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# Covariant Faddeev calculation of $N\text{-}\Delta(1232)$ form factors

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[Craig Roberts (ANL)]

*Electromagnetic  $N\text{-}N^*$  Transition Form Factors Workshop*

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# Themes

## ❖ Themes

- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

- Demonstrate utility of Nambu–Jona-Lasinio model
  - ❖ Nucleon, Delta,  $N \rightarrow \Delta$  form factors
  - ❖ quark distributions, etc
- Highlight important challenges in modelling form factors
  - ❖ in particular the pion cloud
- Results: Nucleon, Delta form factors
- Preliminary Results:  $N \rightarrow \Delta$  Transition Form Factors
- Examine Off-Shell Form Factors

# Nambu–Jona-Lasinio Model

❖ Themes

❖ NJL model

❖ Baryons . . .

❖  $N \rightarrow \Delta$  FFs

❖ Constituents

❖ Nucleon FFs

❖ Scalar Diquark FF

❖ Axial-Vector FF

❖ Nucleon FFs

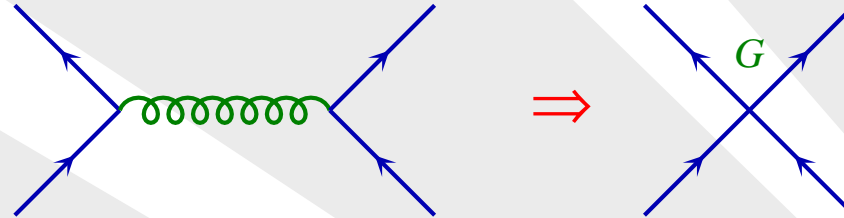
❖ Delta FFs

❖  $N \rightarrow \Delta$  FFs

❖ Off-Shell

❖ Conclusion

- Low energy chiral effective theory of QCD



- Investigate the role of quark degrees of freedom.

- Much in common with DSE

- Lagrangian has same symmetries as QCD:

- ❖ Importantly chiral symmetry and  $D_\chi$ SB,
  - Dynamically generated quark masses,
  - Non-zero chiral condensate.

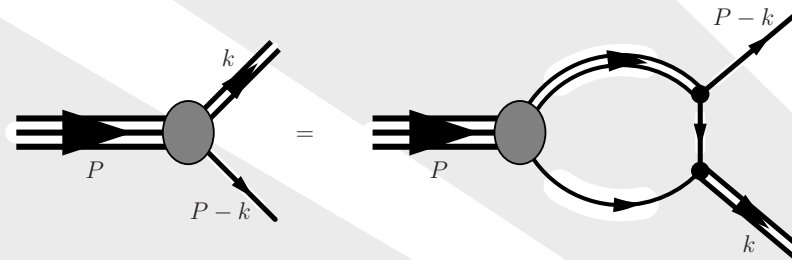
- Lagrangian ( $\Gamma =$  Dirac, colour, isospin matrices)

$$\mathcal{L}_{NJL} = \bar{\psi} (i\cancel{D} - m) \psi + G (\bar{\psi}\Gamma\psi)^2$$

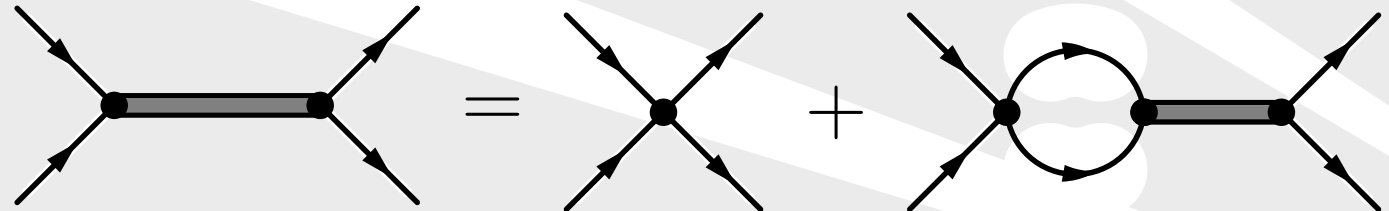
# Baryons in the NJL model

- ❖ Themes
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- ❖ Off-Shell
- ❖ Conclusion

- Baryons approximated as quark-diquark bound states.
- Use relativistic Faddeev approach:



- Diquark - bound state of two quarks:
- Solve Bethe-Salpeter equation for diquark.



- We include scalar and axial-vector diquarks.
- Static Approximation:  $S_{ex}(k) \rightarrow -1/M$ .

# Regularization

- ❖ Themes
- ❖ NJL model
- ❖ **Baryons . . .**
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
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- ❖ Scalar Diquark FF
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## ● Proper-time regularization

$$\frac{1}{X^n} = \frac{1}{(n-1)!} \int_0^\infty d\tau \tau^{n-1} e^{-\tau X}$$
$$\longrightarrow \frac{1}{(n-1)!} \int_{1/(\Lambda_{UV})^2}^{1/(\Lambda_{IR})^2} d\tau \tau^{n-1} e^{-\tau X}.$$

- $\Lambda_{IR}$  eliminates unphysical thresholds for the nucleon to decay into quarks:  $\rightarrow$  simulates confinement.

❖ G. Hellstern, R. Alkofer and H. Reinhardt, Nucl. Phys. A **625**, 697 (1997).

- Needed for: nuclear matter saturation,  $\Delta$  baryon.

❖ W. Bentz, A.W. Thomas, Nucl. Phys. A **696**, 138 (2001)

# Model Parameters

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
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- ❖ Nucleon FFs
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- Free Parameters:

$$\Lambda_{IR}, \Lambda_{UV}, M_0, G_\pi, G_s, G_a, G_\omega \text{ and } G_\rho$$

- Constraints:

- ❖  $f_\pi = 93 \text{ MeV}, m_\pi = 140 \text{ MeV} \quad \& \quad M_N = 940 \text{ MeV}$

- ❖  $\int_0^1 dx (\Delta u_v(x) - \Delta d_v(x)) = g_A = 1.267$

- ❖  $(\rho, E_B/A) = (0.16 \text{ fm}^{-3}, -15.7 \text{ MeV})$

- ❖  $a_4 = 32 \text{ MeV}$

- ❖  $\Lambda_{IR} = 240 \text{ MeV}$

- We obtain [MeV]:

- ❖  $\Lambda_{UV} = 644$

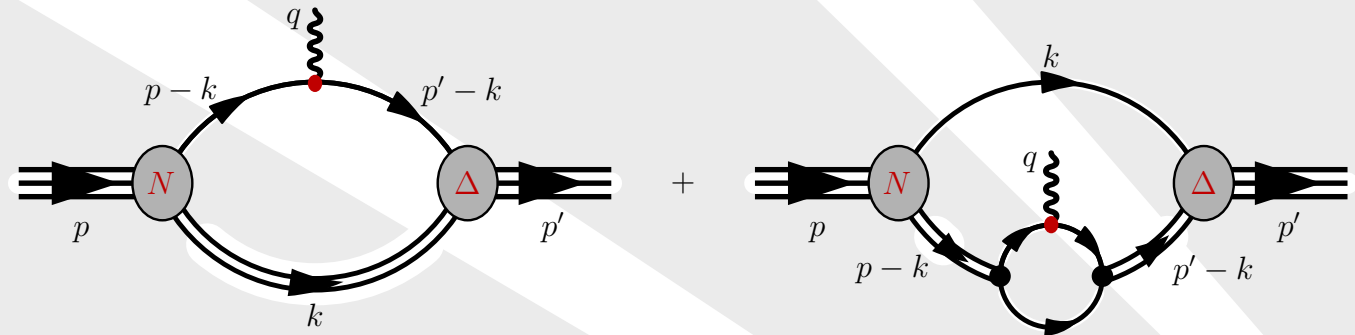
- ❖  $M_0 = 400, \quad M_s = 690, \quad M_a = 990, \quad \dots$

- Can now model a very large array of observables

# Nucleon-Delta Transition Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
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- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
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- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

## ● Transition Form Factor Feynman diagrams



$$\color{red}{\blacklozenge} \quad \bullet = \gamma^\mu F_{1q}(Q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2M} F_{2q}(Q^2)$$

- Approach is **completely covariant**
- **No frame is assumed** & Current is conserved
- Diagrams are expressed in form (**Jones & Scadron**):

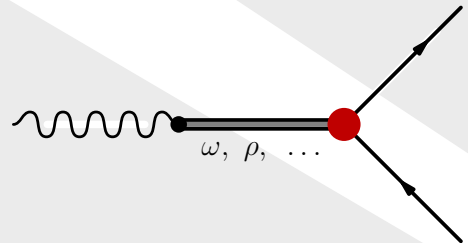
$$J^\mu = \bar{u}_{\Delta,\alpha} \left[ H_1^{\alpha\mu} G_1(Q^2) + H_2^{\alpha\mu} G_2(Q^2) + H_3^{\alpha\mu} G_3(Q^2) \right] u_N$$

- Many Model ingredients necessary

# Constituent Quark Form Factors

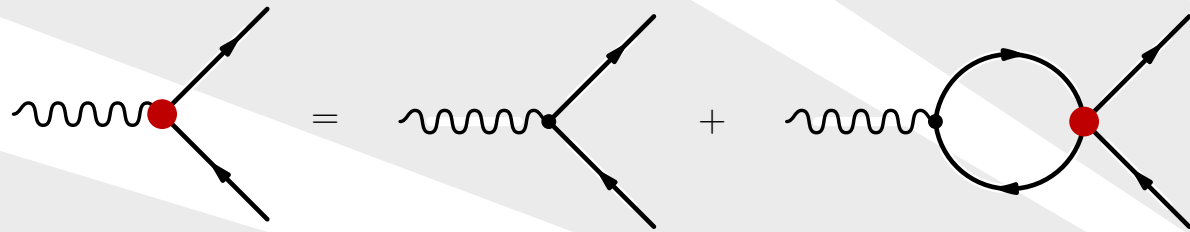
- ❖ Themes
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- Vector Meson Dominance – traditional view



$$\propto \frac{M_\rho^2}{M_\rho^2 + Q^2}$$

- We solve integral equation for vertex



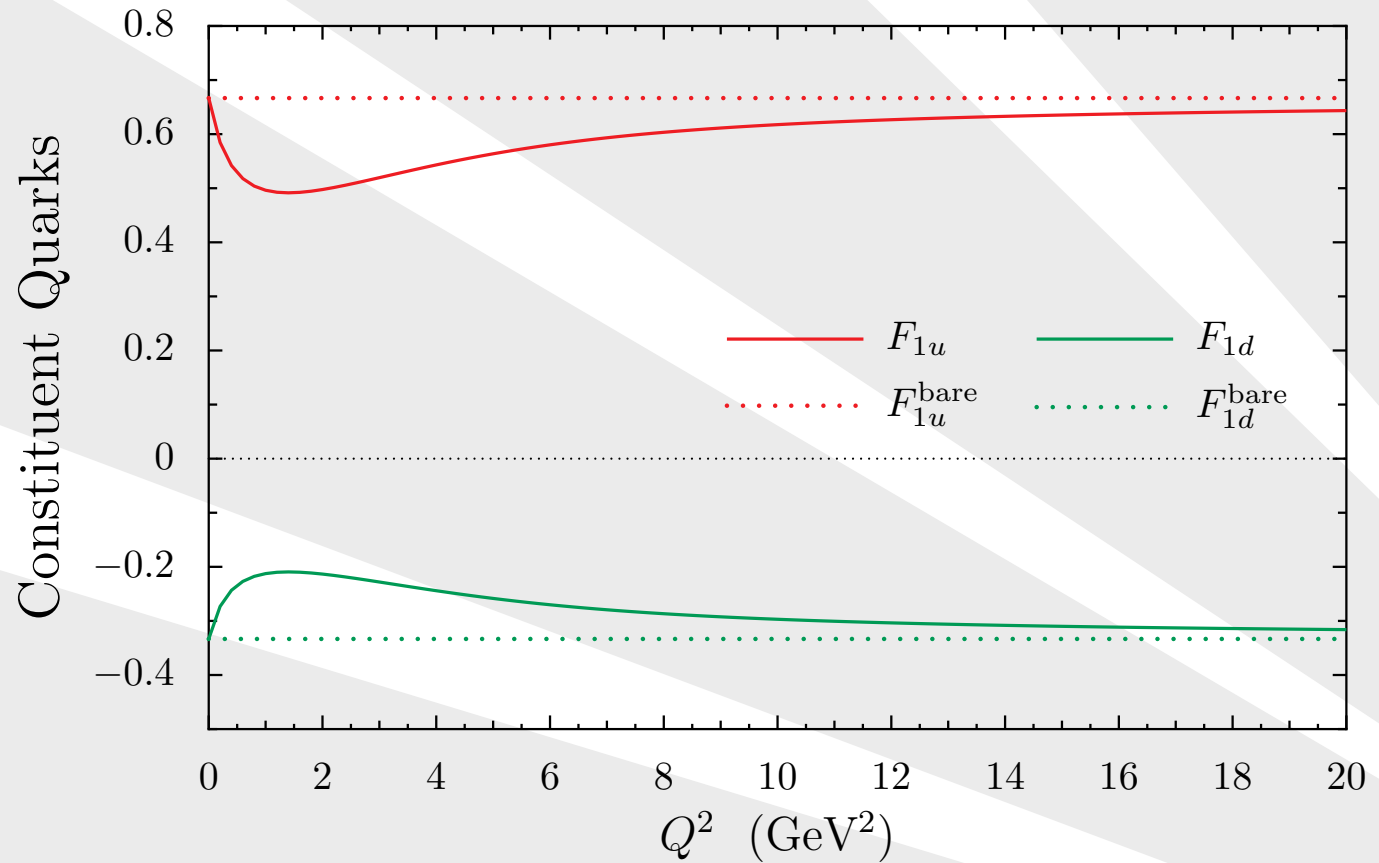
- Vertex becomes

$$\left( \frac{1}{6} + \frac{\tau_3}{2} \right) \gamma^\mu \rightarrow \left[ \frac{1}{6} F_\omega + \frac{\tau_3}{2} F_\rho \right] \gamma^\mu$$



# VMD Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ **Constituents**
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
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- ❖ Conclusion

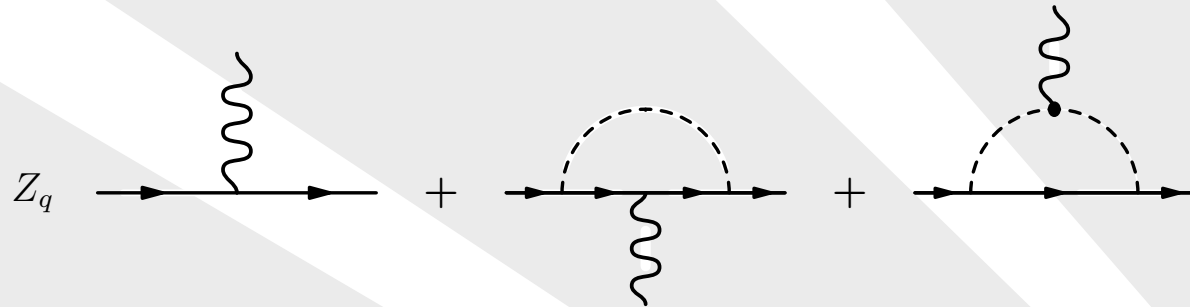


- Expanding about  $Q^2 = 0$  gives

$$\left[ \frac{1}{6} F_\omega + \frac{\tau_3}{2} F_\rho \right] \gamma^\mu \sim \left[ \frac{1}{6} \frac{M_\omega^2}{M_\omega^2 + Q^2} + \frac{\tau_3}{2} \frac{M_\rho^2}{M_\rho^2 + Q^2} \right] \gamma^\mu$$

# Constituent Quarks – Pion

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ **Constituents**
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



- **Probability** for find bare quark:  $Z_q = 1 + \frac{\partial \Sigma_q}{\partial \not{p}}$
- **Pion cloud**  $\rightarrow$  **anomalous m.m** for constituent quarks.

