

Lattice Calculations of $N-N^*$ Form Factors

Huey-Wen Lin

The Jefferson Lab logo consists of the words 'Jefferson Lab' in a bold, black, sans-serif font. A red swoosh underline is positioned under the 'e' in 'Jefferson'. Below the main text is the full name 'Thomas Jefferson National Accelerator Facility' in a smaller, black, sans-serif font, preceded by a small red dot.
Thomas Jefferson National Accelerator Facility



Oct. 14, 2008

In collaboration with

Saul Cohen, Robert Edwards, and David Richards (JLab)



Kostas Orginos
(W&M; JLab)

Lattice Baryon Form Factors

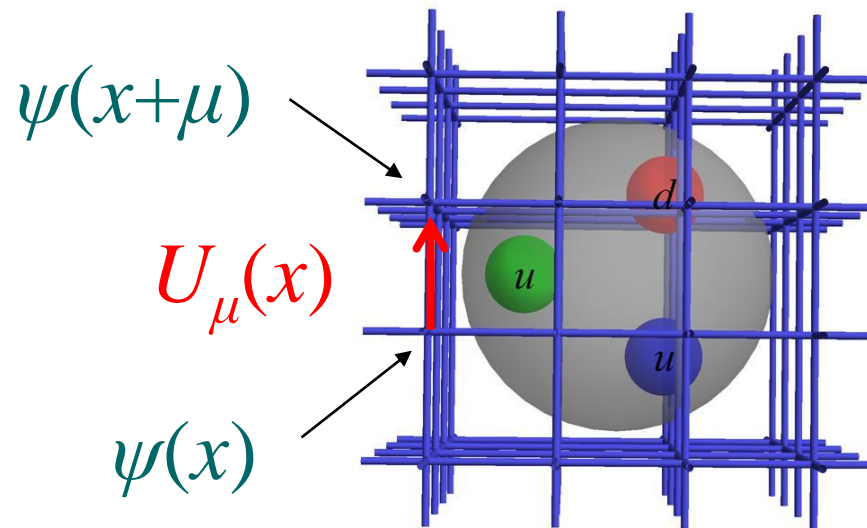
- ◆ Lattice QCD group at JLab
 - ◆ main physics directions in support of (JLab) hadronic physics experimental program
- ◆ Baryon form factors proposals
 - ◆ Strange baryon (transition) form factors
 - ◆ Radically excited transition form factors
 - ◆ Challenge: never been done in lattice calculations before
 - ◆ Starting with first excited state of nucleon: Roper(?)
 - ◆ Experiments at Jefferson Laboratory (**CLAS**), MIT-Bates, LEGS, Mainz, Bonn, GRAAL, and Spring-8
 - ◆ Long term goal: extend calculation to more N^*

Lattice QCD

- ◆ Physical observables are calculated from the path integral

$$\langle 0|O(\bar{\psi}, \psi, A)|0\rangle = \frac{1}{Z} \int [dA][d\bar{\psi}][d\psi] O(\bar{\psi}, \psi, A) e^{i \int d^4x \mathcal{L}^{\text{QCD}}(\bar{\psi}, \psi, A)}$$

- ◆ Strong-coupling regions: expansions no longer converge
- ◆ Lattice QCD is a discrete version of continuum QCD theory



- ◆ Numerical integration to calculate the path integral
- ◆ Take $a \rightarrow 0$ and $V \rightarrow \infty$ in the continuum limit

Lattice Challenge

- ◆ Euclidean space:

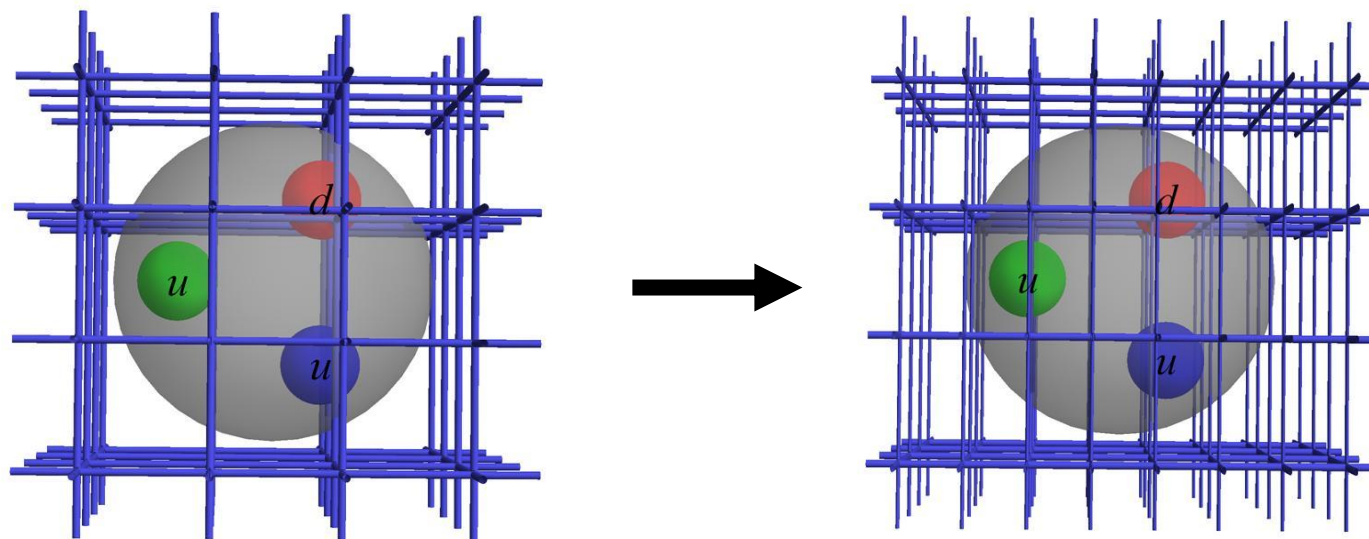
obtain correlators with time-dependent form

$$\sum_n Z_{n,B} e^{-E_n(\vec{P})t}$$

- ◆ Signal falls exponentially with time

dominated by ground state; challenge for excited states

- ◆ Solution: increase resolution

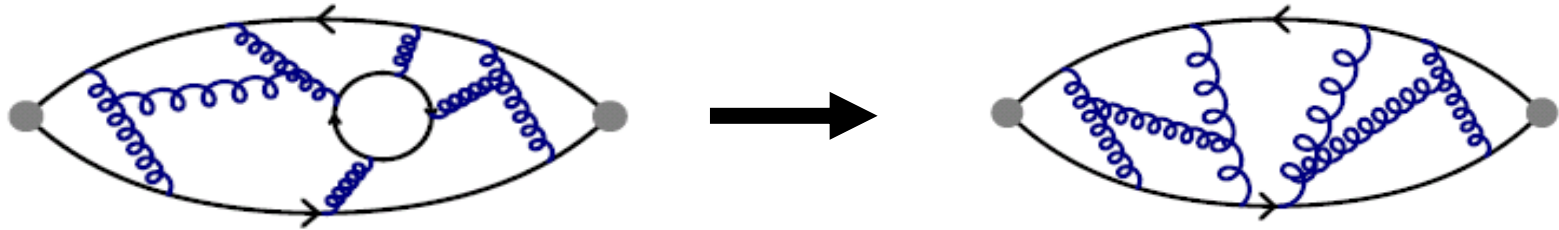


Anisotropic lattice

more complications to generate 2+1f lattice ensembles

Lattice Setup

- ◆ “Quenched” for exploratory study



no sea quark contributions

- ◆ Bad: Uncontrollable systematic error
- ◆ Good?
 - ◆ Preserve nice features:
confining, asymptotically free, spontaneously broken chiral symmetry
 - ◆ Cheap exploratory studies to develop new methods
- ◆ Some detailed Parameters
 - ◆ $16^3 \times 64$ anisotropic lattice, $\xi = 3$
 - ◆ Wilson gauge action + clover fermion action
 - ◆ $a_t^{-1} \approx 6$ GeV and $a_s \approx 0.125$ fm ($L < 2$ fm)
 - ◆ $m_\pi \approx 720$ (480 and 1100) MeV
 - ◆ 200 configurations

Green Functions

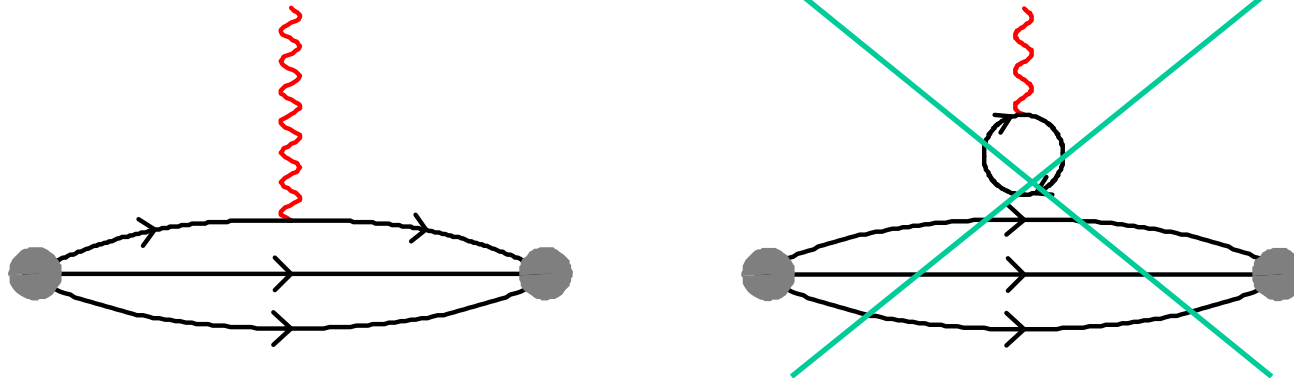
- ◆ Three-point function with interpolation operator J

$$C_{3\text{pt}}^{\Gamma, \mathcal{O}}(\vec{p}, t, \tau) = \sum_{\alpha, \beta} \Gamma^{\alpha, \beta} \langle J_{\beta}(\vec{p}, t) \mathcal{O}(\tau) \bar{J}_{\alpha}(\vec{p}, 0) \rangle$$

