

Two-Photon Exchange in Electron-Proton Elastic Scattering

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Exclusive Reactions at High Momentum Transfer

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Plan of talk

Radiative corrections for elastic electron scattering

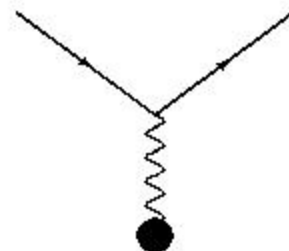
- Model-independent and model-dependent; soft and hard photons
- Refined bremsstrahlung calculations

Two-photon exchange effects in the process $e+p \rightarrow e+p$

- Models for two-photon exchange
 - Cross sections
 - Polarization transfer
 - Single-spin asymmetries



Elastic Nucleon Form Factors



- Based on one-photon exchange approximation

$$M_{fi} = M_{fi}^{1\gamma}$$

$$M_{fi}^{1\gamma} = e^2 \bar{u}_e \gamma_\mu u_e \bar{u}_p (F_1(t) \gamma_\mu - \frac{\sigma_{\mu\nu} q_\nu}{2m} F_2(t)) u_p$$

- Two techniques to measure

$$\sigma = \sigma_0 (G_M^2 \tau + \epsilon \cdot G_E^2) \quad : \text{Rosenbluth technique}$$

$$\frac{P_x}{P_z} = -\frac{A_x}{A_z} = -\frac{G_E \sqrt{\tau} \sqrt{2\epsilon(1-\epsilon)}}{G_M \tau \sqrt{1-\epsilon^2}} \quad : \text{Polarization technique}$$

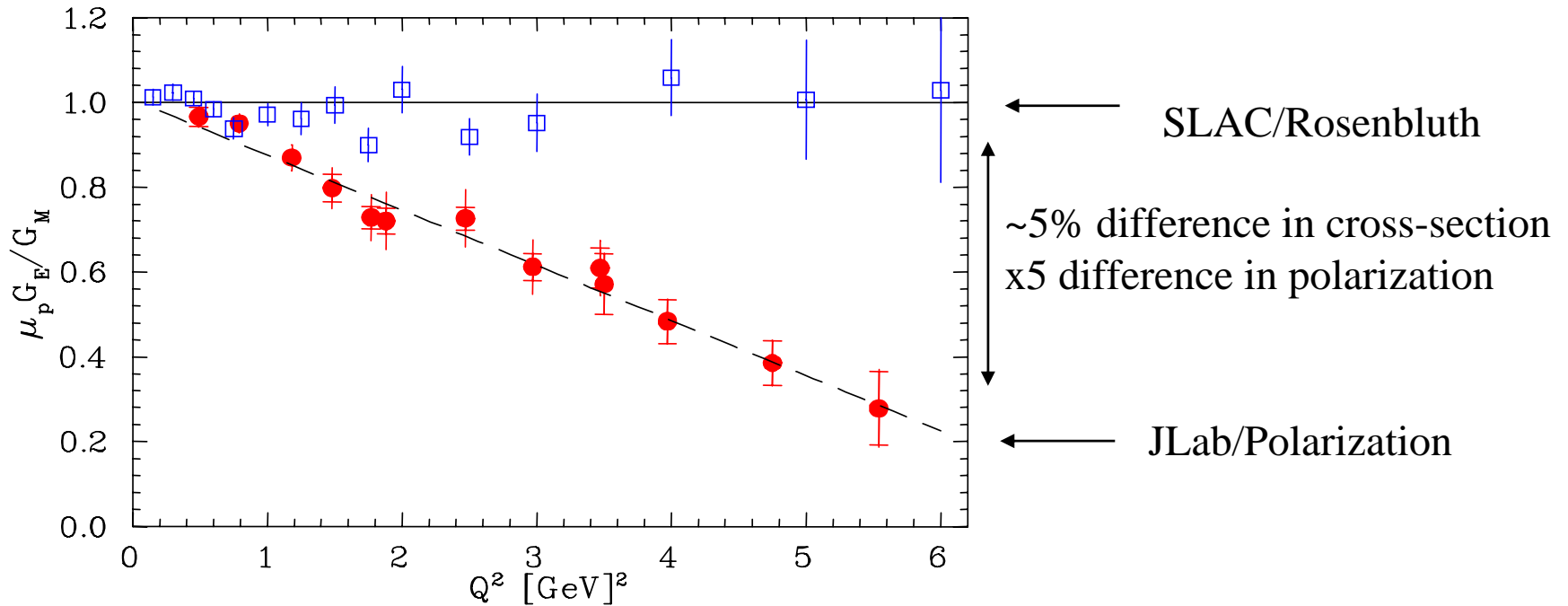
$$G_E = F_1 - \tau F_2, \quad G_M = F_1 + F_2$$

$$(P_y = 0)$$

Latter due to: Akhiezer, Rekalov; Arnold, Carlson, Gross



Do the techniques agree?



- Both early SLAC and Recent JLab experiments on (super)Rosenbluth separations followed $G_E/G_M \sim \text{const}$, see I.A. Quattan et al., Phys.Rev.Lett. 94:142301,2005
- JLab measurements using polarization transfer technique give different results (Jones'00, Gayou'02)

Radiative corrections, in particular, a short-range part of 2-photon exchange is a likely origin of the discrepancy



Basics of QED radiative corrections



(First) Born approximation

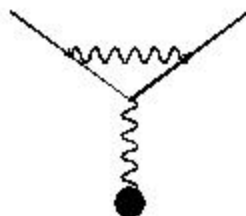


Initial-state radiation



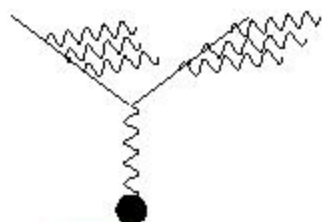
Final-state radiation

Cross section $\sim d\omega/\omega \Rightarrow$ integral diverges logarithmically: **IR catastrophe**



Vertex correction \Rightarrow cancels divergent terms; Schwinger (1949)

$$\sigma_{\text{exp}} = (1 + \delta)\sigma_{\text{Born}}, \quad \delta = \frac{-2\alpha}{\pi} \left\{ \left(\ln \frac{E}{\Delta E} - \frac{13}{12} \right) \left(\ln \frac{Q^2}{m_e^2} - 1 \right) + \frac{17}{36} + \frac{1}{2} f(\theta) \right\}$$

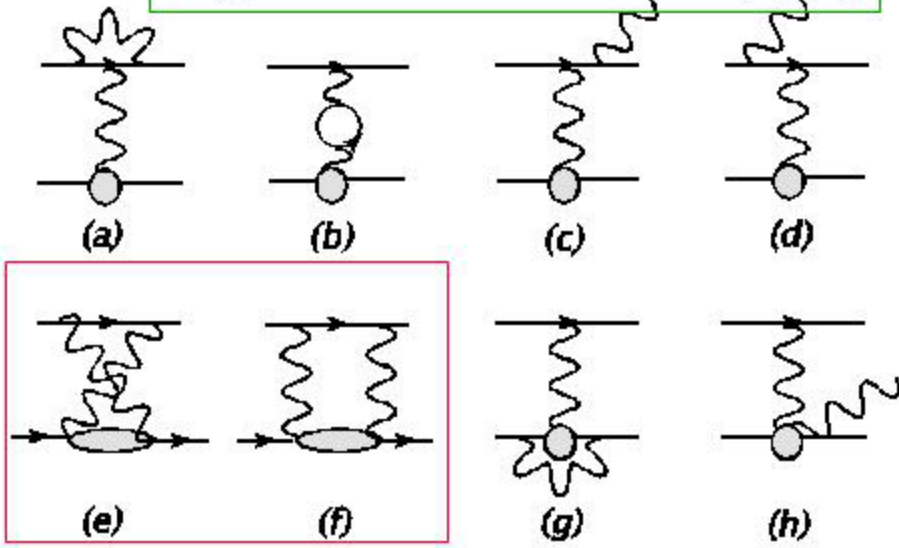
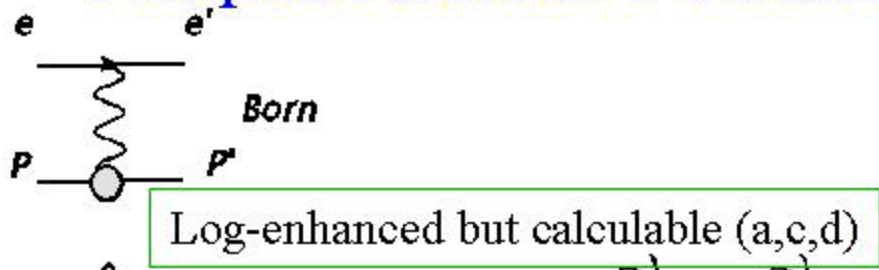


Multiple soft-photon emission: solved by exponentiation, Yennie-Frautschi-Suura (YFS), 1961

$$(1 + \delta) \rightarrow e^{\delta}$$



Complete radiative correction in $O(\alpha_{em})$



Radiative Corrections:

- Electron vertex correction (a)
- Vacuum polarization (b)
- Electron bremsstrahlung (c,d)
- Two-photon exchange (e,f)
- Proton vertex and VCS (g,h)
- Corrections (e-h) depend on the nucleon structure
- Meister & Yennie; Mo & Tsai
- Further work by Bardin & Shumeiko; Maximon & Tjon; AA, Akushevich, Merenkov;
- Guichon & Vanderhaeghen '03:
Can (e-f) account for the Rosenbluth vs. polarization experimental discrepancy? Look for ~3% ...

