

Quark models
of N^* electrocouplings at high Q^2
and their *connection to QCD*

Nucleon Resonance Structure in Exclusive
Electroproduction at High Photon Virtualities with
the CLAS 12 Detector

Jefferson Lab, May 16th, 2011

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Some observations

CQMs

- not **derived** from QCD (at least for the moment)
- **approximated** approaches to QCD dynamics
- **necessary** (as other models) QCD equations not yet solved

- LQCD results: very important, not yet systematic

- evolution of CQMs: benefits from QCD aspects

quark models **before QCD**

First formulation of Quark Model (Gell-Mann, Zweig 1964)

- realizations of SU(3) symmetry
- internal indexes useful for quantum numbers

Constituent Quark (Morpurgo 1965)

- hadrons as bound states of quarks
- binding from a potential
- high quark mass (up to 10 GeV)

Non Relativistic model

Excited baryon wave function with h.o.
symmetric model working well (Horgan-Dalitz 1973)

Birth of QCD

Quarks have been observed

Bjorken scaling:

Proton composed of **pointlike** objects

spin 1/2 (partons)

confined

asymptotically free

QCD

Colour **SU(3) gauge theory**
quark and gluons
current quarks
massless

≠ Constituent
Quarks

Some immediate benefits for CQM

$$\Psi_{3q} = \theta_{\text{colour}} \times \chi_{\text{spin}} \times \phi_{\text{fl}} \times \psi_{\text{space}}$$

$\text{SU}(3)_c$ $\text{SU}(2)$ $\text{SU}(3)_f$ $\text{O}(3)$

θ_{colour} is antisymmetric

the rest is SYMMETRIC

Confinement

h.o. is acceptable

but non realistic

spectrum

degeneracy of levels

Q^2 behavior of ff

strongly damped

Proton (and baryon) structure

quarks ,gluons, quark-antiquark pairs
the interaction increases at low Q^2

Two types of quarks

valence quarks

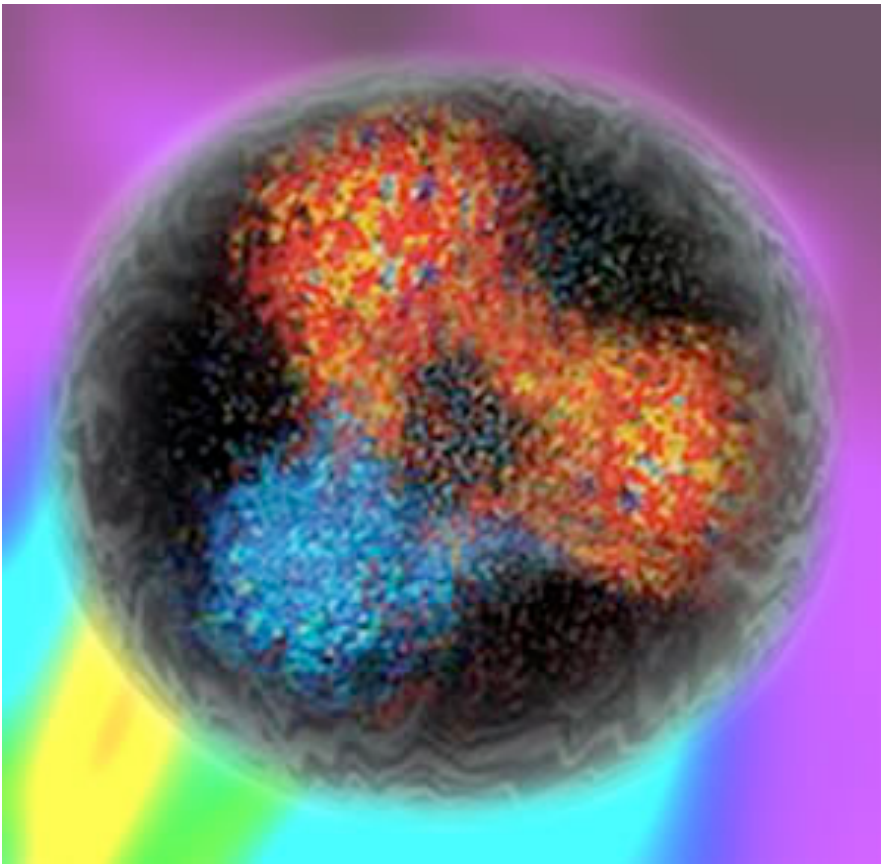
(determine quantum numbers)

sea quarks

High Q^2 regime

Possible picture for proton structure

**three valence quarks,
each surrounded by quark-antiquark pairs and gluons
plus ...**



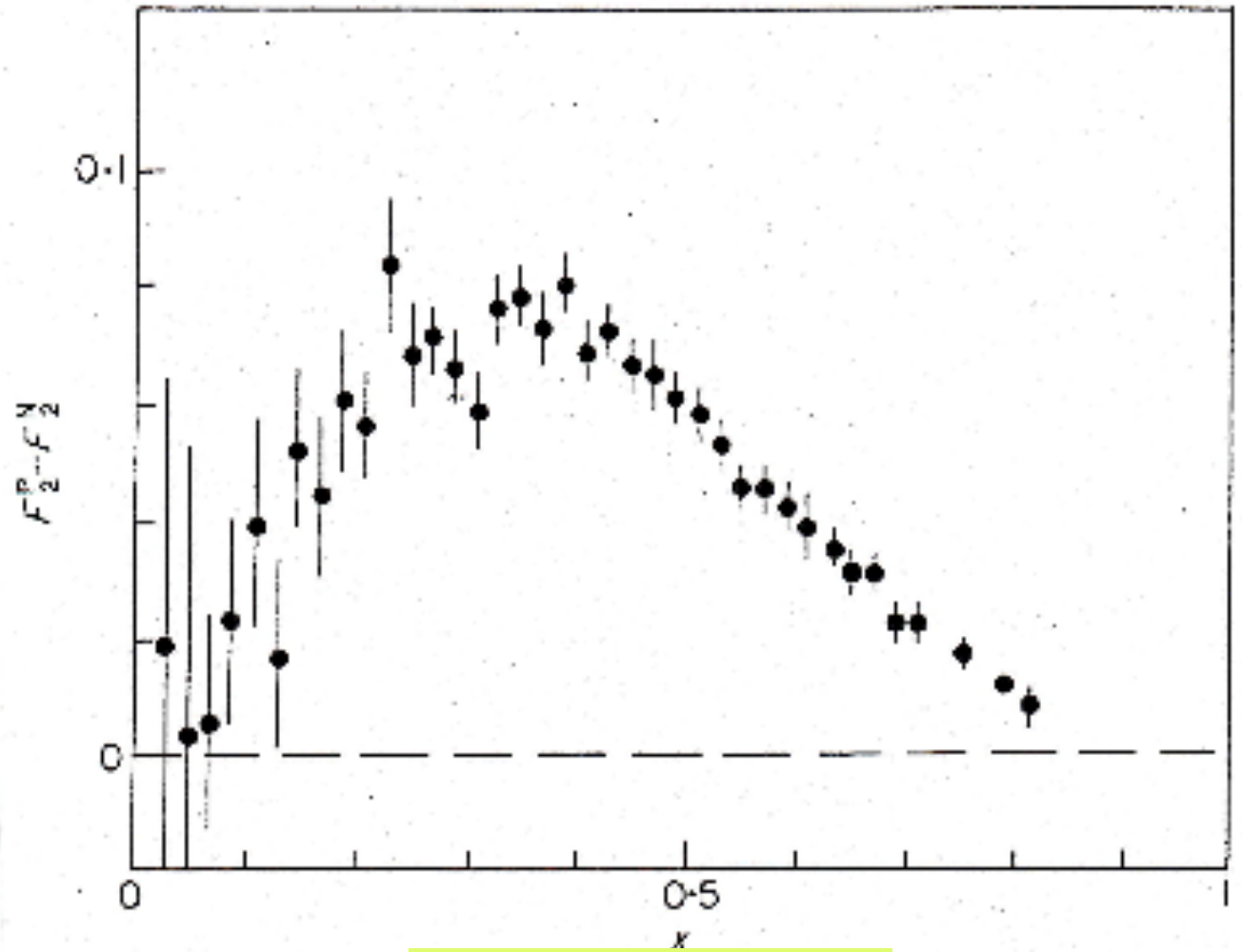
Constituent Quarks

At variance with QCD quarks:

CQ acquire mass &
size
carrier of the proton spin

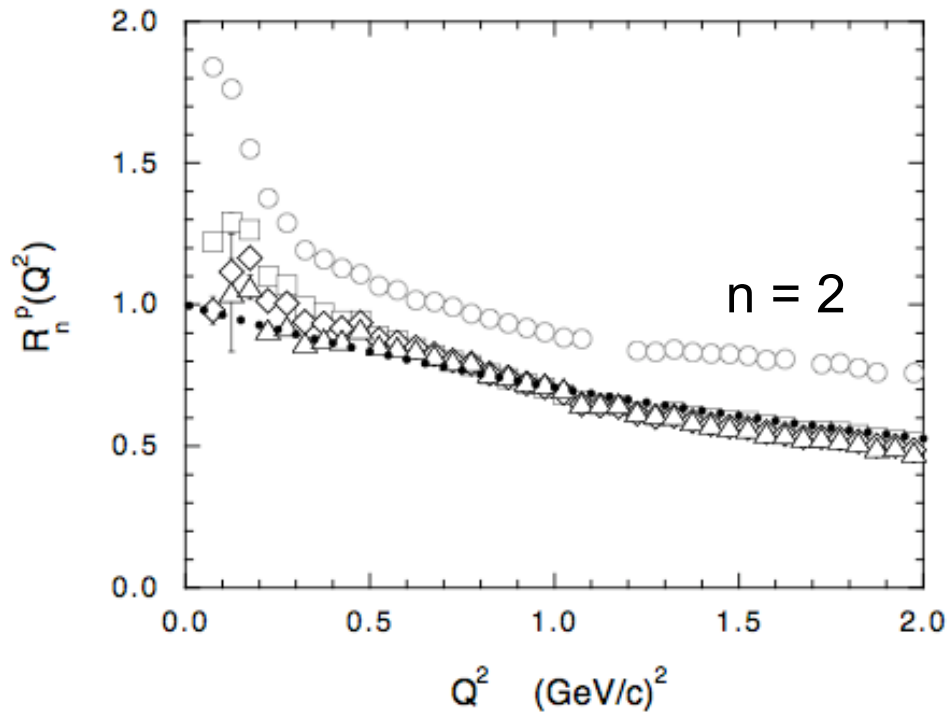
Medium-low Q^2

Further support 1



Peak at $x \gg 1/3$

Further support 2



Ratio between
proton Nachtmann moments &
CQ distribution

Bloom-Gilman duality

Inelastic proton scattering as elastic scattering on CQ

(approximate) scaling function \longrightarrow square of CQ ff

$$F(Q^2) = 1/(1 + 1/6 r_{\text{CQ}}^2 Q^2)$$

$$r_{\text{CQ}} \cong 0.2 \text{ fm}$$

Ricco et al., PR **D67**, 094004 (2003)

Emerging picture

At medium-low Q^2 CQ are suitable internal degrees of freedom they have mass, size and account for all quantum numbers

Including SPIN

spin crisis?

