

$$\phi_{KS} = 120x_1x_2x_3(8.40 - 26.88x_1 - 35.28x_3 + 35.28x_1^2 + 37.80x_3^2 + 30.24x_1x_3)$$

M.Gari and N.G.Stefanis,Phys.Lett.B175,462(86);PRD35,1074(87)

$$\phi_{GS} = 120x_1x_2x_3(6.040 - 16.775x_1 - 34.985x_3 - 1.027x_1^2 + 12.307x_3^2 + 111.320x_1x_3)$$

Asymptotic DA

$$\phi_{ASY} = 120x_1x_2x_3$$

AdS/CFT?

It is possible to write the helicity amplitude in the form

$$M_{hh'}^{\lambda\lambda'} = \frac{4}{9} (4\pi\alpha_{em})(4\pi\alpha_s)^2 \left(\frac{120f_N}{8\sqrt{6}} \right)^2 \sum_d \sum_{m,n} C^{(d)}(m_1, m_3, n_1, n_3) I^{(d)}(m_1, m_3, n_1, n_3)$$

where

$$I^{(d)}(m_1, m_3, n_1, n_3) = \int_0^1 dx_1 dx_2 dx_3 dy_1 dy_2 dy_3 \tilde{T}^{(d)} x_1^{m_1+1} x_2 x_3^{m_3+1} y_1^{n_1+1} y_2 y_3^{n_3+1}$$

m_1, m_3, n_1, n_3 are powers of momentum fractions (from ϕ_i)

$C^{(d)}$

is a coefficient that sums contributions for 3 Fock states.

It is dependent on $Z_i^{(d)}$ (product of charges of struck quarks)

and the coefficients appearing in ϕ_i

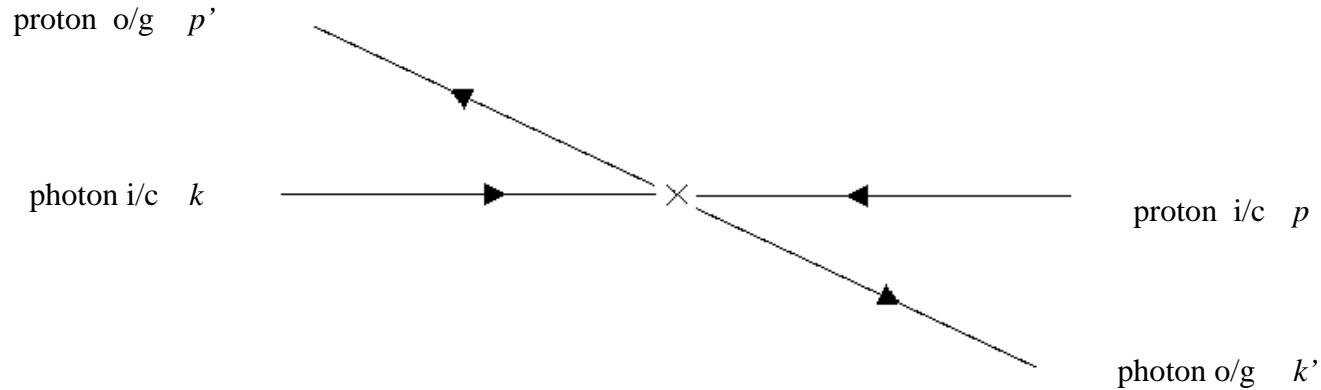
f_N

is a normalization constant (taken as 0.0052 GeV^2)

$\tilde{T}^{(d)}$

is the color/flavor independent part of $T_i^{(d)}$

Kinematics



$$p = P(1,0,0,1)$$

$$k = P(\sqrt{1-Q^2/P^2}, 0, 0, -1)$$

$$p' = E(1, \sin \theta, 0, \cos \theta)$$

$$k' = E(1, -\sin \theta, 0, -\cos \theta)$$

- Define $S = (2E)^2$ (Mandelstam invariant)

$$R = 1 + Q^2 / S \quad (\text{Virtuality parameter})$$

- For real photon: $R=1$, for DVCS: $R=2$

- $\tilde{T}^{(d)}$ calculated as a function of S and R : e.g.

$$\tilde{T}_{A16}^{\uparrow\uparrow} = \frac{8}{S^2} \frac{R^{3/2} s^4}{c} \frac{1 - Rx_3}{\langle \bar{y}_1, x_3 \rangle \langle \bar{y}_1, \bar{x}_1 \rangle \langle y_3, x_3 \rangle (1 - R\bar{x} + i\varepsilon)}$$

$$\langle y, x \rangle = y(1 - Rs^2x) - Rc^2x + i\varepsilon$$

$$s = \sin(\theta/2), c = \cos(\theta/2), \bar{y} = 1 - y$$

- See details of calculation for A51 in Appendix A of Thomson, Pang and Ji, PRD73, 054023 (2006).
- The pole expressions have been expanded into real and imaginary parts using

$$\frac{1}{\langle y, x \rangle + i\varepsilon} = P\left(\frac{1}{\langle y, x \rangle}\right) - i\pi\delta(\langle y, x \rangle)$$

where P means principal value.

