

Extraction of Total Amplitudes from Complete Photoproduction Experiments

Frank Tabakin

Department of Physics and Astronomy,
University of Pittsburgh, Pittsburgh, PA 15260

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- **Observables & Helicity Amplitudes**
- Observables & Transversity Amplitudes
- Bilinear Form and SU(4) Γ Matrices
- Discrete Ambiguities
- Fierzing The Observables
- Fierz Theorem Consequences
- Sample case: $\pi - Nucleon$
- BDS versus CT Theorems
- Limitations
- Conclusions

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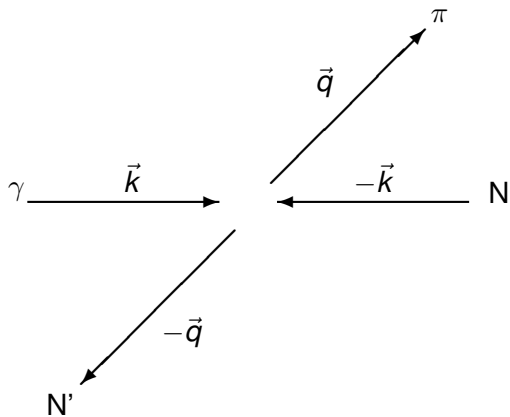
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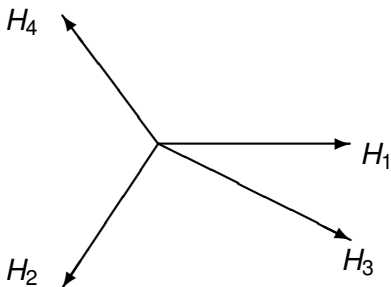
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$$\sigma(\theta, \phi) = \frac{q}{k} \mathcal{I}(\theta, \phi)$$



The Four Complex Amplitudes In Helicity Space

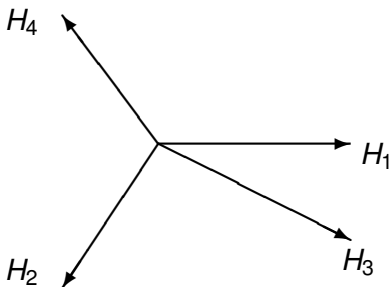


$$H_i(E, \theta, \phi) = |H_i| e^{i\phi_i}$$

$$4 + 3 = 7$$

Discrete ambiguities 7 + ?

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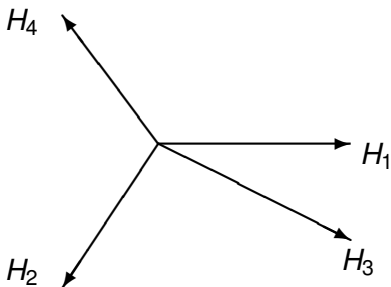


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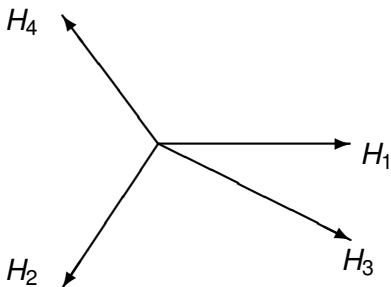


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$$H_i(E, \theta, \phi) = |H_i| e^{i\phi_i}$$

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$$\begin{aligned}\check{\Omega}^\alpha &= \Omega^\alpha \mathcal{I}(\theta) \\ &= \frac{1}{2} H_i^* \Gamma_{ij}^\alpha H_j \\ &\equiv \frac{1}{2} \langle H | \Gamma^\alpha | H \rangle ,\end{aligned}$$

$$\alpha = 1, \dots \quad 16$$

Spin Observables: Sixteen spin observables are expressed in helicity representation and BHP forms.

Classified into four sets:

- S for the differential cross section and single spin observables,
- BT beam-target,
- BR beam-recoil,
- TR target-recoil .

Helicity Representation \mathcal{S}

Spin Observable	Helicity Representation	BHP
$\check{\Omega}^1 \equiv \mathcal{I}(\theta)$	$\frac{1}{2}(H_1 ^2 + H_2 ^2 + H_3 ^2 + H_4 ^2)$	$\frac{1}{2}\langle H \Gamma^1 H\rangle$
$\check{\Omega}^4 \equiv \check{\Sigma}$	$\text{Re}(-H_1H_4^* + H_2H_3^*)$	$\frac{1}{2}\langle H \Gamma^4 H\rangle$
$\check{\Omega}^{10} \equiv -\check{T}$	$\text{Im}(H_1H_2^* + H_3H_4^*)$	$\frac{1}{2}\langle H \Gamma^{10} H\rangle$
$\check{\Omega}^{12} \equiv \check{P}$	$\text{Im}(-H_1H_3^* - H_2H_4^*)$	$\frac{1}{2}\langle H \Gamma^{12} H\rangle$