**Consistent inclusion of quark-antiquark pairs
in the Unquenched quark model**

	val	sea	total
$\Delta\Sigma$	0.378	0.298	0.676
ΔL	0.000	0.324	0.324
$2J$	0.378	0.622	1.000

main contribution coming
from quark-antiquark pairs
and their orbital angular
momentum

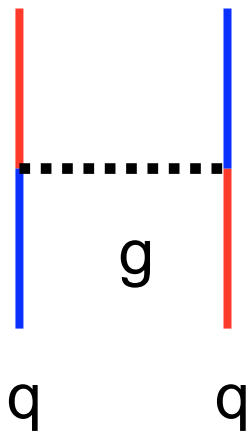
Bijker-Santopinto
Phys.Rev. C80 (2009) 065210

Introducing dynamics

LQCD (De Rújula, Georgi, Glashow, 1975)

the quark interaction contains

- a long range **spin-independent** confinement
→ SU(6) configurations
- a short range **spin dependent** term



One Gluon Exchange

$$V_{\text{OGE}} = -a/r + \text{Hyperfine interaction}$$

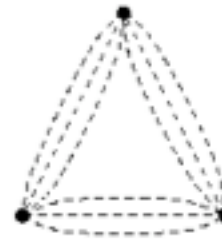
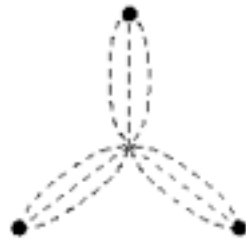
- QCD fundamental mechanism

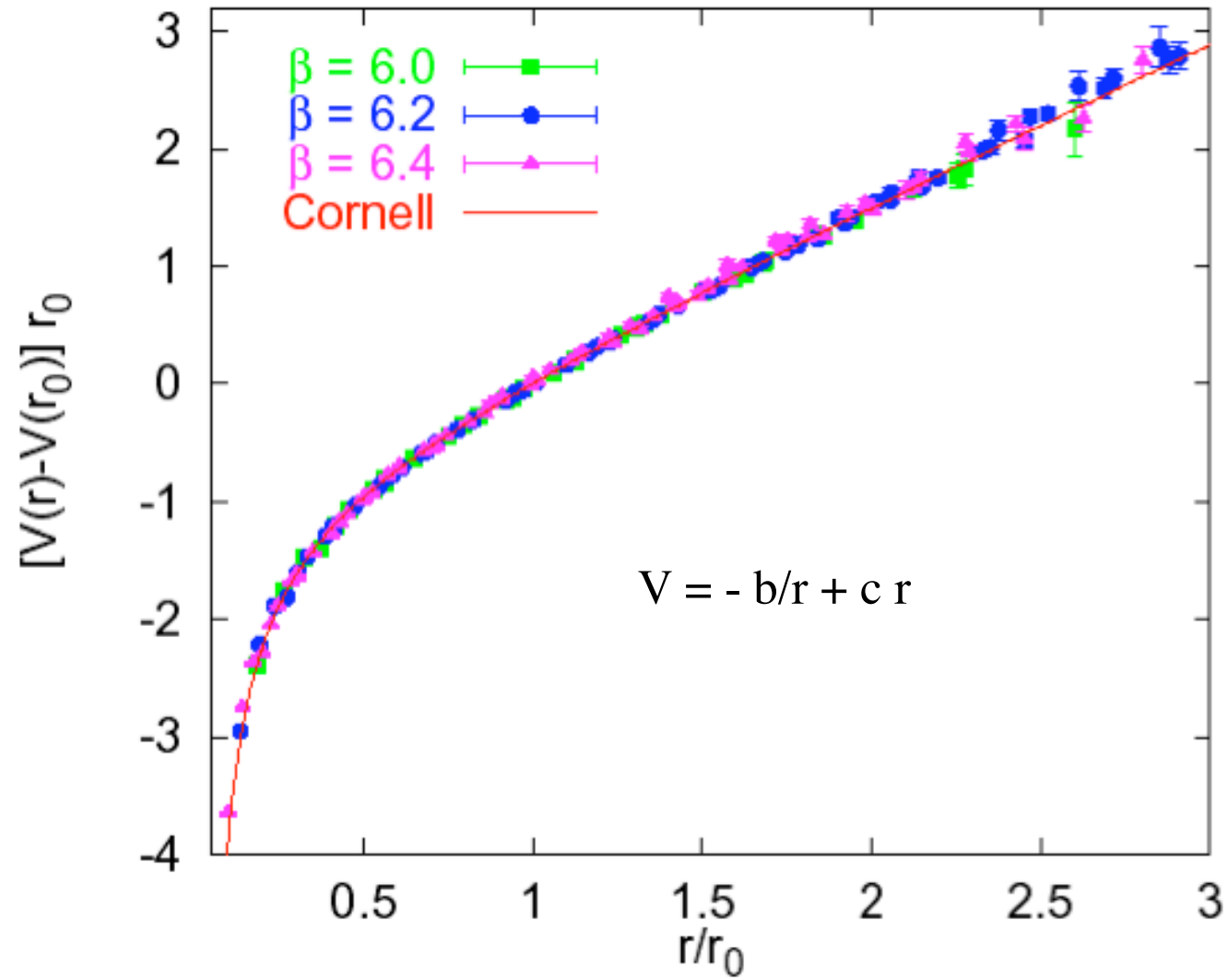


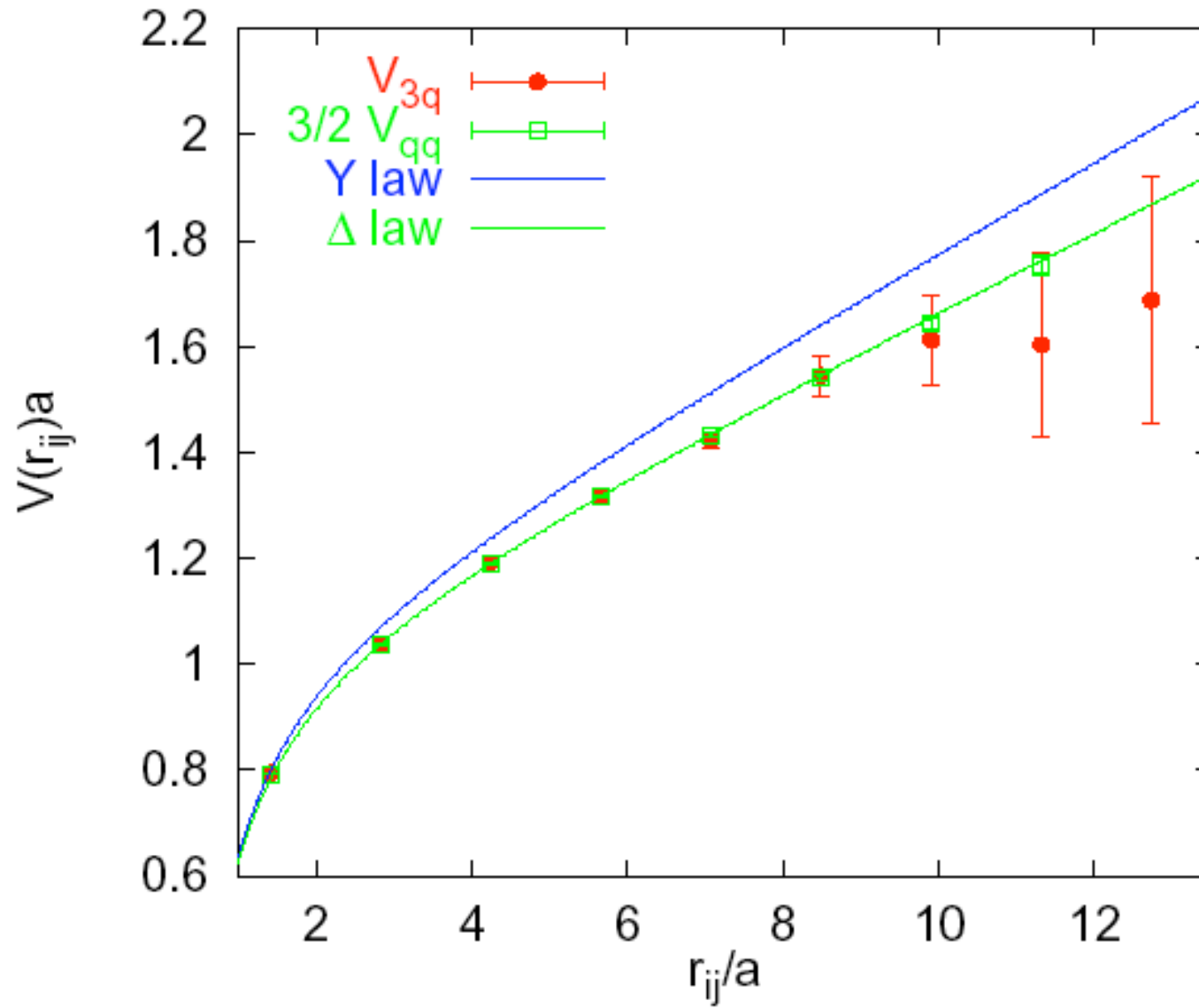
3-body forces

Carlson et al, 1983
Capstick-Isgur 1986
hCQM 1995

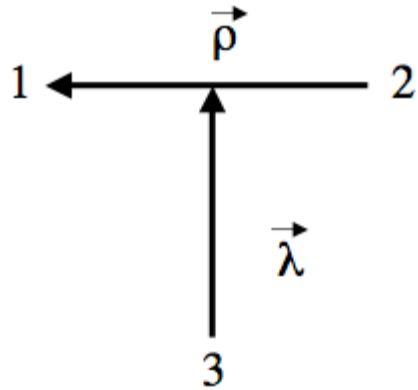
- Flux tube model







Jacobi coordinates



$$L^2(\Omega)Y_{[\gamma]}(\Omega) = -\gamma(\gamma + 4)Y_{[\gamma]}(\Omega)$$

γ grand angular quantum number

$$\sum_{i < j} V(\mathbf{r}_{ij}) \approx V(\mathbf{x}) + \dots$$

Hyperspherical Coordinates

$$(\rho, \Omega_\rho, \lambda, \Omega_\lambda) \Rightarrow (\mathbf{x}, t, \Omega_\rho, \Omega_\lambda)$$

$$x = \sqrt{\rho^2 + \lambda^2}$$

$$t = \text{arctg} \frac{\rho}{\lambda}$$

$$L^2(\Omega) \Leftrightarrow C_2(O(6))$$

$Y_{[\gamma]}(\Omega)$

Hyperspherical harmonics

Hasenfratz et al. 1980:
 $\Sigma V(\mathbf{r}_i, \mathbf{r}_j)$ is approximately hypercentral

