Low Energy Precision Tests of Supersymmetry



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M.R-M & S. Su, hep-ph/0612057 J. Erler & M.R-M, PPNP 54, 351 (2005)

Outline

- I. Motivation: Why New Symmetries ? Why Low Energy Probes ?
- II. Prime Suspect: Supersymmetry
- **III.** Low Energy Precision Tests
 - Weak Decays
 - PVES

I. Motivation

Why New Symmetries ? Why Low Energy Probes ?









There must have been additional symmetries in the earlier Universe to

- Unify all matter, space, & time
- Stabilize the weak scale
- Produce all the matter that exists
- Account for neutrino properties
- Give self-consistent quantum gravity

Supersymmetry, GUT's, extra dimensions...

What are the new fundamental symmetries?

Two frontiers in the search

Collider experiments (pp, e^+e^- , etc) at higher energies (E >> M_Z)

High energy physics

Indirect searches at lower energies (E < M_Z) but high precision



Particle, nuclear & atomic physics

Precision Probes of New Symmetries

New Symmetries

- 1. Origin of Matter
- 2. Unification & gravity
- 3. Weak scale stability
- 4. Neutrinos













Precision Probes of New Symmetries



Probing Fundamental Symmetries beyond the SM:

Use precision lowenergy measurements to probe virtual effects of new symmetries & compare with collider results

Stunning SM Success

Precision, low energy measurements can probe for new symmetries in the desert

Precision ~ Mass Scale

$$\delta_{NEW} = \frac{\Delta O^{NEW}}{O^{SM}} \approx \frac{\alpha}{\pi} \left(\frac{M}{\tilde{M}}\right)^2$$

 $M = m_{\mu} \qquad \delta \sim 2 \times 10^{-9} \\ \delta^{exp} \sim 1 \times 10^{-9} \\ M = M_{W} \qquad \delta \sim 10^{-3}$

Interpretability

- Precise, reliable SM predictions
- Comparison of a variety of observables
- Special cases: SM-forbidden or suppressed processes

II. Prime suspect: Supersymmetry

SUSY: a candidate symmetry of the early Universe

3 of 4
Yes
Maybe so
Maybe
Probably necessary

Couplings unify with SUSY



SUSY protects G_F from shrinking



SUSY may help explain observed abundance of matter

Cold Dark Matter Candidate

 χ^0 Lightest SUSY particle

Baryonic matter: electroweak phase transition



SUSY: a candidate symmetry of the early Universe

Supersymmetry

Fermions

 $e_{L,R}$, $q_{L,R}$

 W, Z, γ, g

Bosons

H

105 new parameters: masses, mixing angles, CPV phases (40)

$$\mathcal{L}_{\text{soft}} = -\frac{1}{2}(M_3 \bar{g} \bar{g} + M_2 \bar{W} \bar{W} + M_1 \bar{B} \bar{B}) + c.c.$$

$$-(\bar{u} \mathbf{a}_u \bar{Q} H_u - \bar{d} \mathbf{a}_d \bar{Q} H_d - \bar{e} \mathbf{a}_e \bar{L} H_d) + c.c.$$

$$-\bar{Q}^{\dagger} \mathbf{m}_Q^2 \bar{Q} - \bar{L}^{\dagger} \mathbf{m}_L^2 \bar{L} - \bar{u} \mathbf{m}_Q^2 \bar{u}^{\dagger} - \bar{d} \mathbf{m}_d^2 \bar{d}^{\dagger} - \bar{e} \mathbf{m}_e^2 \bar{e}^{\dagger} - m_{H_u}^2 H_u^* H_u - m_{H_d}^2 H_d^* H_d$$

$$-(bH_u H_d + c.c.) \qquad (16)$$

Models: relate weak scale parameters to each other at high scales ("hidden sector")



SUSY and R Parity

$$P_{R} = (-1)^{3(B-L)} (-1)^{2S}$$

If nature conserves $P_R \rightarrow vertices$ have even number of superpartners

Consequences

• Lightest SUSY particle
$$\left(\widetilde{\chi}^0
ight)$$
 is stable \longrightarrow viable dark matter candidate

Proton is stable

Superpartners appear only in loops

R-Parity Violation (RPV)



RPV : Four-fermion Operators



III. SUSY & Weak Decays

Weak Decays & SUSY

 μ^{-}

 μ^{-}

$$d \rightarrow u \ e^{-} \ \overline{v}_{e}$$

$$s \rightarrow u \ e^{-} \ \overline{v}_{e}$$

$$b \rightarrow u \ e^{-} \ \overline{v}_{e}$$

$$u \ c \ v_{ud} \ V_{us} \ V_{ub}$$

$$V_{ud} \ V_{us} \ V_{ub}$$

$$V_{ub} \ V_{ub}$$

$$V_{ub} \ V_{ub}$$

$$V_{ub} \ V_{ub}$$



Weak Decays & SUSY

$$d \rightarrow u e^{-} \overline{v}_{e}$$
$$s \rightarrow u e^{-} \overline{v}_{e}$$
$$b \rightarrow u e^{-} \overline{v}_{e}$$





Kurylov, R-M, Su

R Parity Violation

No long-lived LSP or SUSY DM

Weak decays







Weak decays & new physics Correlations

$$\begin{array}{cccc}
d \rightarrow u \ e^{-} \ \overline{v}_{e} \\
s \rightarrow u \ e^{-} \ \overline{v}_{e} \\
b \rightarrow u \ e^{-} \ \overline{v}_{e}
\end{array}$$

$$\begin{pmatrix}
u \ c \ t
\end{pmatrix}
\begin{pmatrix}
V_{ud} \ V_{us} \ V_{ub} \\
V_{cd} \ V_{cs} \ V_{cb} \\
V_{td} \ V_{ts} \ V_{tb}
\end{pmatrix}
\begin{pmatrix}
d \\
s \\
b
\end{pmatrix}$$

$$\begin{array}{c}
\psi_{u} \ v_{us} \ v_{ub} \\
\psi_{cd} \ v_{cs} \ v_{cb} \\
\psi_{b} \\
\psi_{cd} \ v_{ts} \ v_{tb}
\end{pmatrix}$$

 ${oldsymbol{v}}_{\mu}$

Weak decays & SUSY : Correlations



Pion leptonic decay & SUSY



Pion leptonic decay & SUSY

New TRIUMF, PSI





 $\boldsymbol{v}_{\scriptscriptstyle M}$

Raidal, Santamaria; Cirigliano, Kurylov, R-M, Vogel







IV. SUSY & PVES

Q_W and SUSY Radiative Corrections

Tree Level

$$Q_W^f = g_V^f g_A^e$$

Radiative Corrections

Flavor-dependent

$$Q_W^f = \rho_{PV} \left(2I_3^f - 4Q_f \kappa_{PV} \sin^2 \theta_W \right) + \lambda_f$$

Constrained by Z-pole precision observables

Scale-dependent effective weak mixing

Flavor-independent

SUSY Radiative Corrections









$$\frac{\delta Q_W^e}{Q_W^e} \approx -30 \Delta_{12k}(\tilde{e}_R^k) \approx -45 \left(\frac{100 \, GeV}{m_{\tilde{e}_R^k}}\right)^2 |\lambda_{12k}|^2$$

 $\lambda_{12k} \sim 0.3$ for $m_{SUSY} \sim 1~TeV$ & $\delta Q_W{}^e/Q_W{}^e \sim 5\%$





Comparing A_d^{DIS} and $Q_w^{p,e}$



Low Energy Probes of SUSY



We're making progress...

...won't leave until the job is done...

...and open to new ideas.