Spin Observable	Helicity Representation	BHP
$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(H_1 H_4^* - H_3 H_2^*)$	$\frac{1}{2} \langle H \Gamma^3 H \rangle$
$\check{\Omega}^5 \equiv \check{H}$	$\text{Im}(-H_2 H_4^* + H_1 H_3^*)$	$\frac{1}{2} \langle H \Gamma^5 H \rangle$
$\check{\Omega}^9 \equiv \check{E}$	$\frac{1}{2} (H_1 ^2 - H_2 ^2 + H_3 ^2 - H_4 ^2)$	$\frac{1}{2} \langle H \Gamma^9 H \rangle$
$\check{\Omega}^{11} \equiv \check{F}$	$\text{Re}(-H_2 H_1^* - H_4 H_3^*)$	$\frac{1}{2} \langle H \Gamma^{11} H \rangle$

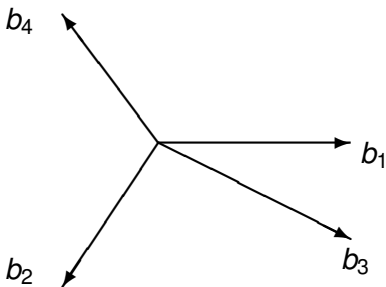
Helicity Representation \mathcal{BR}

Spin Observable	Helicity Representation	BHP
$\check{\Omega}^{14} \equiv \check{O}_x$	$\text{Im}(-H_2 H_1^* + H_4 H_3^*)$	$\frac{1}{2} \langle H \Gamma^{14} H \rangle$
$\check{\Omega}^7 \equiv -\check{O}_z$	$\text{Im}(H_1 H_4^* - H_2 H_3^*)$	$\frac{1}{2} \langle H \Gamma^7 H \rangle$
$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Re}(H_2 H_4^* + H_1 H_3^*)$	$\frac{1}{2} \langle H \Gamma^{16} H \rangle$
$\check{\Omega}^2 \equiv -\check{C}_z$	$\frac{1}{2} (H_1 ^2 + H_2 ^2 - H_3 ^2 - H_4 ^2)$	$\frac{1}{2} \langle H \Gamma^2 H \rangle$

Helicity Representation \mathcal{TR}

Spin Observable	Helicity Representation	BHP
$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-H_1 H_4^* - H_2 H_3^*)$	$\frac{1}{2} \langle H \Gamma^6 H \rangle$
$\check{\Omega}^{13} \equiv -\check{T}_z$	$\text{Re}(-H_1 H_2^* + H_4 H_3^*)$	$\frac{1}{2} \langle H \Gamma^{13} H \rangle$
$\check{\Omega}^8 \equiv \check{L}_x$	$\text{Re}(H_2 H_4^* - H_1 H_3^*)$	$\frac{1}{2} \langle H \Gamma^8 H \rangle$
$\check{\Omega}^{15} \equiv \check{L}_z$	$\frac{1}{2}(- H_1 ^2 + H_2 ^2 + H_3 ^2 - H_4 ^2)$	$\frac{1}{2} \langle H \Gamma^{15} H \rangle$

The Four Complex Amplitudes In Transversity Space

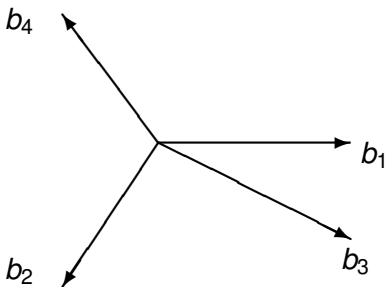


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$$4 + 3 = 7$$

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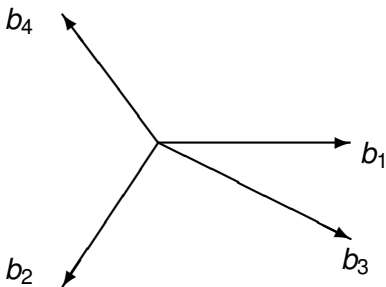


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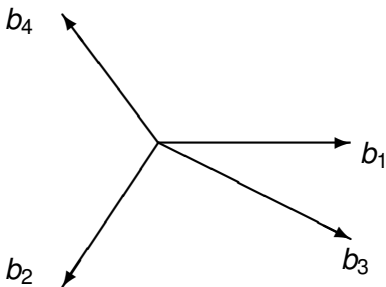


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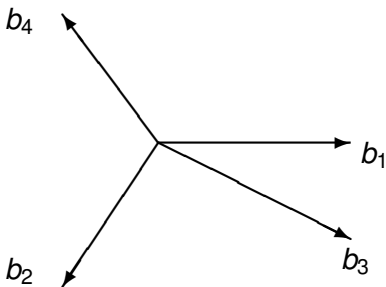