Main issue: Corrections dependent on nucleon structure

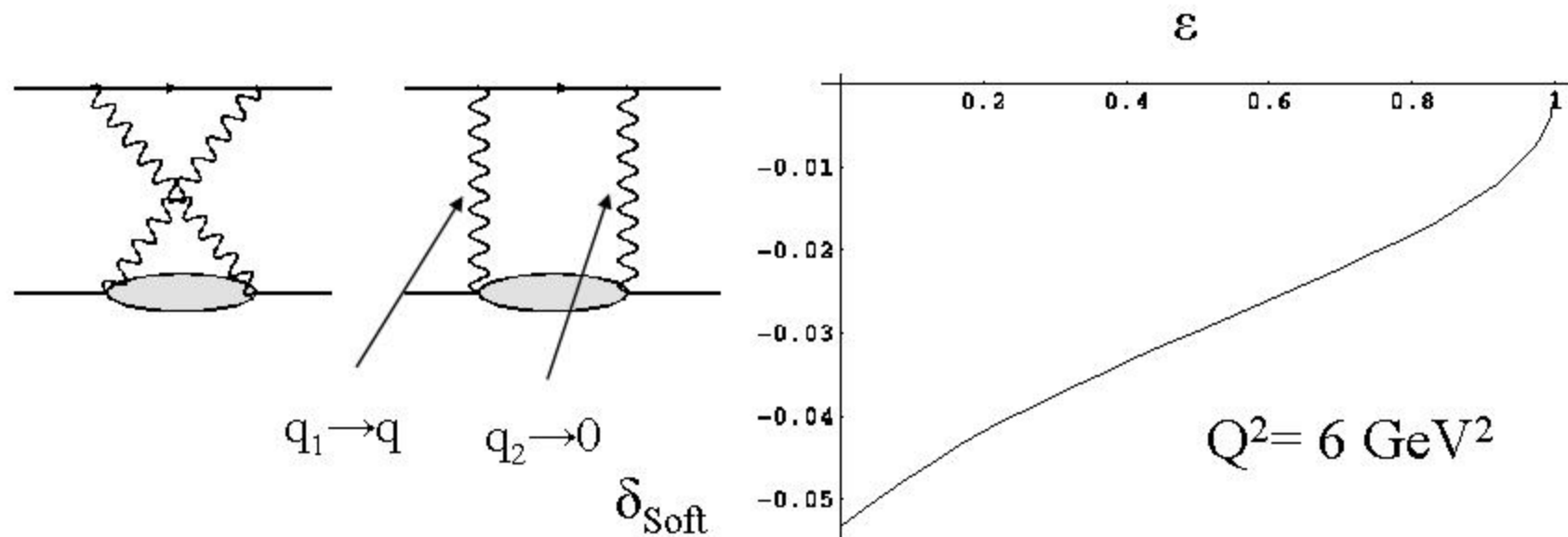
Model calculations:

- Blunden, Melnitchouk, Tjon, Phys.Rev.Lett. **91**:142304, 2003
- Chen, AA, Brodsky, Carlson, Vanderhaeghen, Phys.Rev.Lett. **93**:122301, 2004



Separating *soft* 2-photon exchange

- Tsai; Maximon & Tjon ($k \rightarrow 0$); similar to Coulomb corrections at low Q^2
- Grammer & Yennie prescription PRD 8, 4332 (1973) (also applied in QCD calculations)
- Shown is the resulting (soft) QED correction to [cross section](#)
- **Already included in experimental data analysis**
- **NB:** Corresponding effect to polarization transfer and/or asymmetry is zero



What is missing in the calculation?

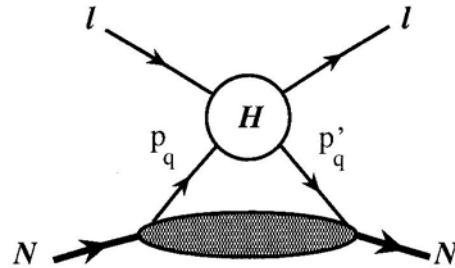
- 2-photon exchange contributions for non-soft intermediate photons
 - Can estimate based on a text-book example from *Berestetsky, Lifshitz, Pitaevsky: Quantum Electrodynamics*
 - Double-log asymptotics of electron-quark backward scattering

$$\delta = -\frac{e_q e}{8\pi^3} \log^2 \frac{s}{m_q^2}$$

- Negative sign for backward ep-scattering; zero for forward scattering → Can (at least partially) mimic the electric form factor contribution to the Rosenbluth cross section
- Numerically ~3-4% (for SLAC kinematics and $m_q \sim 300$ MeV)
- **Motivates a more detailed calculation of 2-photon exchange at quark level**

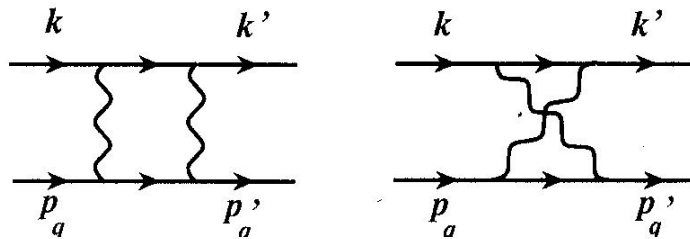


Calculations using Generalized Parton Distributions



Model schematics:

- Hard eq-interaction
- GPDs describe quark emission/absorption
- Soft/hard separation
 - Use Grammer-Yennie prescription



Hard interaction with a quark

AA, Brodsky, Carlson, Chen, Vanderhaeghen,
Phys.Rev.Lett.**93**:122301,2004; Phys.Rev.D**72**:013008,2005



Short-range effects; on-mass-shell quark (AA, Brodsky, Carlson, Chen, Vanderhaeghen)

Two-photon probe directly interacts with a (massless) quark
Emission/reabsorption of the quark is described by GPDs

$$A_{eq \rightarrow eq}^{2\gamma} = \frac{e_q^2}{t} \frac{\alpha_{em}}{2\pi} (V_\mu^e \otimes V_\mu^q \times f_V + A_\mu^e \otimes A_\mu^q \times f_A),$$

$$V_\mu^{e,q} = \bar{u}_{e,q} \gamma_\mu u_{e,q}, \quad A_\mu^{e,q} = \bar{u}_{e,q} \gamma_\mu \gamma_5 u_{e,q}$$

$$f_V = -2[\log(-\frac{u}{s}) + i\pi] \log(-\frac{t}{\lambda^2}) - \frac{t}{2} [\frac{1}{s} (\log(\frac{u}{t}) + i\pi) - \frac{1}{u} \log(-\frac{s}{t})] +$$

$$+ \frac{(u^2 - s^2)}{4} [\frac{1}{s^2} (\log^2(\frac{u}{t}) + \pi^2) + \frac{1}{u^2} \log(-\frac{s}{t}) (\log(-\frac{s}{t}) + i2\pi)] + i\pi \frac{u^2 - s^2}{2su}$$

$$f_A = -\frac{t}{2} [\frac{1}{s} (\log(\frac{u}{t}) + i\pi) + \frac{1}{u} \log(-\frac{s}{t})] +$$

$$+ \frac{(u^2 - s^2)}{4} [\frac{1}{s^2} (\log^2(\frac{u}{t}) + \pi^2) - \frac{1}{u^2} \log(-\frac{s}{t}) (\log(-\frac{s}{t}) + i2\pi)] + i\pi \frac{t^2}{2su}$$