- ◆ Baryon interpolating field

$$J_{\alpha}(\vec{p}, t) = \sum_{\vec{x}, a, b, c} e^{i\vec{p}\cdot\vec{x}} \epsilon^{abc} [u_a^T(y_1, t) C \gamma_5 d_b(y_2, t)] u_{c, \alpha}(y_3, t) \phi(y_1 - x) \phi(y_2 - x) \phi(y_3 - x)$$

- ◆ Two contraction categories:



- ◆ We use only the “connected” construction for this work
- ◆ Ongoing investigation into “disconnected” contribution

Form Factors

- ◆ The form factors are buried in the amplitudes

$$\begin{aligned} \Gamma_{\mu,AB}^{(3),T}(t_i, t, t_f, \vec{p}_i, \vec{p}_f) &= a^3 \sum_n \sum_{n'} \frac{1}{Z_j} \frac{Z_{n',B}(p_f) Z_{n,A}(p_i)}{4E'_n(\vec{p}_f) E_n(\vec{p}_i)} e^{-(t_f-t)E'_n(\vec{p}_f)} e^{-(t-t_i)E_n(\vec{p}_i)} \\ &\times \sum_{s,s'} T_{\alpha\beta} u_{n'}(\vec{p}_f, s')_{\beta} \langle N_{n'}(\vec{p}_f, s') | j_{\mu}(0) | N_n(\vec{p}_i, s) \rangle \bar{u}_n(\vec{p}_i, s)_{\alpha} \end{aligned}$$

- ◆ Nucleon form factor ($n = n' = 0$)

$$\langle N | V_{\mu} | N \rangle(q) = \bar{u}_N(p') \left[\gamma_{\mu} F_1(q^2) + \sigma_{\mu\nu} q_{\nu} \frac{F_2(q^2)}{2m} \right] u_N(p) e^{-iq \cdot x}$$

- ◆ Nucleon-Roper form factor ($n = 0, n' = 1$ or $n = 1, n' = 0$)

$$\langle N_2 | V_{\mu} | N_1 \rangle_{\mu}(q) = \bar{u}_{N_2}(p') \left[F_1(q^2) \left(\gamma_{\mu} - \frac{q_{\mu}}{q^2} \not{q} \right) + \sigma_{\mu\nu} q_{\nu} \frac{F_2(q^2)}{M_{N_1} + M_{N_2}} \right] u_{N_1}(p) e^{-iq \cdot x}$$

- ◆ Need best possible input from two-point correlators

Variational Method

◆ Generalized eigenvalue problem:

[C. Michael, Nucl. Phys. B 259, 58 (1985)]

[M. Lüscher and U. Wolff, Nucl. Phys. B 339, 222 (1990)]

◆ Construct the matrix

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t)^\dagger \mathcal{O}_j(0) | 0 \rangle$$

◆ Solve for the generalized eigensystem of

$$C(t_0)^{-1/2} C(t) C(t_0)^{-1/2} v = \lambda(t, t_0) v$$

with eigenvalues

$$\lambda_n(t, t_0) = e^{-(t-t_0)E_n} (1 + \mathcal{O}(e^{-|\delta E|(t-t_0)}))$$

Now the original correlator matrix can be approximated by

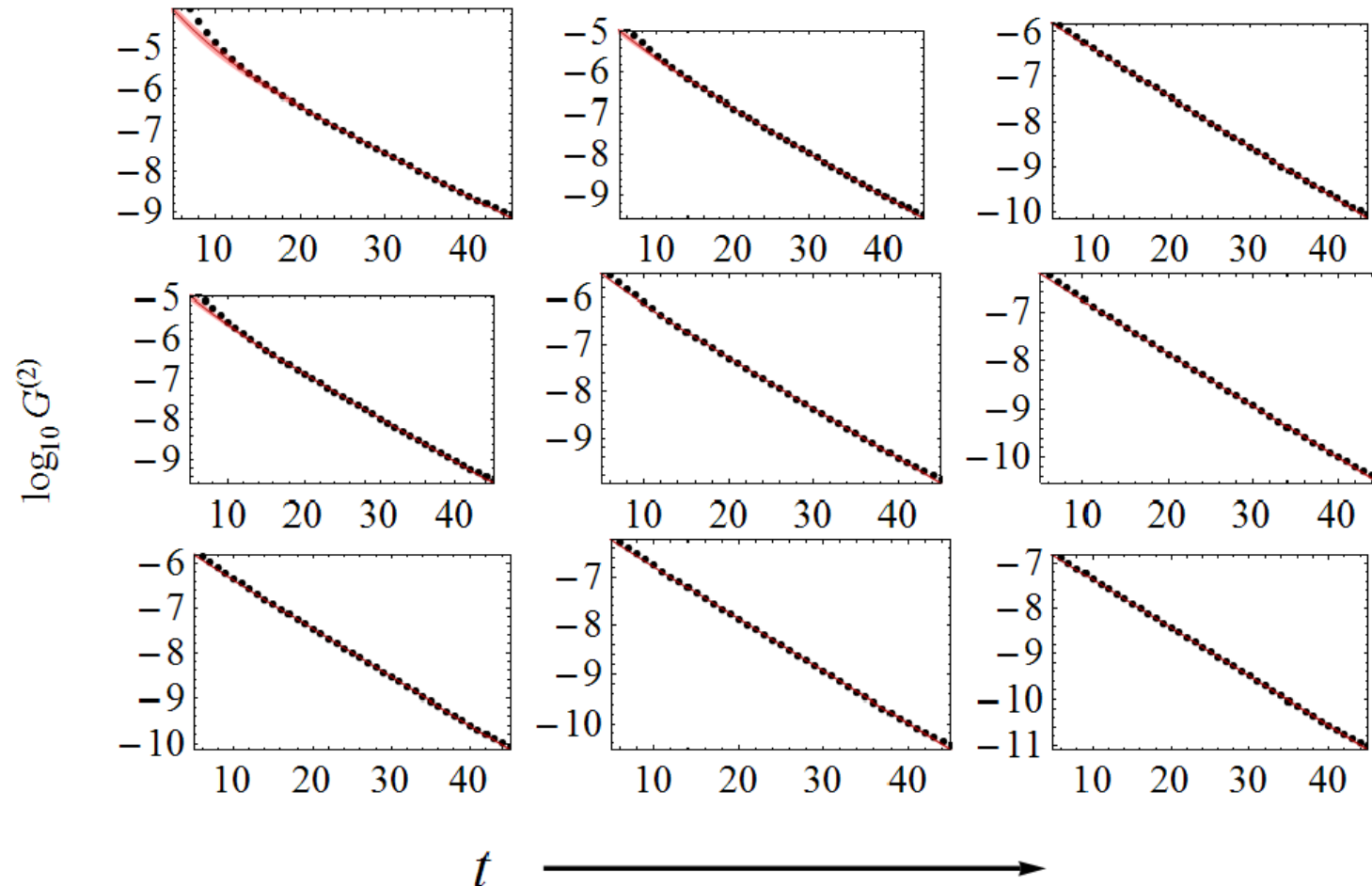
$$C_{ij} = \sum_{n=1}^r (C(t_0)^{1/2} v_n^*)_i (v_n C(t_0)^{1/2})_j \lambda_n(t, t_0) = \sum_n \frac{E_n + m}{2E_n} Z_{i,n} Z_{j,n} e^{-E_n t}$$