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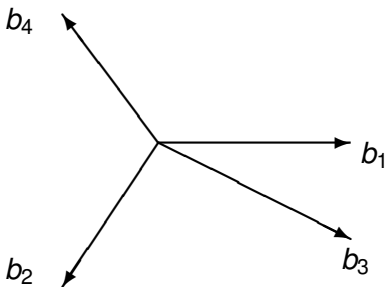


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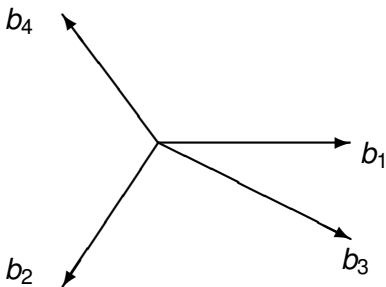


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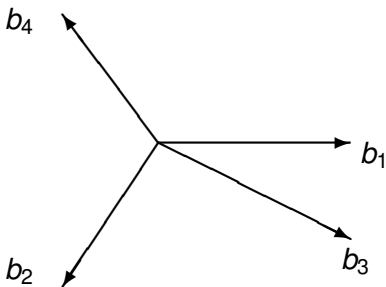


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The Four Complex Amplitudes In Transversity Space



$$b_i(E, \theta, \phi) = |b_i| e^{i\phi_i}$$

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Transversity Representation \mathcal{S}

Spin Observable	Transversity Representation	BTP
$\check{\Omega}^1 \equiv \mathcal{I}(\theta)$	$\frac{1}{2}(b_1 ^2 + b_2 ^2 + b_3 ^2 + b_4 ^2)$	$\frac{1}{2}\langle b \tilde{\Gamma}^1 b\rangle$
$\check{\Omega}^4 \equiv \check{\Sigma}$	$\frac{1}{2}(b_1 ^2 + b_2 ^2 - b_3 ^2 - b_4 ^2)$	$\frac{1}{2}\langle b \tilde{\Gamma}^4 b\rangle$
$\check{\Omega}^{10} \equiv -\check{T}$	$\frac{1}{2}(- b_1 ^2 + b_2 ^2 + b_3 ^2 - b_4 ^2)$	$\frac{1}{2}\langle b \tilde{\Gamma}^{10} b\rangle$
$\check{\Omega}^{12} \equiv \check{P}$	$\frac{1}{2}(- b_1 ^2 + b_2 ^2 - b_3 ^2 + b_4 ^2)$	$\frac{1}{2}\langle b \tilde{\Gamma}^{12} b\rangle$

Spin Observable	Transversity Representation	BTP
$\check{\Omega}^3 \equiv \check{G}$	$\text{Im}(-b_1 b_3^* - b_2 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^3 b \rangle$
$\check{\Omega}^5 \equiv \check{H}$	$\text{Re}(b_1 b_3^* - b_2 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^5 b \rangle$
$\check{\Omega}^9 \equiv \check{E}$	$\text{Re}(b_1 b_3^* + b_2 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^9 b \rangle$
$\check{\Omega}^{11} \equiv \check{F}$	$\text{Im}(b_1 b_3^* - b_2 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^{11} b \rangle$

Spin Observable	Transversity Representation	BTP
$\check{\Omega}^{14} \equiv \check{O}_x$	$\text{Re}(-b_1 b_4^* + b_2 b_3^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^{14} b \rangle$
$\check{\Omega}^7 \equiv -\check{O}_z$	$\text{Im}(-b_1 b_4^* - b_2 b_3^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^7 b \rangle$
$\check{\Omega}^{16} \equiv -\check{C}_x$	$\text{Im}(b_1 b_4^* - b_2 b_3^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^{16} b \rangle$
$\check{\Omega}^2 \equiv -\check{C}_z$	$\text{Re}(b_1 b_4^* + b_2 b_3^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^2 b \rangle$

Spin Observable	Transversity Representation	BTP
$\check{\Omega}^6 \equiv -\check{T}_x$	$\text{Re}(-b_1 b_2^* + b_3 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^6 b \rangle$
$\check{\Omega}^{13} \equiv -\check{T}_z$	$\text{Im}(b_1 b_2^* - b_3 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^{13} b \rangle$
$\check{\Omega}^8 \equiv \check{L}_x$	$\text{Im}(-b_1 b_2^* - b_3 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^8 b \rangle$
$\check{\Omega}^{15} \equiv \check{L}_z$	$\text{Re}(-b_1 b_2^* - b_3 b_4^*)$	$\frac{1}{2} \langle b \tilde{\Gamma}^{15} b \rangle$