Hypercentral Hypothesis

$$V = V(x)$$

Factorization

$$\psi(x, t, \Omega_\rho, \Omega_\lambda) = \underbrace{\psi_{\nu\gamma}(x)}_{\text{("dynamics")}} \underbrace{Y_{[\gamma, l_\rho, l_\lambda]}}_{\text{("geometry")}}$$

Only one differential equation in x (hyperradial equation)

Hypercentral Model

Genoa group, 1995

$$V(x) = -\tau/x + \alpha x$$

Hypercentral approximation of

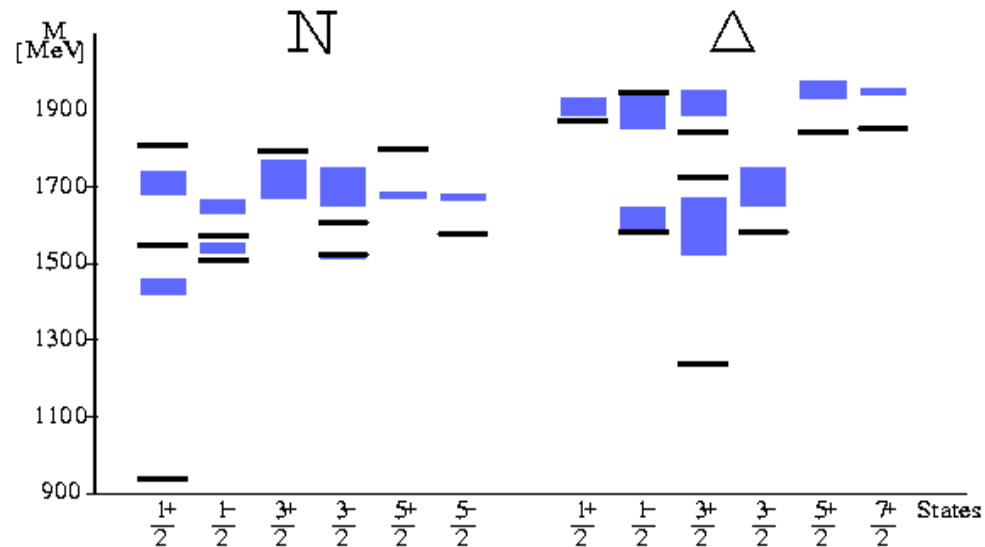
$$V = -b/r + c r$$

Hypercentral Model (1)

$$H_{3q} = 3m + \sum_{i=1}^3 \frac{\mathbf{p}_i^2}{2m} + V(\mathbf{x}) + H_{hyp}$$

M. Ferraris, M. M. Giannini, M. Pizzo, E. Santopinto, L. Tiator, Phys. Lett. B364 (1995), 231

- $V(\mathbf{x}) = -\frac{\tau}{x} + \alpha x$; $H_{hyp} = A \left[\sum_{i < j} V^S(\mathbf{r}_i, \mathbf{r}_j) \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j + \text{tensor} \right]$
- 3 parameters τ α $A \leftarrow$ fixed to the spectrum, $m = \frac{M}{3}$



$$\tau = 4.59$$

$$\alpha = 1.61 \text{ fm}^{-1}$$

$$A \leftarrow (N - \Delta)$$

$$x = \sqrt{\rho^2 + \lambda^2}$$

hyperradius

Results (predictions)
with the Hypercentral Constituent
Quark Model

for

- Helicity amplitudes
- Elastic nucleon form factors

HELICITY AMPLITUDES

Definition

$$A_{1/2} = \langle N^* J_z = 1/2 | H_{em}^T | N J_z = -1/2 \rangle^* \xi \quad \S$$

$$A_{3/2} = \langle N^* J_z = 3/2 | H_{em}^T | N J_z = 1/2 \rangle^* \xi \quad \S$$

$$S_{1/2} = \langle N^* J_z = 1/2 | H_{em}^L | N J_z = 1/2 \rangle^* \xi$$

N, N^* nucleon and resonance as 3q states

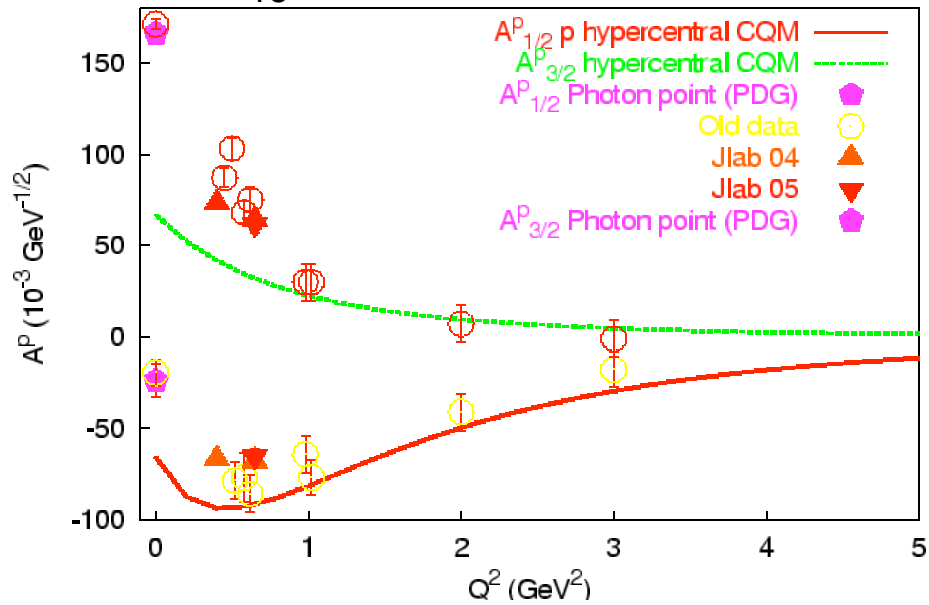
H_{em}^T, H_{em}^L model transition operator

ξ overall sign \rightarrow problem

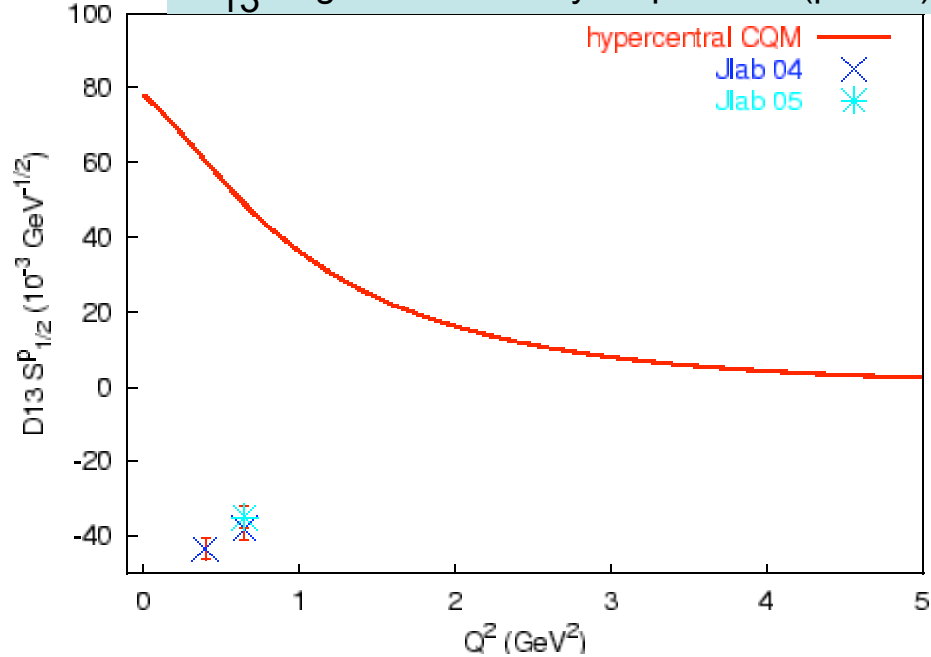
\S results for the negative parity resonances:

M. Aiello et al. J. Phys. G24, 753 (1998)

D_{13} transverse helicity amplitudes (proton)

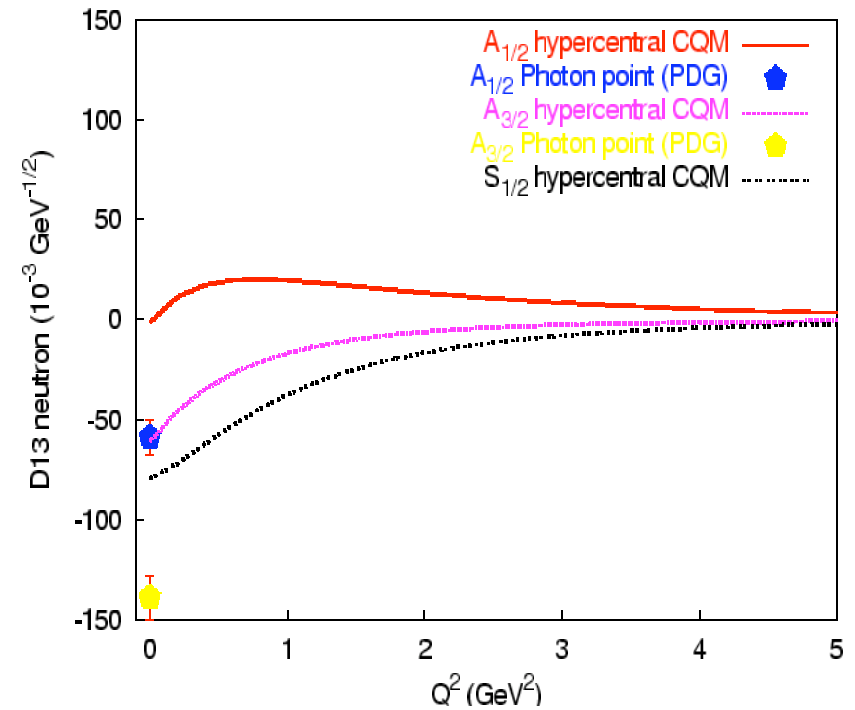


D_{13} longitudinal helicity amplitudes (proton)