Note the additional effective (axial-vector)² interaction; absence of mass terms



'Hard' contributions to generalized form factors

GPD integrals

$$A \equiv \int_{-1}^1 \frac{dx}{x} \frac{[(\hat{s} - \hat{u}) \tilde{f}_1^{hard} - \hat{s}\hat{u}\tilde{f}_3]}{(s-u)} \sum_q e_q^2 (H^q + E^q),$$

$$B \equiv \int_{-1}^1 \frac{dx}{x} \frac{[(\hat{s} - \hat{u}) \tilde{f}_1^{hard} - \hat{s}\hat{u}\tilde{f}_3]}{(s-u)} \sum_q e_q^2 (H^q - \tau E^q)$$

$$C \equiv \int_{-1}^1 \frac{dx}{x} \tilde{f}_1^{hard} \operatorname{sgn}(x) \sum_q e_q^2 \tilde{H}^q$$

Two-photon-exchange form factors from GPDs

$$\delta \tilde{G}_M^{hard} = C$$

$$\delta \tilde{G}_E^{hard} = -\left(\frac{1+\varepsilon}{2\varepsilon}\right) (A - C) + \sqrt{\frac{1+\varepsilon}{2\varepsilon}} B$$

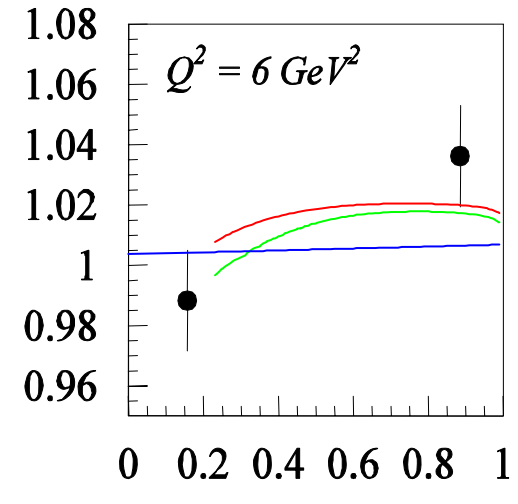
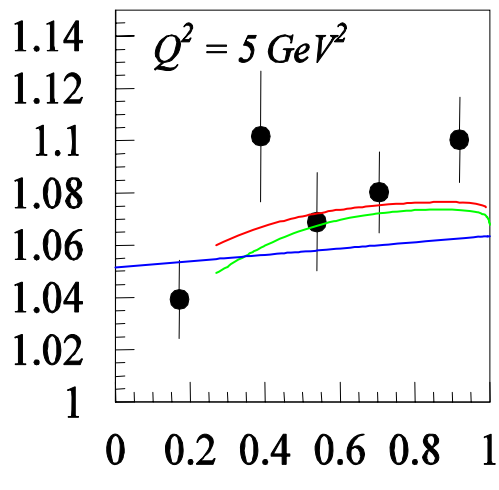
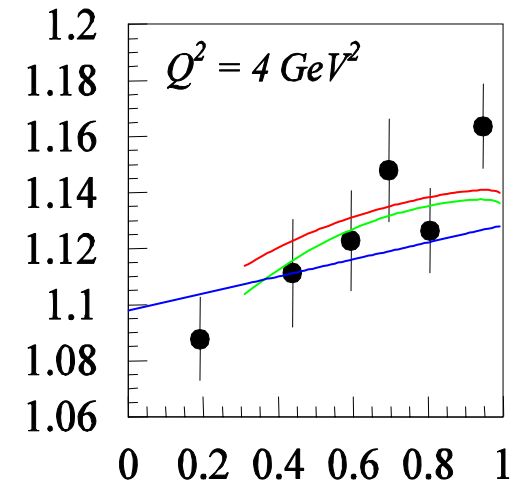
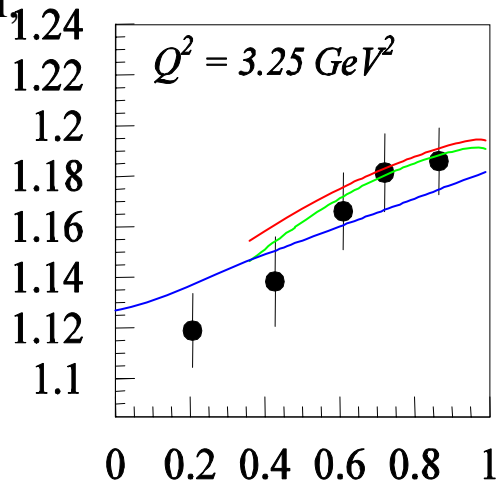
$$\tilde{F}_3 = \frac{M^2}{\nu} \left(\frac{1+\varepsilon}{2\varepsilon}\right) (A - C)$$



Two-Photon Effect for Rosenbluth Cross Sections

$$\sigma_R / (\mu_p G_D)^2$$

- Data shown are from Andivahis et al. PRD 50, 5491 (1994)
- Included GPD calculation of two-photon-exchange effect
- Qualitative agreement with data:
- Discrepancy likely reconciled



ϵ

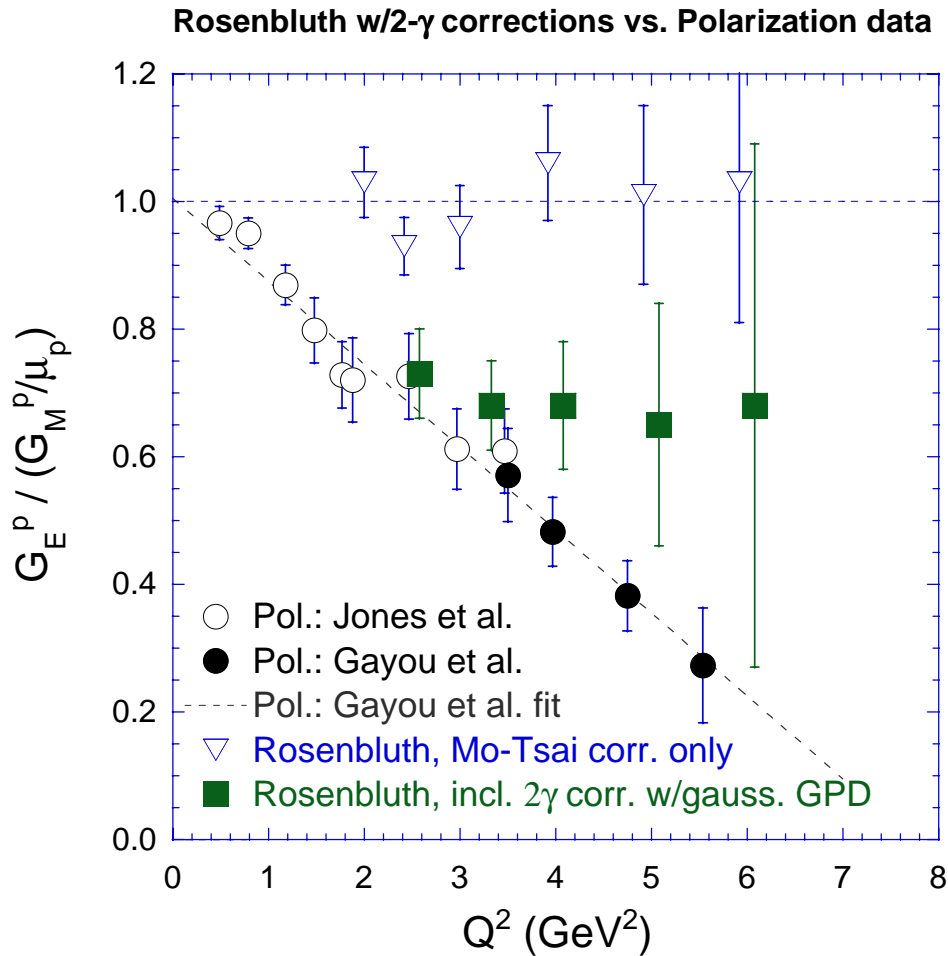
ϵ



Updated Ge/Gm plot

AA, Brodsky, Carlson, Chen, Vanderhaeghen,

Phys.Rev.Lett.93:122301, 2004; Phys.Rev.D72:013008, 2005



Full Calculation of Bethe-Heitler Contribution

Additional work by AA et al., using MASCARAD (*Phys.Rev.D64:113009,2001*)
 Full calculation including soft and hard bremsstrahlung