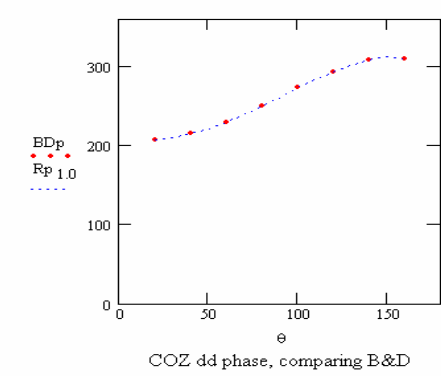
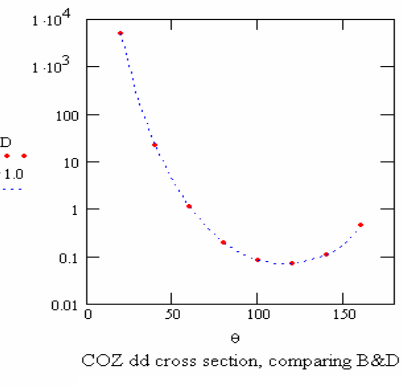
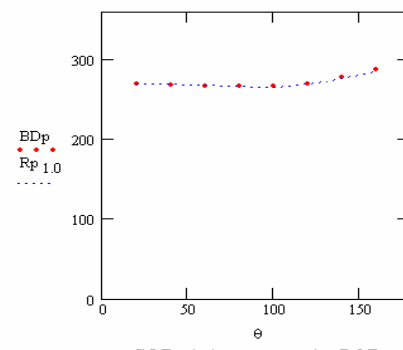
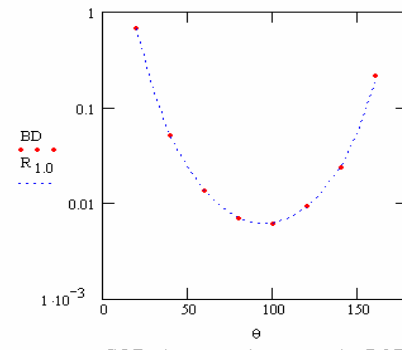
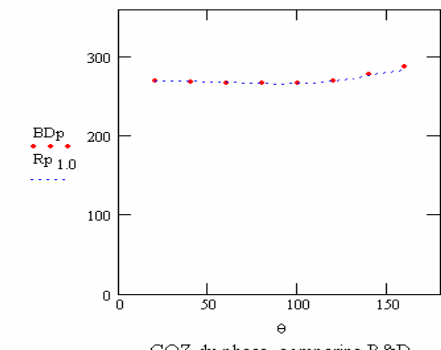
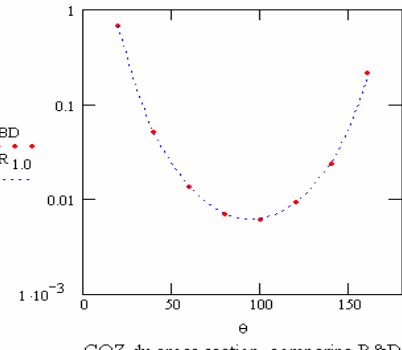
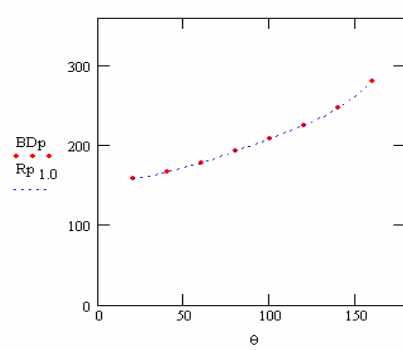
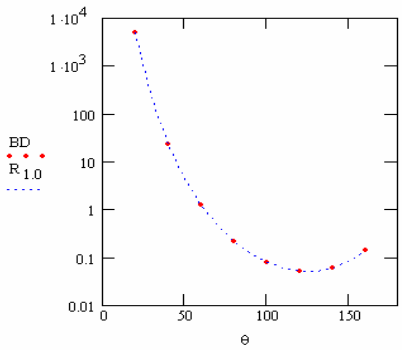
- Delta function integrals have been evaluated explicitly.
- The remaining principal value integrals have been transformed using the ‘folding method’ of Kronfeld & Nizic which renders the integrand finite over the range of integration.
- Brooks and Dixon used a different method to do the integrations (contour deformation)
- The integration has then been completed numerically using a Monte Carlo algorithm in Fortran.

Real and imaginary parts of all helicity amplitudes were tabulated in the basis of Appell polynomial expansion up to A_3 :

Thomson, Pang and Ji, hep-ph/0602164v2

for the computation of cross sections and phase calculations performed in the range of R values (1.0, 1.25, 1.5, 1.75, 2.0) and angles (20° to 160°).