◆ Three smearings (i, j) are chosen for this work

- ◆ 2nd excited state is contaminated by remaining states

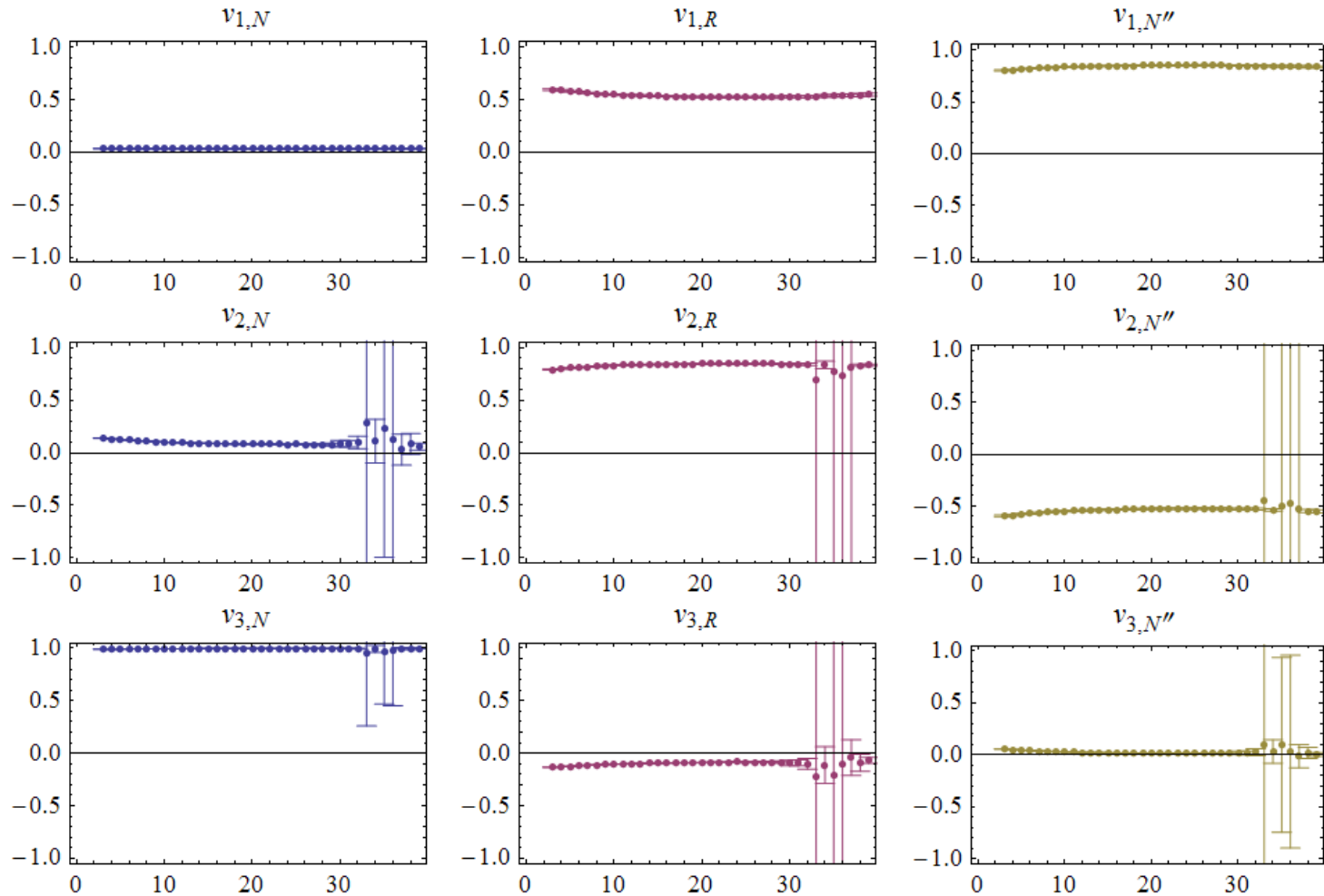
Variational Method

- ◆ Reconstruct two-point correlators from Z and λ



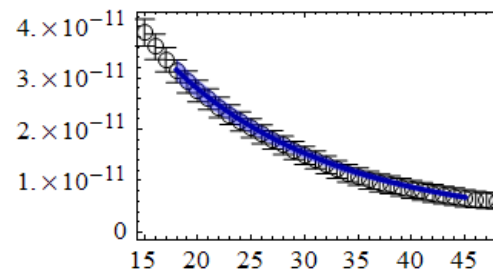
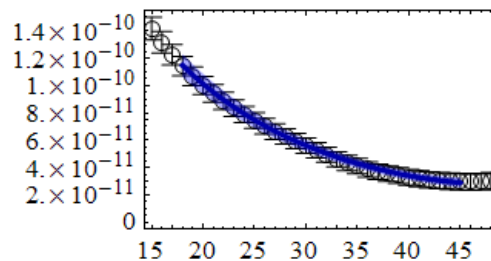
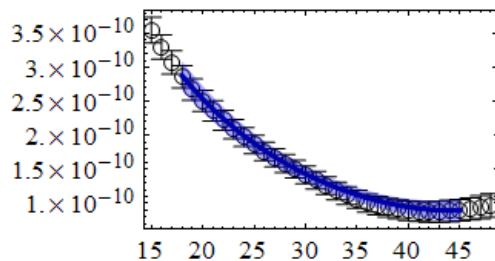
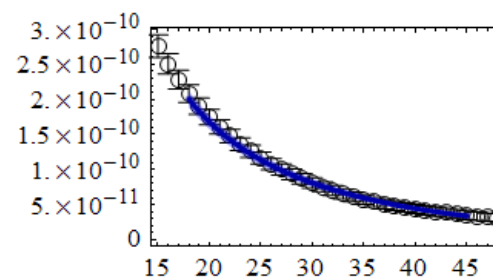
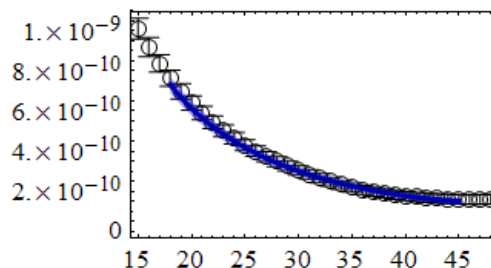
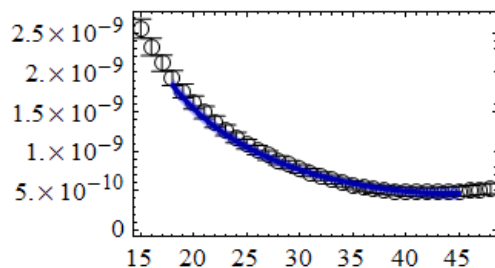
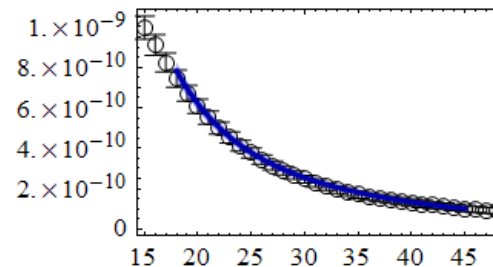
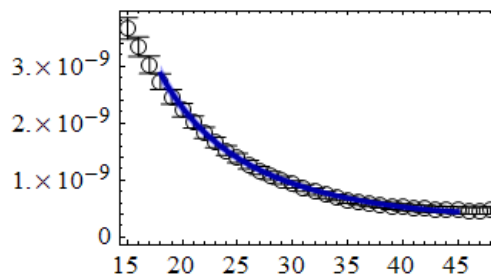
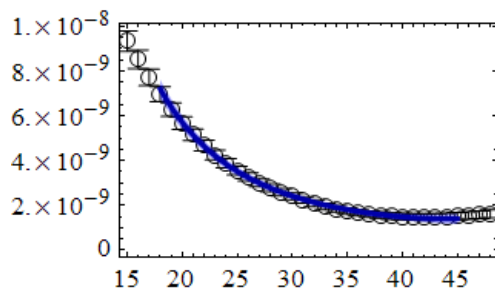
Variational Method

- ◆ Eigenvectors (at $p = 0$) show overlap of smearings with states



Three-Point Fitting

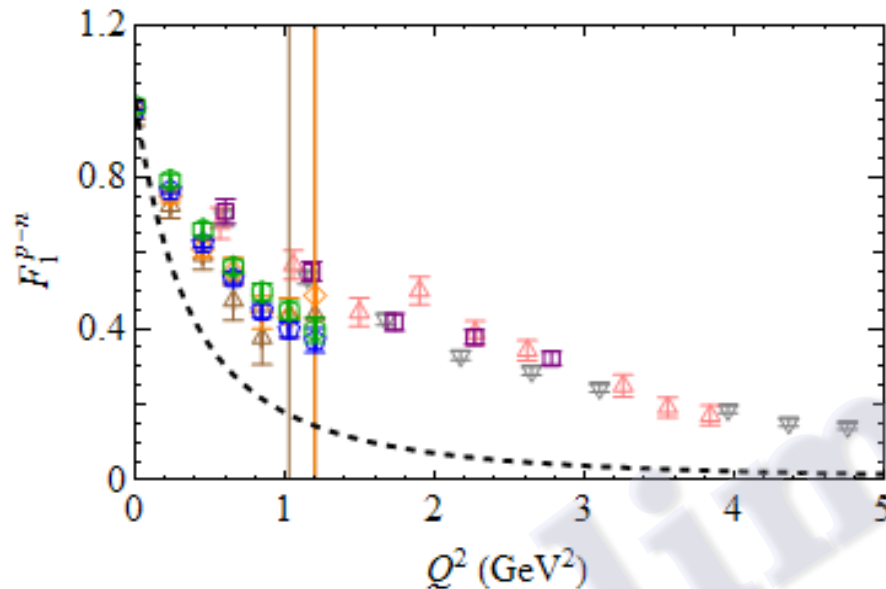
◆ Example: $P_f = \{0,0,0\}$, $P_i = \{0,1,1\}$, V_4