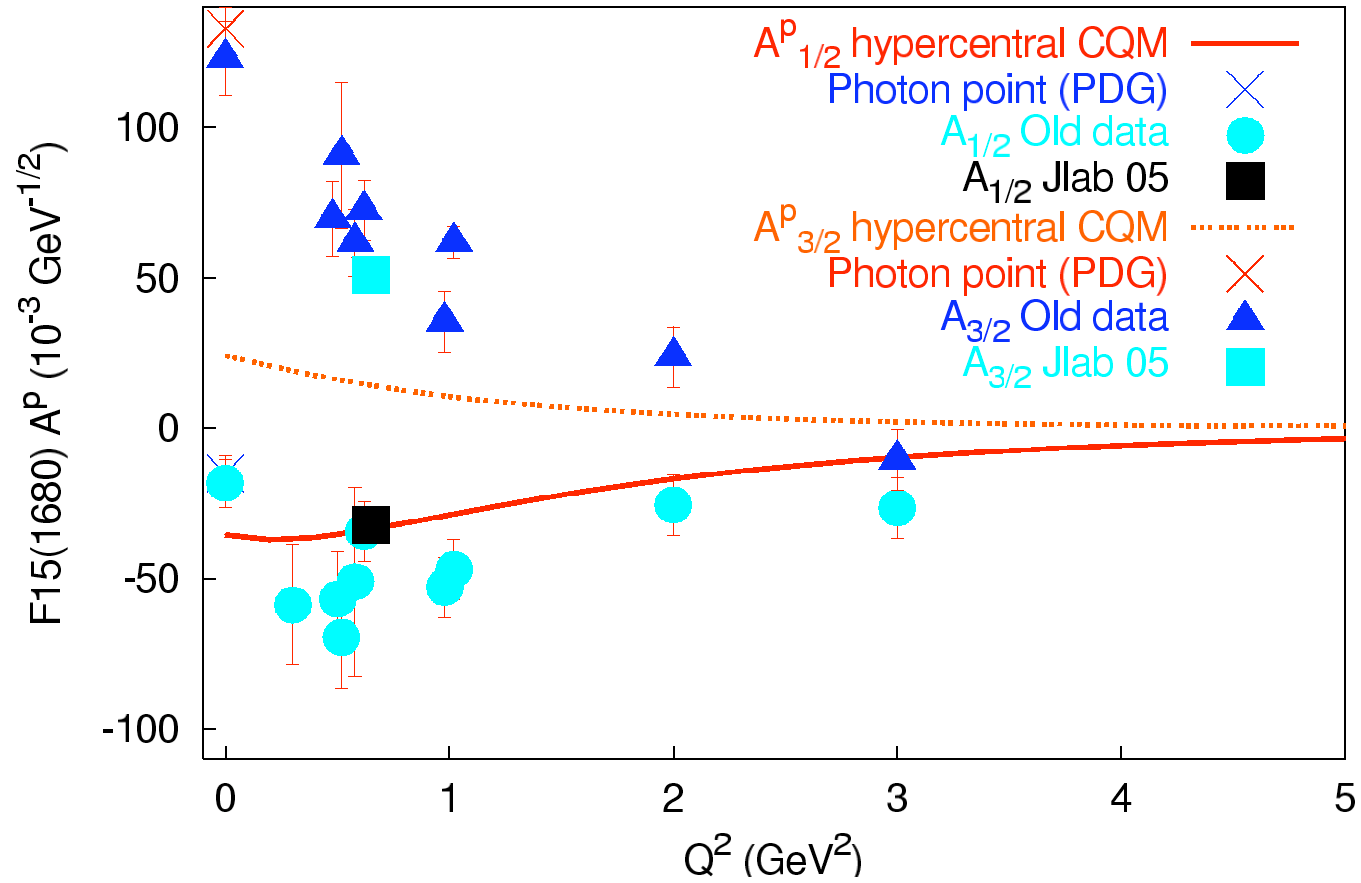


D_{13}

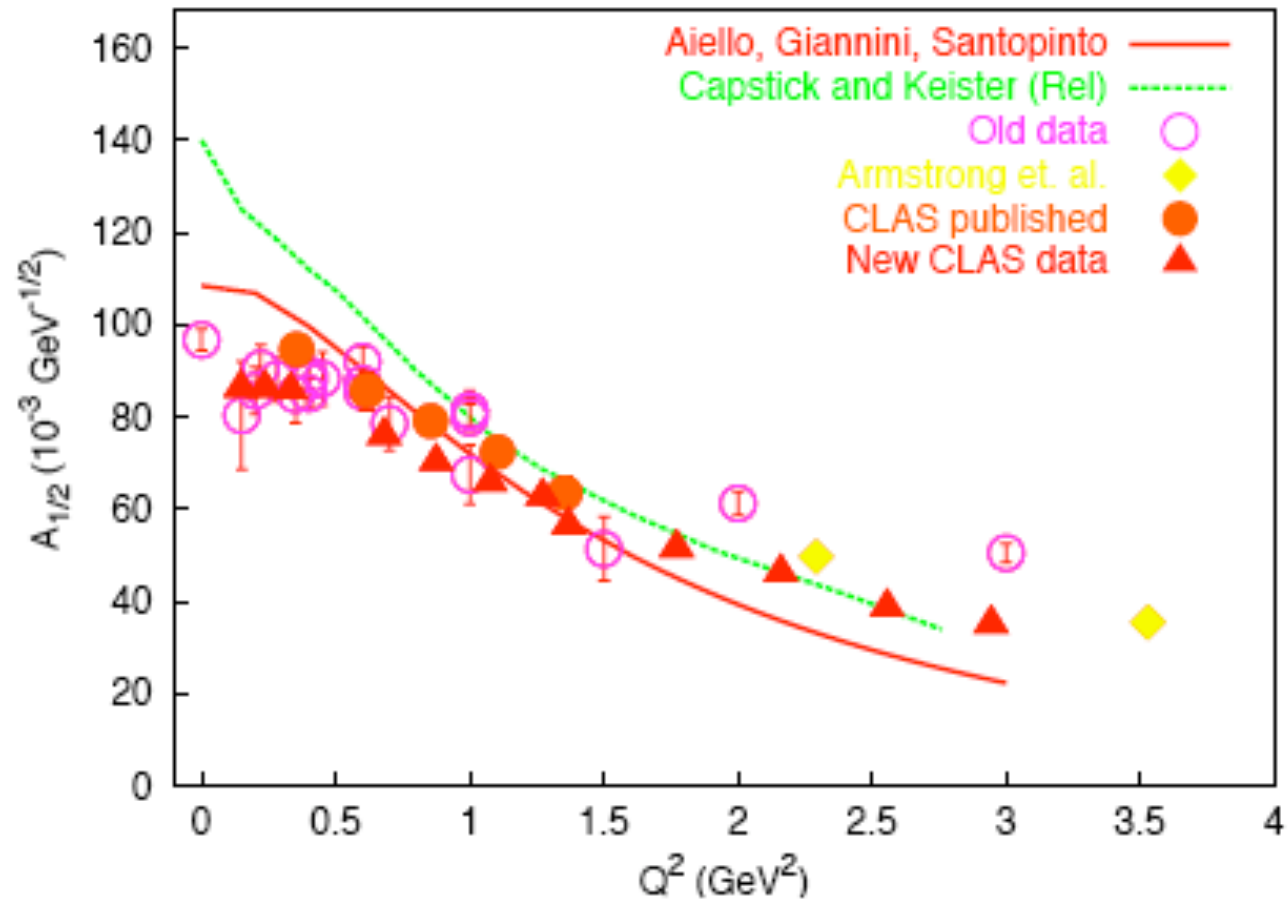
D_{13} helicity amplitudes (neutron)



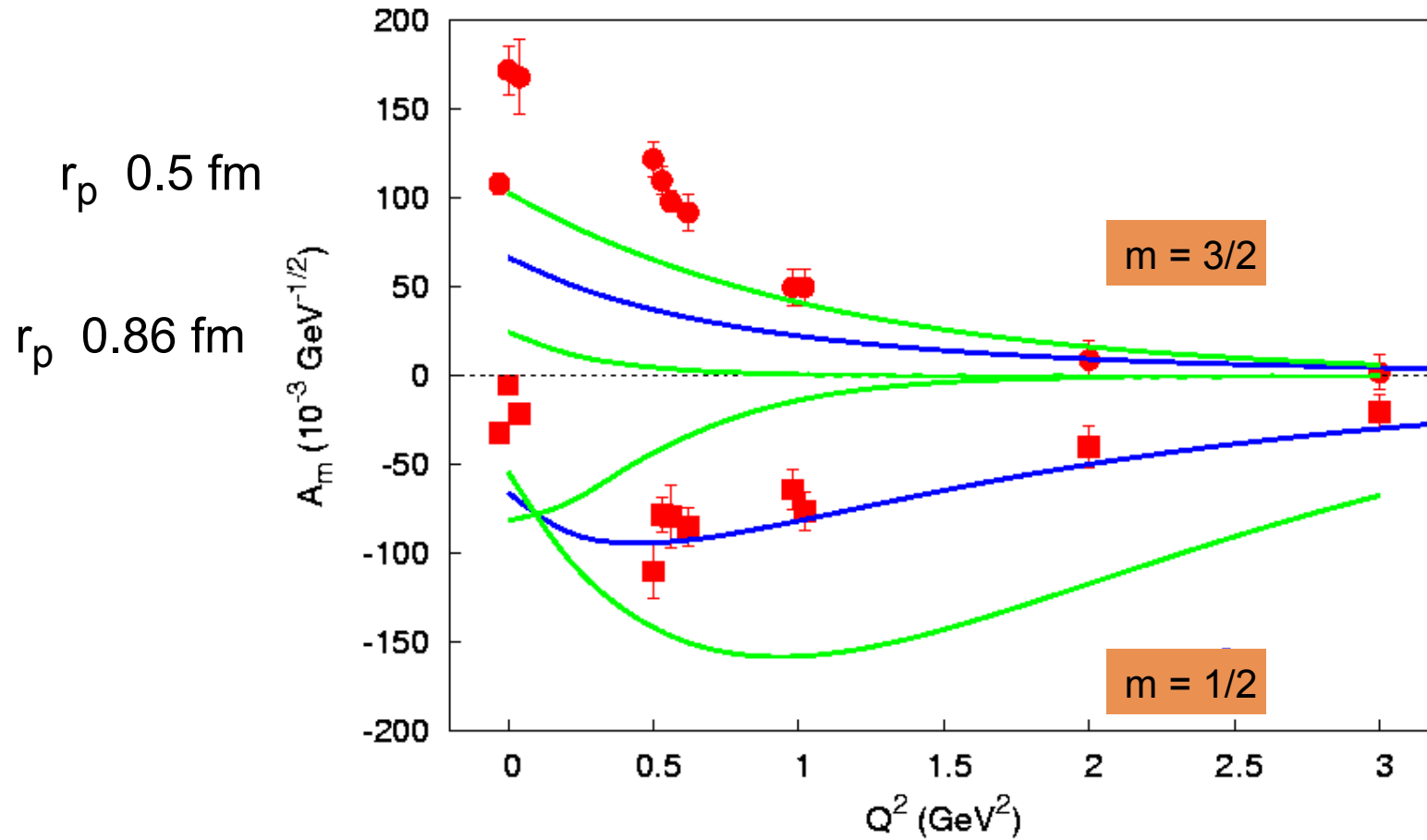
F15 transverse helicity amplitudes



S11(1535)