Comparison with previous work for real photon



• T. Brooks and L. Dixon, PRD62, 114021 (2000)

... R. Thomson, A. Pang and C. Ji, PRD73, 054023 (2006)

Initial-State Helicity Correlation

$$A_{LL}(\text{or } K_{LL}) = \frac{\frac{d\sigma_+^+}{dt} - \frac{d\sigma_+^-}{dt}}{\frac{d\sigma_+^+}{dt} + \frac{d\sigma_+^-}{dt}} = \frac{|M_{\uparrow\uparrow}^{\uparrow\uparrow}|^2 + |M_{\uparrow\uparrow}^{\uparrow\downarrow}|^2 - |M_{\uparrow\uparrow}^{\downarrow\uparrow}|^2 - |M_{\uparrow\uparrow}^{\downarrow\downarrow}|^2}{|M_{\uparrow\uparrow}^{\uparrow\uparrow}|^2 + |M_{\uparrow\uparrow}^{\uparrow\downarrow}|^2 + |M_{\uparrow\uparrow}^{\downarrow\uparrow}|^2 + |M_{\uparrow\uparrow}^{\downarrow\downarrow}|^2}$$

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

- T. Brooks and L. Dixon, PRD62, 114021 (2000)
- M. Vanderhaeghen, P. Guichon and J. Van de Wiele, NPA622, 144 (1997)
- A. S. Kronfeld and B. Nizic, PRD44, 3445 (1991)
- G. R. Farrar and H. Zhang, PRD41, 3348 (1990);42, 2413(E) (1990)

JLab Data: $0.678 \pm 0.083 \pm 0.04$ at 120°
 $s=6.9 \text{ GeV}^2$, $t=-4.0 \text{ GeV}^2$, $u=-1.1 \text{ GeV}^2$
PRL94,242001(2005)



QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Handbag: M.Diehl, T.Feldmann, R.Jakob and P.Kroll, PLB460,204(1999)

Diquark: P.Kroll, M.Schurmann and W.Schwiger, IJMPA6,4107(1991)

Handbag approach to wide-angle RCS using GPD

- It is argued, at currently accessible kinematics, the Compton scattering amplitude is dominated by soft overlap contributions which can be described in light front gauge via the handbag diagram.
- In this approach, the helicity amplitude is given by (for example)

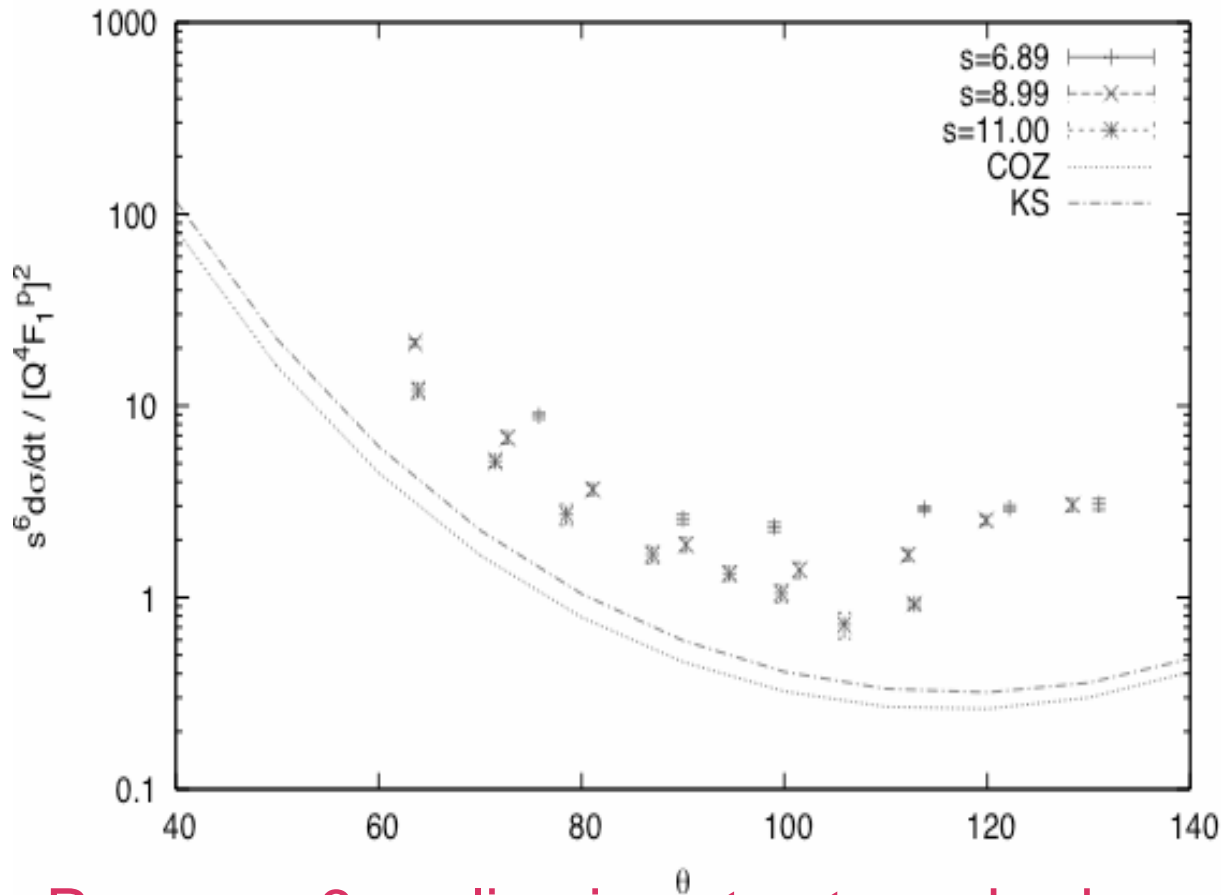
$$M_{++}^{\lambda\lambda'}(s,t) = 2\pi\alpha_{em} \left(H_{++}^{\lambda\lambda'}(s,t)(R_V(t) + R_A(t)) + H_{--}^{\lambda\lambda'}(s,t)(R_V(t) - R_A(t)) \right)$$