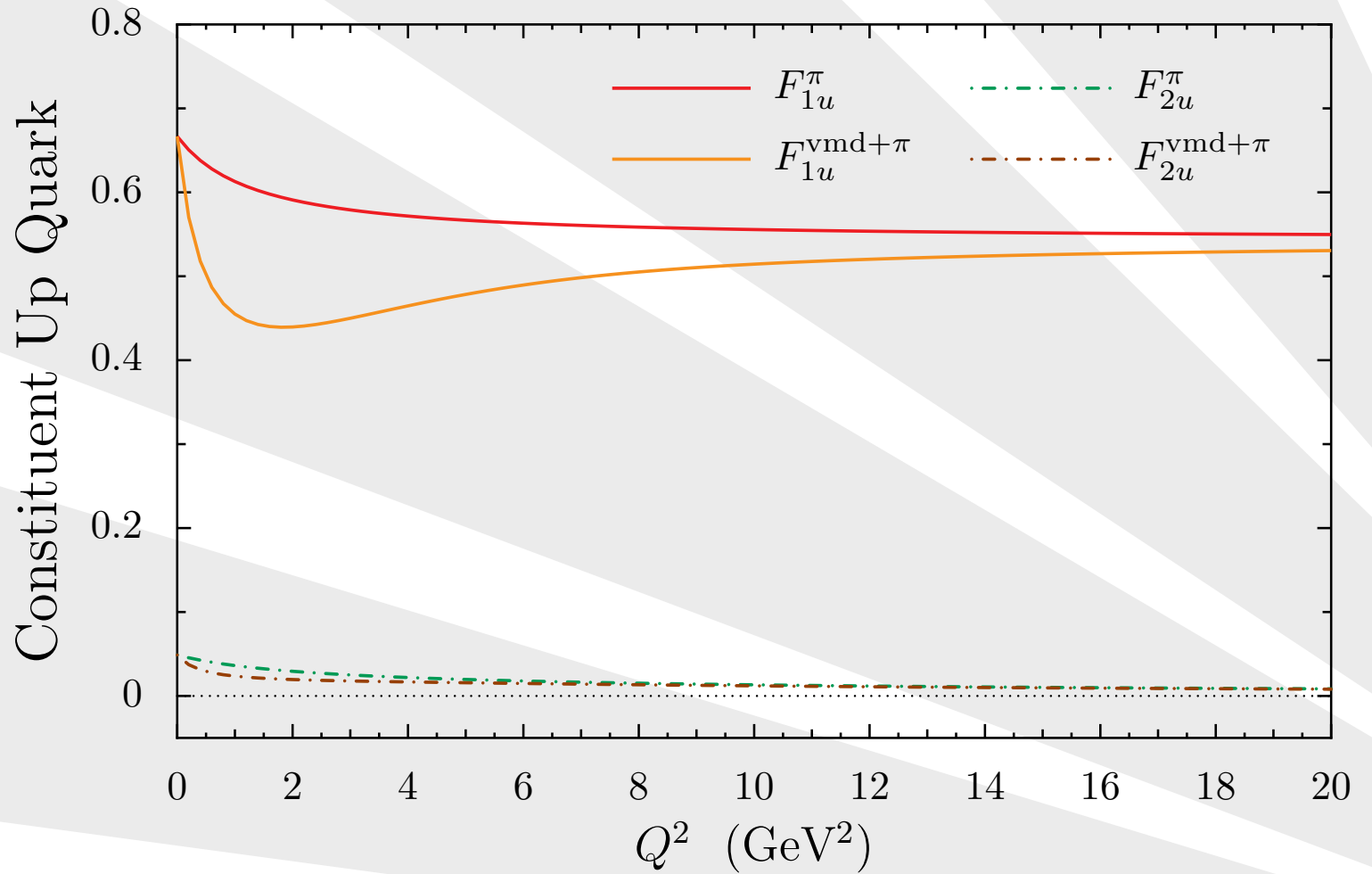
$$F_{1q}(Q^2) = Z_q \left( \frac{1}{6} F_\omega + \frac{1}{2} \tau_3 F_\rho \right) + (F_\omega - \tau_3 F_\rho) F_{1q}^{(q)} + \tau_3 F_\rho F_{1q}^{(\pi)}$$

$$F_{2q}(Q^2) = (F_\omega - \tau_3 F_\rho) F_{2q}^{(q)} + \tau_3 F_\rho F_{2q}^{(\pi)}$$

- **Self-consistent pion cloud**
- However **no pion exchange** between quarks
- **Better** to add pion at **nucleon level**

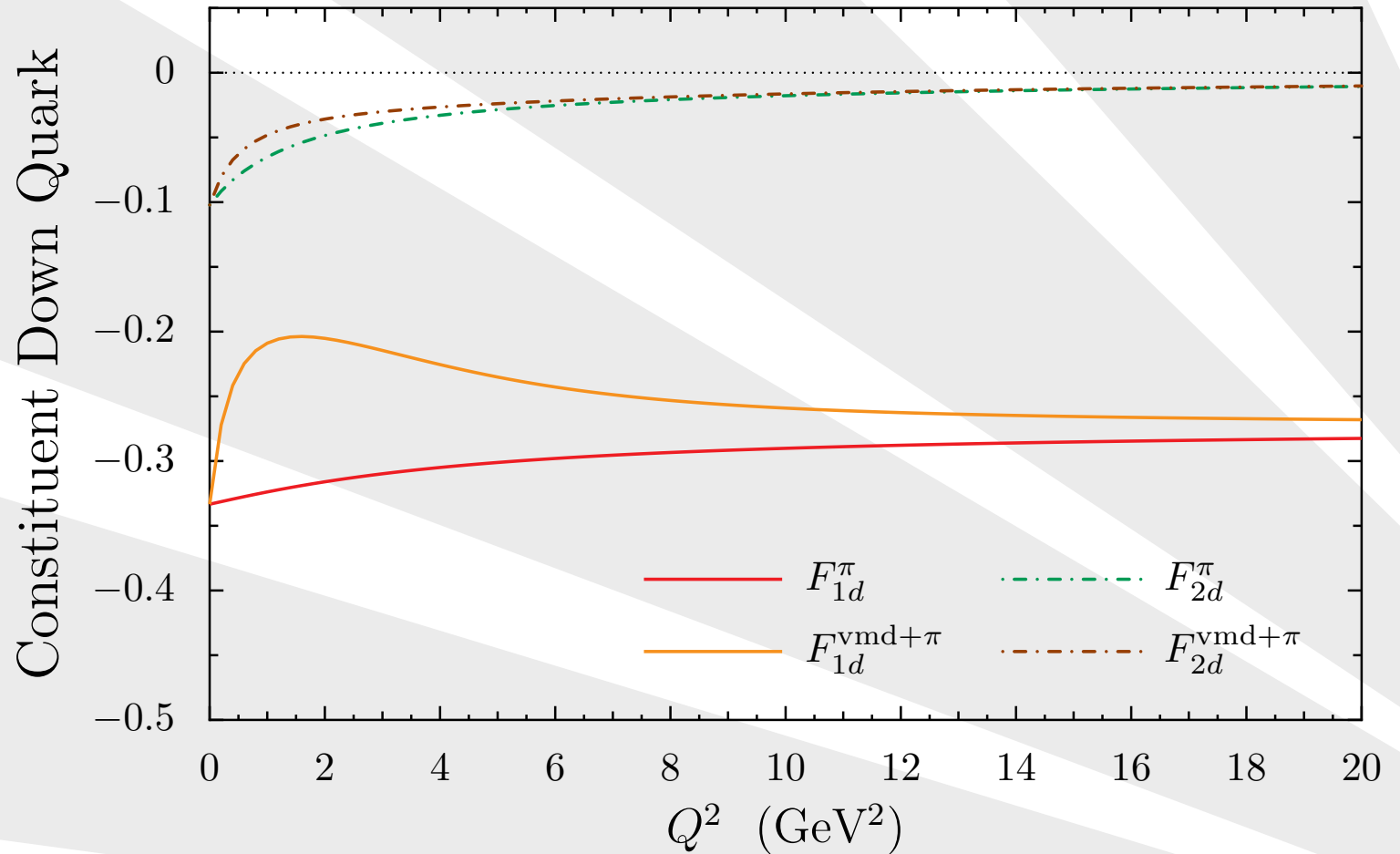
# Constituent Up Quark Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents**
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
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- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
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- ❖ Conclusion



# Constituent Down Quark Results

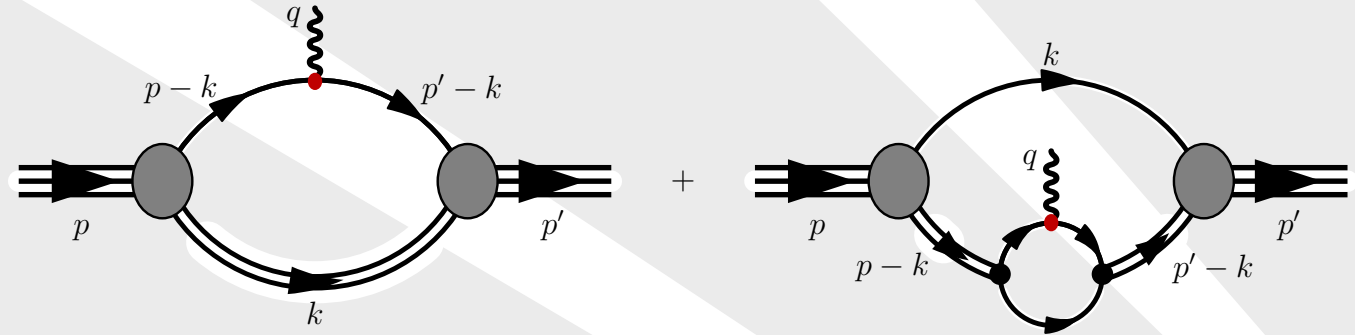
- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
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# Nucleon Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ **Nucleon FFs**
- ❖ Scalar Diquark FF
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## ● Form Factor Feynman diagrams



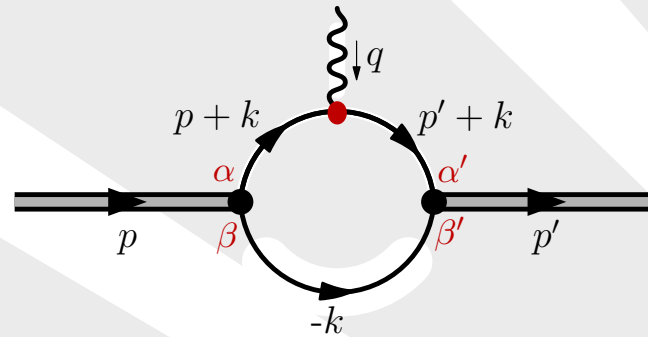
$$\bullet = \gamma^\mu F_{1q}(Q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2M} F_{2q}(Q^2)$$

- Approach is **completely covariant**
- **No frame is assumed**
- Charge is conserved automatically
- Diagrams are expressed in form:

$$\langle J^\mu \rangle = \bar{u}_N(p') \left[ \gamma^\mu F_{1N}(Q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2M_N} F_{2N}(Q^2) \right] u_N(p)$$

# Scalar Diquark & Pion Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ **Scalar Diquark FF**
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



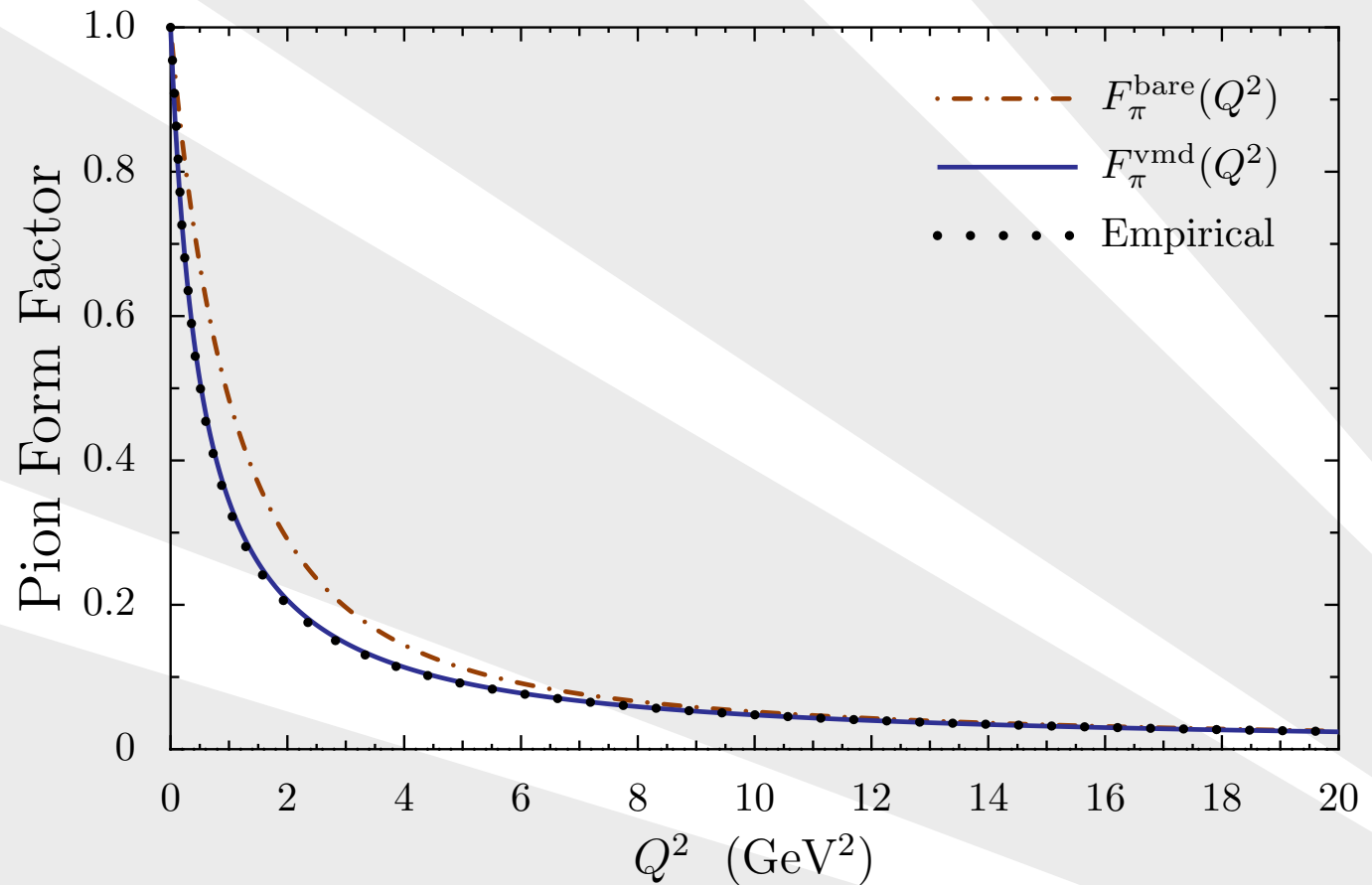
- **Scalar diquark:**  $(\gamma_5 C \tau_2 \beta^{A'}) (C^{-1} \gamma_5 \tau_2 \beta^A)$ , **Pion:**  $(\tau \gamma_5)$
- Form Factor expressions same:  $g_\pi \leftrightarrow g_s$ ,  $m_\pi \leftrightarrow M_s$
- Two form factors in general

$$\langle J_\pi^\mu \rangle = (p' + p)^\mu F_\pi(Q^2) + \underbrace{(p' - p)^\mu F_\pi^{\text{OS}}(Q^2)}_{\rightarrow 0}$$

- **Result charge radius:**  $\langle r_E^2 \rangle_\pi = 0.46 \text{ fm}^2$
- **Experiment:**  $\langle r_E^2 \rangle_\pi = 0.45 \pm 0.01 \text{ fm}^2$