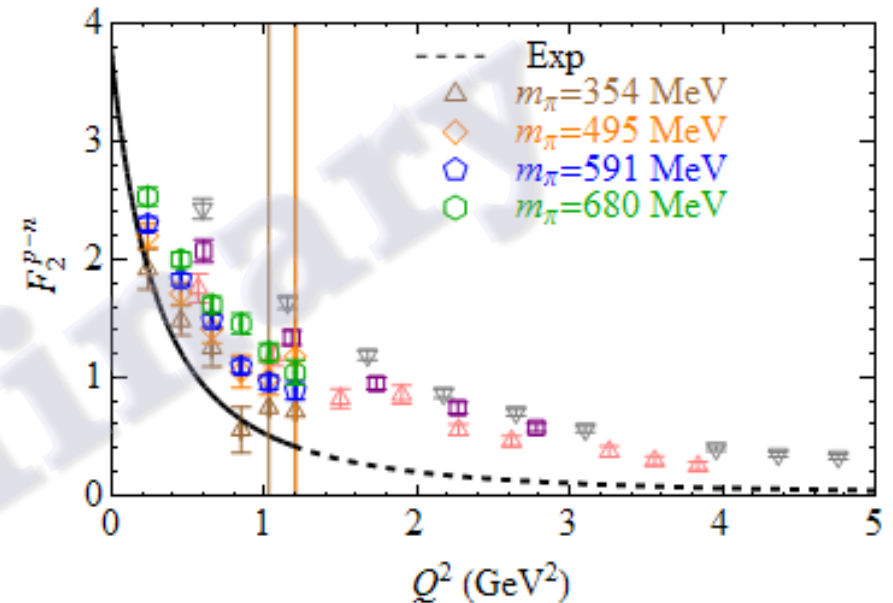
Nucleon Form Factors

- ◆ Pion masses around 480, 720 and 1100 MeV

Isovector F_1



Isovector F_2

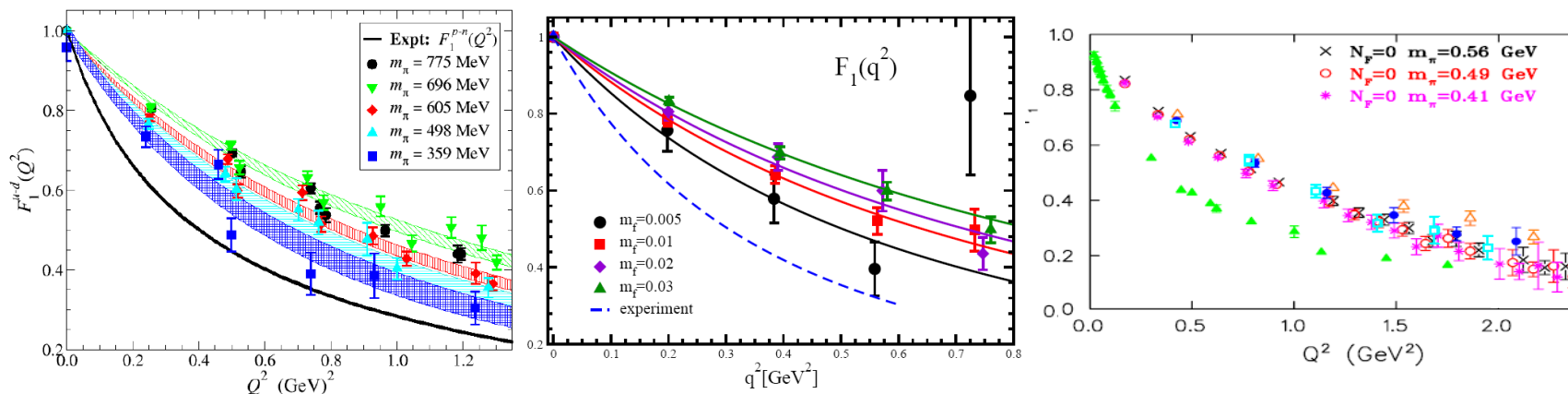


- ◆ Compare with $N_f = 2+1$ mixed action (DWF+ asqtad) calculation with conventional approach
- ◆ **Clean signal for momentum region of $3 < Q^2 < 5$ GeV²**

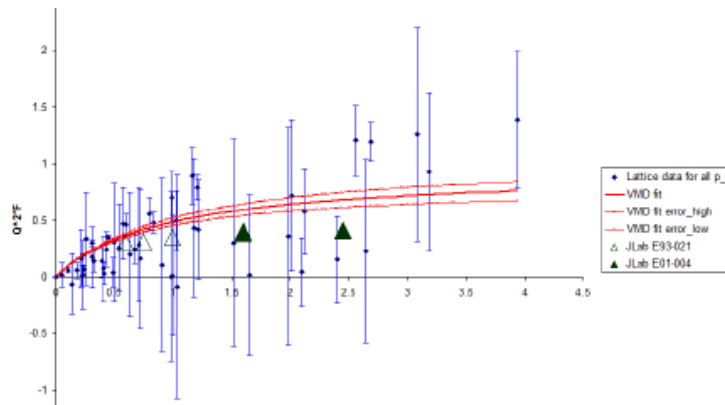
Lattice Form Factors

◆ Large- Q^2 calculations

◆ Typical Q^2 range for nucleon form factors is $< 3.0 \text{ GeV}^2$



◆ Higher- Q^2 calculations suffer from poor noise-to-signal ratios even for pion



Nucleon-Roper Form Factors

- ◆ Converting experimental data ($\gamma^* N_1 \rightarrow N_2$)

$$A_{1/2}(Q^2)/\kappa_A = G_M(Q^2) = F_1^*(Q^2) + F_2^*(Q^2)$$

$$S_{1/2}(Q^2)/\kappa_S = G_E(Q^2) = F_1^*(Q^2) - F_2^*(Q^2) Q^2 / (M_{N_1} + M_{N_2})^2$$

with

$$k_A(Q^2) \equiv \sqrt{2\pi\alpha} \frac{Q^2 + (M_{N_1} - M_{N_2})^2}{M_{N_1} (M_{N_1}^2 - M_{N_2}^2)}$$

$$k_S(Q^2) \equiv k_A(Q^2) \frac{M_{N_1} + M_{N_2}}{2\sqrt{2}Q^2 M_{N_2}} \sqrt{Q^2 + (M_{N_1} - M_{N_2})^2} \\ \times \sqrt{Q^2 + (M_{N_1} + M_{N_2})^2}$$

- ◆ Use CLAS analysis and PDG results to solve for $F_{1,2}^*(Q^2)$

I. Aznauryan et al., arXiv:0804.0447[nucl-ex], arXiv:0711.1120[nucl-th];

K. Park et al., Phys. Rev. C77, 015208 (2008);

V. I. Mokeev et. al, AIP Conf. Proc. 842, 339 (2006).

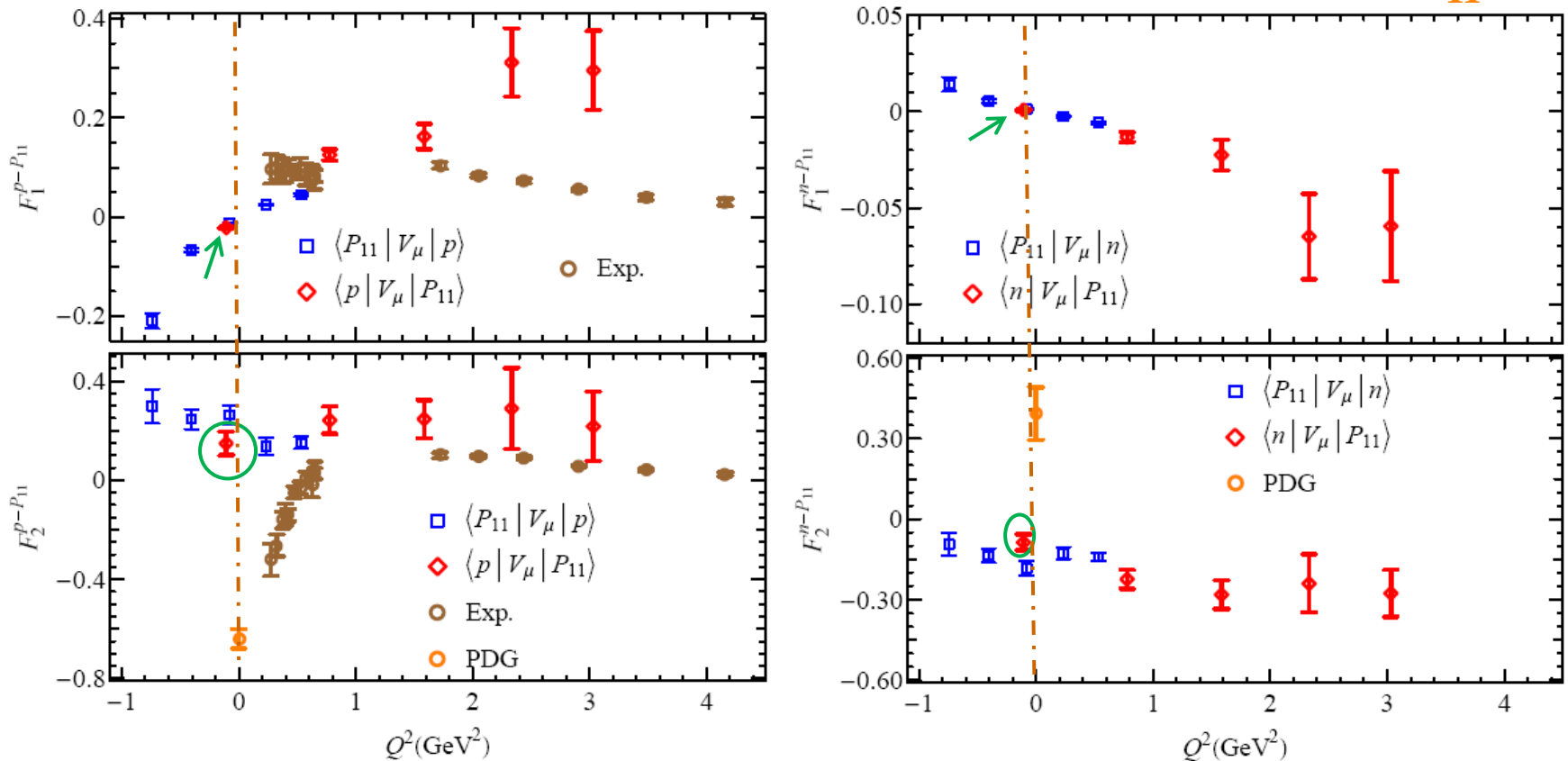
Nucleon-Roper Form Factors

- Completed exploratory study on quenched lattices [arXiv:0803.3020](https://arxiv.org/abs/0803.3020)

Proton- P_{11}

720 MeV Pion

Neutron- P_{11}



- Possible decaying state (circled above)
- 200 configurations give us reasonable signal
- Lower pion mass will shift the time-like region to space-like region