where $H_{++}^{\lambda\lambda'}(s,t)$ denotes the amplitude for the sub-process $\gamma q \rightarrow \gamma q$ and $R_A(t)$, $R_V(t)$ are soft form factors which can be expressed in terms of GPDs. For example, summing over flavors a

$$R_V(t) = \sum_a e_a^2 \int_{-1}^1 \frac{d\bar{x}}{x} H^a(\bar{x}, \xi = 0, t)$$

Scaled Unpolarized RCS Cross Section

Hall A Collab, PRL98, 152001(2007)

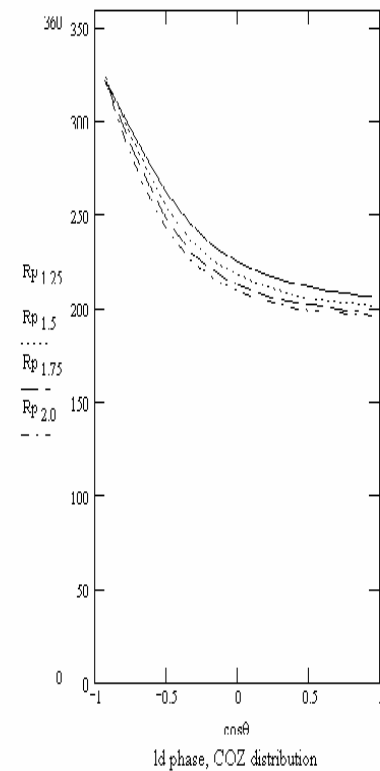
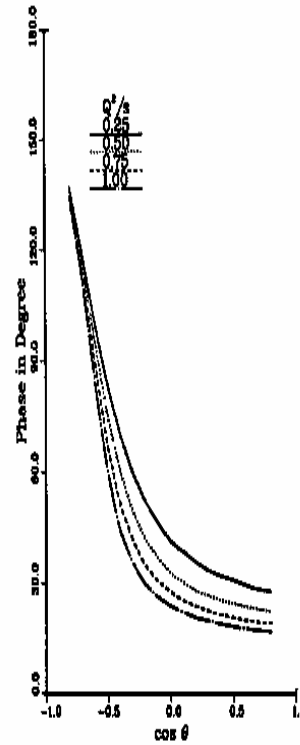
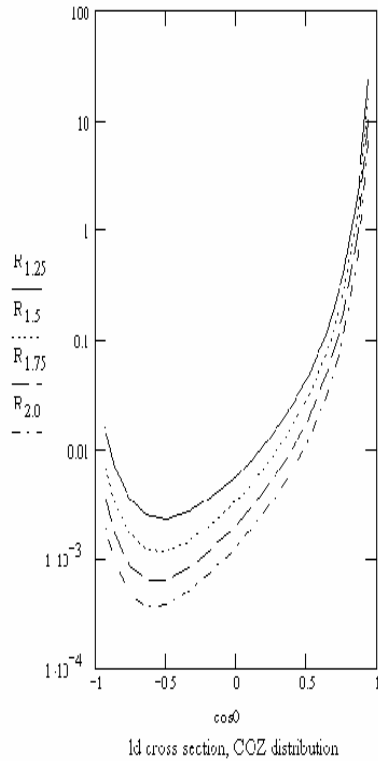
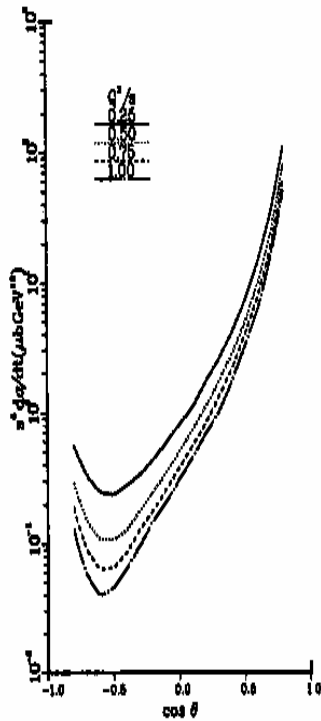


Suppress the
Uncertainty in
Normalization:

$$0.19 < \alpha_s < 0.39,$$
$$0.75 < \phi_{\text{evol}} < 1.33,$$
$$f_N = (5.2 \pm 0.3) \times 10^{-3} \text{ GeV}^2.$$