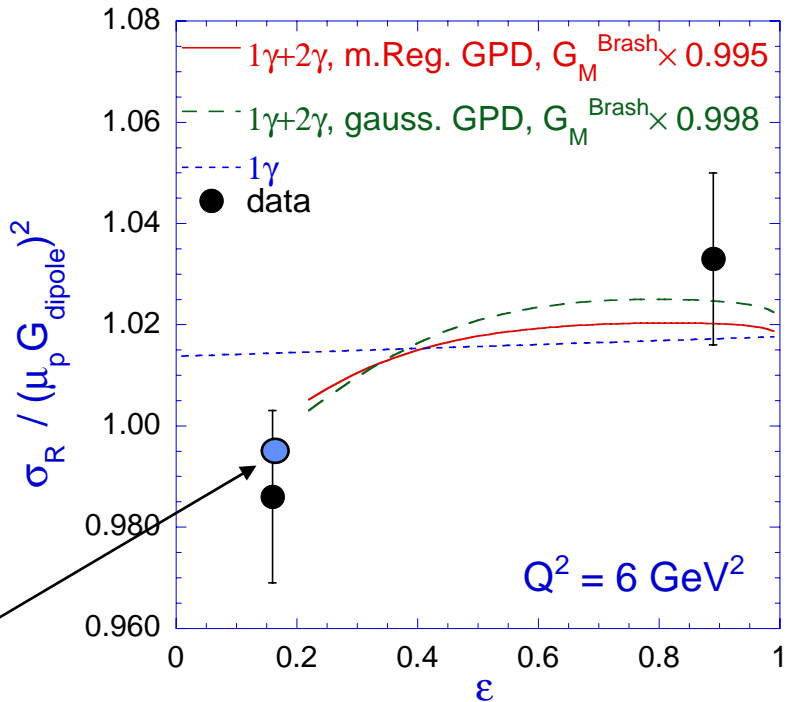
Radiative leptonic tensor in full form
 AA et al, *PLB 514, 269 (2001)*

$$L'_{\mu\nu} = -\frac{1}{2} \text{Tr}(\hat{k}_2 + m) \Gamma_{\mu\alpha} (1 + \gamma_5 \hat{\xi}_e) (\hat{k}_1 + m) \bar{\Gamma}_{\alpha\nu}$$

$$\Gamma_{\mu\alpha} = \left(\frac{k_{1\alpha}}{k \cdot k_1} - \frac{k_{2\alpha}}{k \cdot k_2} \right) \gamma_\mu \left[\frac{\gamma_\mu \hat{k} \gamma_\alpha}{2k \cdot k_1} - \frac{\gamma_\alpha \hat{k} \gamma_\mu}{2k \cdot k_2} \right]$$

$$\Gamma_{\alpha\nu} = \left(\frac{k_{1\alpha}}{k \cdot k_1} - \frac{k_{2\alpha}}{k \cdot k_2} \right) \gamma_\nu \left[\frac{\gamma_\alpha \hat{k} \gamma_\nu}{2k \cdot k_1} - \frac{\gamma_\nu \hat{k} \gamma_\alpha}{2k \cdot k_2} \right]$$

Cross section for ep elastic scattering

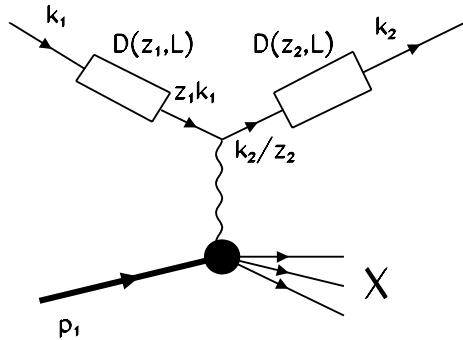


Additional effect of full soft+hard brem \rightarrow +1.2% correction to ϵ -slope
Resolves additional ~25% of Rosenbluth/polarization discrepancy!

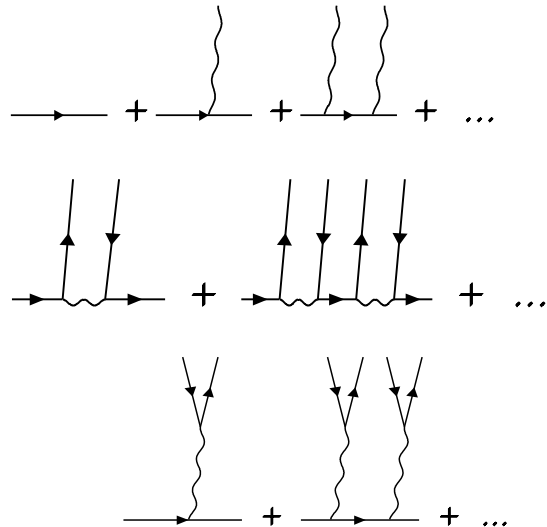


Electron Structure Functions

- For polarized $ep \rightarrow e'X$ scattering, AA et al, JETP 98, 403 (2004); elastic ep: AA et al. PRD 64, 113009 (2001).



- Resummation technique for collinear photons (=peaking approx.)
- Difference <0.5% from previous calculation including hard brem

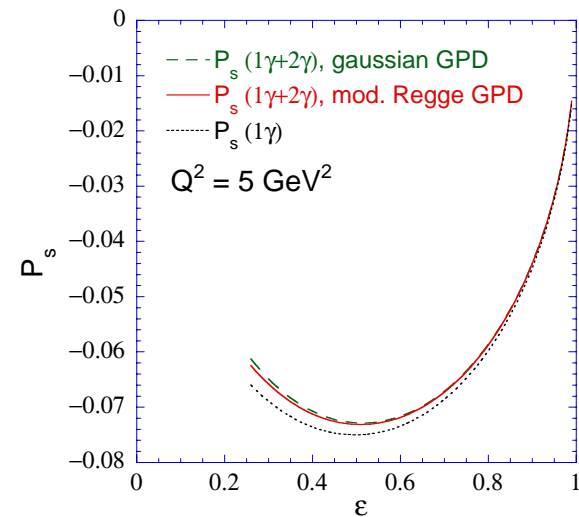
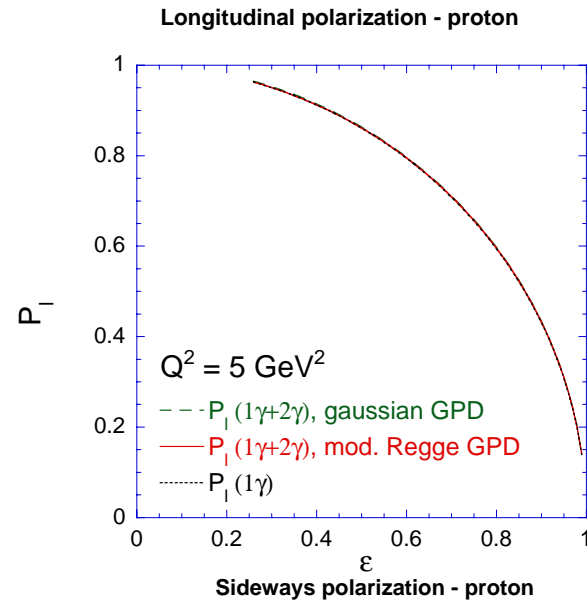
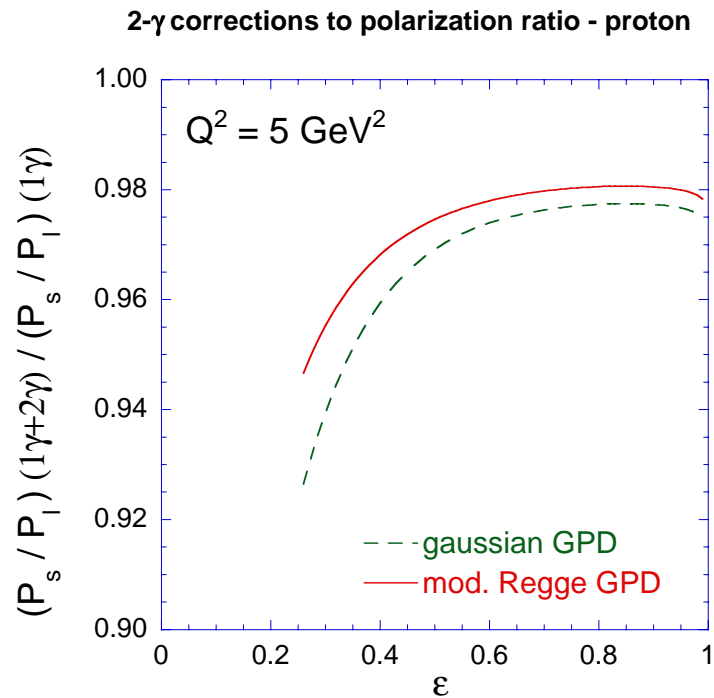


- Bystritskiy, Kuraev, Tomasi-Gustafson (2007) claimed this approach resolves Rosenbluth vs polarization discrepancy... **but used incorrect energy cutoff $\Delta E/E$ of 3% (instead of e.g. 1.5%)** => **miscalculated rad.correction by ~5% (absolute)**