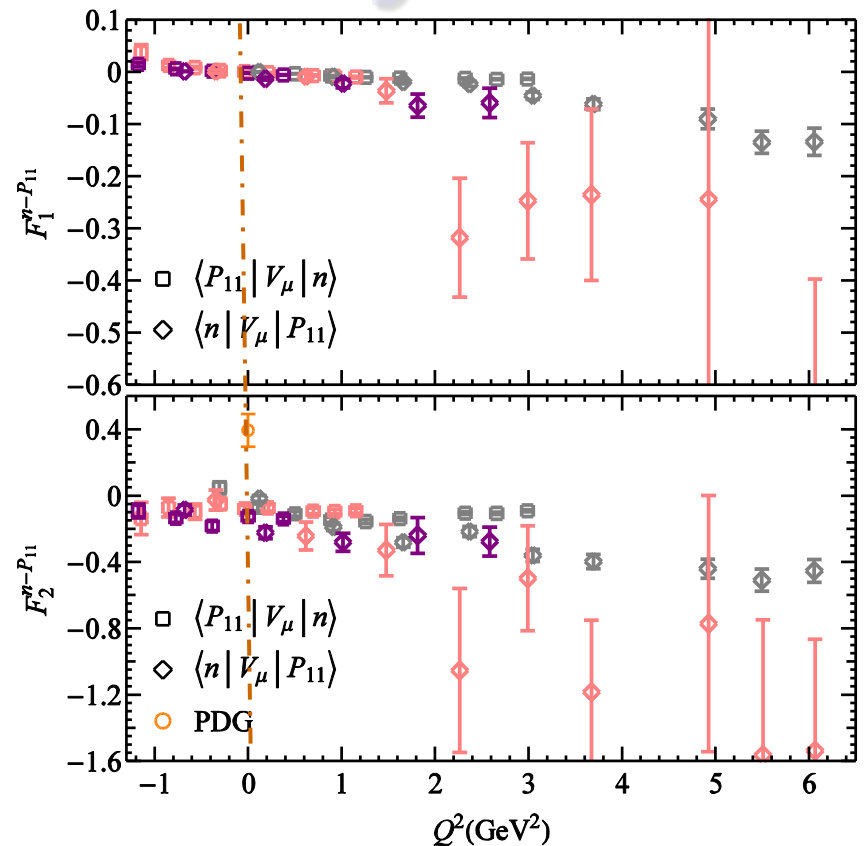
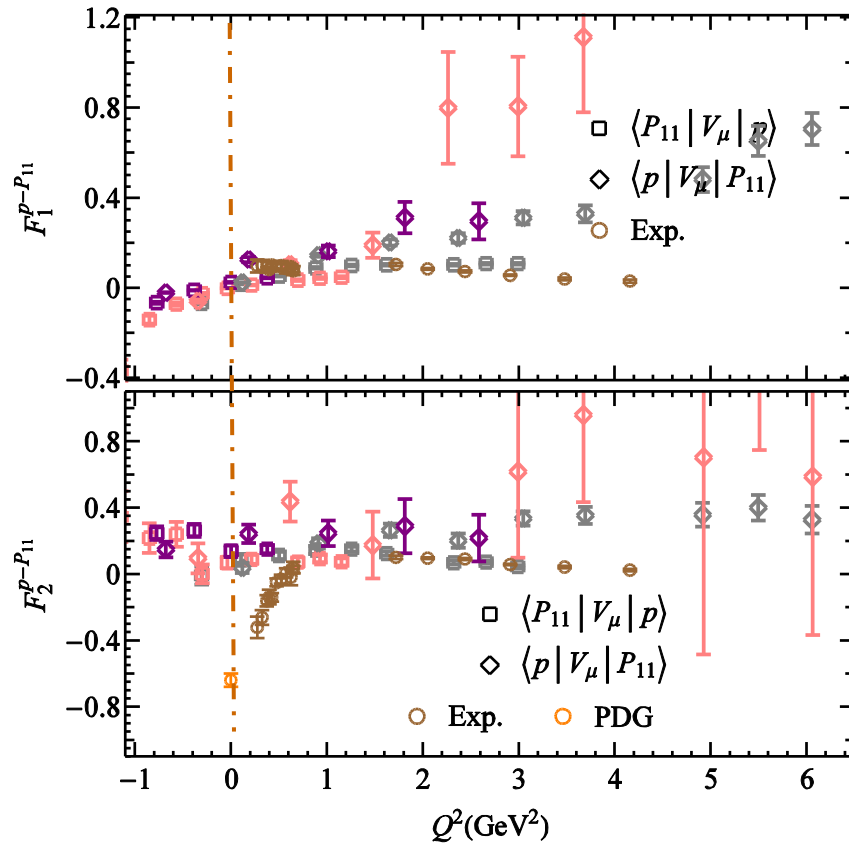
Nucleon-Roper Form Factors

- ◆ Add two more mass points at $m_\pi \sim 480$ and 1100 MeV

Proton- P_{11}

Neutron- P_{11}

Preliminary



- ◆ Need to remove the points with potential decay kinematics
- ◆ May want to calculate on a second volume

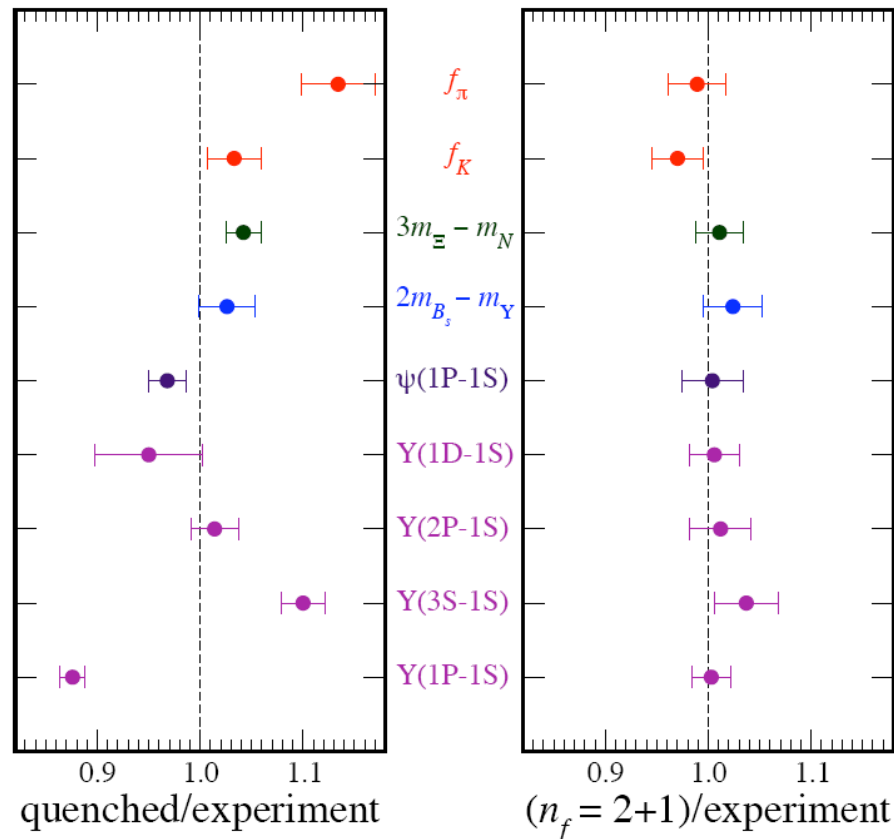
Full-QCD Ensembles

Why Dynamical?

Lattice QCD spectrum

- ◆ Successfully calculates many ground states (Nature, ...)

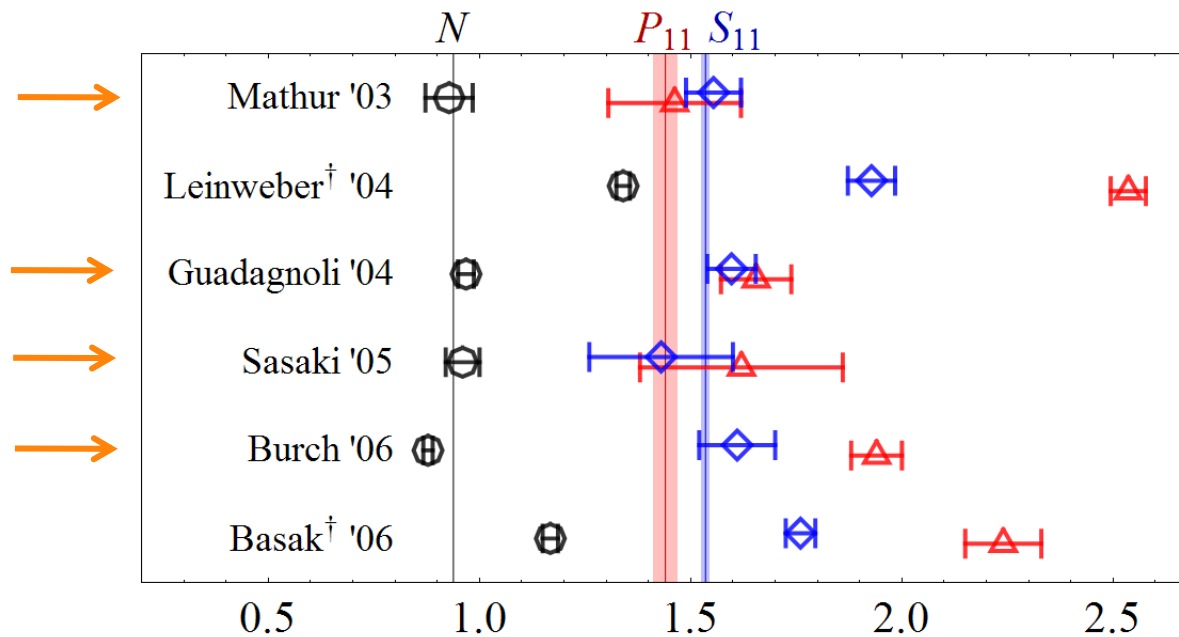
HPQCD



- ◆ Predictions: B_c mass, D and D_s decay constants, $D \rightarrow Kl\nu$ form factors