$A_m^P N(1520) D13$



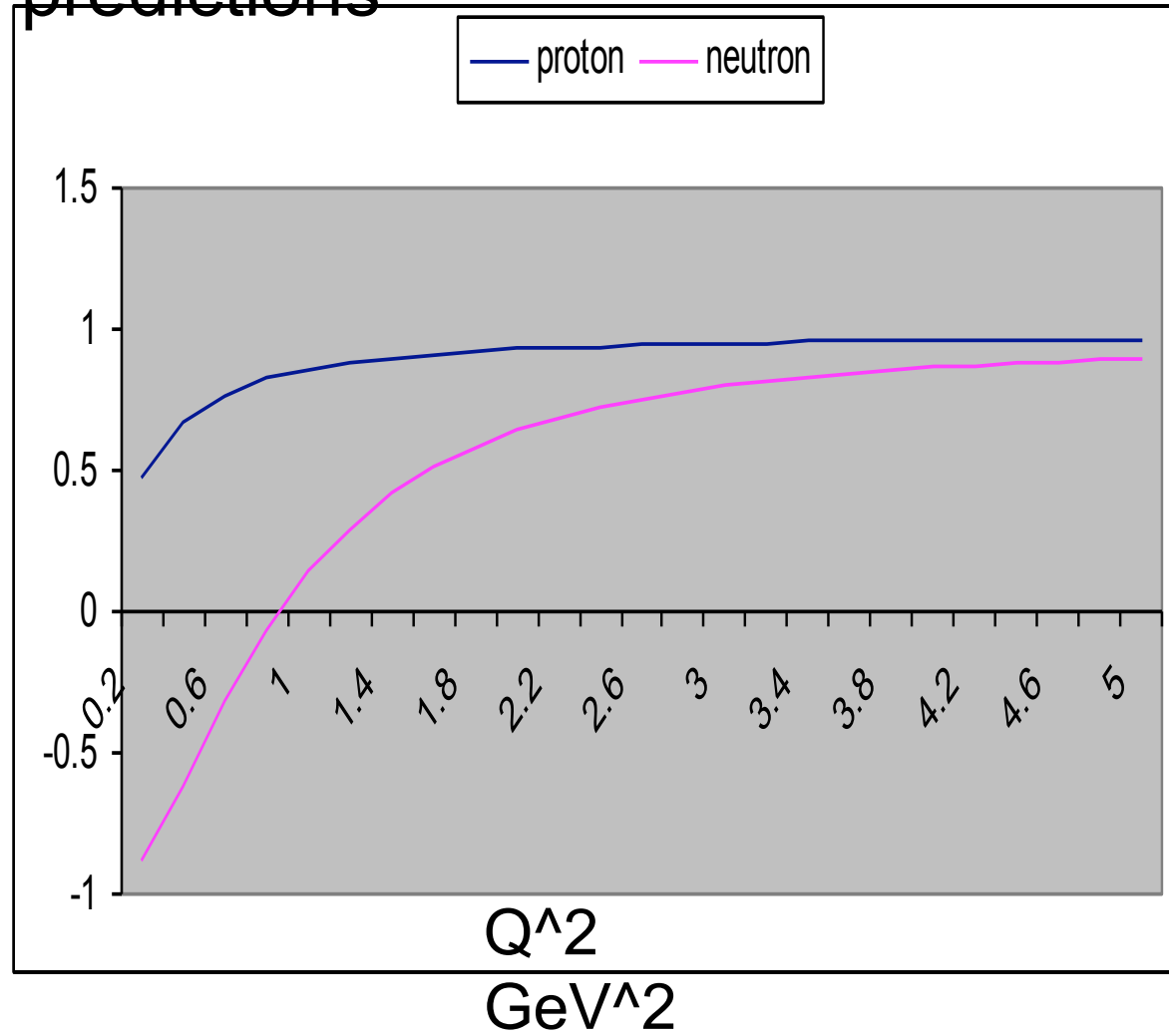
Green curves H.O.

Blue curves hCQM

D13

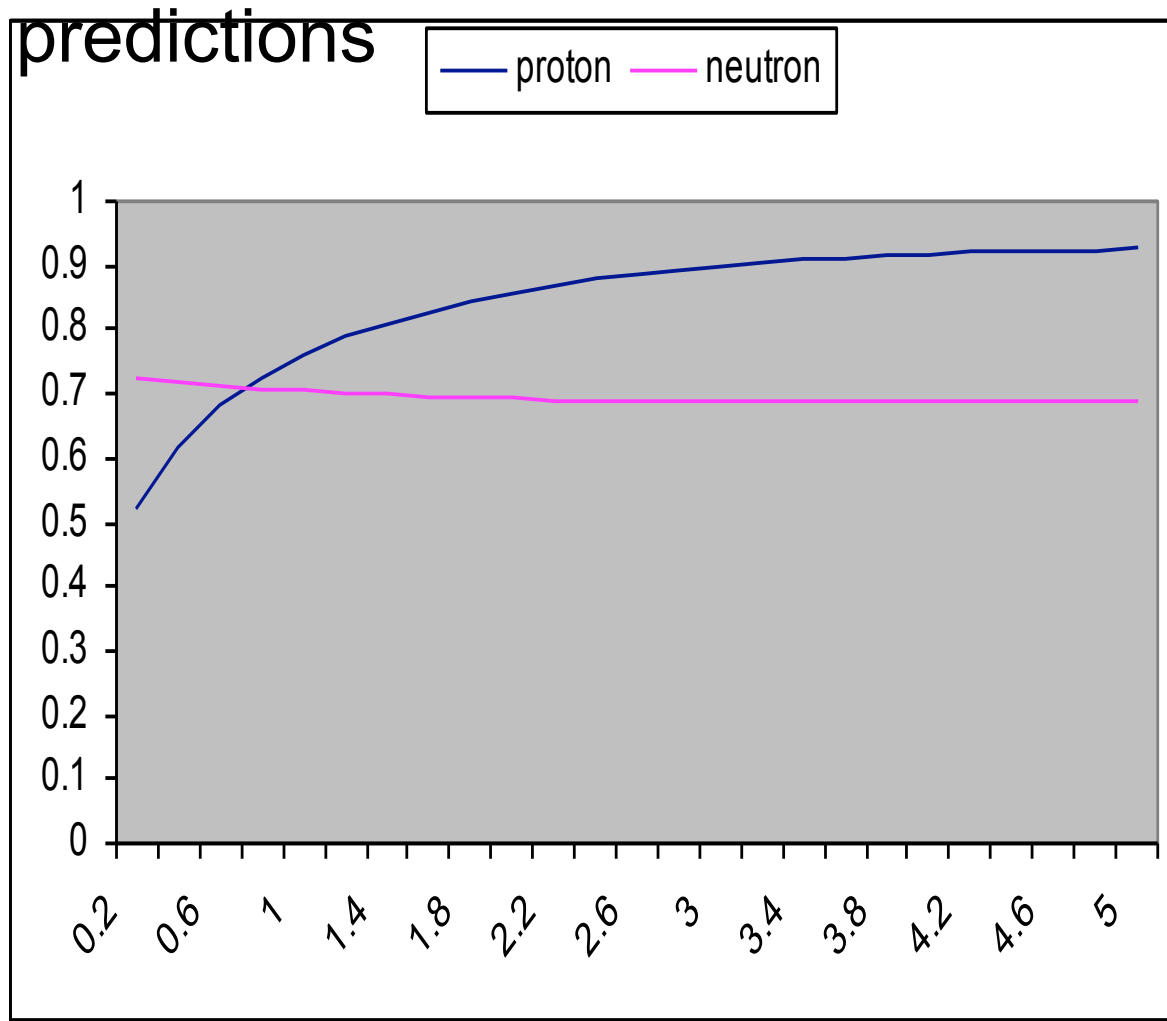
hCQM

predictions



F15

hCQM



Q^2
GeV²

Helicity ratio

	proton	neutron
P33	≈ -0.5	
D13	ok	ok
F15	ok	≈ 0.7
D13*	ok	0.96
D33	ok	
D15	1/3	≈ 0.32
F35	-0.82	
F37	-0.32	
P13	ok	ok

Structure effects ?