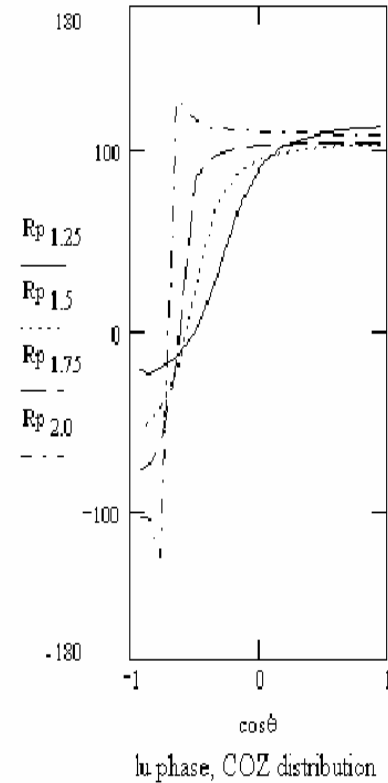
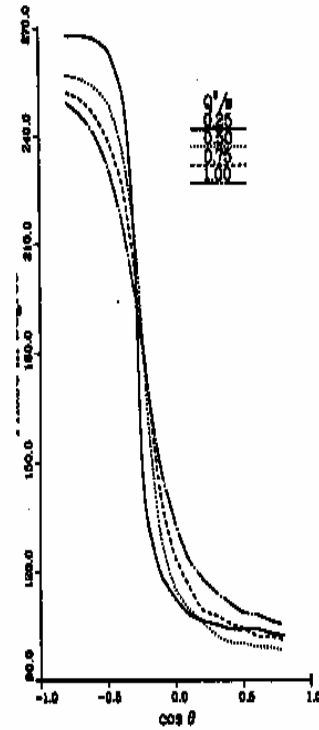
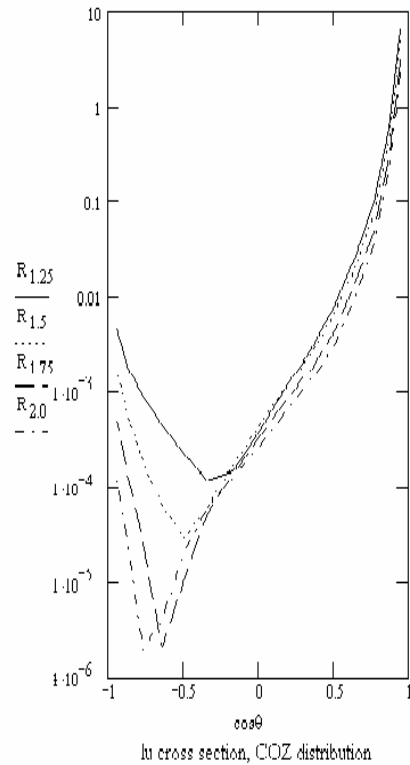
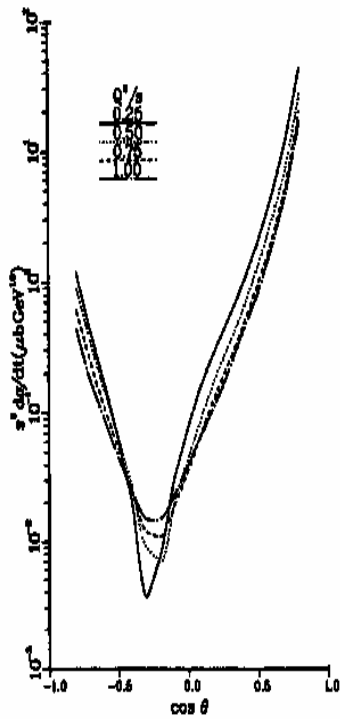
Power $n=6$ scaling is not yet reached.
12 GeV upgrade is anticipated.

Virtual photon comparison: longitudinal polarization (l to d)

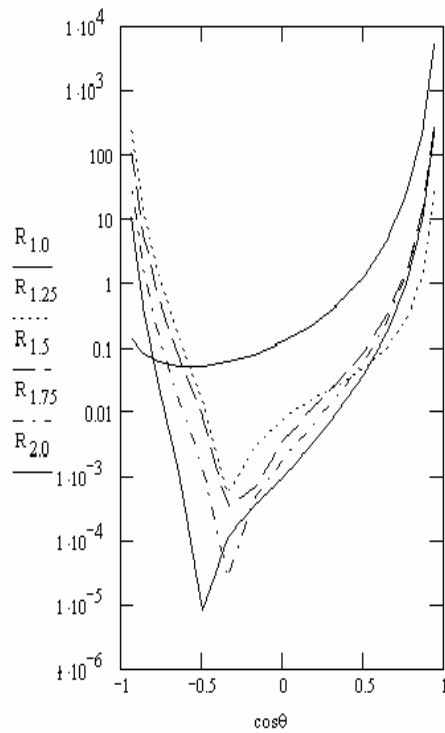
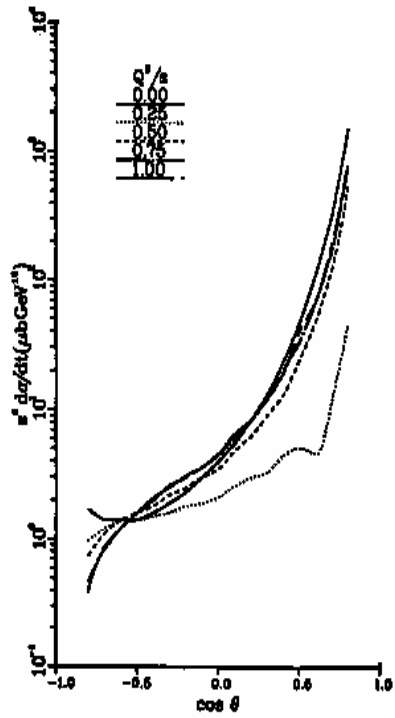


The results have been compared with Farrar and Zhang, PRD41,3348(1990).