# Pion Form Factor

- ❖ Themes
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- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ **Scalar Diquark FF**
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

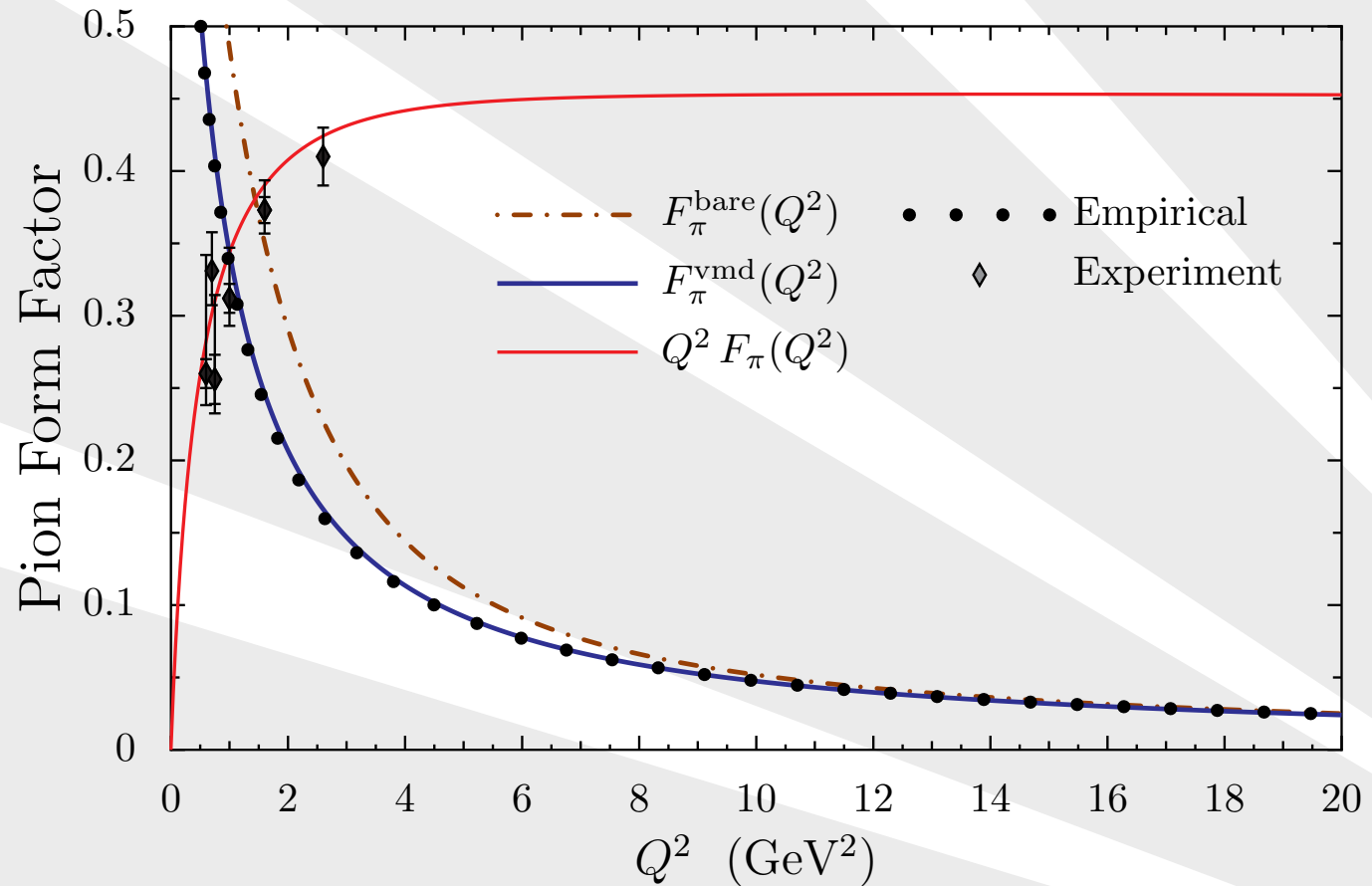


●  $F_\pi(Q^2) = [1 + Q^2/\Lambda^2]^{-1} \quad \Lambda^2 = 0.5 \text{ GeV}^2$

● No pion dressing on quarks  $\leftrightarrow \rho$  excitation

# Pion Form Factor

- ❖ Themes
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- ❖ Off-Shell
- ❖ Conclusion



● V. Tadevosyan *et al.* [Jefferson Lab F( $\pi$ ) Collaboration], Phys. Rev. C **75**, 055205 (2007)

●  $Q^2 F_\pi(Q^2) \rightarrow 16\pi f_\pi^2 \alpha_s(Q^2) \Rightarrow \alpha_{NJL} = 0.94 \Rightarrow Q^2 \sim 0.46 \text{ GeV}^2$



# Axial-Vector Diquark & Rho Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ **Axial-Vector FF**
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

- **AV diquark:**  $(\gamma^\beta C \tau_i \tau_2 \beta^{A'}) (C^{-1} \gamma^\alpha \tau_2 \tau_j \beta^A)$ , **Rho:**  $(\tau_j \gamma^\mu)$
- Form Factor expressions same:  $g_\rho \leftrightarrow g_a$ ,  $m_\rho \leftrightarrow M_a$
- 3 on shell form factors

$$J_\rho^\mu = \left[ g^{\alpha\beta} F_1(Q^2) - \frac{q^\alpha q^\beta}{2M_a^2} F_2(Q^2) \right] (p+p')^\mu - (q^\alpha g^{\mu\beta} - q^\beta g^{\mu\alpha}) F_3(Q^2)$$

- **Sachs Form Factors**

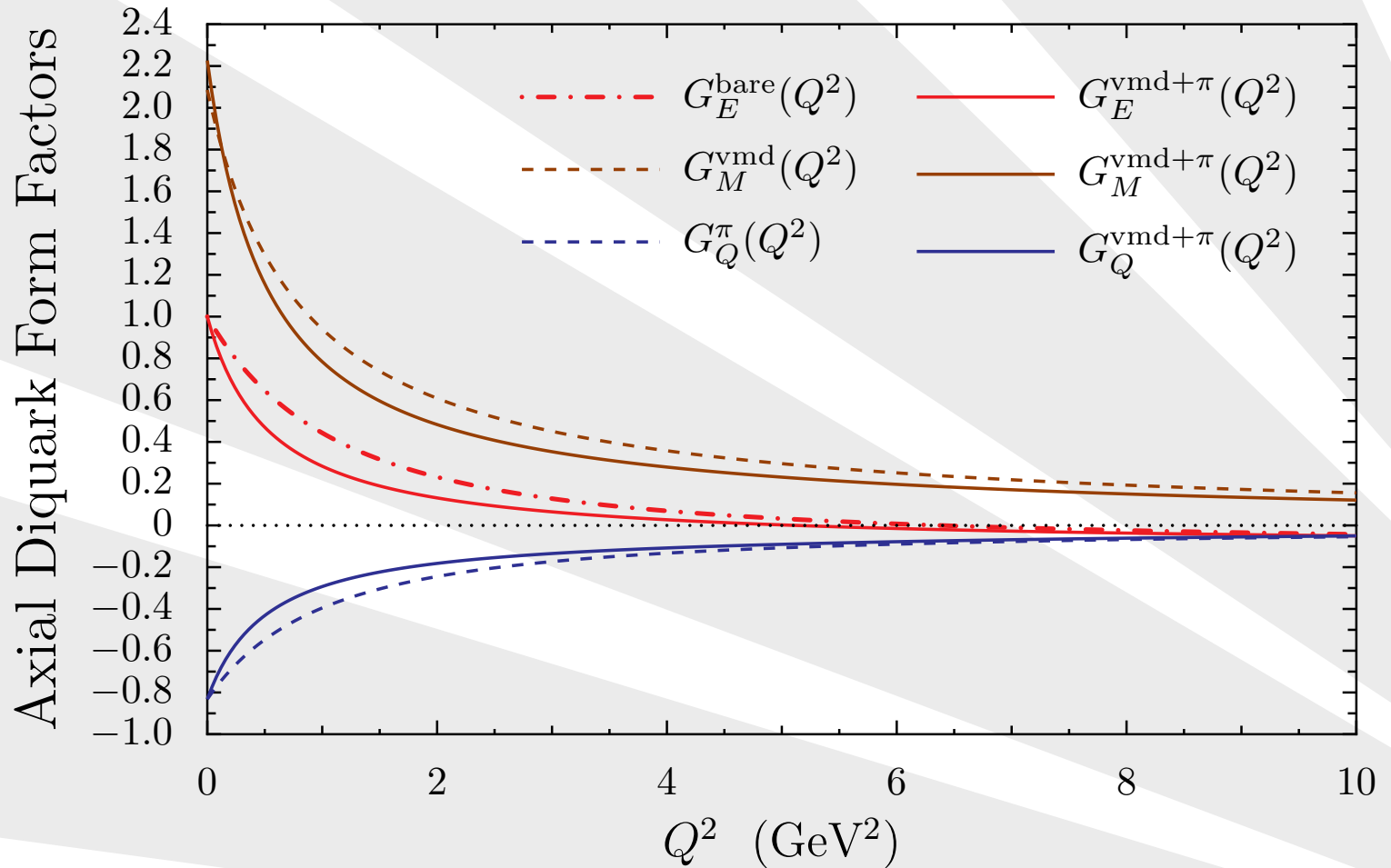
$$G_C(Q^2) = F_1(Q^2) + \frac{2}{3} \frac{Q^2}{4M^2} G_Q(Q^2), \quad G_M(Q^2) = F_3(Q^2)$$
$$G_Q(Q^2) = F_1(Q^2) + \left( 1 + \frac{Q^2}{4M^2} \right) F_2(Q^2) - F_3(Q^2)$$

- **NJL Results:**  $\langle r_E^2 \rangle_\rho = 0.52$ ,  $\mu_\rho = 2.08$ ,  $Q_\rho = -0.52$
- **DSE Results:**  $\langle r_E^2 \rangle_\rho = 0.54$ ,  $\mu_\rho = 2.01$ ,  $Q_\rho = -0.41$

❖ M. S. Bhagwat and P. Maris, Phys. Rev. C **77**, 025203 (2008)

# Axial-Vector Diquark Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



● For large  $Q^2$  we find  $G_Q \rightarrow G_E$ .

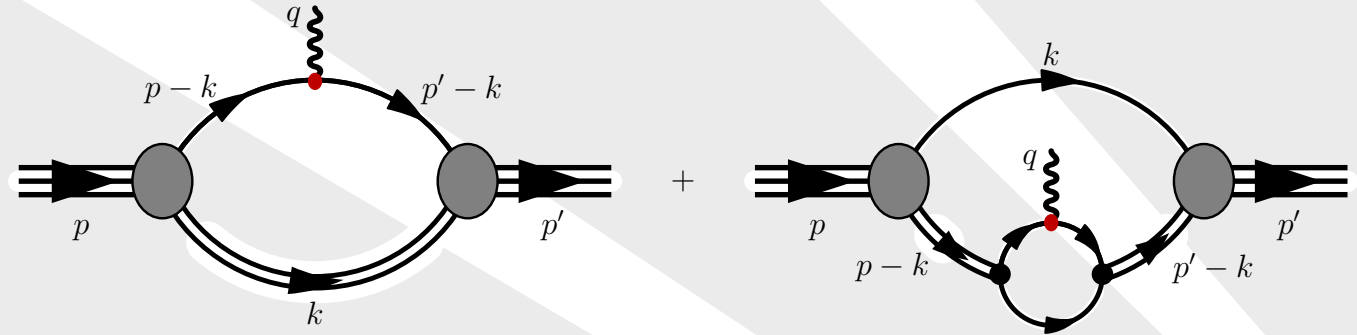
# Nucleon Form Factors – Recall

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF

## ❖ Nucleon FFs

- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

### ● Form Factor Feynman diagrams



$$❖ \bullet = \gamma^\mu F_{1q}(Q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2M} F_{2q}(Q^2)$$

### ● Diagrams are expressed in form:

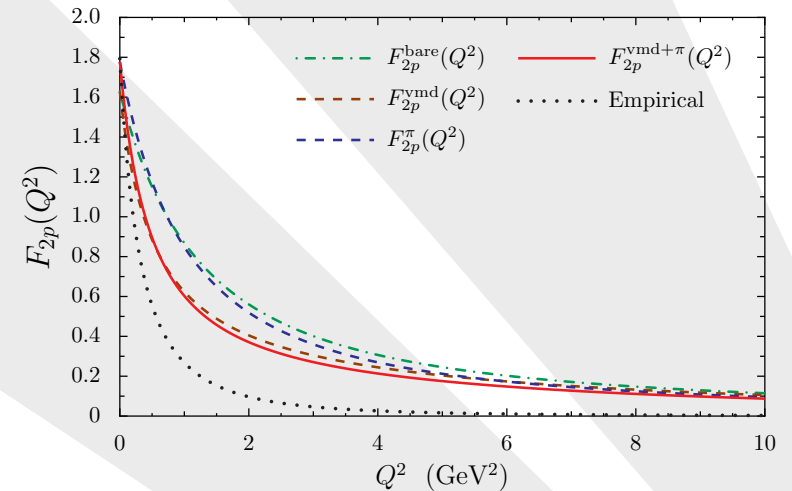
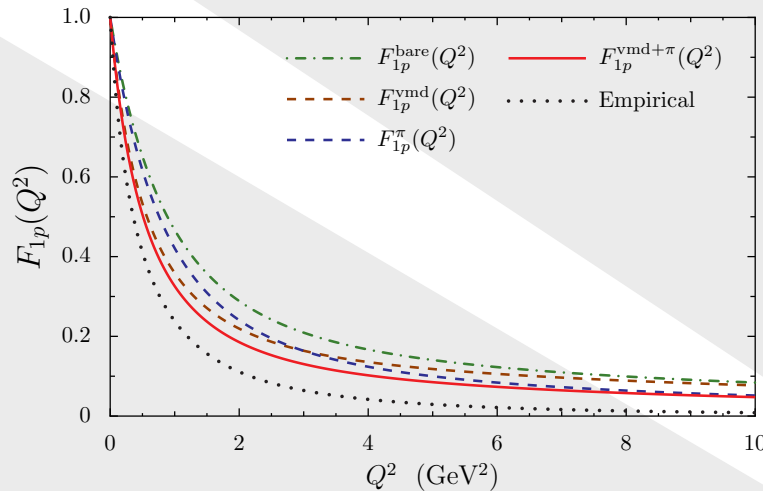
$$\langle J^\mu \rangle = \bar{u}_N(p') \left[ \gamma^\mu F_{1N}(Q^2) + \frac{i \sigma^{\mu\nu} q_\nu}{2M_N} F_{2N}(Q^2) \right] u_N(p)$$

# Proton Form Factors: Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF

## ❖ Nucleon FFs

- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



● Arrington, et al, Phys. Rev. C **76**, 035205 (2007)

● NJL:  $\kappa_p = 1.77$ , Experiment:  $\kappa_p = 1.79$

● NJL:  $\langle r_E^2 \rangle_p = 0.58 \text{ fm}^2$ , Experiment:  $\langle r_E^2 \rangle_p = 0.72 \text{ fm}^2$

● NJL:  $\langle r_M^2 \rangle_p = 0.56 \text{ fm}^2$ , Experiment:  $\langle r_M^2 \rangle_p = 0.71 \text{ fm}^2$

● NJL<sub>bare</sub>:  $\kappa_p = 1.61$ ,  $\langle r_E^2 \rangle_p = 0.36 \text{ fm}^2$ ,  $\langle r_M^2 \rangle_p = 0.38 \text{ fm}^2$