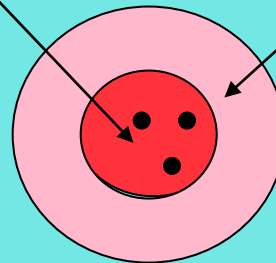
observations

- the **calculated** proton radius is about **0.5 fm**
(value previously obtained by fitting the helicity amplitudes)
- the medium Q^2 behaviour is fairly well reproduced (**$1/x$ potential**)
- there is lack of strength at **low** Q^2 (outer region) in the e.m. transitions

specially for the $A_{3/2}$ amplitudes

- emerging picture: quark core (**0.5 fm**) plus (meson or sea-quark) **cloud**

Quark-antiquark pairs effects are important for the low Q^2 behavior



Elastic form factors

The non relativistic calculation gives values higher than data since the calculated radius is 0.5 fm

a relativistic calculation is needed

Performed by variuos groups

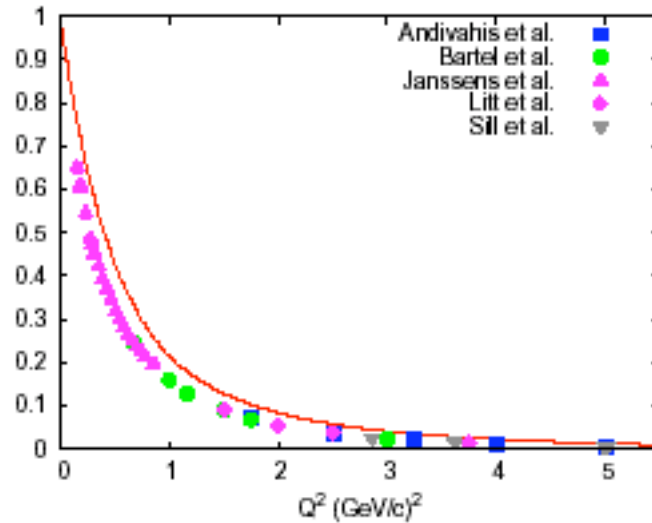
Rome,	Pavia-Graz,	Bonn,	Genova
LF	PF	BS	PF

Calculated values!

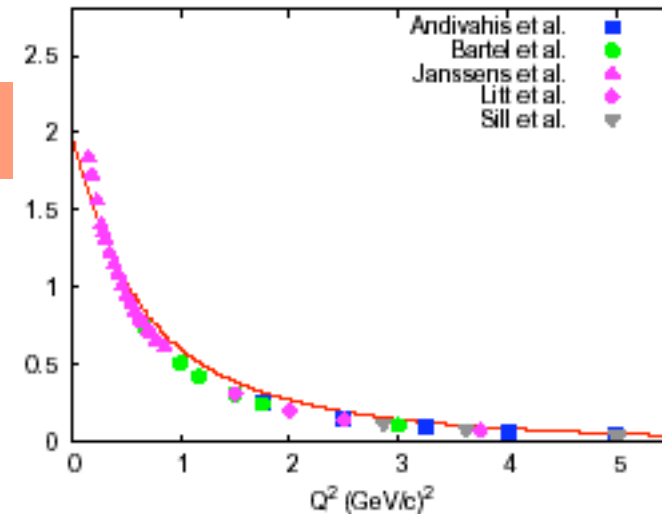
- Boosts to initial and final states
- Expansion of current to any order
- Conserved current

Relativistic theory
in point form

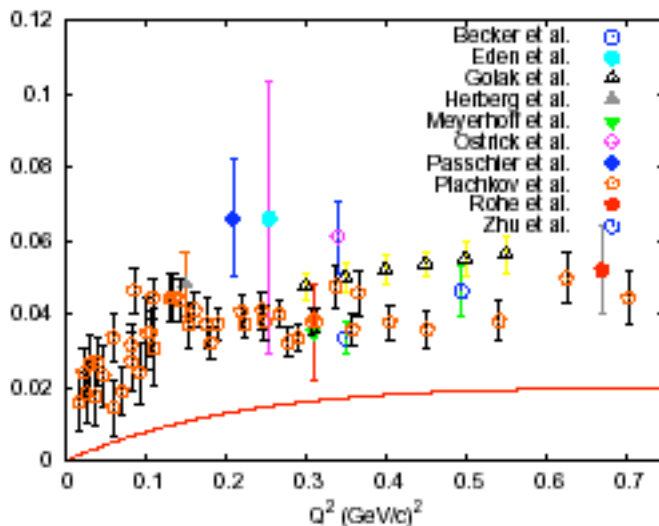
G_E^p



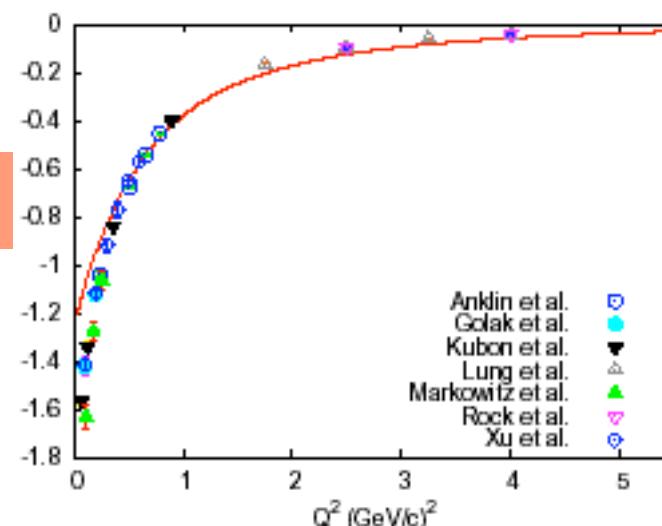
G_M^p

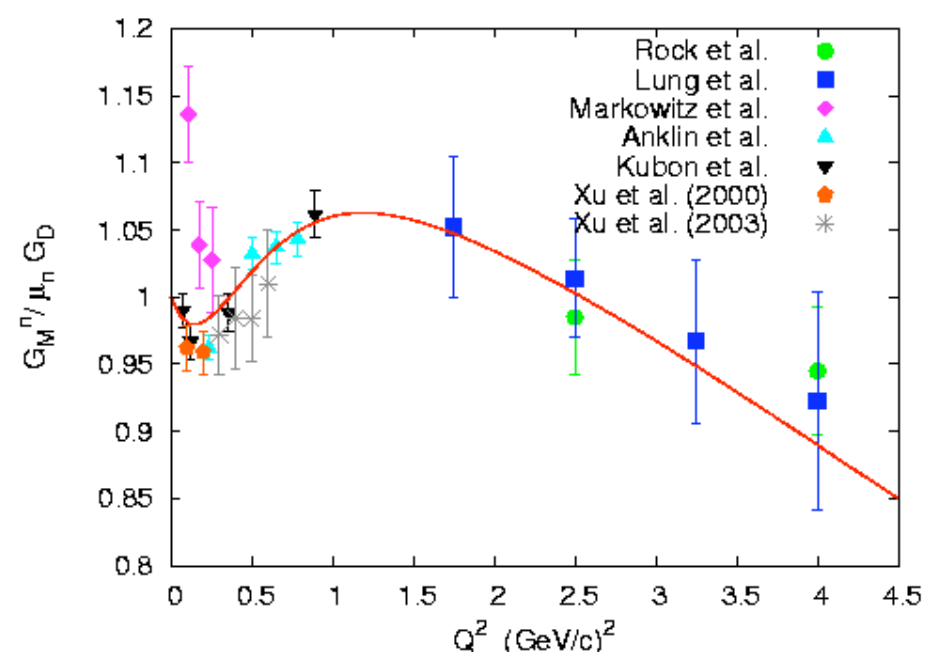
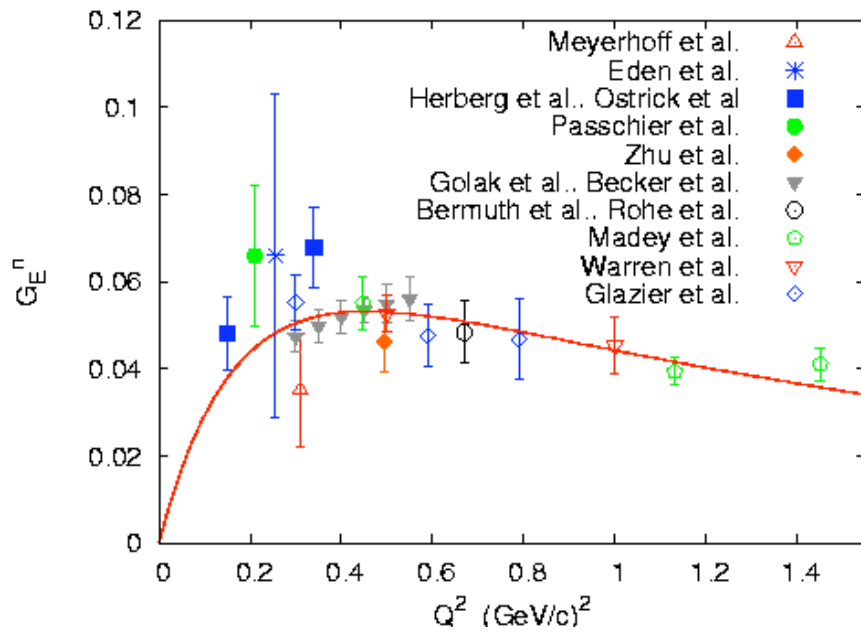
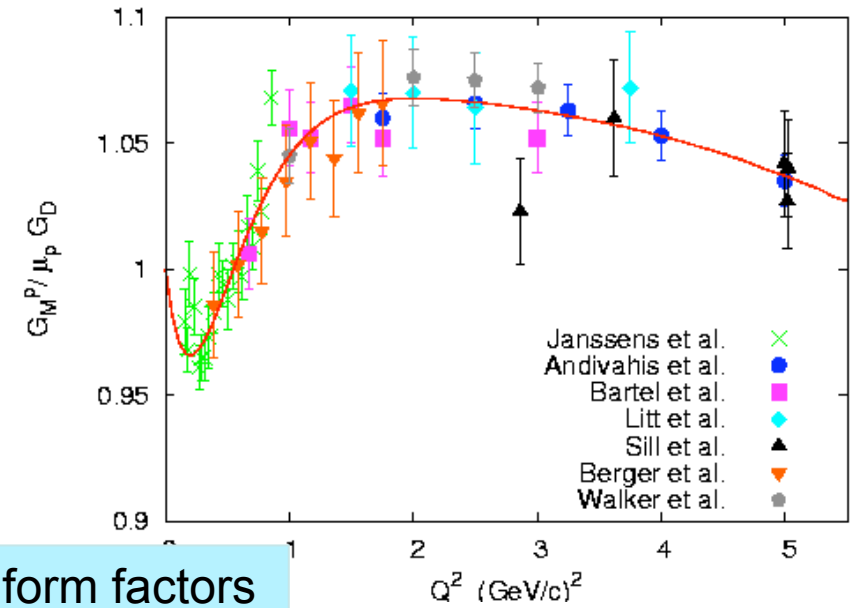
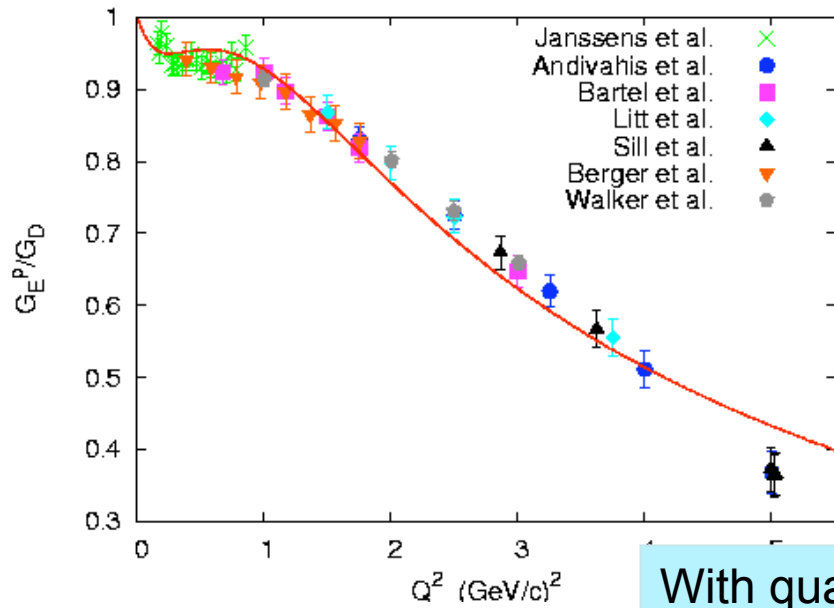