Polarization transfer

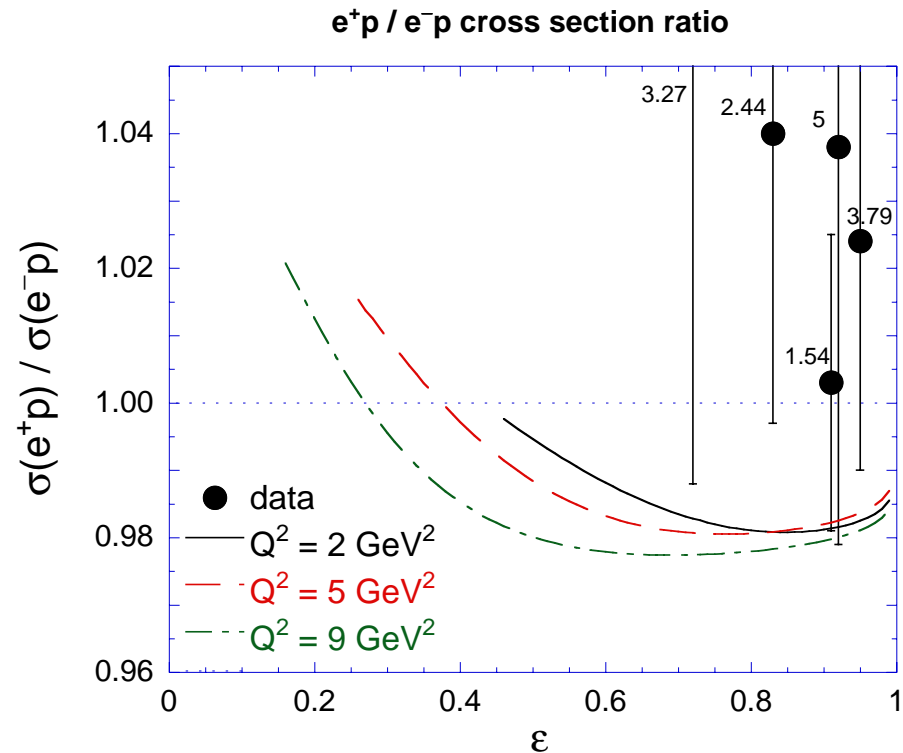
- Also corrected by two-photon exchange, but with little impact on G_{ep}/G_{mp} extracted ratio



Charge asymmetry

- Cross sections of electron-proton scattering and positron-proton scattering are equal in one-photon exchange approximation
 - Different for two- or more photon exchange

$$\sigma_{\pm} = \sigma_{1\gamma} \pm \sigma_{2\gamma}$$



Measured in JLab Experiment 04-116,
 Spokepersons: W. Brooks, AA, J. Arrington, K.Joo, B.Raue, L.Weinstein
 First run scheduled for Fall 2006



Single-Spin Asymmetries in Elastic Scattering

Parity-conserving

- Observed spin-momentum correlation of the type:

$$\vec{s} \cdot \vec{k}_1 \times \vec{k}_2$$

where $k_{1,2}$ are initial and final electron momenta, s is a polarization vector of a target OR beam

- For elastic scattering asymmetries are due to *absorptive part* of 2-photon exchange amplitude

Parity-Violating

$$\vec{s} \cdot \vec{k}_1$$



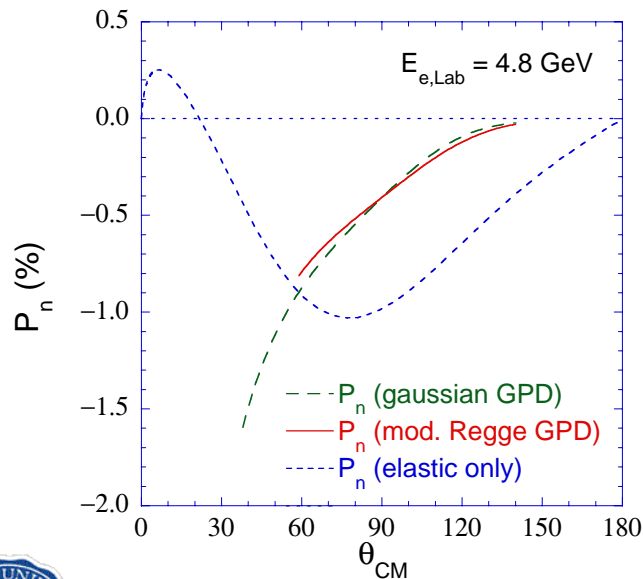
Quark+Nucleon Contributions to A_n

- Single-spin asymmetry or polarization normal to the scattering plane
- Handbag mechanism prediction for single-spin asymmetry/polarization of elastic ep-scattering on a polarized proton target

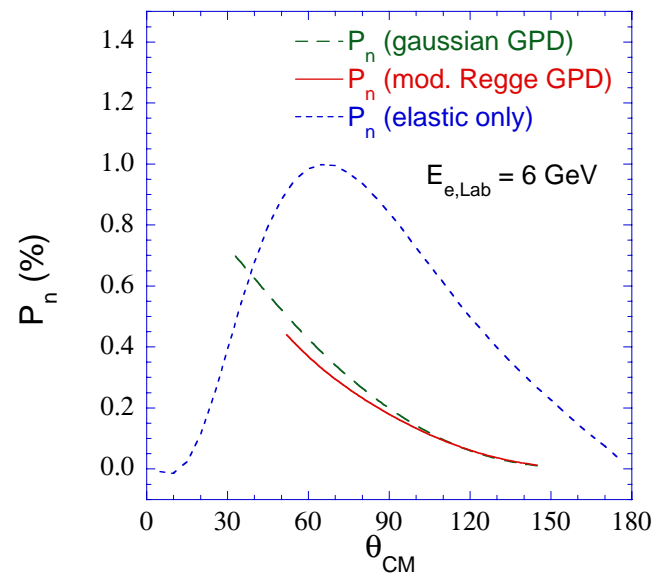
$$A_n = \sqrt{\frac{2\varepsilon(1+\varepsilon)}{\tau}} \frac{1}{\sigma_R} \left[G_E \operatorname{Im}(A) - \sqrt{\frac{1+\varepsilon}{2\varepsilon}} G_M \operatorname{Im}(B) \right] \quad \textit{Only minor role of quark mass}$$

No dependence on GPD \tilde{H}

Normal Polarization or Analyzing Power - Neutron



Normal Polarization or Analyzing Power - Proton



Proton Mott Asymmetry at Higher Energies

Spin-orbit interaction of electron moving in a Coulomb field

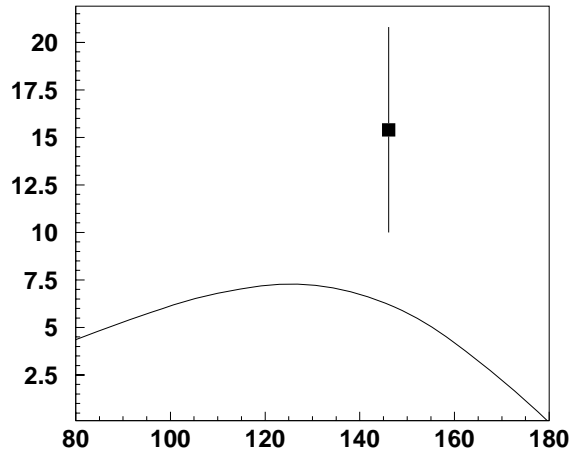
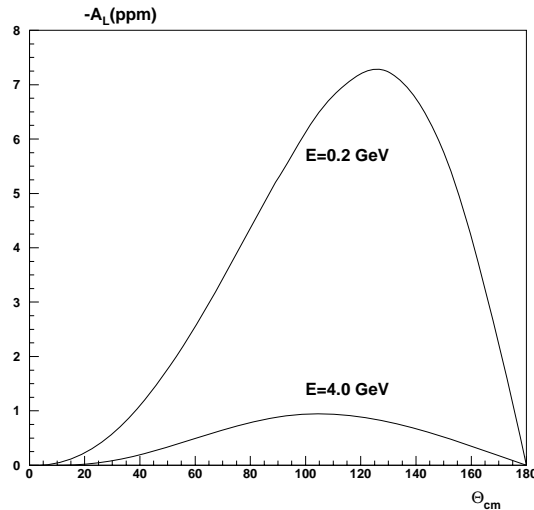
N.F. Mott, Proc. Roy. Soc. London, Set. A **135**, 429 (1932);

BNSA for electron-muon scattering: Barut, Fronsdal, Phys.Rev.120, 1871 (1960);