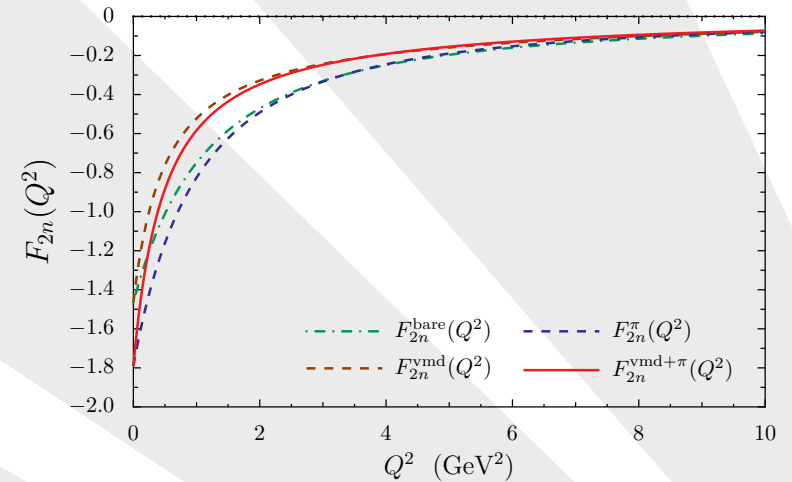
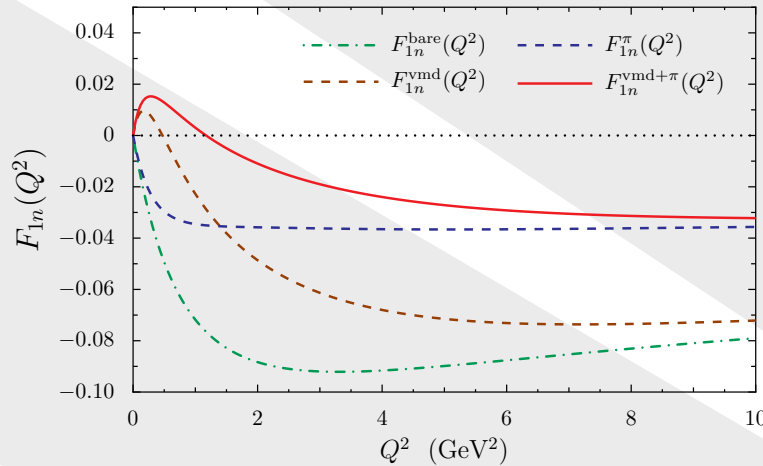
● Need extra  $\sim 1/Q^2$  factor  $\leftrightarrow$  Static Approximation

# Neutron Form Factors: Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
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- ❖ Nucleon FFs
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- ❖ Axial-Vector FF

## ❖ Nucleon FFs

- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
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- ❖ Conclusion

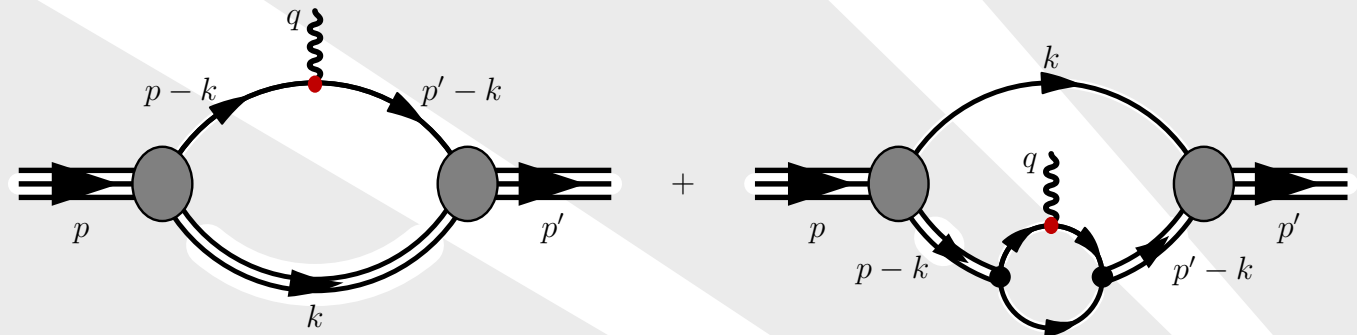


- **NJL:**  $\kappa_n = 1.79$ ,      **Experiment:**  $\kappa_n = 1.91$
- **NJL:**  $\langle r_E^2 \rangle_n = -0.15 \text{ fm}^2$ ,      **Experiment:**  $\langle r_E^2 \rangle_n = -0.12 \text{ fm}^2$
- **NJL:**  $\langle r_M^2 \rangle_n = 0.54 \text{ fm}^2$ ,      **Experiment:**  $\langle r_M^2 \rangle_n = 0.79 \text{ fm}^2$
- **bare:**  $\kappa_n = -1.46$ ,  $\langle r_E^2 \rangle_n = -0.07 \text{ fm}^2$ ,  $\langle r_M^2 \rangle_p = 0.38 \text{ fm}^2$
- **Need extra  $\sim 1/Q^2$  factor  $\leftrightarrow$  Static Approximation**

# Delta Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

## ● Form Factor Feynman diagrams



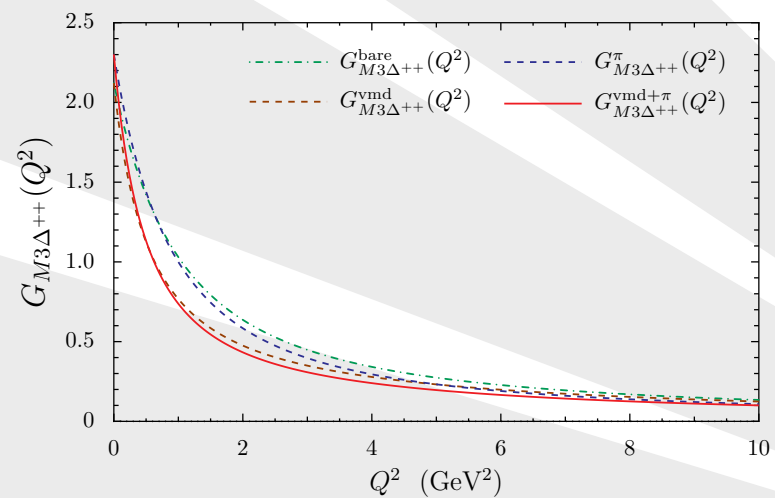
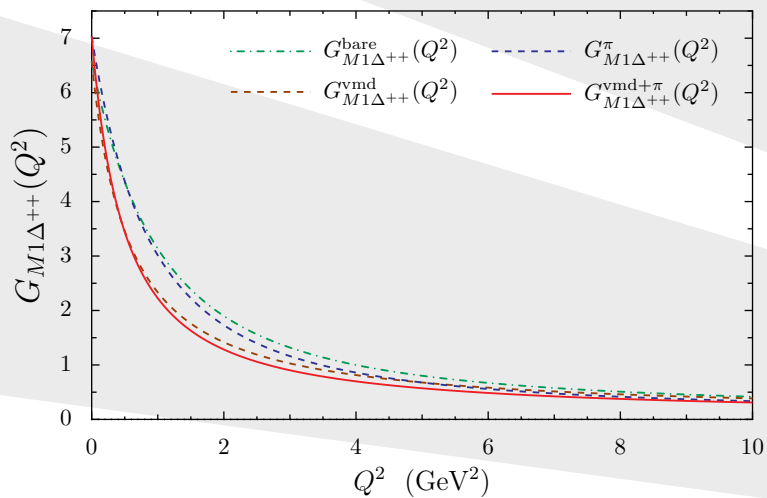
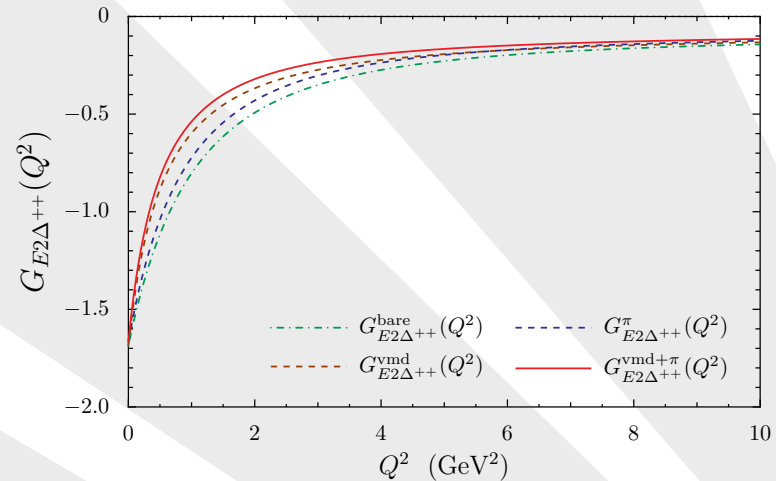
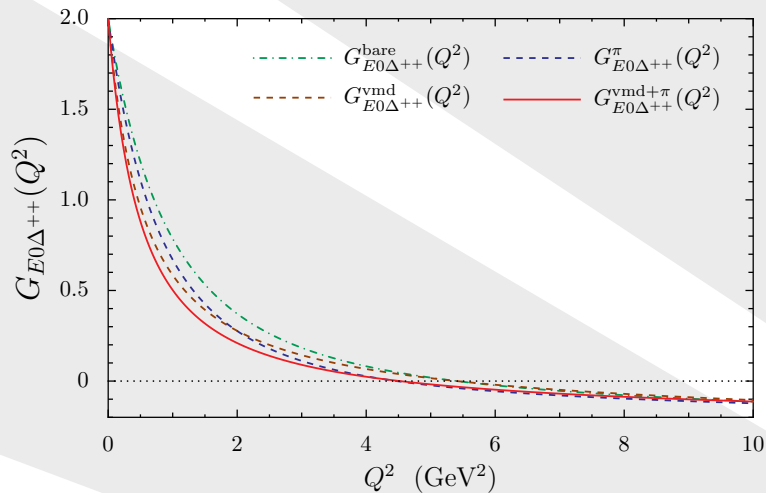
- Only axial-vector diquarks
- completely covariant, current conservation, no frame
- Diagrams are expressed in form:

$$\Gamma^{\mu,\alpha\beta} = H_1^{\alpha\beta,\mu} a_1 + H_2^{\alpha\beta,\mu} a_2 + H_3^{\alpha\beta,\mu} c_1 + H_4^{\alpha\beta,\mu} c_2$$

- Multipole Form Factors:  $G_{E0}, G_{E2}, G_{M1}, G_{M3}$

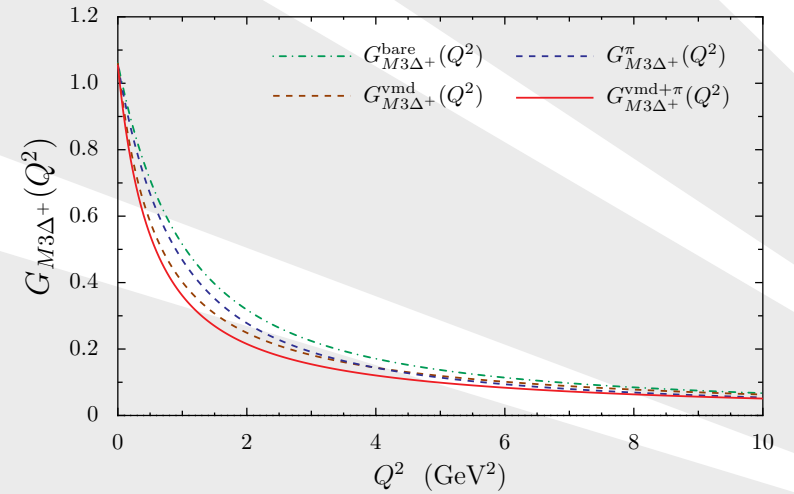
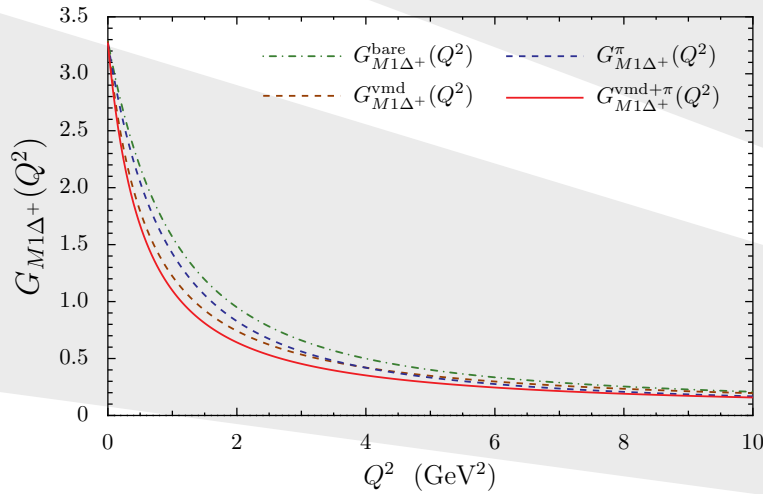
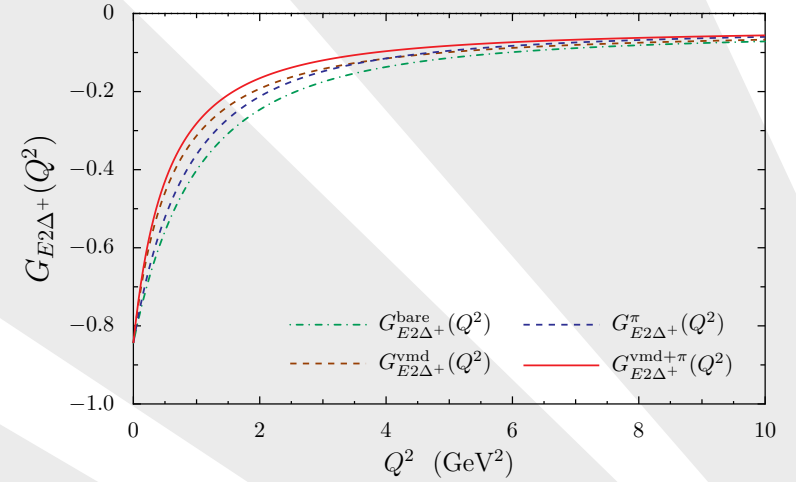
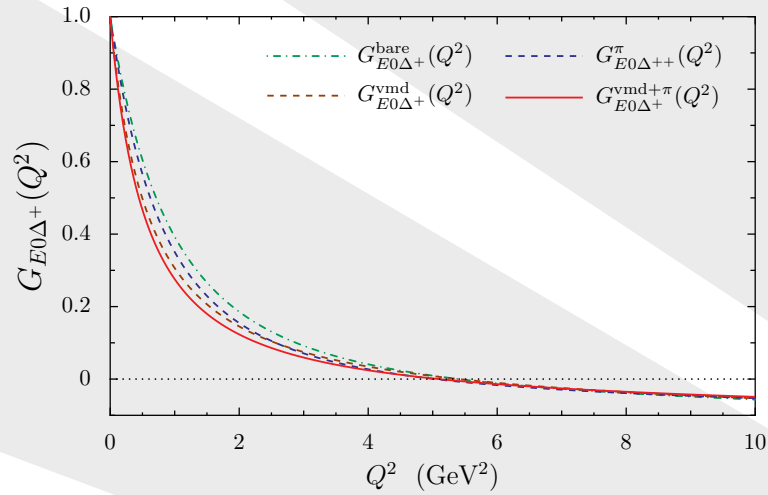
# $\Delta^{++}$ Form Factors Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