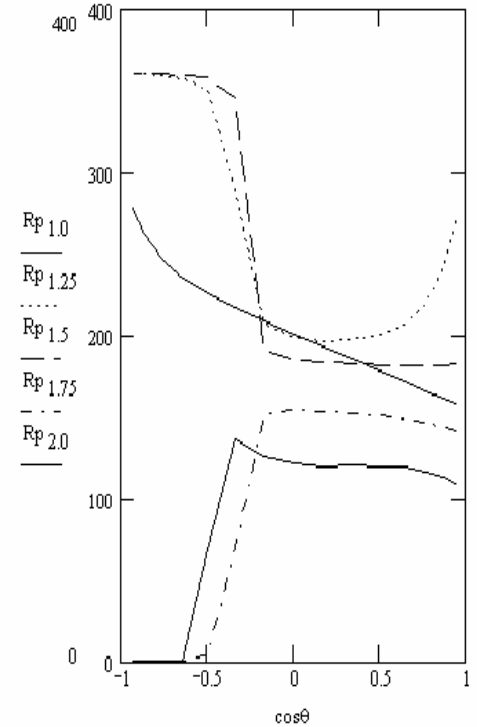
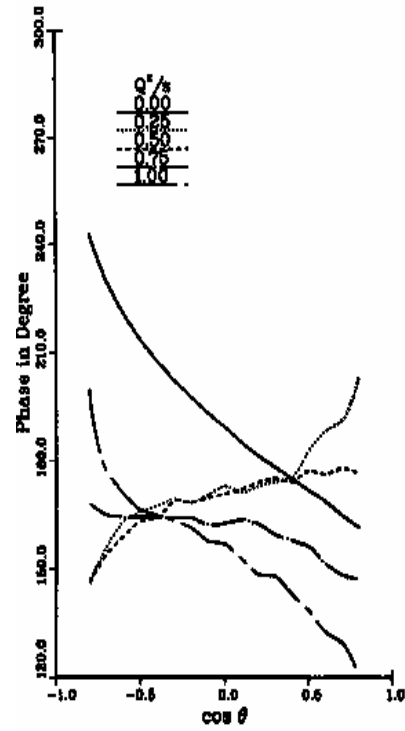
Virtual photon comparison: longitudinal polarization (l to u)



Virtual photon comparison: up to up

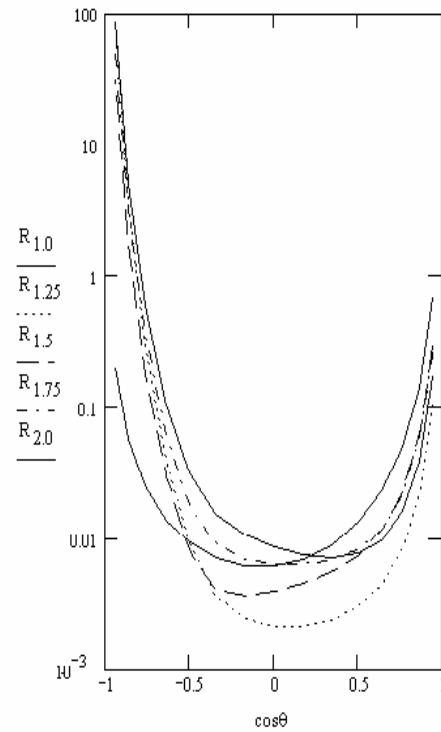
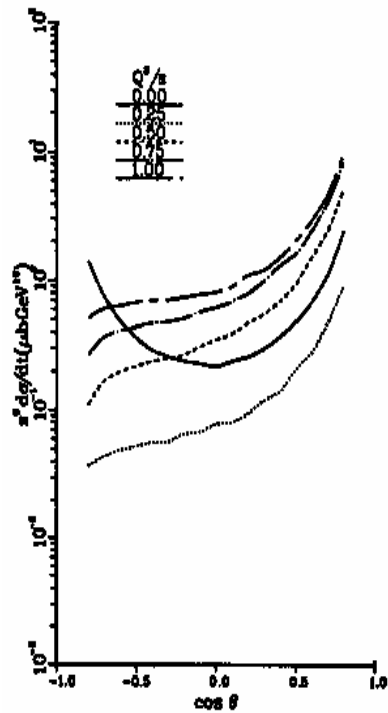