BNSA for electron-proton scattering: Afanasev, Akushevich, Merenkov, hep-ph/0208260

Transverse beam SSA, units are parts per million

Figures from AA et al, hep-ph/0208260



- **Due to absorptive part of two-photon exchange amplitude; shown is elastic contribution**
- **Nonzero effect observed by SAMPLE Collaboration (S.Wells et al., PRC63:064001,2001) for 200 MeV electrons**

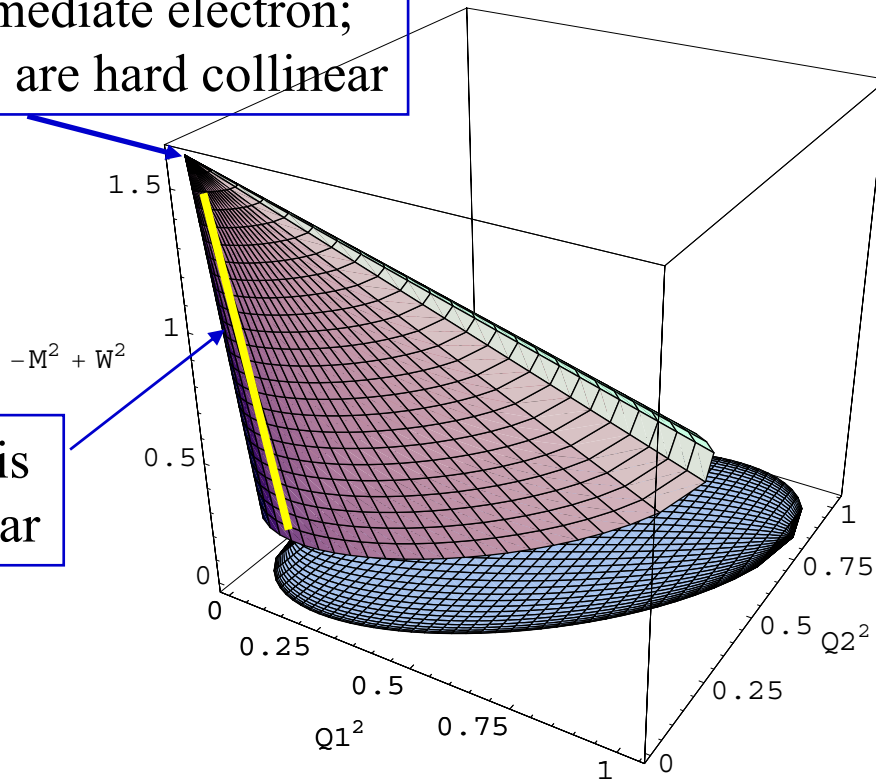


Phase Space Contributing to the absorptive part of 2γ -exchange amplitude

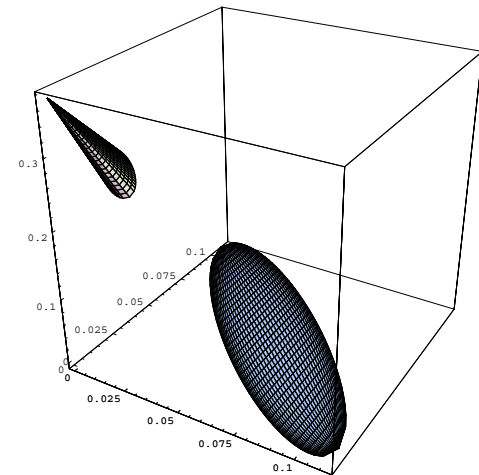
- 2-dimensional integration (Q_1^2, Q_2^2) for the elastic intermediate state
- 3-dimensional integration (Q_1^2, Q_2^2, W^2) for inelastic excitations

'Soft' intermediate electron;
Both photons are hard collinear

One photon is
Hard collinear

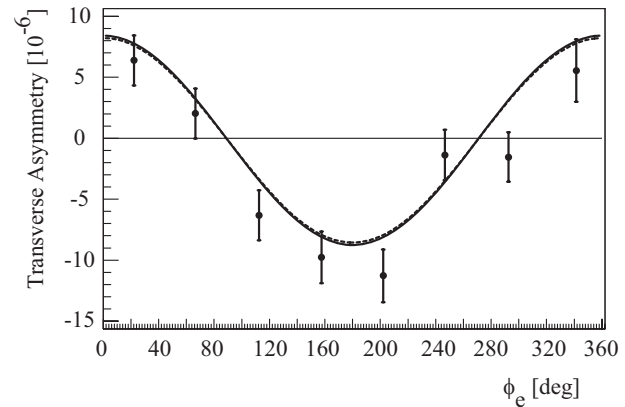


Examples: MAMI A4
E= 855 MeV
 $\Theta_{cm}= 57$ deg;
SAMPLE, E=200 MeV



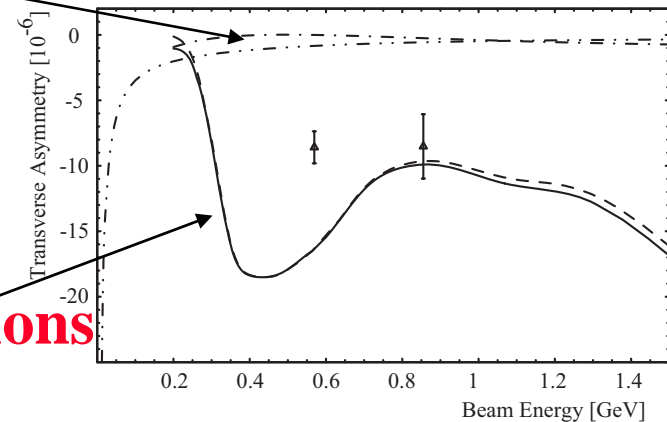
MAMI data on Mott Asymmetry

- F. Maas et al., [MAMI A4 Collab.]
Phys.Rev.Lett.**94**:082001, 2005
- Pasquini, Vanderhaeghen:
Phys.Rev.C**70**:045206,2004
Surprising result: Dominance of inelastic
intermediate excitations

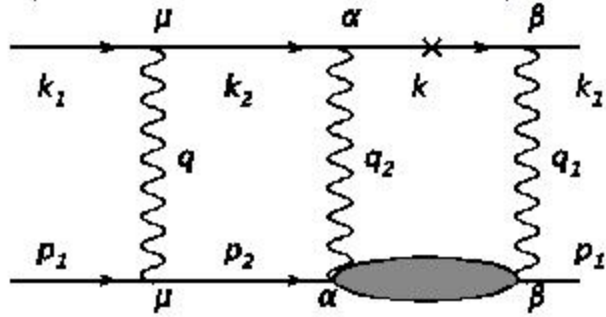


Elastic intermediate
state

Inelastic excitations
dominate



Beam Normal Asymmetry (=Mott asymmetry) (AA, Merenkov)



$$A_n^{e,p} = -\frac{\alpha Q^2}{\pi^2 D(s, Q^2)} \text{Im} \int \frac{d^3 \mathbf{k}}{2k_0} \cdot \frac{L_{\mu\alpha\beta} H_{\mu\alpha\beta}}{Q_1^2 Q_2^2}$$

$$L_{\mu\alpha\beta} = \frac{1}{4} \text{Tr}(\hat{\mathbf{k}}_2 + m_e) \gamma_\mu (\hat{\mathbf{k}}_1 + m_e) (1 - \gamma_5 \hat{\xi}_e) \gamma_\beta (\hat{\mathbf{k}} + m_e) \gamma_\alpha$$

$$H_{\mu\alpha\beta} = \frac{1}{4} \text{Tr}(\hat{\mathbf{p}}_2 + M) \Gamma_\mu (\hat{\mathbf{p}}_1 + M) (1 - \gamma_5 \hat{\xi}_p) T_{\beta\alpha}$$

$$\hat{\mathbf{a}} \equiv \mathbf{a}_{\mu} \gamma_\mu$$

$$L_{\mu\alpha\beta} \mathbf{q}_\mu = L_{\mu\alpha\beta} \mathbf{q}_{2\alpha} = L_{\mu\alpha\beta} \mathbf{q}_{1\beta} = H_{\mu\alpha\beta} \mathbf{q}_\mu = H_{\mu\alpha\beta} \mathbf{q}_{2\alpha} = H_{\mu\alpha\beta} \mathbf{q}_{1\beta} = 0$$

EM gauge invariance important for cancellation of collinear singularity for target asymmetry

$$L_{\mu\alpha\beta} H_{\mu\alpha\beta} \rightarrow 0 \quad \text{if} \quad Q_1^2 \text{ and/or } Q_2^2 \rightarrow 0$$