# $\Delta^+$ Form Factors Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

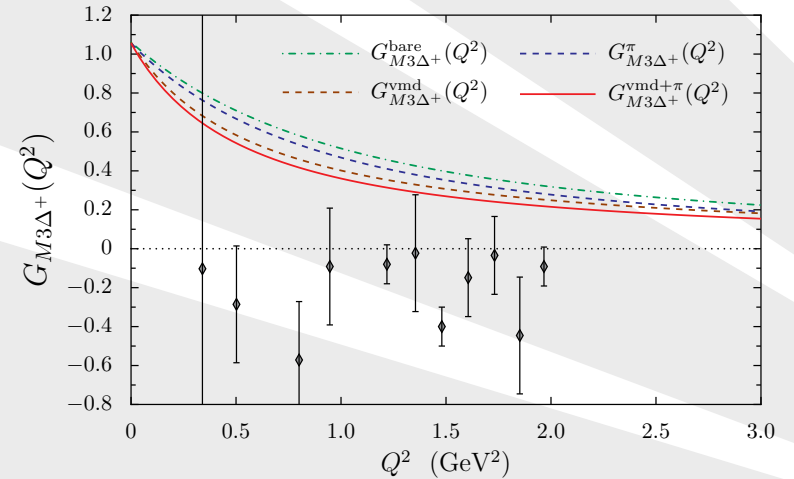
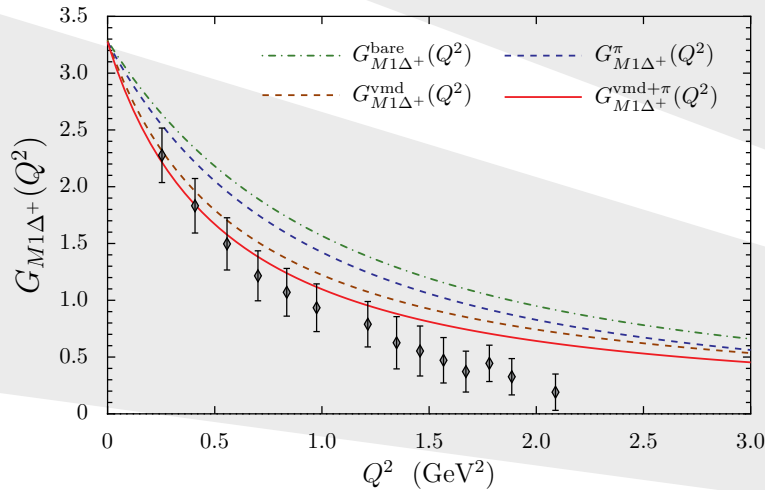
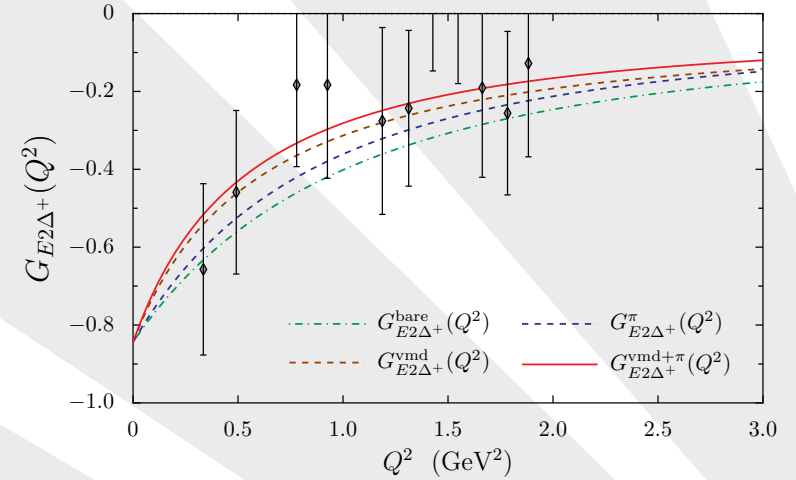
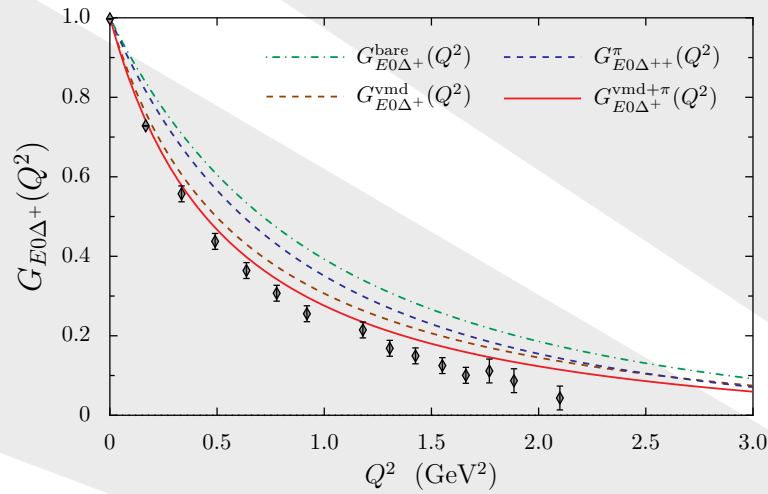


●  $G_{\Delta^+} \sim \frac{1}{2} G_{\Delta^{++}}$



# $\Delta^+$ Form Factors with Lattice Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



● C. Alexandrou, *et al*, PoS LAT2007, 149 (2007)

# Delta Moments

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ **Delta FFs**
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

## ● Largely unexplored experimentally: Summary

$$\blacklozenge \mu_{\Delta^{++}} = 4.52 \pm 0.51 \pm 0.45 \mu_N \text{ Or } \mu_{\Delta^{++}} = 6.14 \pm 0.51 \mu_N$$

Bosshard, 1991 Lopez Castro, 2002

$$\blacklozenge \mu_{\Delta^+} = 2.7_{-1.3}^{+1.0}(\text{stat.}) \pm 1.5(\text{syst.}) \pm 3(\text{theory}) \mu_N \text{ Kotulla, 2002}$$

## ● Chiral Extrapolated Lattice QCD Cloët, 2003

$$\blacklozenge \mu_{\Delta^{++}} = 4.99 \pm 0.56 \mu_N, \mu_{\Delta^+} = 2.49 \pm 0.29 \mu_N$$

$$\blacklozenge \mu_{\Delta^0} \sim 0.06 \mu_N, \mu_{\Delta^-} = 2.45 \pm 0.27 \mu_N$$

## ● NJL results ( $\mu_N$ )

$$\blacklozenge \text{Pion: } \mu_{\Delta^{++}} = 5.59, \mu_{\Delta^+} = 2.60, \mu_{\Delta^0} = -0.38, \mu_{\Delta^-} = -3.36$$

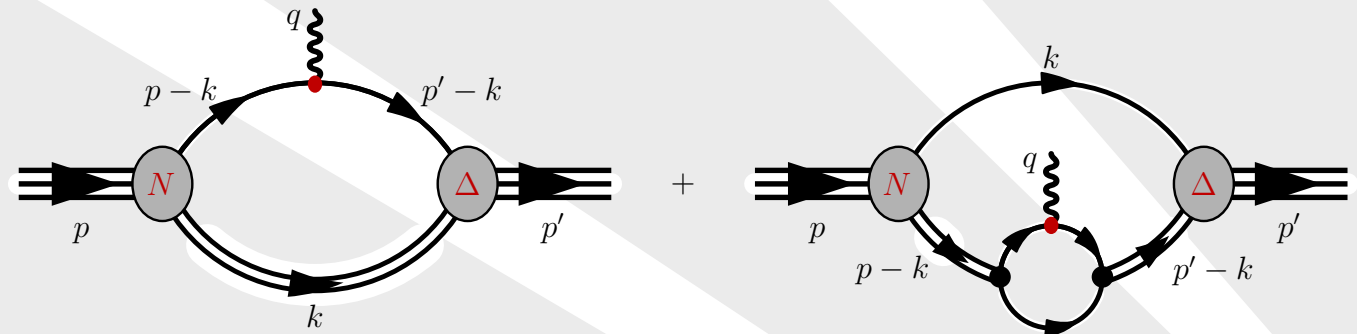
$$\blacklozenge \text{Bare: } \mu_{\Delta^{++}} = 5.39, \mu_{\Delta^+} = 2.61, \mu_{\Delta^0} = 0.0, \mu_{\Delta^-} = -2.61$$

## ● Lattice: $\mu_{\Delta^+} < \mu_p$ , NJL agrees.

# $N \rightarrow \Delta$ Transition Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

## ● Form Factor Feynman diagrams



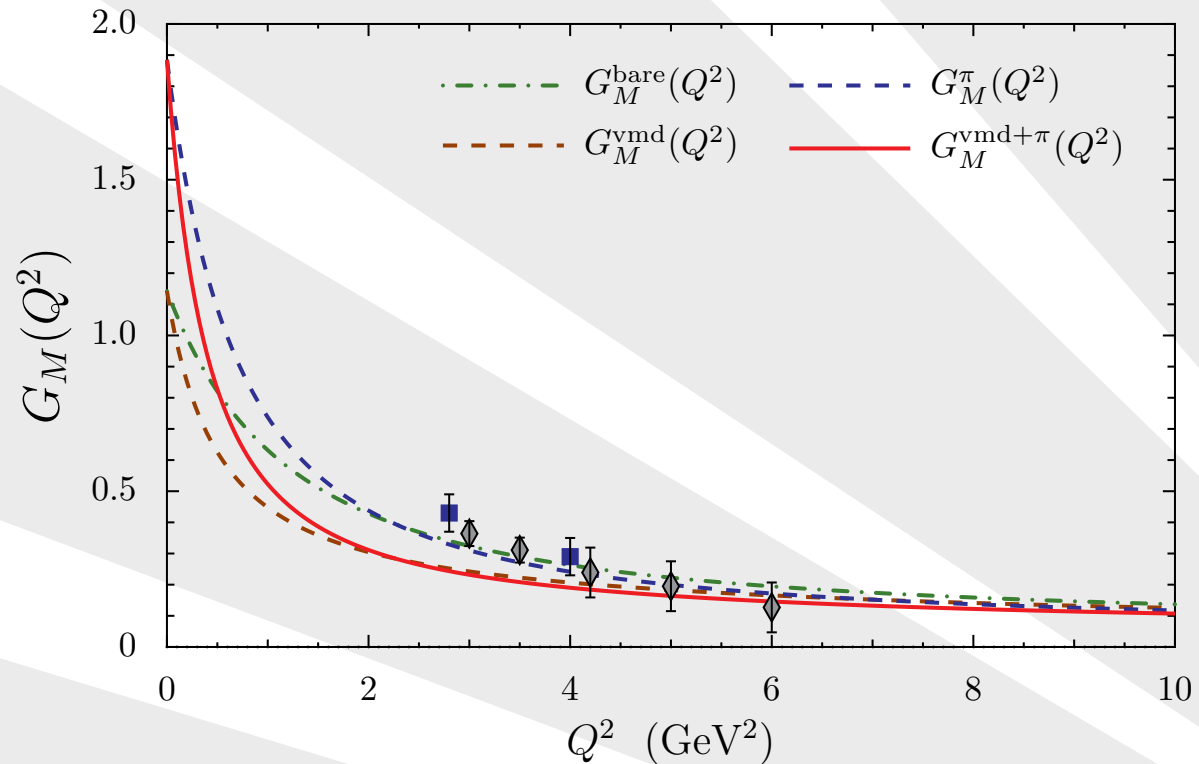
- Necessary ingredients are in place
- Scalar and axial-vector diquarks
- completely covariant, current conservation, no frame
- Diagrams are expressed in form:

$$J^\mu = \bar{u}_{\Delta,\alpha} \left[ H_1^{\alpha\mu} G_1(Q^2) + H_2^{\alpha\mu} G_2(Q^2) + H_3^{\alpha\mu} G_3(Q^2) \right] u_N$$

- Jones & Scadron multipole form factors:  $G_E, G_M, G_C$

# Form Factor Results: $G_M$

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



● V.V. Frolov, *et al.*, Phys. Rev. Lett. 82 (1999) 45.

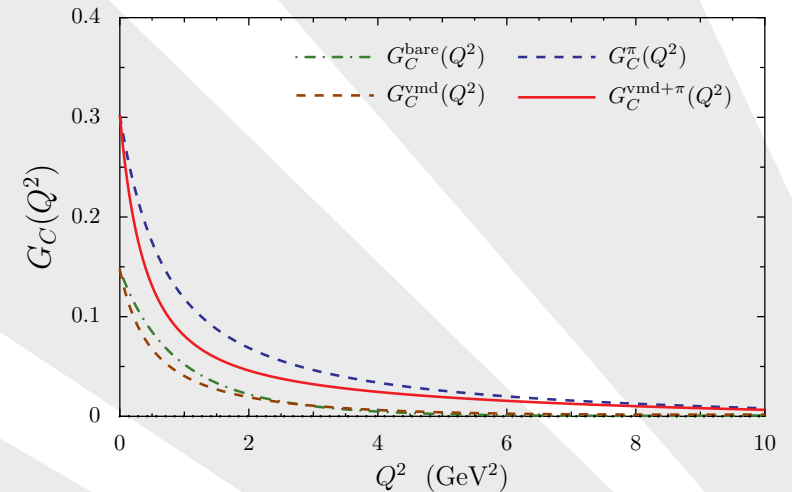
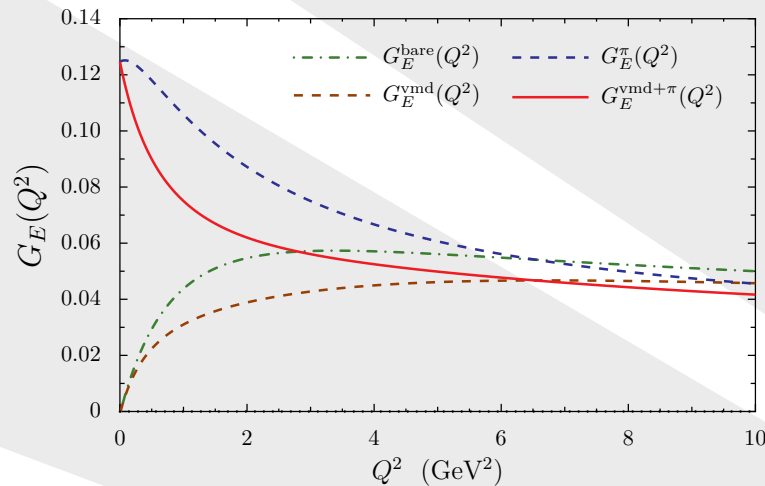
M. Ungaro, *et al.*, Phys. Rev. Lett. 97 (2006) 112003.

● Pion playing very important role

● We have a lot of missing strength at  $G_M(Q^2) \sim 0$ .

# Form Factor Results: $G_E$ & $G_C$

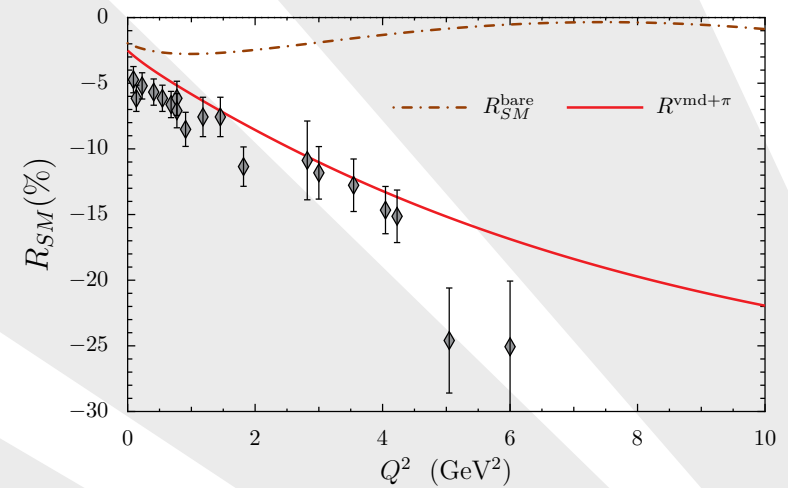
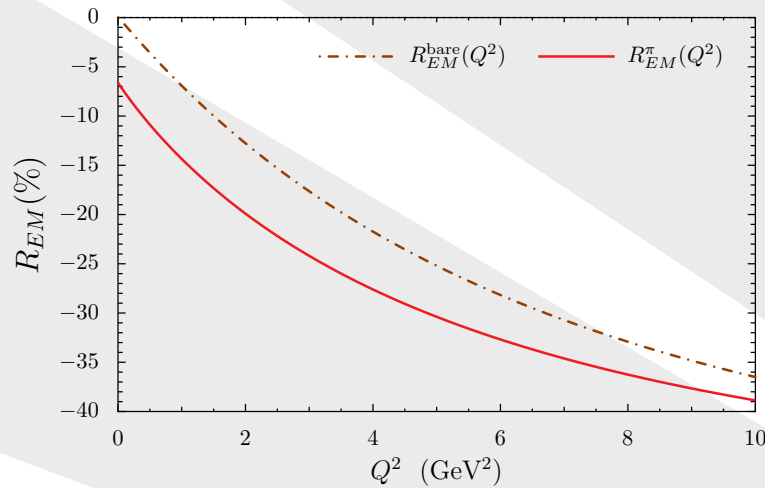
- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



- Form factors are small
- Pion cloud playing important large role
- Significant pion cloud effects for  $G_E$  at large  $Q^2$ .