Roper Resonance on the Lattice

◆ Mostly done in “quenched” approx. N , P_{11} , S_{11} spectrum

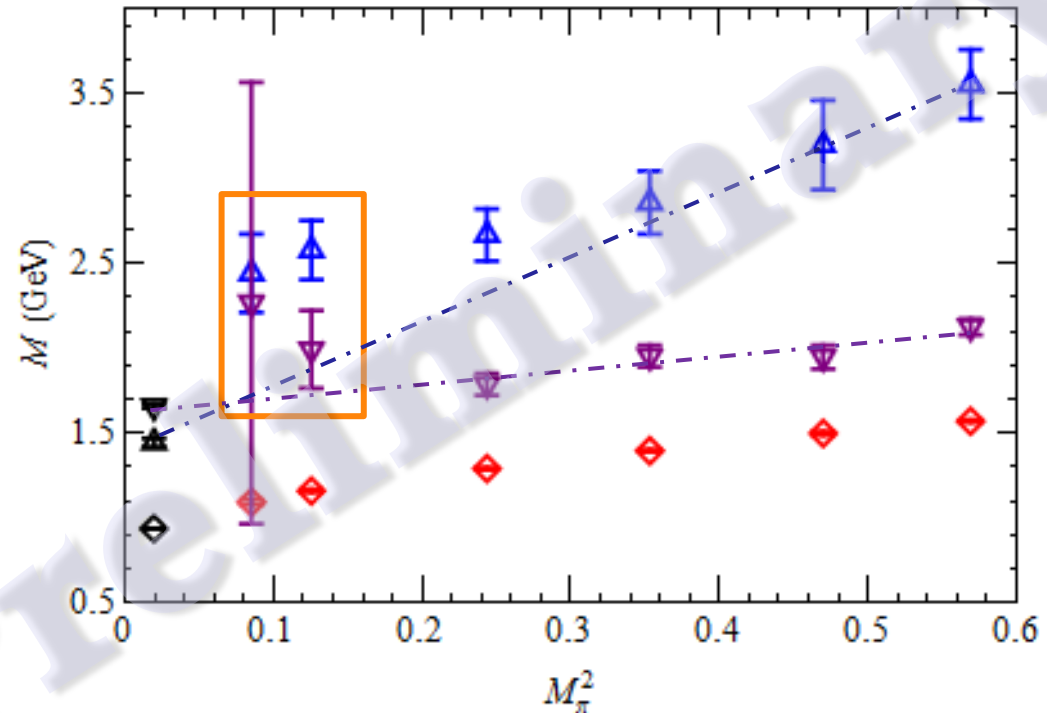


Group	N_f	S_f	a_t^{-1} (GeV)	M_π (GeV)	L (fm)	Method	Extrapolation
Basak et al. [12]	0	Wilson	6.05	0.49	2.35	VM	N/A
Burch et al. [11]	0	CIDO	1.68, 1.35	0.35–1.1	2.4	VM	$a + bm_\pi^2$
Sasaki et al. [9]	0	Wilson	2.1	0.61–1.22	1.5, 3.0	MEM	$\sqrt{a + bm_\pi^2}$
Guadagnoli et al. [7]	0	Clover [13]	2.55	0.51–1.08	1.85	SBBM	$a + bm_\pi^2 + cm_\pi^4$
Leinweber et al. [8]	0	FLIC	1.6	0.50–0.91	2.0	VM	N/A
Mathur et al. [6]	0	Overlap [14]	1.0	0.18–0.87	2.4, 3.2	CCF	$a + bm_\pi + cm_\pi^2$

Roper in Full QCD

- ◆ $N_f = 2+1$ mixed action (DWF+ asqtad) calculation ($L \sim 2.5$ fm)
- ◆ Symbols: J^P

◆ $1/2^+$ N
◆ $1/2^-$ S_{11}
◆ $1/2^+$ P_{11}

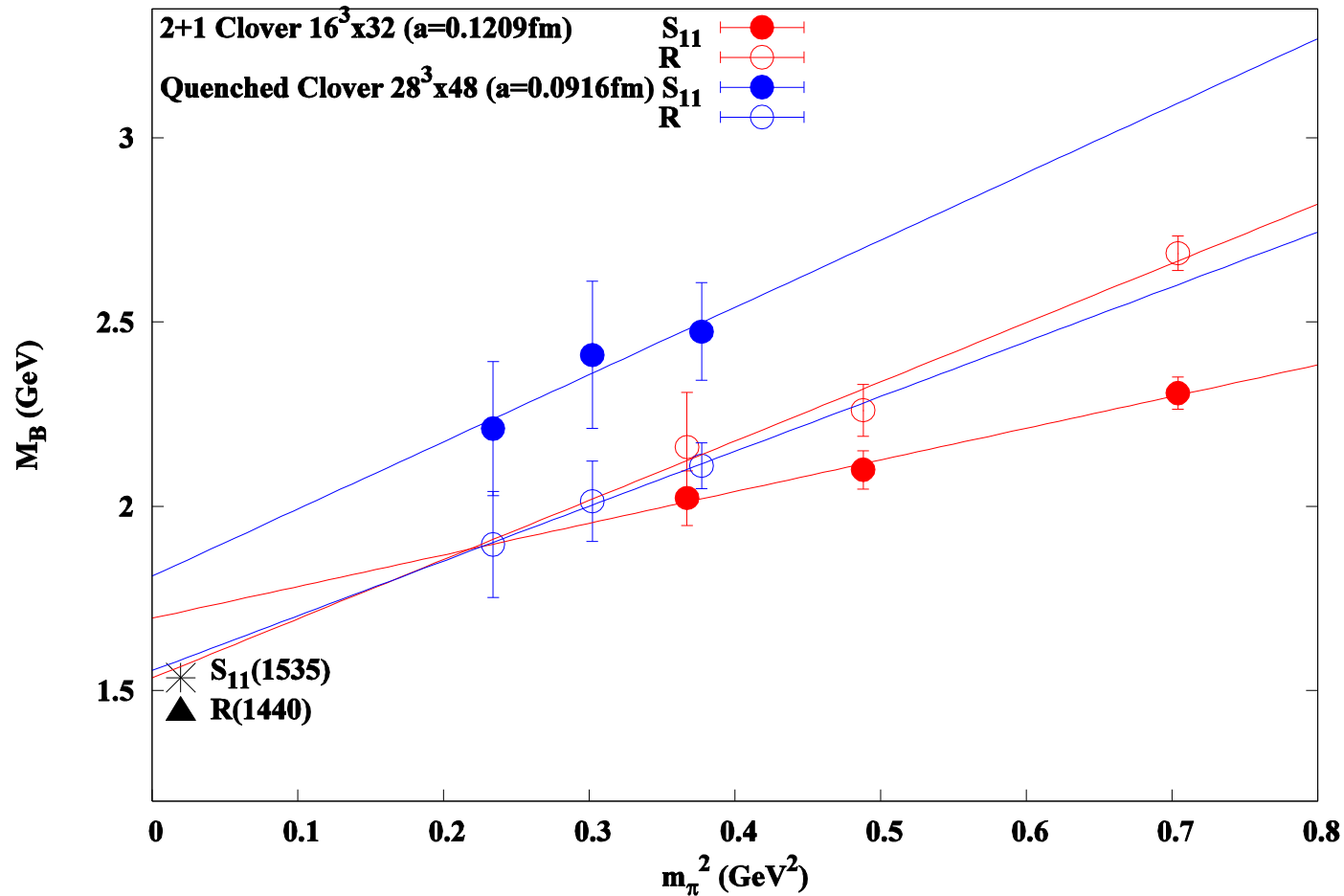


- ◆ Finite-volume effects starting at 350 MeV pion?
- ◆ Not a crazy possibility (see the hand-drawn extrapolation lines)
- ◆ Stay tuned for future $N_f = 2+1$ lattice calculations

Roper in Full QCD

◆ $N_f = 2+1$ isotropic clover action calculation ($L \sim 2$ fm)

[K.F. Liu, LAT2008]



Dynamical Anisotropic Lattices

◆ $N_f = 2+1$ anisotropic clover action available ensemble list

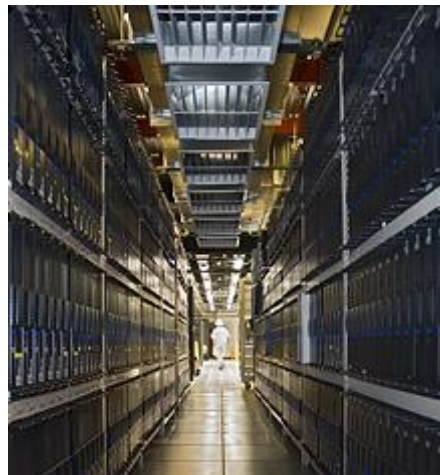
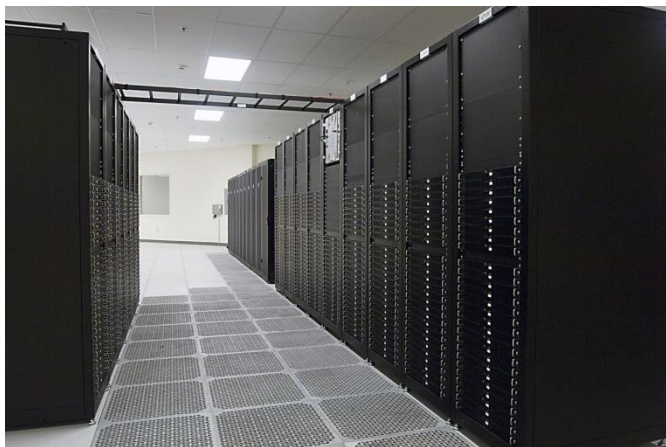
[R. Edwards et al., Phys. Rev. D 78, 014505 (2008)]