uu cross section, COZ distribution

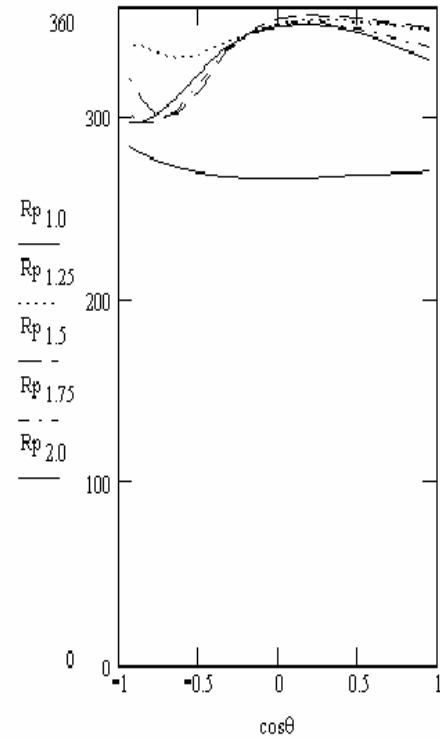
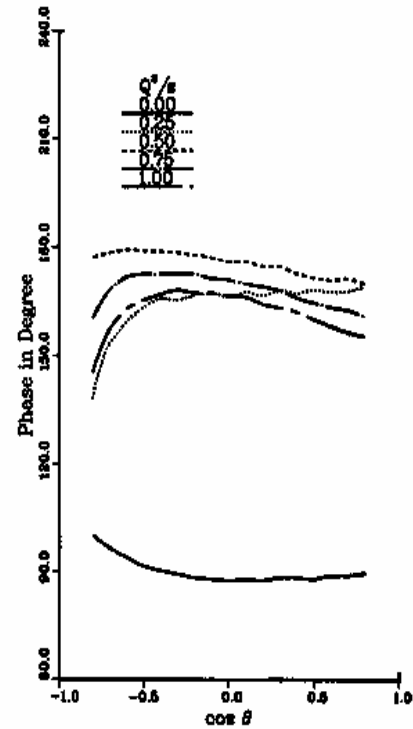


uu phase, COZ distribution

Virtual photon comparison: up to down

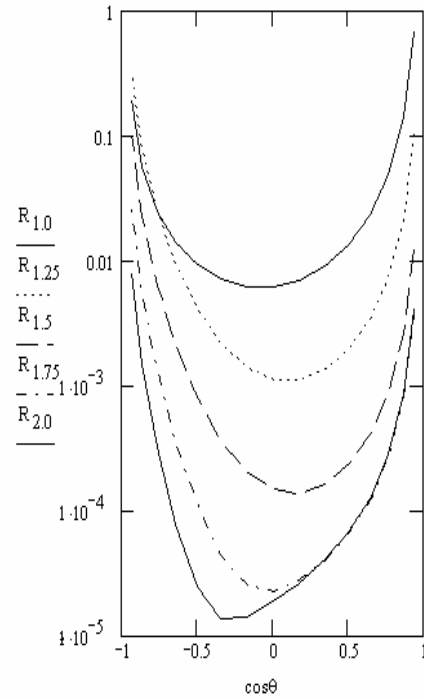
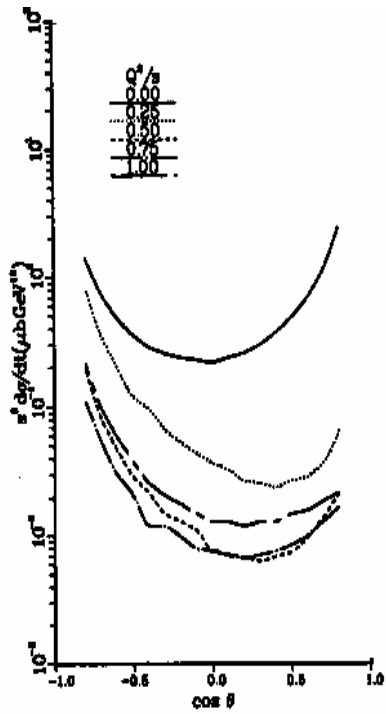


ud cross section, COZ distribution

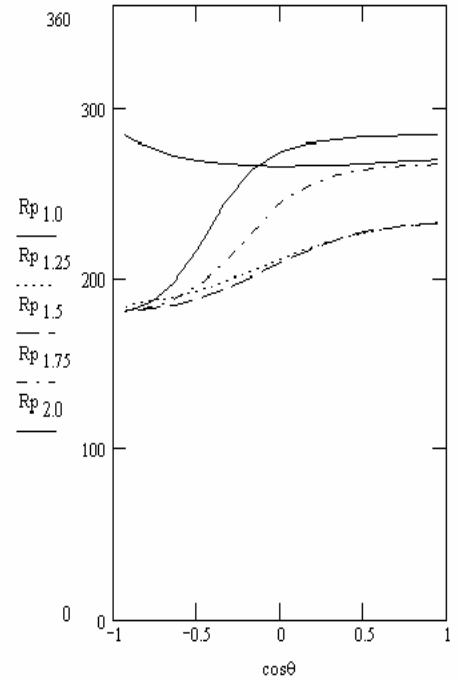
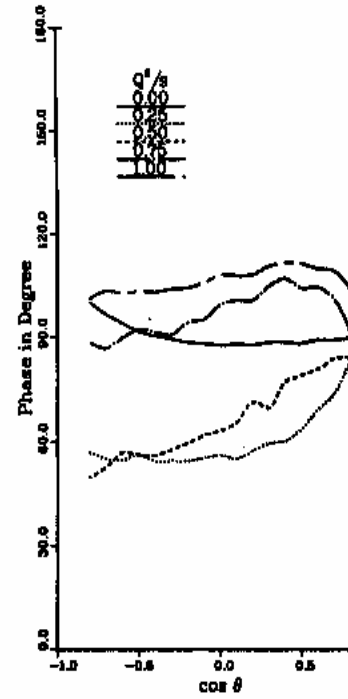