Feature of the normal beam asymmetry: After m_e is factored out, the remaining expression is singular when virtuality of the photons reach zero in the loop integral.

$$L_{\mu\alpha\beta} H_{\mu\alpha\beta} \rightarrow m_e \cdot \text{const} \quad \text{if} \quad Q_1^2 \text{ and/or } Q_2^2 \rightarrow 0 \Rightarrow A \sim m_e \log^2 \frac{Q^2}{m_e^2}, m_e \log \frac{Q^2}{m_e^2}$$

Also calculations by Vanderhaeghen, Pasquini (2004); Gorchtein (2005); Borisyuk&Kobushkin (2005) confirm *quasi-real photon exchange* enhancement of inelastic intermediate excitations



Special property of Mott asymmetry

AA, Merenkov, Phys.Lett.B599:48,2004, Phys.Rev.D70:073002,2004;
+Erratum (hep-ph/0407167v2)

- Reason for the unexpected behavior: hard collinear quasi-real photons
 - Intermediate photon is collinear to the parent electron
 - It generates a dynamical pole and logarithmic enhancement of inelastic excitations of the intermediate hadronic state
 - For $s \gg -t$ and above the resonance region, the asymmetry is given by:

$$A_n^e = \sigma_{\gamma p} \frac{(-m_e)\sqrt{Q^2}}{8\pi^2} \cdot \frac{F_1 - \tau F_2}{F_1^2 + \tau F_2^2} \left(\log\left(\frac{Q^2}{m_e^2}\right) - 2 \right)$$

In addition multiply by a standard diffractive factor $\exp(-BQ^2)$, where $B=3.5-4 \text{ GeV}^{-2}$
Compare with no-inelastic-excitation asymmetry at small $\theta \Rightarrow$ may differ by orders of magnitude depending on kinematics of the measurement

$$A_n^e \propto \alpha \frac{m_e}{\sqrt{s}} \theta^3$$

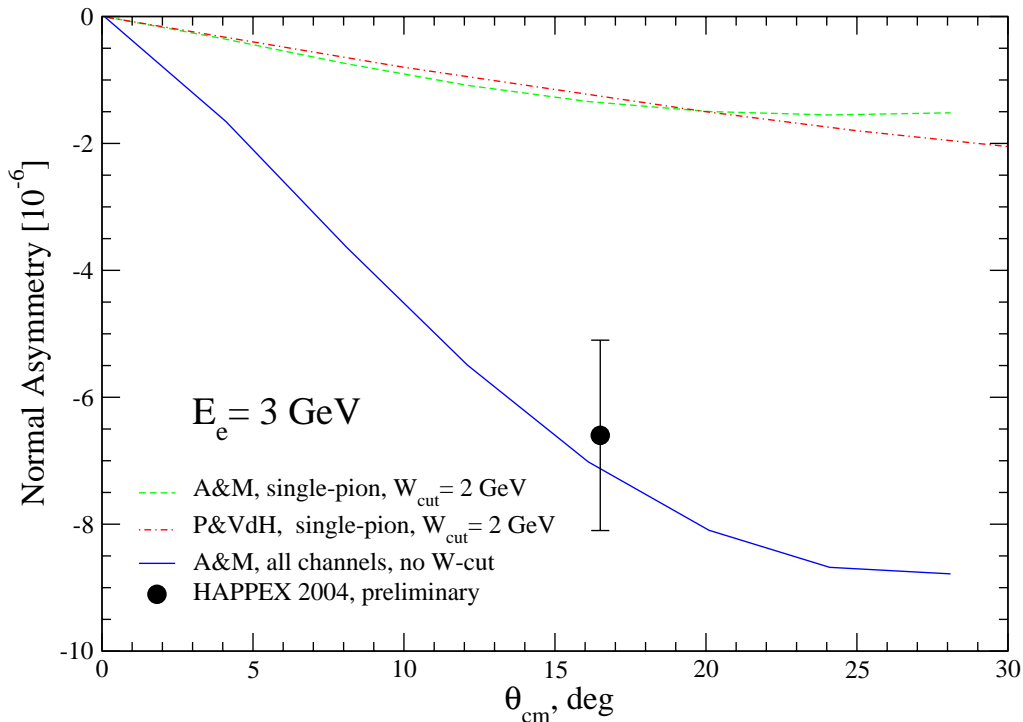


Predictions vs experiment for Mott asymmetry

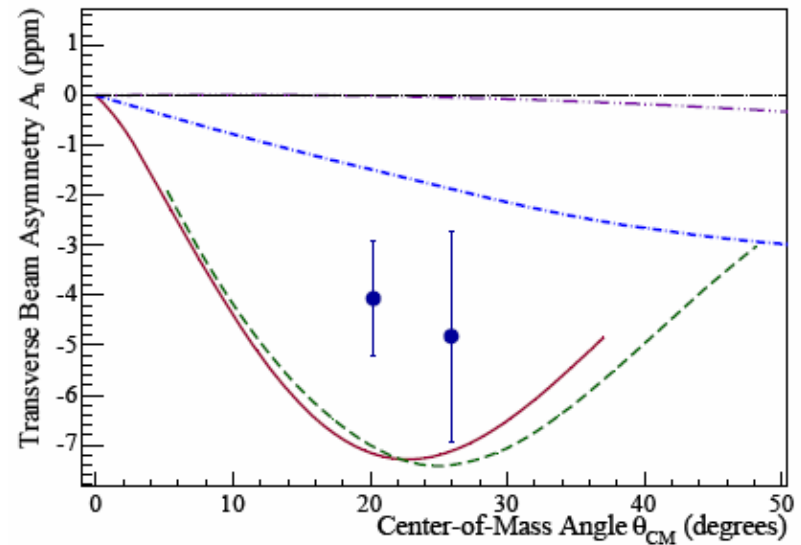
Use fit to experimental data on $\sigma_{\gamma p}$ and exact 3-dimensional integration over phase space of intermediate 2 photons

Normal beam asymmetry for elastic ep-scattering

Unitarity-based model predictions

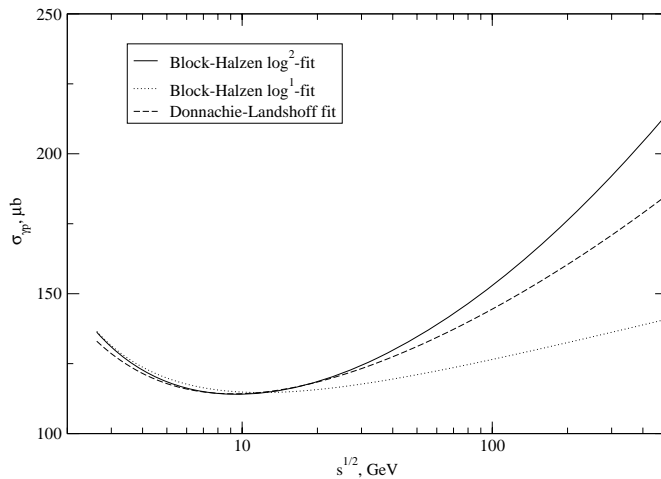
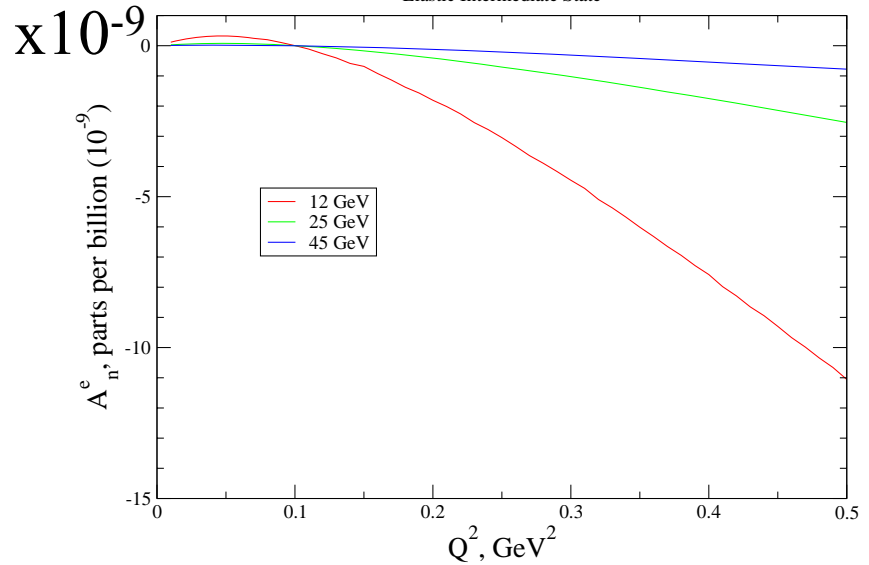
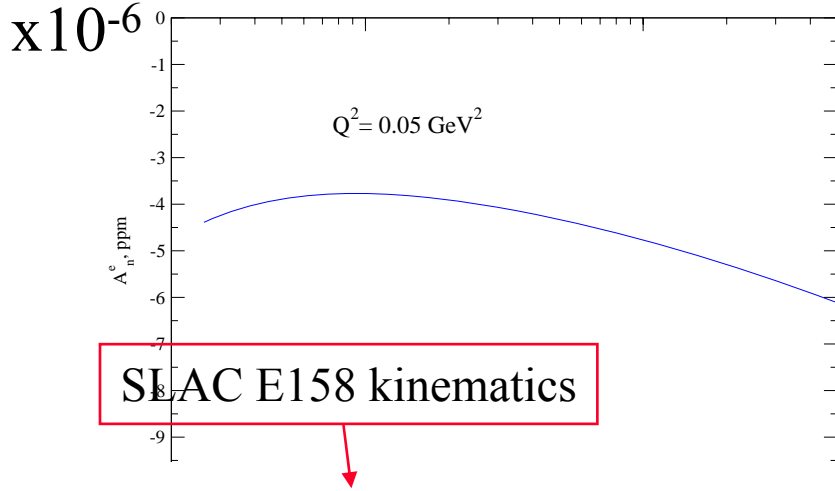