[R. Edwards and M. Peardon, LAT2008]

L_x	L_t	m_l	m_s	L (fm)	m_π L	m_π (MeV)
12	96	-0.0540	-0.0540	1.44	11.6	~ 1600
12	96	-0.0699	-0.0540	1.44	8.3	
12	96	-0.0794	-0.0540	1.44	5.9	
12	96	-0.0826	-0.0540	1.44	9.6	
16	96	-0.0826	-0.0540	1.92	6.3	~ 660
12	96	-0.0618	-0.0618	1.44	9.7	~ 1340
16	128	-0.0743	-0.0743	1.92	8.1	~ 850
16	128	-0.0808	-0.0743	1.92	5.6	
16	128	-0.0830	-0.0743	1.92	4.5	
16	128	-0.0840	-0.0743	1.92	3.8	
24	128	-0.0840	-0.0743	2.88	5.7	~ 360

Computational Resources

USQCD facilities: JLab, Fermilab, BNL



Non-lattice resources open to USQCD: ORNL, LLNL, ANL



NSF supercomputer and world-wide increase in computation facilities

Gauge Generation

- ◆ Most of the major 2+1-flavor gauge ensembles:

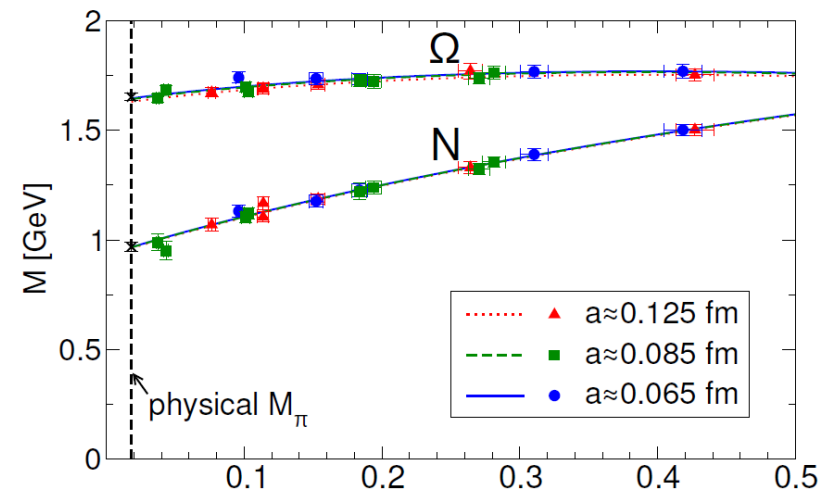
$M_\pi < 300 \text{ MeV}$ (Most of them have multiple lattice spacings and volumes)

- ◆ MILC (staggered): $M_\pi \sim 217 \text{ MeV}$
- ◆ PACS-CS (Clover action): $M_\pi \sim 156 \text{ MeV}$ (but small volume)
- ◆ The Budapest-Marseille-Wuppertal (BMW) Collaboration: $M_\pi \sim 193 \text{ MeV}$
- ◆ RBC/UKQCD (DWF), $M_\pi \sim 210 \text{ MeV}$ (on-going)
- ◆ The Hadron Spectrum Collaboration (anisotropic clover):
 $M_\pi \sim 175 \text{ MeV}$ (on-going)

- ◆ Small range of chiral extrapolation:

significantly reduces the
systematic uncertainties

Example: BMW Collaboration,
LAT2008

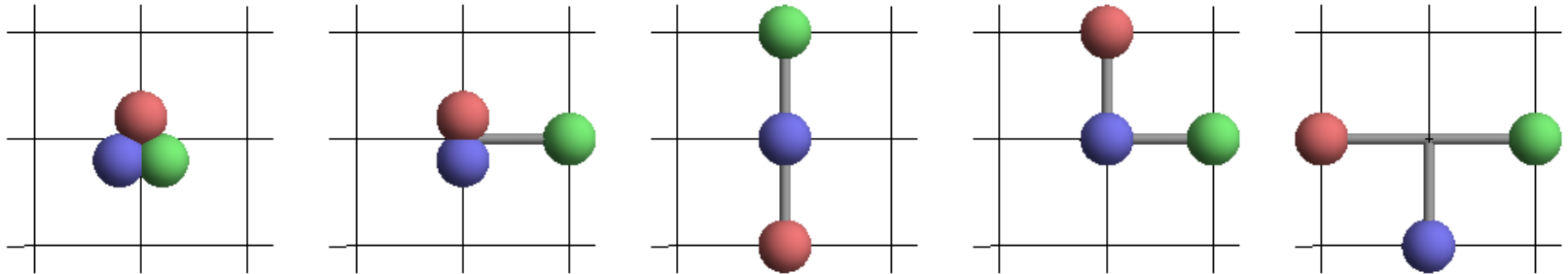


Extended Operators

on-going project

Orthogonal Operators

◆ Baryon field $\Phi_{\alpha\beta\gamma,ijk}^{ABC}(x) = \epsilon_{abc}[\tilde{D}_i^{(3)}\tilde{\psi}]_{Aa\alpha}(x)[\tilde{D}_j^{(3)}\tilde{\psi}]_{Bb\beta}(x)[\tilde{D}_k^{(3)}\tilde{\psi}]_{Cc\gamma}(x)$



D. Richards, this workshop

- ◆ Classify states according to symmetry properties
- ◆ Projection onto irreducible representations of finite groups
- ◆ Number of operators:

N^+ Operator type	G_{1g}	H_g	G_{2g}
Single-Site	3	1	0
Singly-Displaced	24	32	8
Doubly-Displaced-I	24	32	8
Doubly-Displaced-L	64	128	64
Triply-Displaced-T	64	128	64
Total	179	321	144

S. Basak et al., Phys. Rev. D72, 094506 (2005)

Summary and Outlook

Lattice QCD calculations of N - P_{11} form factors...

- ◆ Test case is in a small “quenched” box with large pion mass
- ◆ We demonstrate a method to determine N - N^* form factors with reasonable signal on anisotropic lattices
- ◆ Clean signal for large Q^2 momentum N - N form factors

Further along our roadmap...

- ◆ $g_{\pi NN^*}$ from axial coupling and Goldberger-Treiman relation
- ◆ Moving forward to full-QCD anisotropic lattice calculations
- ◆ Implement group theory operators for baryons
- ◆ Other N - N^* form factors. The methodology developed can be applied to many other excited-nucleon form factors.

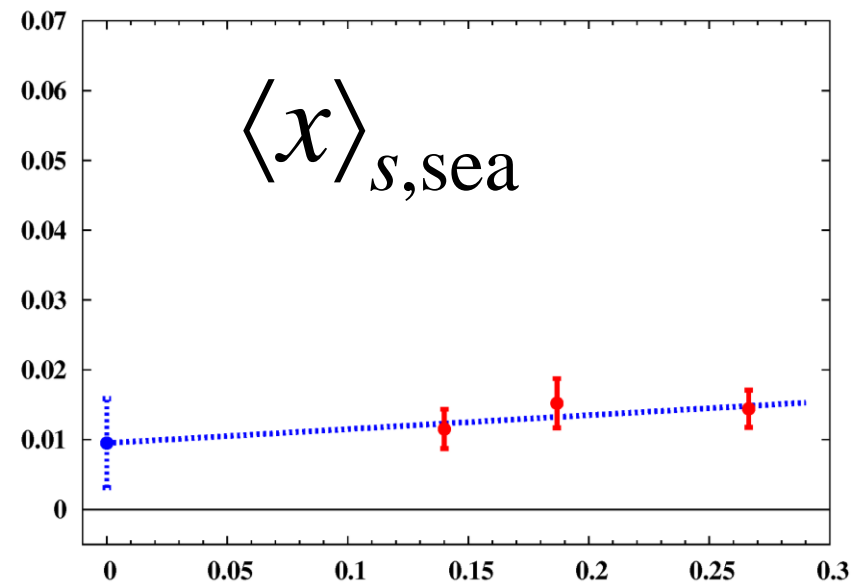
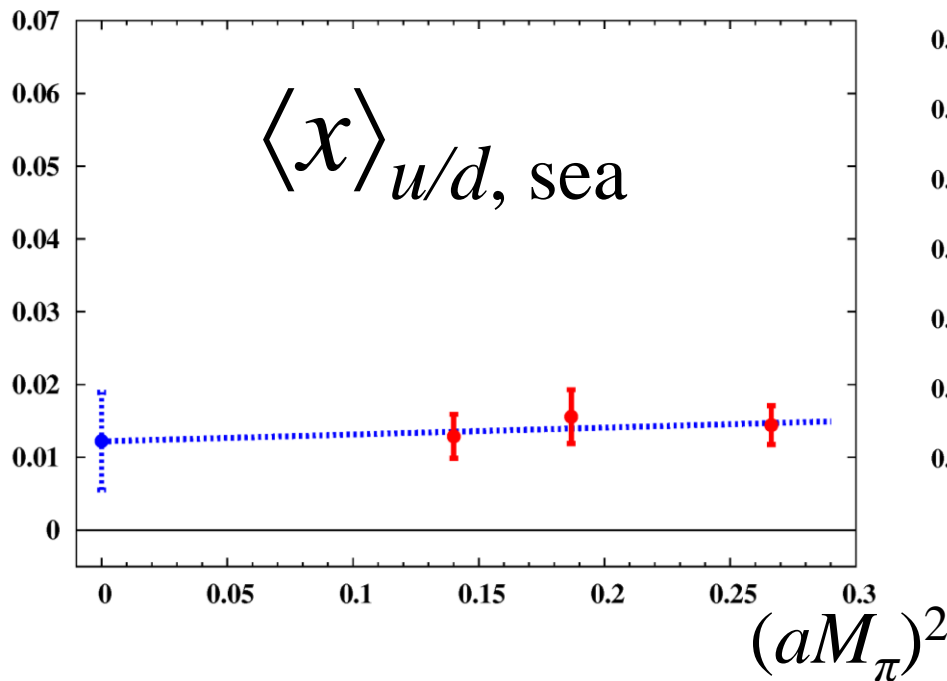
Backup Slides