ud phase, COZ distribution

Virtual photon comparison: down to up

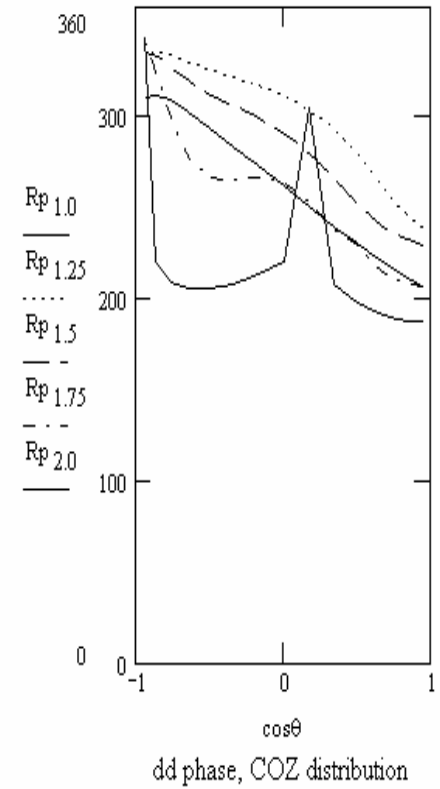
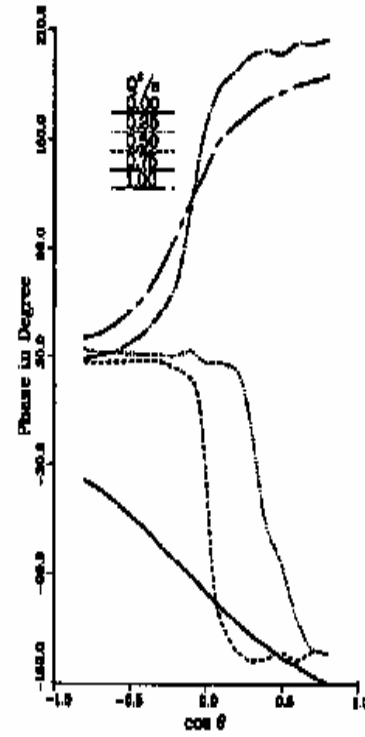
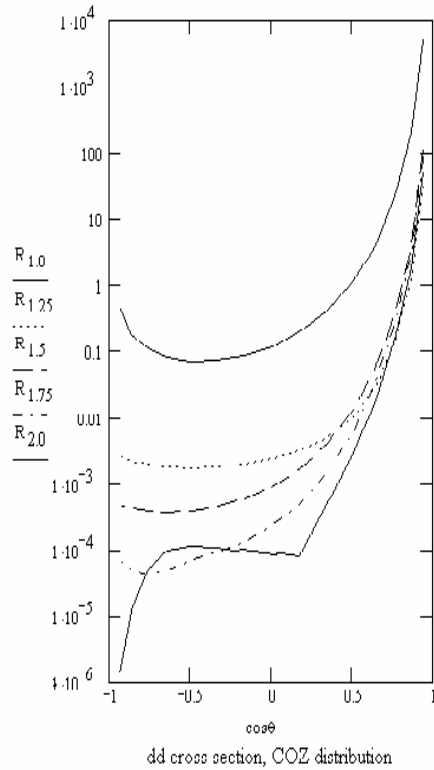
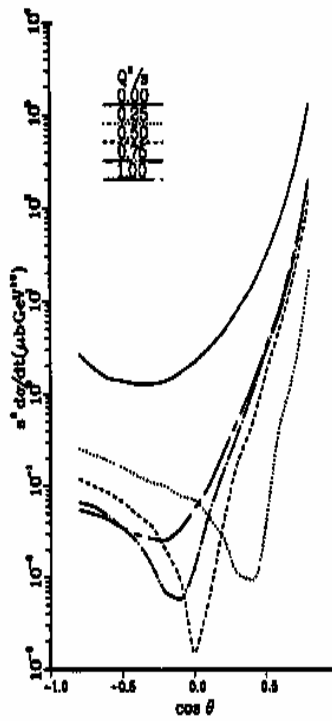


du cross section, COZ distribution



du phase, COZ distribution

Virtual photon comparison: down to down



Checking Handbag Dominance in PQCD



- FeynComp converted to work in Light Front gauge.
- Hand calculation of a few sample diagrams has been used to check the output of FeynComp.
- Explicit calculation using FeynComp shows the gauge invariance:

$$\sum_{\text{LightFront}} \text{amplitudes} = \sum_{\text{Feynman}} \text{amplitudes}$$

where the sum is over diagrams with photons attached to the same quarks.

Work in progress ...

Conclusions

- Agreement with Brooks & Dixon in RCS results gives more confidence on our VCS calculation over the previous one by Farrar & Zhang.
- JLab 12 GeV upgrade is highly desirable to shed some light on the validity of PQCD in exclusive processes.
- GPDs can be expressed in terms of DAs using PQCD at large momentum transfer and the previous results are currently under investigation.
- Handbag dominance check in PQCD analysis is also in progress.