# Form Factor Ratios

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

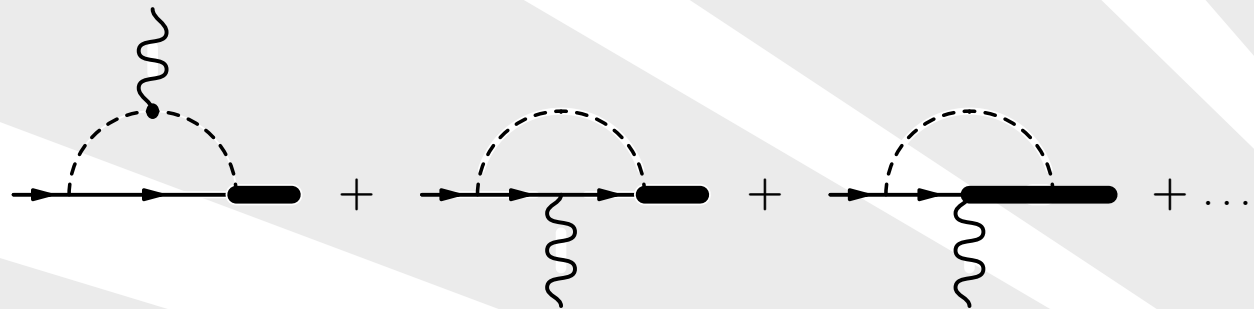


- $R_{EM}$  is roughly a fact 10 too large
- $G_E$  is very sensitive to cancellations between diagrams
  - ❖ compare neutron charge radius
- Nice result for  $R_{SM}$ 
  - ❖ Data from MAMI, LEGS, MIT-Bates and JLab

# What's Missing

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

- Need Faddeev calculation without static approx.
- Detailed description of pion cloud
- Need diagrams like:



- Important to quantify these effects
- Possible to calculate diagrams self-consistently in NJL

# Off-Shell Effects

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

- Largely unexplored in literature
- Relax constraint:  $p'^2 = p^2 = M^2$ .
- Very difficult, or impossible, in many model approaches
  - ❖ Dynamical model in terms of elementary d.o.f.
  - ❖ Solve dynamical equations.
- Potentially important: in-medium quark distributions, form factors, etc
- Pion has two off-shell form factors

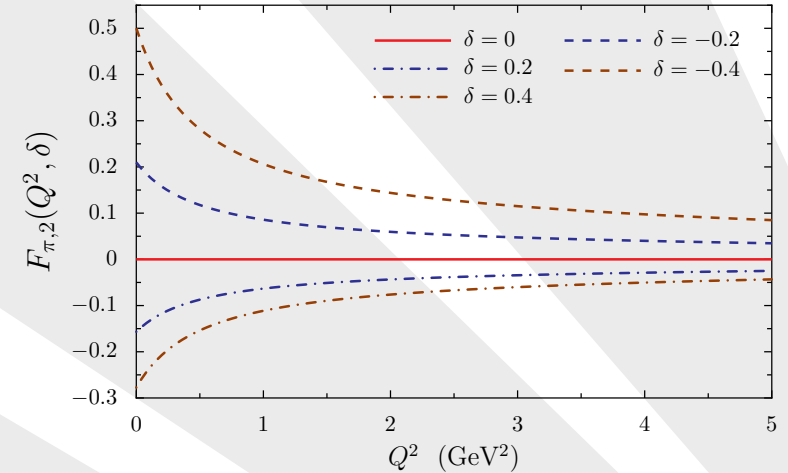
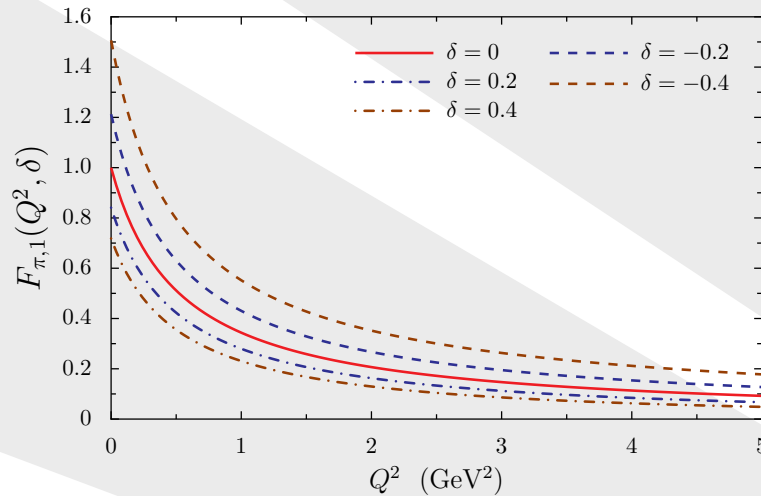
$$j_{\pi}^{\mu} = (p'^{\mu} + p^{\mu}) F_{\pi,1} + (p'^{\mu} - p^{\mu}) F_{\pi,2}$$

- For  $p'^2 = p^2 = M^2$ , we have  $F_{\pi,2} = 0$  &  $F_{\pi,1} \rightarrow F_{\pi}$ .



# Pion Off-Shell Form Factors

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ **Off-Shell**
- ❖ Conclusion



- Define  $\delta = p'^2 - p^2$ , put final state on-shell, i.e.  $p'^2 = m_\pi^2$
- Effects maybe large
- Off-shell effects may be important for experimental extraction of  $N \rightarrow \Delta$  transition.
- Also for in-medium form factors

# Conclusion & Outlook

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell

## ❖ Conclusion

- Covariant, Confining, Faddeev formalism:
  - ❖ Form Factors, Quark Distributions, etc
- Good description of “ $Q^2 = 0$ ” observables
  - ❖ e.g. magnetic moments, quark distributions
  - ❖ has been long standing problem
- Transition form factors require detailed understanding of pion contributions
  - ❖ Important challenge for quark models
- Size of off-shell effects may be studied
- Results so far suggest further investigation warranted.

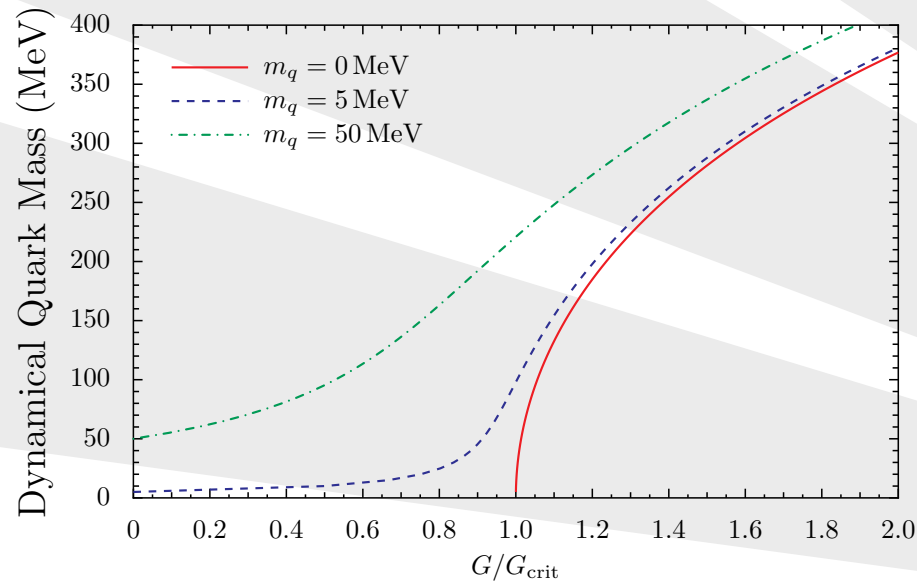
# Gap Equation

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



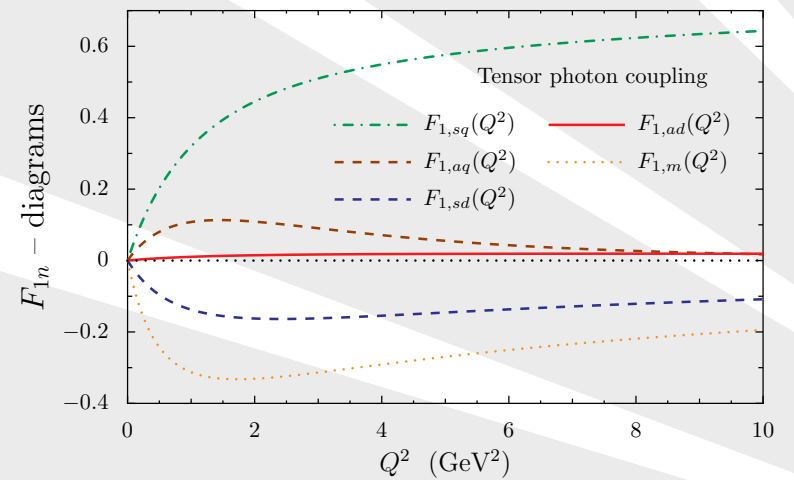
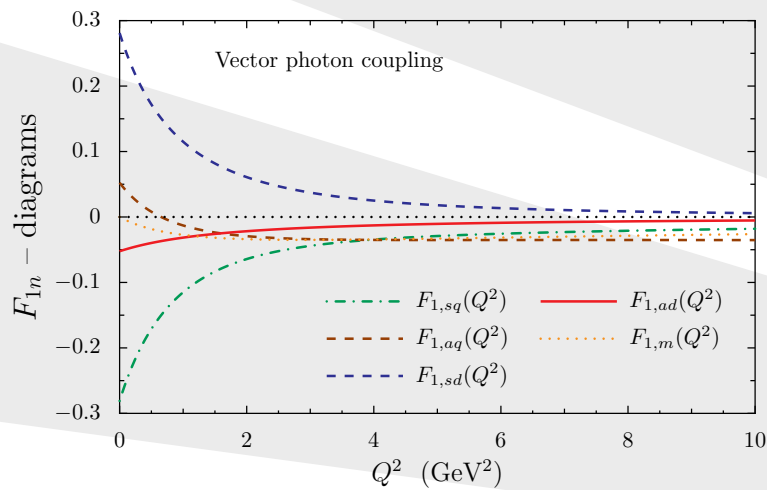
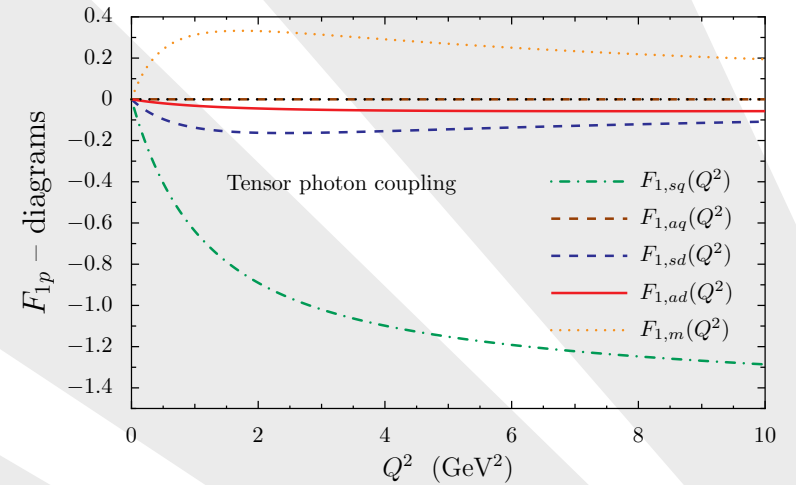
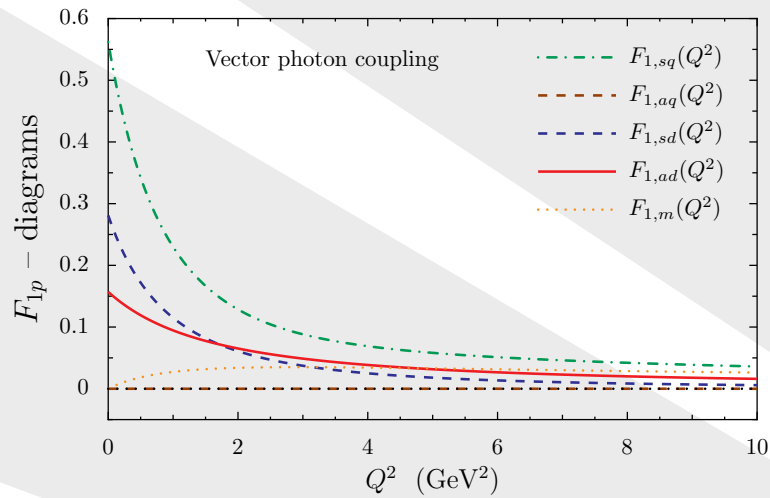
- Self-consistent solution – gives Quark Propagator

$$\frac{1}{\not{p} - m + i\epsilon} \longrightarrow \frac{1}{\not{p} - M + i\epsilon}$$