Including Disconnected Diagrams

- ◆ Example: Sea-quark contribution to (unrenormalized) quark momentum fraction

- ◆ 2+1f clover lattices, $M_\pi \sim 610\text{--}840$ MeV

T. Doi (χ QCD), Lattice 2008



Challenges for the Future

As one goes to lighter pion-mass regions...

- ◆ Need more statistics to get sufficient single-to-noise ratio

$$\begin{aligned} \frac{\text{Signal}}{\text{Noise}} &= \frac{\langle J(t)J(0) \rangle}{\frac{1}{\sqrt{N}} \sqrt{\langle |J(t)J(0)|^2 \rangle - \langle J(t)J(0) \rangle^2}} \\ &\sim \frac{Ae^{-Mnt}}{\frac{1}{\sqrt{N}} \sqrt{Be^{-3m_\pi t} - Ce^{-2Mnt}}} \\ &\sim \sqrt{N} D e^{-(M_n - \frac{3}{2}m_\pi)t} \end{aligned}$$

Challenges for the Future

As one goes to lighter pion-mass regions...

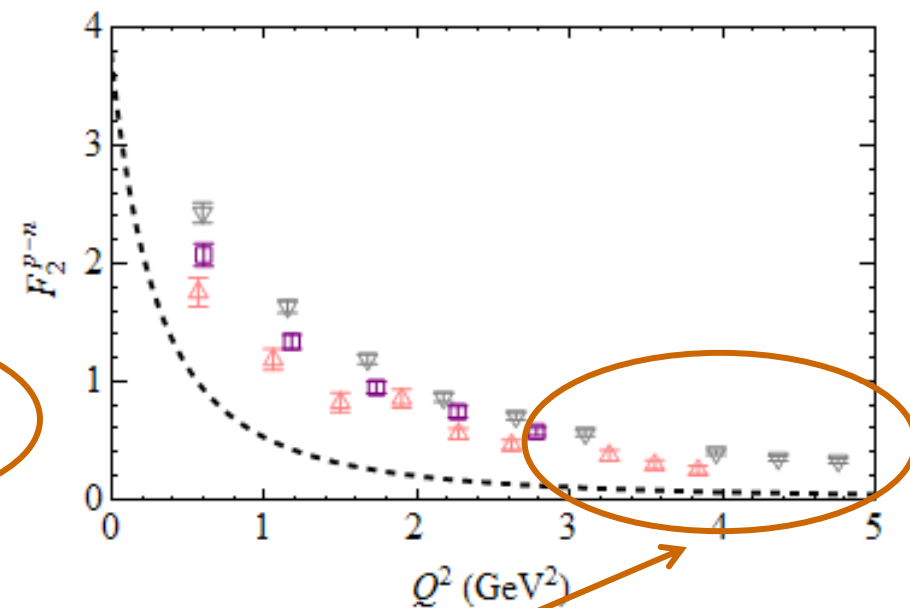
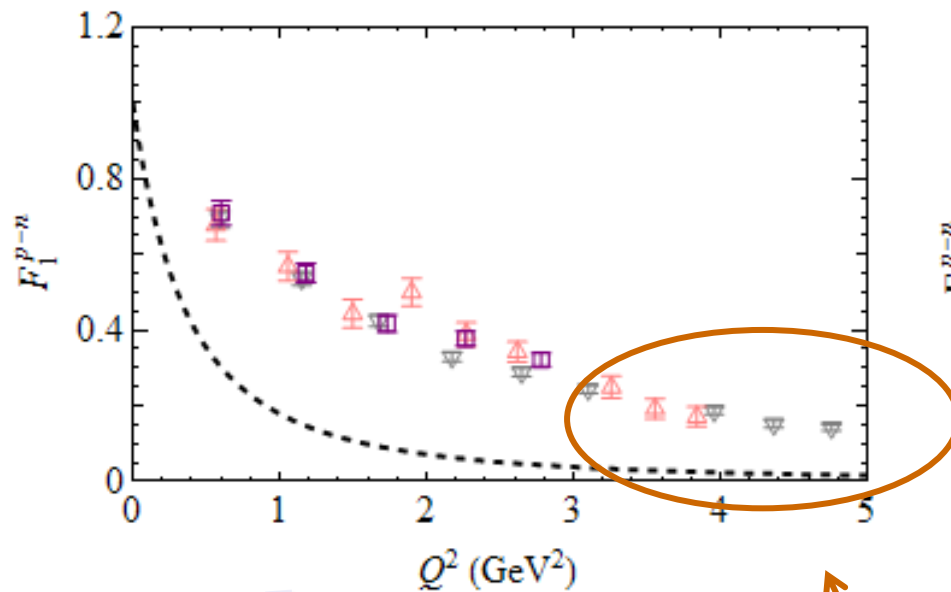
- ◆ Need more statistics to get sufficient single-to-noise ratio
- ◆ Decay channels open up...
 - ◆ Conservative approach: using multiple volumes to identify single- and multiple-particle states
- ◆ Further improvement:
 - ◆ Symmetry-breaking on the lattice:
 - ◆ construct more operators that would generate the same quantum numbers in $a = 0$ world
 - ◆ More data input
 - ◆ Better discrimination among different states

Nucleon Form Factors

- ◆ Pion masses around 480, 720 and 1100 MeV

Isovector F_1

Isovector F_2



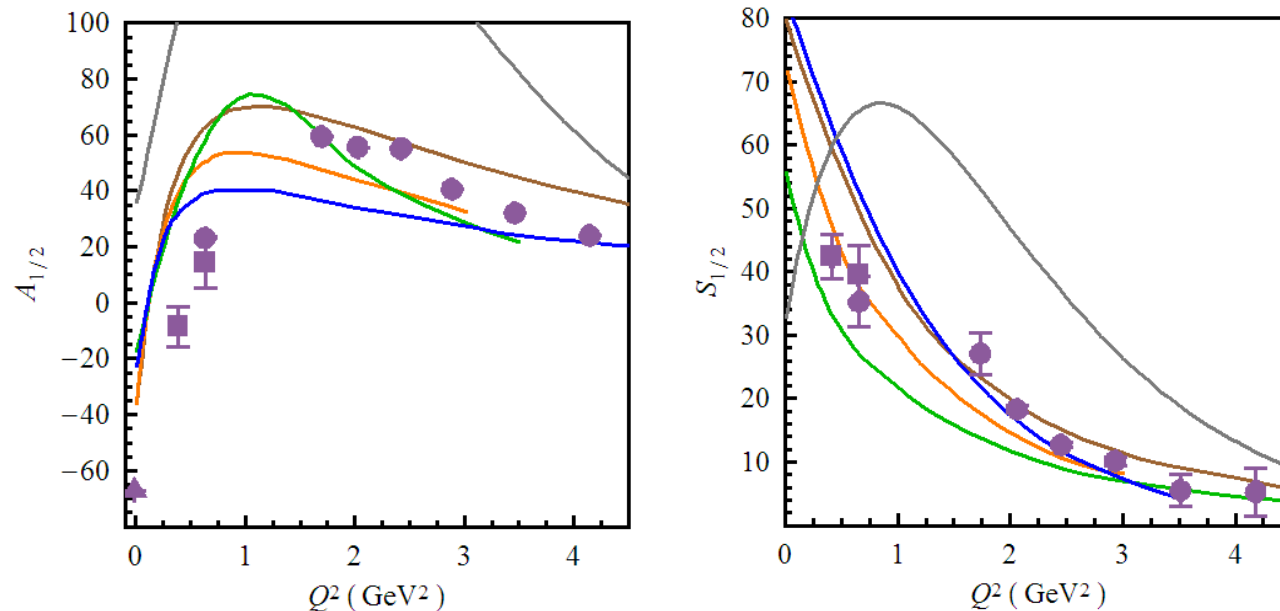
Preliminary

Clean signal for momentum region of $3 < Q^2 < 5$ GeV²

- ◆ Working on larger- Q^2 regions with non-zero p_f
- ◆ $N_f = 3$ and 2+1 data are on the way

$N-P_{11}$ Form Factor

- ◆ Experiments at Jefferson Laboratory (**CLAS**), MIT-Bates, LEGS, Mainz, Bonn, GRAAL, and Spring-8
- ◆ **Helicity amplitudes** are measured (in $10^{-3} \text{ GeV}^{-1/2}$ units)
- ◆ Many models disagree (a selection are shown below)



- ◆ If the Roper is the first radially excited state of the nucleon, this is the data to compare with