10 arXiv 0705.1525[nucl-ex]



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No suppression for Mott asymmetry with energy at fixed Q^2



Parts-per-million vs. parts-per-billion scales: a consequence of soft Pomeron exchange, and hard collinear photon exchange

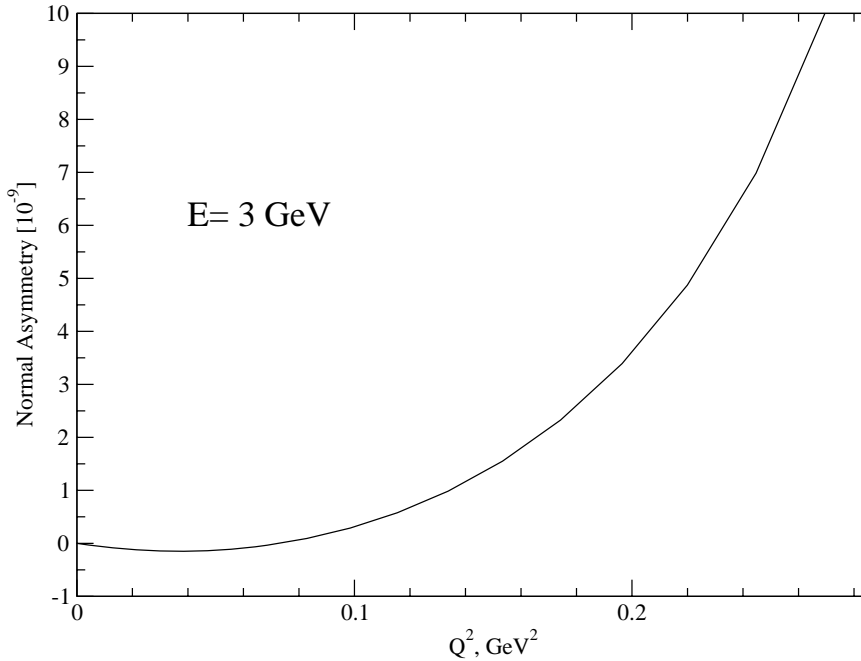


Mott Asymmetry on Nuclei

- Important systematic correction for parity-violation experiments (~ 10 ppm for HAPPEX on ^4He , ~ 5 ppm for PREX on Pb)
 - Measures (integrated) absorptive part of Compton scattering amplitude
- Coulomb distortion: only 10-10 effect (Cooper & Harowitz Phys Rev C72:034602 2005)

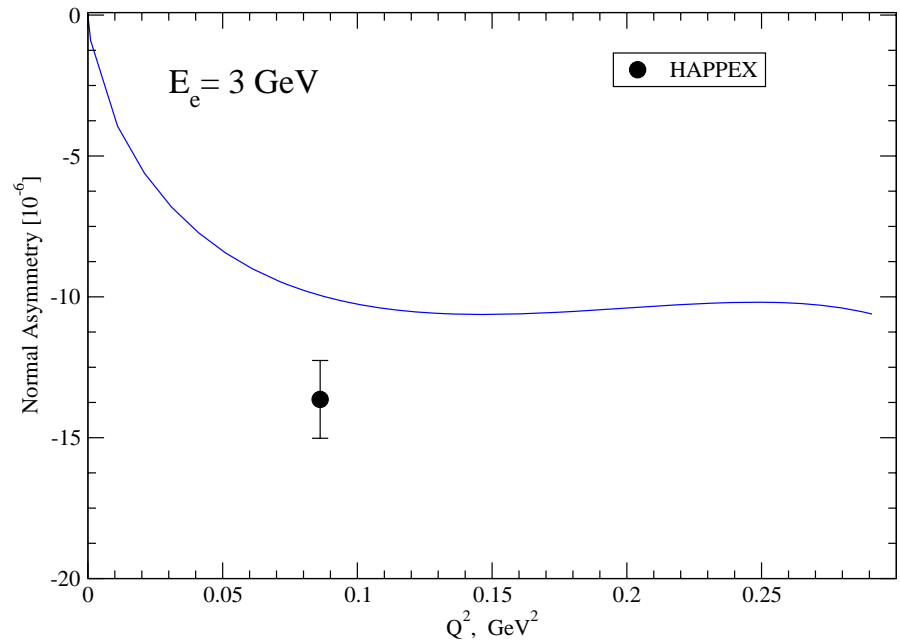
Normal beam asymmetry for elastic $e\text{-}^4\text{He}$ scattering

Contribution of elastic intermediate state



Normal beam asymmetry for elastic $e\text{-}^4\text{He}$ scattering

Unitarity-based model predictions



Five orders of magnitude enhancement in HAPPEX asymmetries due to excitation of inelastic intermediate states in 2γ -exchange (AA, Merenkov; use Compton data from Erevan)



Summary: SSA in Elastic ep-Scattering

- Collinear photon exchange present in (light particle) beam SSA
- **Warning: Models violating EM gauge invariance will suffer collinear divergence** for target SSA
- VCS amplitude in *beam asymmetry* is enhanced in different kinematic regions compared to *target asymmetry*
- Strong-interaction dynamics for Mott asymmetry in small-angle ep-scattering above the resonance region is *soft diffraction*
 - *For the diffractive mechanism A_n is a) not suppressed with beam energy and b) does not grow with Z*
 - *Confirmed experimentally → first observation of diffractive elastic electron-hadron scattering*



Other theoretical developments

- Blunden et al., Phys.Rev.C72:034612, 2005
Approximate proton Compton amplitude by Born terms
- Kondratyuk et al., nucl-th/0506026
Add intermediate Δ -excitation to the above
- Carlson, Vanderhaeghen, Pascalutsa, hep-ph/0509055 (PRL'05)
GPD approach extended to $N \rightarrow \Delta$ transition
- Borisyuk, Kobushkin, Phys.Rev.C72:035207,2005
- Jain et al (2006): Model nonlocal ep-interaction, confirm results of others authors for corrections to Rosenbluth separation
- AA, Strikman, Weiss: Single-spin asymmetries from 2γ -exchange in DIS (see talk by Weiss at this workshop)



Two-photon exchange for electron-nucleon scattering

- Model calculations of 2γ -exchange radiative corrections bring into agreement the results of polarization transfer and Rosenbluth techniques for GEp measurements
- Full treatment of brem corrections removes $\sim 25\%$ of R/P discrepancy **in addition** to 2γ
→ Important to compute conventional corrections accurately
- **Experimental tests of two-photon exchange**
 - C-violation in electron vs positron elastic scattering (**JLab E04-116, E-07-005**)
 - Measurement of nonlinearity of Rosenbluth plot (**JLab E05-017**)
 - Search for deviation of angular dependence of polarization and/or asymmetries from Born behavior at fixed Q^2 (**JLab E04-019**)
 - Elastic single-spin asymmetry or induced polarization (**JLab E05-015**)
 - Extended to inelastic (e, e') in **E-07-013**
 - 2γ normal beam asymmetry measurements parallel to parity-violating experiments (**HAPPEX, G0, PREX**)

These studies provide

- a) Testing precision of the electromagnetic probe and*
- b) Double-virtual VCS studies*