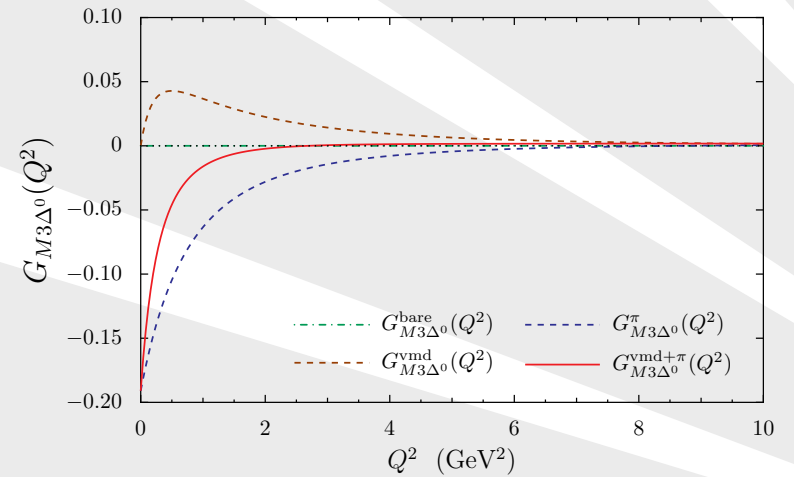
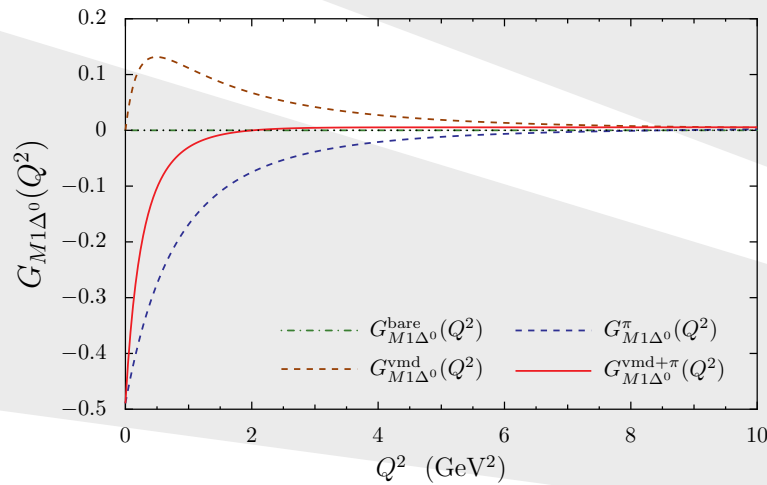
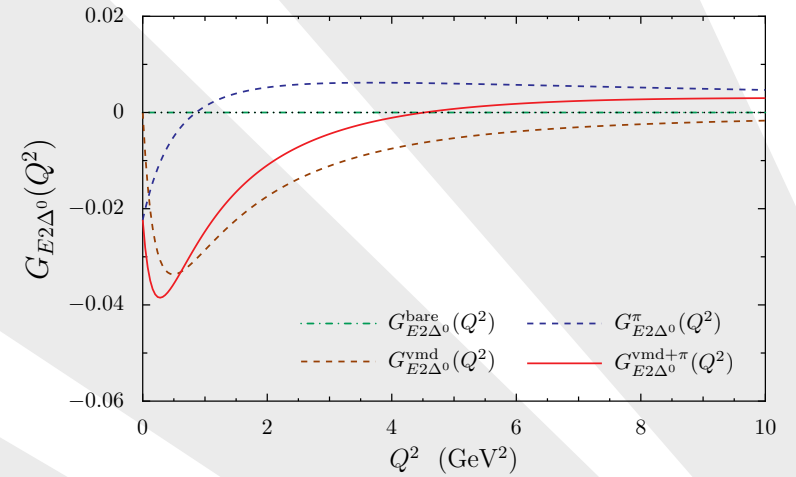
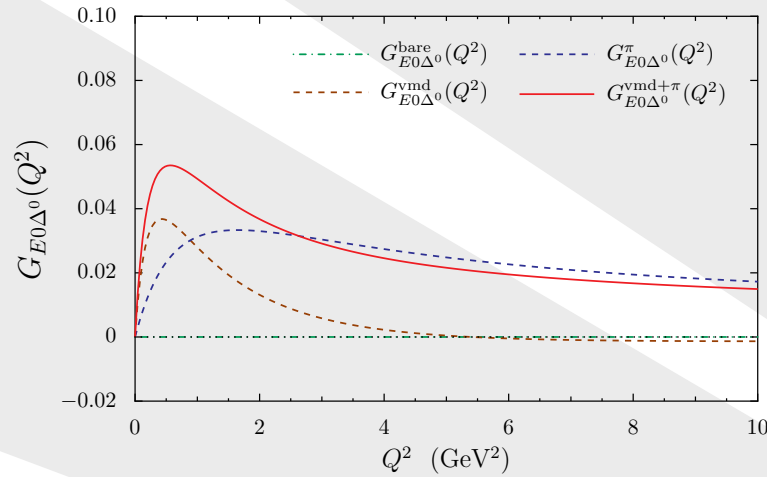
# $F_1$ Form Factor Diagrams

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



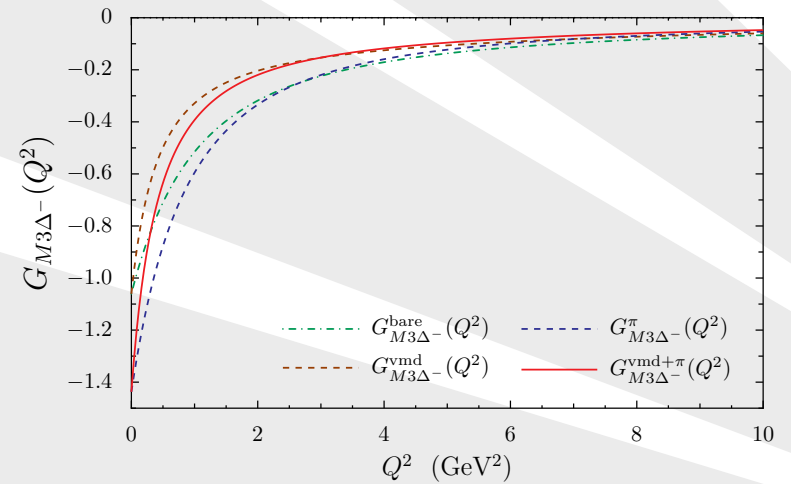
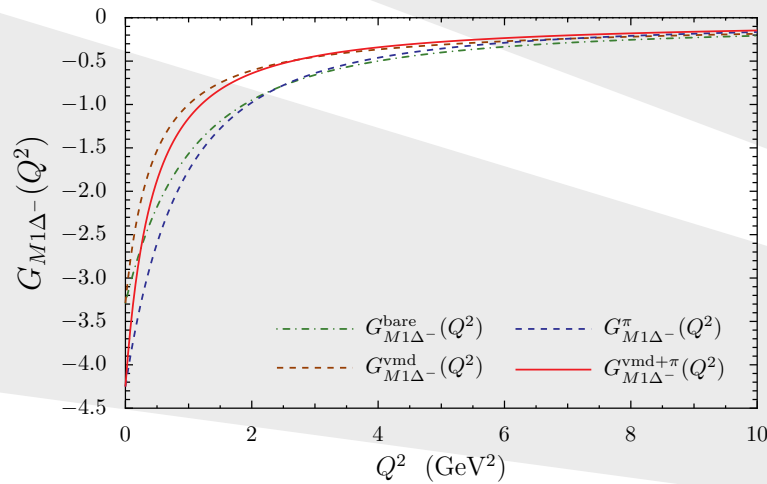
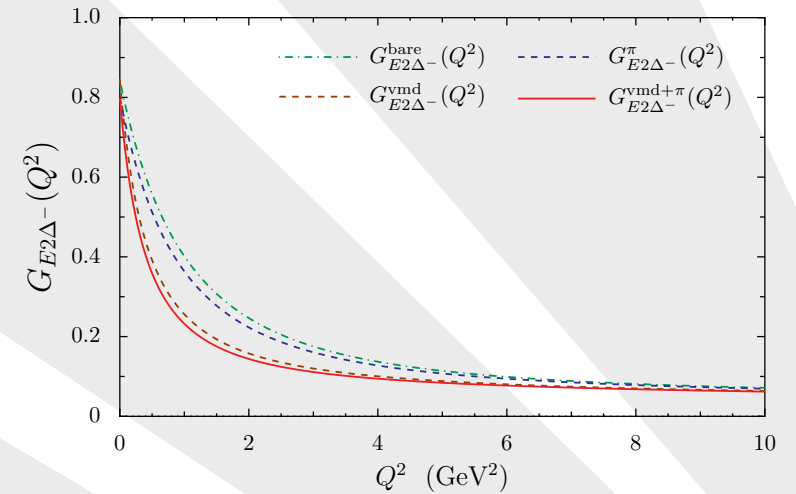
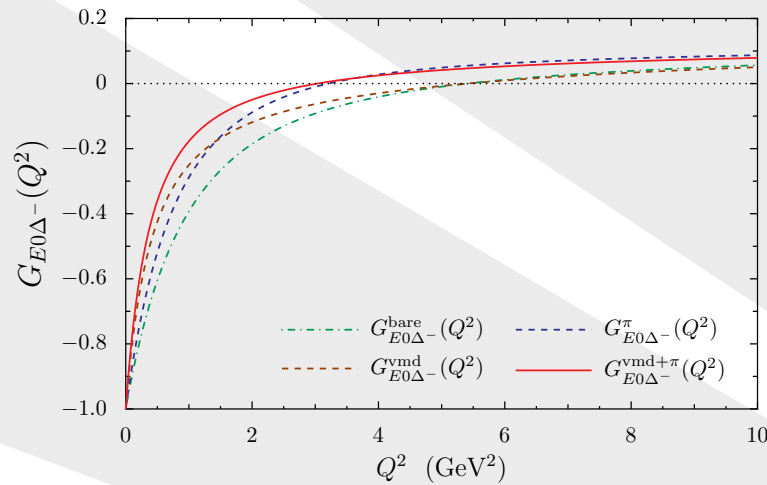
# $\Delta^0$ Form Factors Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



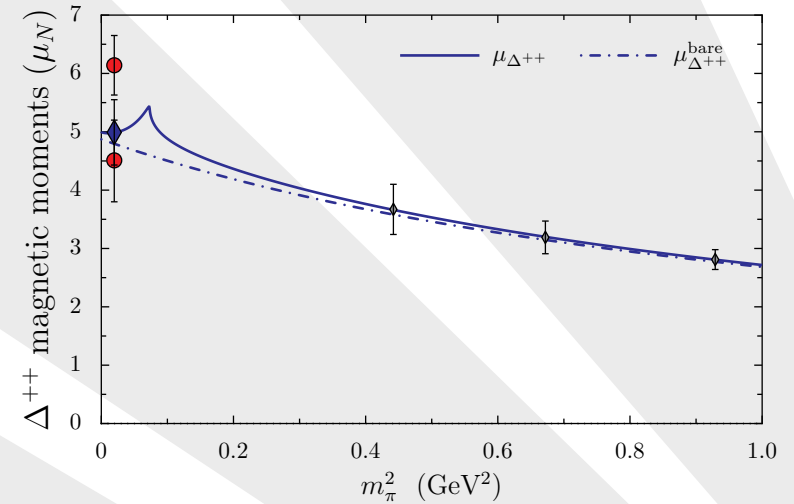
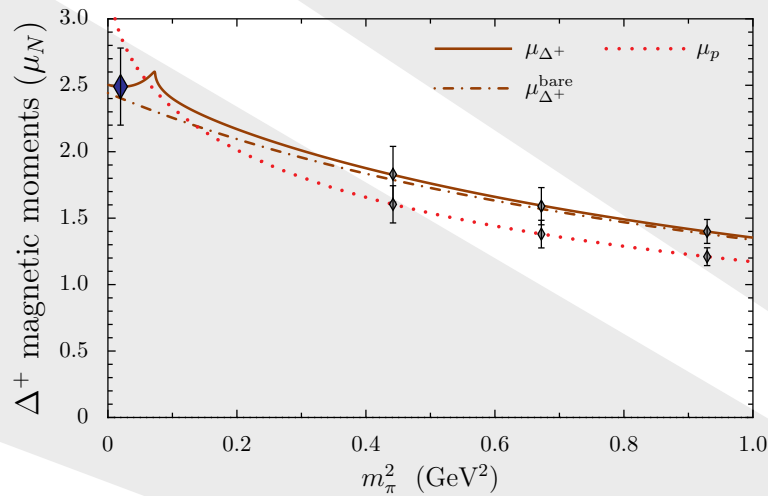
# $\Delta^-$ Form Factors Results

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



# Delta Moments: Pion Cloud

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



- Pion Cloud contribution surprisingly small
- NJL model consistent with this result
- Lattice:  $\mu_{\Delta^+} < \mu_p$ , NJL agrees.

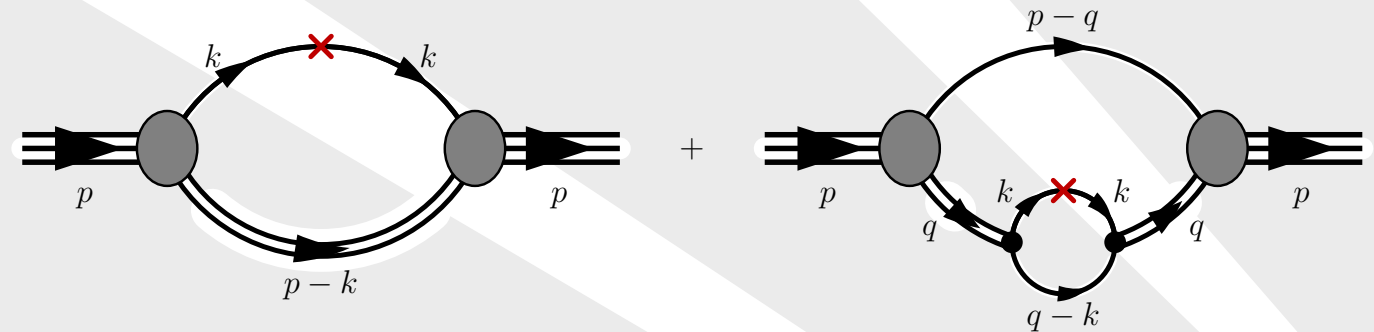
❖ C. Alexandrou, *et al*, Phys. Rev. D **74**, 034508 (2006)

❖ C. Alexandrou, *et al*, PoS **LAT2007**, 149 (2007)

# Nucleon quark distributions

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

- Associated with a Feynman diagram calculation.



$$\diamond [q(x), \Delta q(x), \Delta_T q(x)]$$

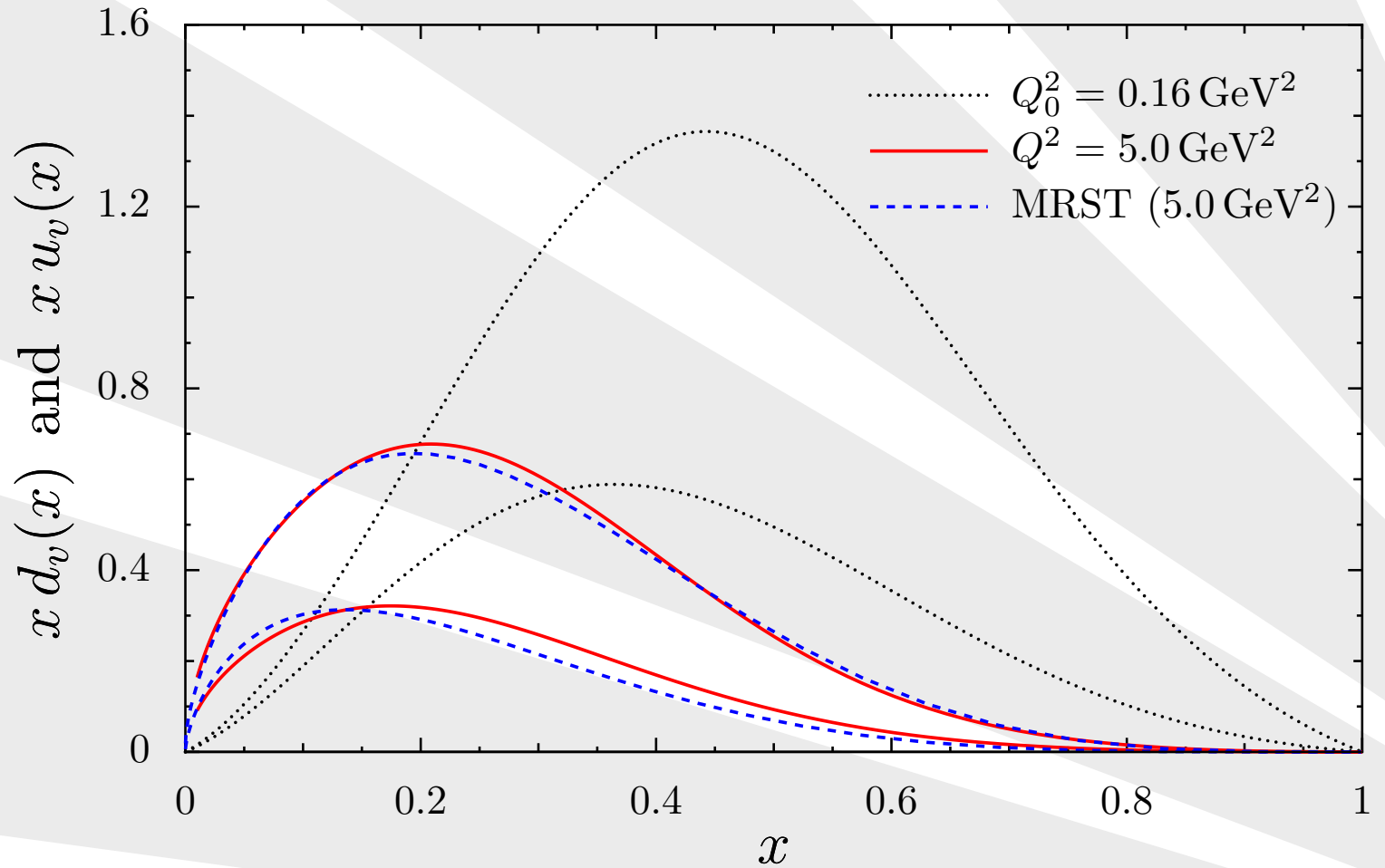
$$\rightarrow \mathbf{X} = \delta(x - \frac{k^+}{p^+}) [\gamma^+, \gamma^+ \gamma_5, \gamma^+ \gamma^1 \gamma_5]$$