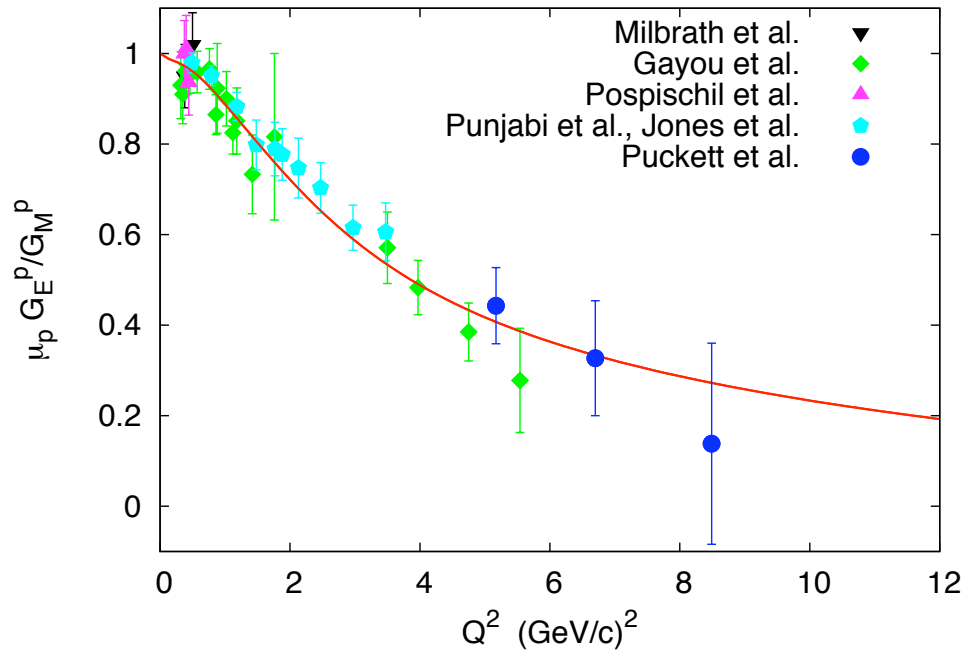
G_E^n



G_M^n

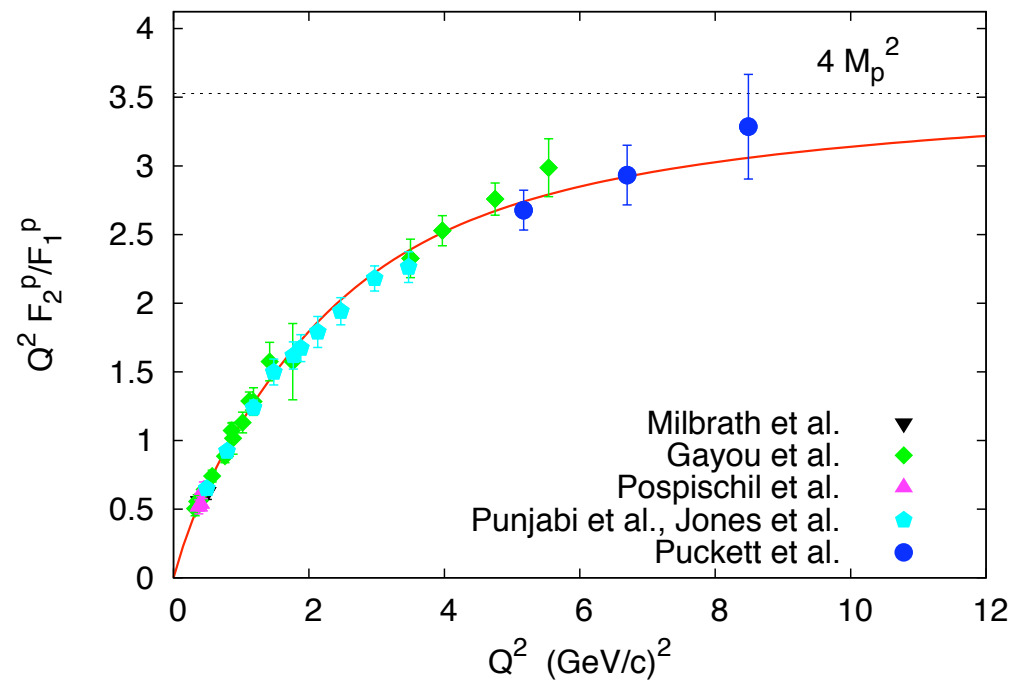






With quark form factors

Santopinto et al.
PR C **82**, 065204 (2010)



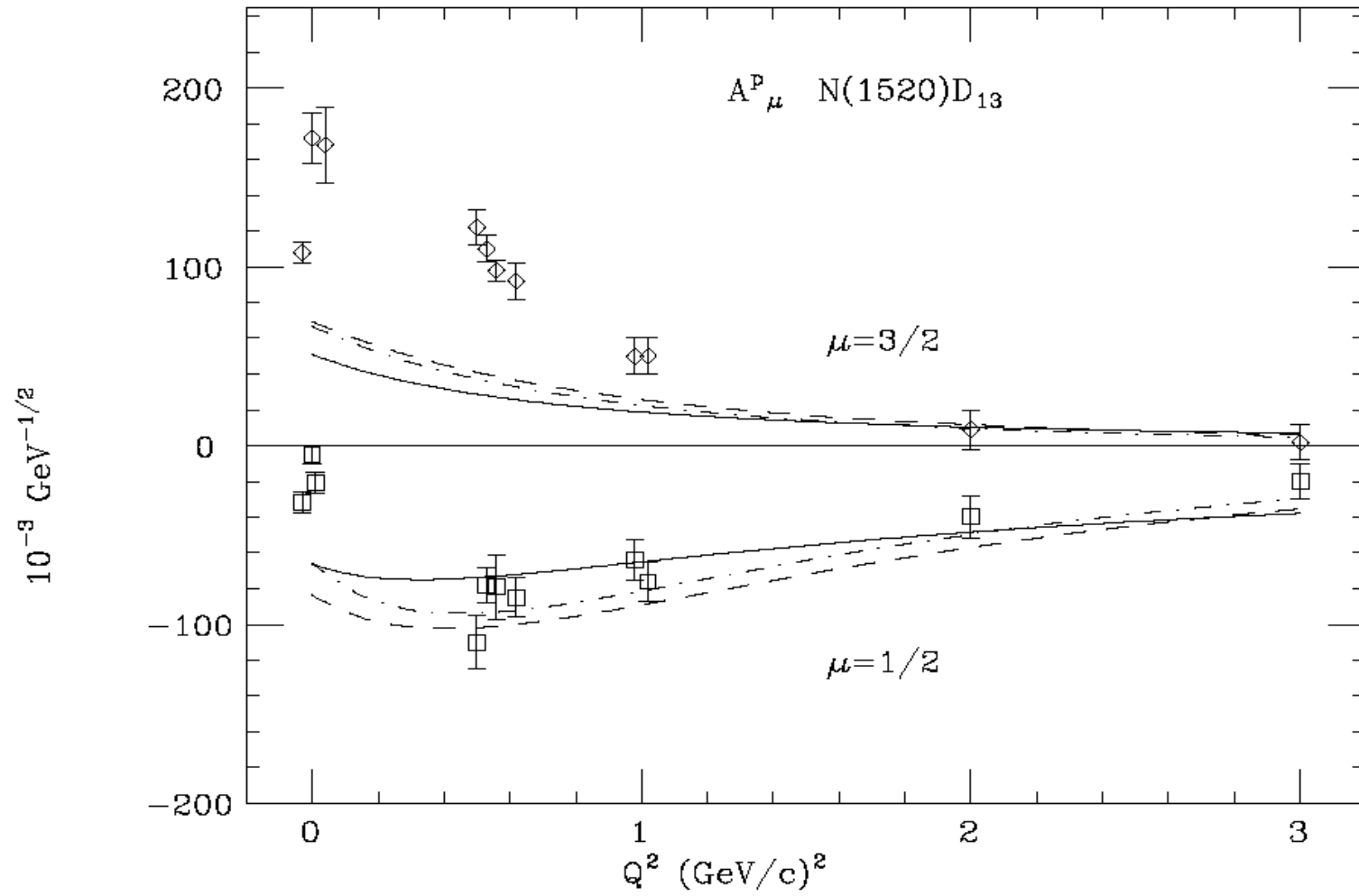
Relativistic corrections to form factors