- Satisfies baryon and momentum sum rules.
- Satisfies positivity constraints and Soffer bound.
- Covariant and gives correct support
- Model testing ground



# $u_v(x)$ and $d_v(x)$ distributions

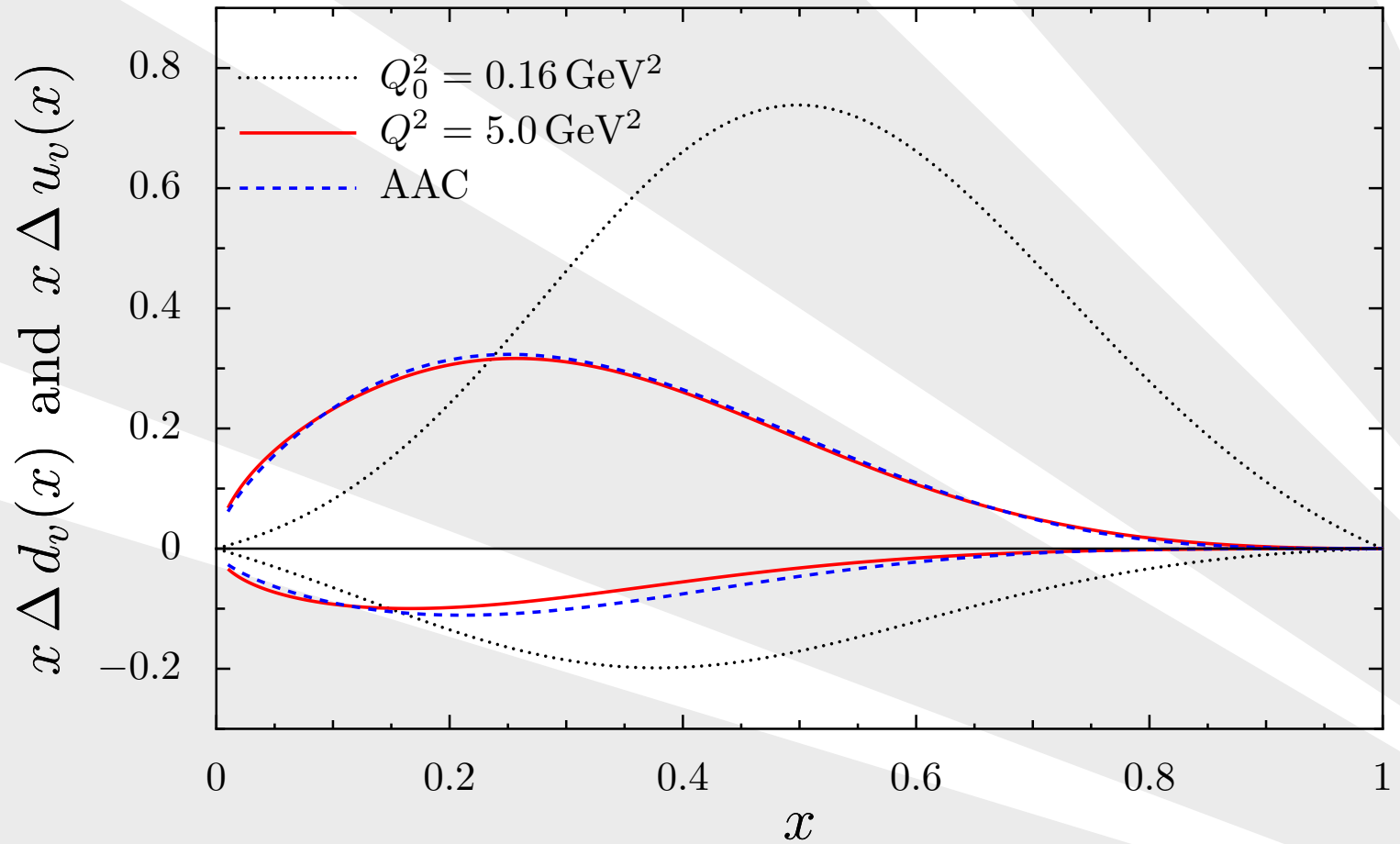
- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



● MRST, Phys. Lett. B **531**, 216 (2002).

# $\Delta u_\nu(x)$ and $\Delta d_\nu(x)$ distributions

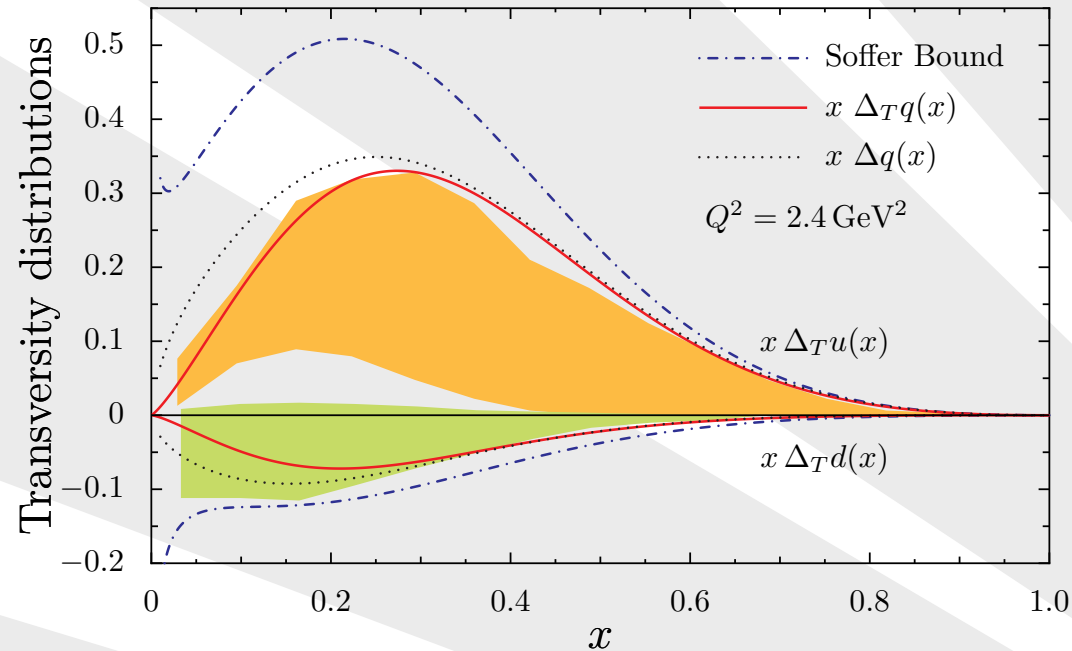
- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
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- ❖ Nucleon FFs
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- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



● M. Hirai, S. Kumano and N. Saito, Phys. Rev. D **69**, 054021 (2004).

# $\Delta_T u_v(x)$ and $\Delta_T d_v(x)$ distributions

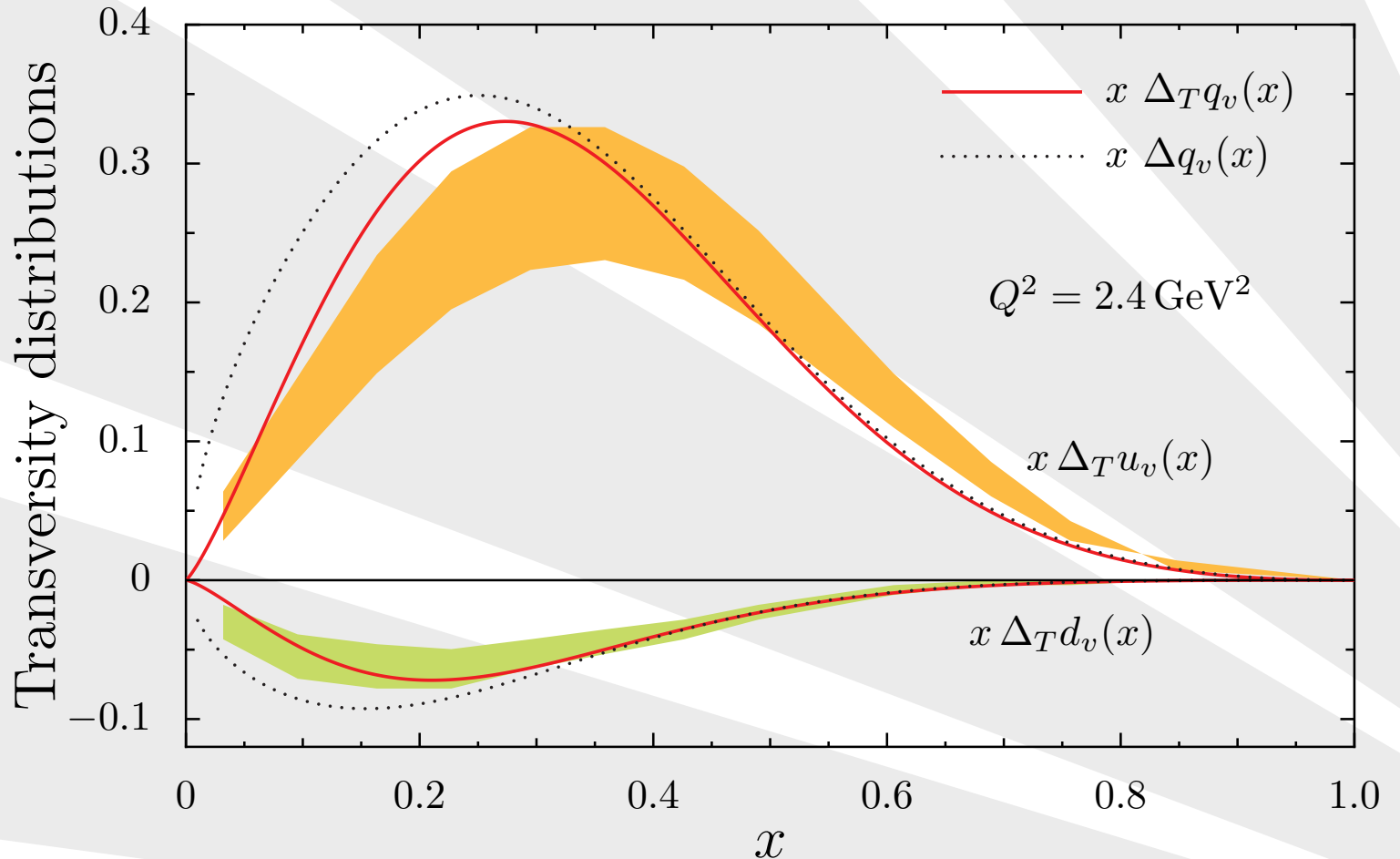
- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



- M. Anselmino *et. al.*, Phys. Rev. D **75**, 054032 (2007).
- Non-relativistic limit:  $\Delta_T q(x) = \Delta q(x)$
- $M \sim 400$  MeV, large relat. corrections unexpected
- Potential problem for models based concept of “constituent quarks” – maybe running mass

# Transversity: Reanalysis

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion

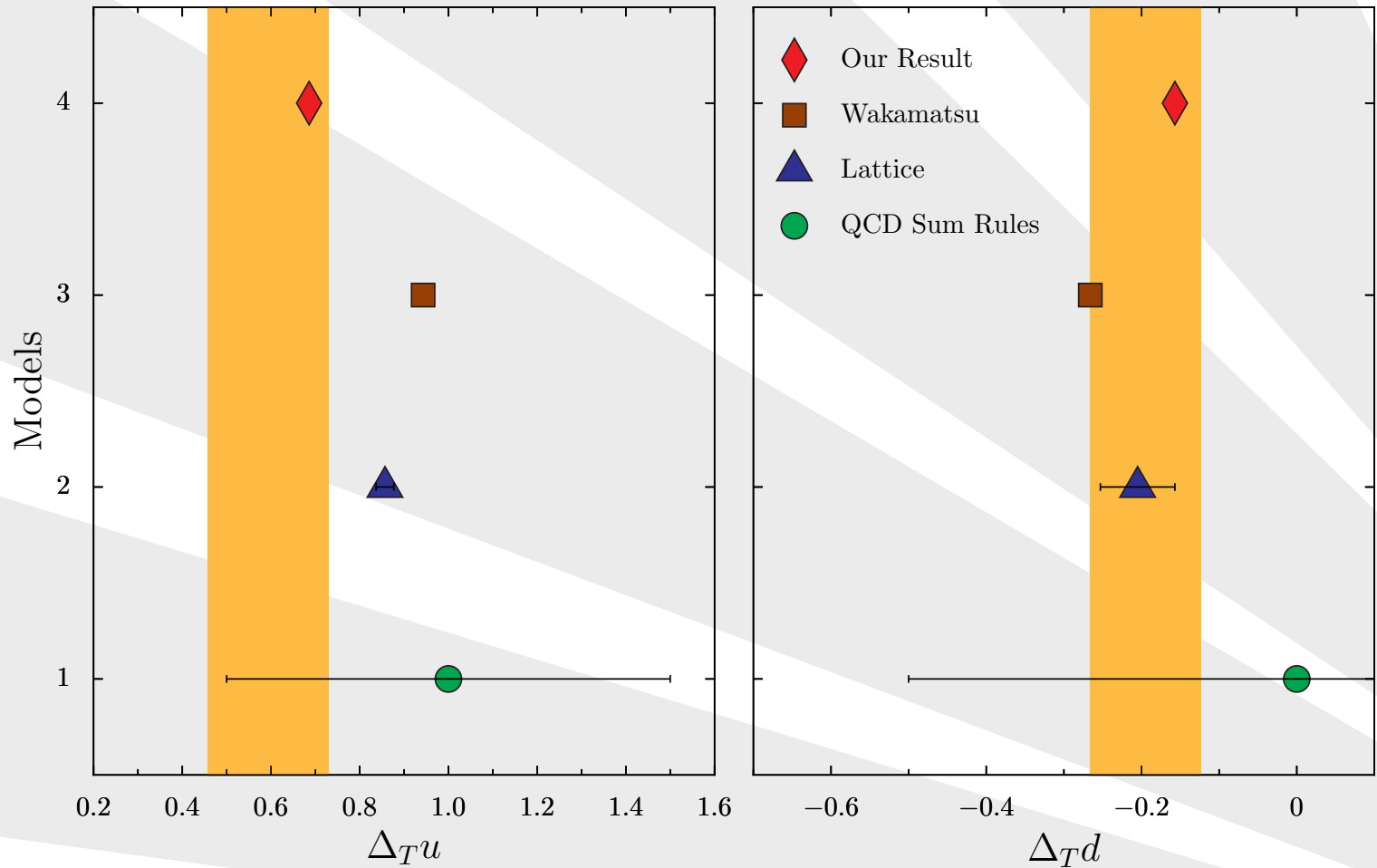


● Anselmino *et al* DIS 08

● I. C. Cloët, W. Bentz and A. W. Thomas, Phys. Lett. B **659**, 214 (2008)

# Transversity Moments

- ❖ Themes
- ❖ NJL model
- ❖ Baryons . . .
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Constituents
- ❖ Nucleon FFs
- ❖ Scalar Diquark FF
- ❖ Axial-Vector FF
- ❖ Nucleon FFs
- ❖ Delta FFs
- ❖  $N \rightarrow \Delta$  FFs
- ❖ Off-Shell
- ❖ Conclusion



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