- Breit frame
- Lorentz boosts applied to the initial and final state
- Expansion of current matrix elements up to first order in quark momentum
- Results

$$A_{\text{rel}}(Q^2) = F A_{\text{n.rel}}(Q_{\text{eff}}^2)$$

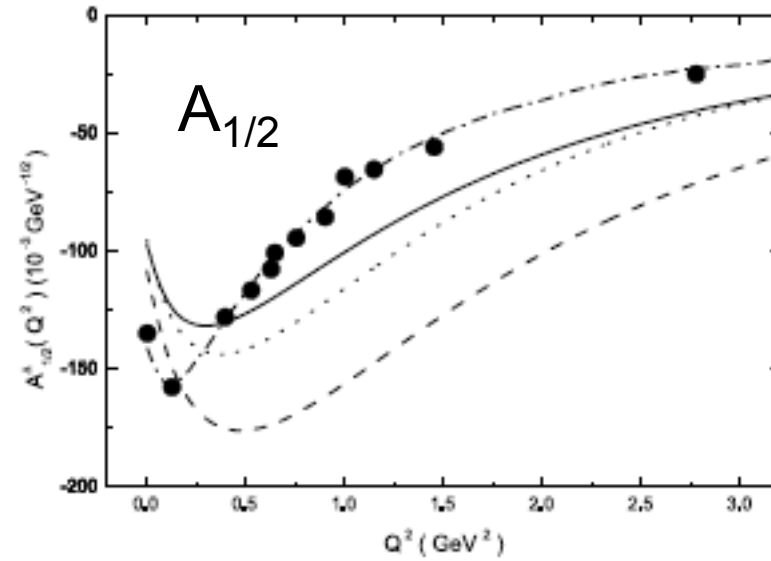
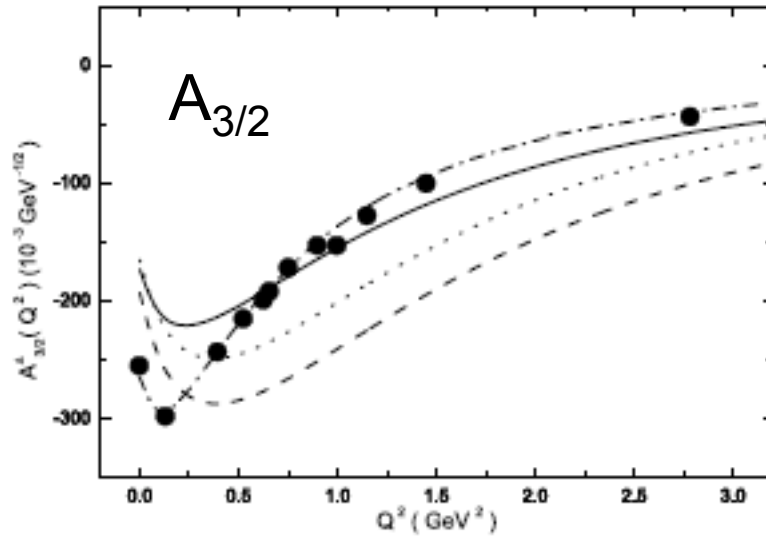
$$F = \text{kin factor} \quad Q_{\text{eff}}^2 = Q^2 (M_N/E_N)^2$$

De Sanctis et al. EPJ 1998



Full curves: hCQM with relativistic corrections
 Dashed curves: hCQM in different frames

Chen, Dong, M.G., Santopinto, Trieste 2006



dot bare
dash dressed
full rel. corr (preliminary calculation)
dash-dot MAID

Relativistic treatment

- elastic form factors: necessary
- helicity amplitudes: probably necessary
exciting higher resonances the
recoil is smaller
- Delta excitation: g.s. in the SU(6) limit
probably more important

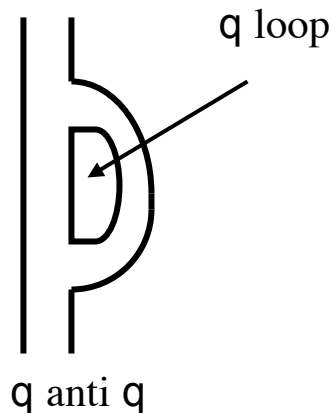
Relativity is an important issue for the description of
elastic and inelastic form factors

but it is not the only important issue

Unquenching the quark model

Mesons

P. Geiger, N. Isgur, Phys. Rev. D41, 1595 (1990)
D44, 799 (1991)

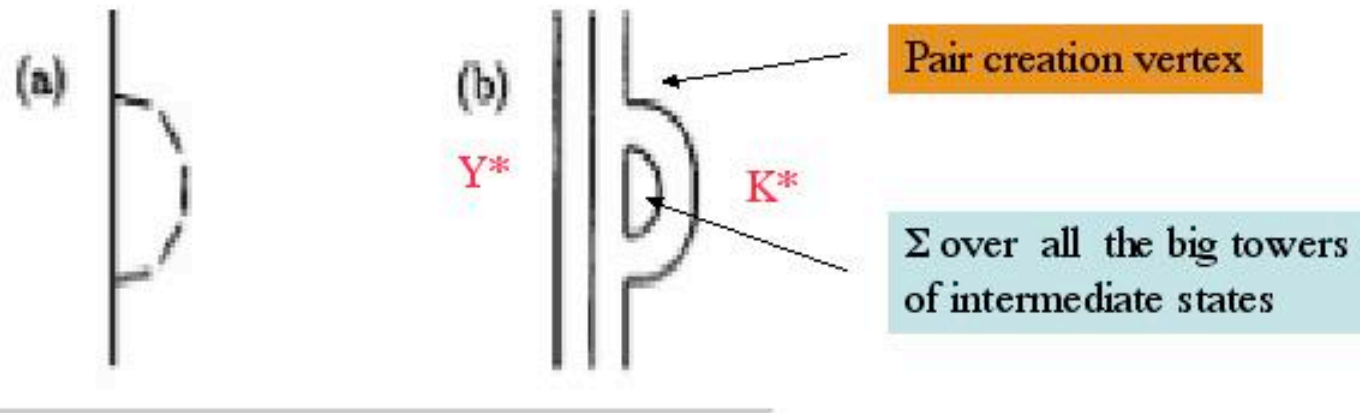


Note:

- sum over all intermediate states
necessary for OZI rule
- linear interaction is preserved
renormalization of the string constant

baryons

The qq-pair creation mechanism is introduced at the microscopical level
→ string-like qq pair creation mechanism



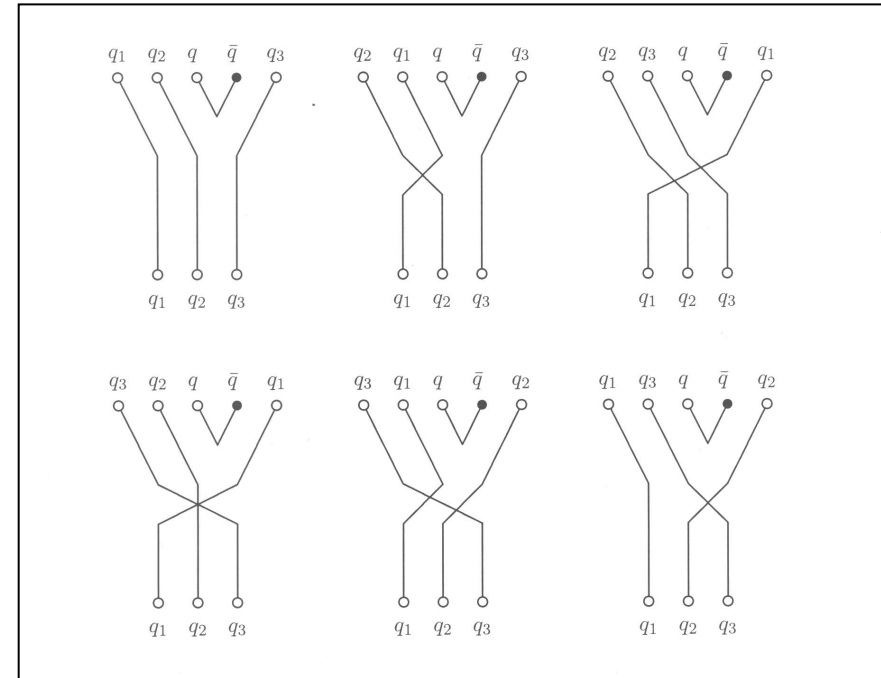
R. Bijker, E. Santopinto,
Phys.Rev.C80:065210,2009

Problems that have been solved for baryons:

- sum over the big tower of intermediate states
- permutational symmetry

(both with group theoretical methods)

- find a quark QCD inspired pair creation mechanism 3P_0
- implementation of the mechanism in such a way to **do not destroy** the good CQM's results



The good magnetic moment results of the CQM are preserved by the UCQM

Bijker, Santopinto, Phys.Rev.C80:065210,2009.

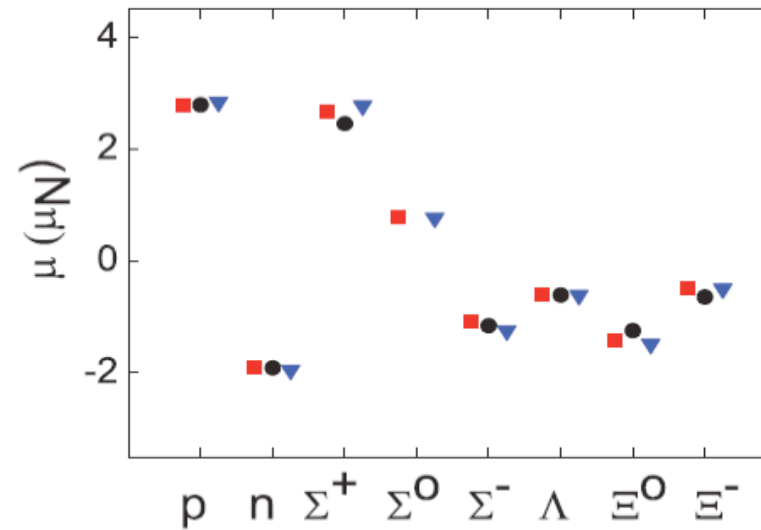
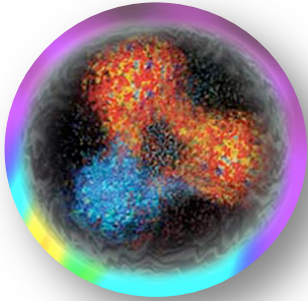


FIG. 3. (Color online) Magnetic moments of octet baryons: experimental values from the Particle Data Group [34] (circles), CQM (squares), and unquenched quark model (triangles).

Possible structure of the nucleon



3-quark core (about 0.5 fm)
+
quark-antiquark pairs
outside and inside the core

Unquenching the CQM:

effects on spectrum

e.m. excitation

consistent evaluation of electroproduction

conclusions

- QCD has affected CQM formulation at various stages
- Important issues for N* high Q^2 electrocouplings:
 - **Relativity:** mainly for ground state
affects Q^2 behaviour
 - **Quark-antiquark pairs**
to be introduced consistently (unquenching)
modify both low and high Q^2 behaviour
 - **Quark form factors**
introduce a Q^2 behaviour of quite different origin
(connected to **fundamental dynamics?**)