The Bilinear Helicity Product (BHP) Summary

$$\begin{array}{ll} S : & (\mathcal{I}, \check{\Sigma}, -\check{T}, \check{P}) & \Gamma^1, \Gamma^4, \Gamma^{10}, \Gamma^{12} \\ BT : & (\check{G}, \check{H}, \check{E}, \check{F}) & \Gamma^3, \Gamma^5, \Gamma^9, \Gamma^{11} \\ BR : & (\check{O}_x, -\check{O}_z, -\check{C}_x, -\check{C}_z) & \Gamma^{14}, \Gamma^7, \Gamma^{16}, \Gamma^2 \\ TR : & (-\check{T}_x, -\check{T}_z, \check{L}_x, \check{L}_z) & \Gamma^6, \Gamma^{13}, \Gamma^8, \Gamma^{15} \end{array}$$

$$\begin{array}{l} \text{Ambiguities:} \\ \Gamma^4, \Gamma^{10}, \Gamma^{12} \\ \Gamma^{15} K_0, \\ \Gamma^6 K_0, \Gamma^8 K_0, \Gamma^{13} K_0 \end{array}$$

Sixteen Hermitian Gamma Matrices

$$\Gamma^{\alpha=1\dots5} = 1, \gamma^0, i\vec{\gamma}$$

$$\Gamma^{\alpha=6\dots11} = \sigma^{0x}, i\sigma^{0y}, i\sigma^{0z}, \\ i\sigma^{xy}, i\sigma^{xz}, i\sigma^{zy}$$

$$\Gamma^{\alpha=12\dots16} = i\gamma^5\gamma^0, \gamma^5\vec{\gamma}, \gamma^5.$$

- Γ^α are Hermitian and unitary.
- $\text{Tr}(\Gamma^\alpha \Gamma^\beta) = 4\delta_{\alpha\beta}$.
- Γ^α are linearly independent. and form a complete set (a basis) for 4×4 matrices. Any 4×4 matrices X can be expanded as $X = \sum_\alpha C_\alpha U^\alpha$ with $C_\alpha = \frac{1}{4}\text{Tr}(\Gamma^\alpha X)$.
- $\sum_\alpha \Gamma_{ba}^\alpha \Gamma_{st}^\alpha = 4\delta_{as}\delta_{bt}$.
- $\Gamma^\alpha \Gamma^\beta = \rho_{\alpha\beta\gamma} \Gamma^\gamma$ with $\rho_{\alpha\beta\gamma} = \frac{1}{4}\text{Tr}(\Gamma^\alpha \Gamma^\beta \Gamma^\gamma)$.
- $\frac{1}{4}\rho_{\alpha\gamma\delta}\rho_{\beta\gamma\eta} = \frac{1}{16}\text{Tr}(\Gamma^\delta \Gamma^\alpha \Gamma^\eta \Gamma^\beta) \equiv C_{\delta\eta}^{\alpha\beta}$: used for the Fierz transformation

Transversity transformation: $|b\rangle = U^4 |H\rangle \quad \tilde{\Gamma}^\alpha = U^4 \Gamma^\alpha U^{\dagger 4}$

$$U^{(4)} = \frac{1}{2} \begin{pmatrix} 1 & -i & i & 1 \\ 1 & i & -i & 1 \\ 1 & i & i & -1 \\ 1 & -i & -i & -1 \end{pmatrix},$$

which involves rotating the helicity quantization axis to the direction normal to the scattering plane. The sixteen spin observables can be expressed in this *transversity basis* by

$$\check{\Omega}^\alpha = \Omega^\alpha \mathcal{I}(\theta) = \frac{1}{2} b_i^* \tilde{\Gamma}_{ij}^\alpha b_j = \frac{1}{2} \langle b | \tilde{\Gamma}^\alpha | b \rangle, \quad \alpha = 1, \dots, 16.$$

The transversity $\tilde{\Gamma}$ matrices form four classes with four members in each class according to their “shape:” diagonal (D); right parallelogram (PR); antidiagonal (AD); and left parallelogram (PL) correspond to S , BT , BR and TR type experiments.

$$\tilde{\Gamma}_D = \tilde{\Gamma}_S = \begin{bmatrix} a & 0 & 0 & 0 \\ 0 & b & 0 & 0 \\ 0 & 0 & c & 0 \\ 0 & 0 & 0 & d \end{bmatrix}; \quad \begin{array}{l} \tilde{\Gamma}_1 \\ \tilde{\Gamma}_4 \\ \tilde{\Gamma}_{10} \\ \tilde{\Gamma}_{12} \end{array} \begin{array}{cccc} a & b & c & d \\ +1 & +1 & +1 & +1 \\ +1 & +1 & -1 & -1 \\ -1 & +1 & +1 & -1 \\ -1 & +1 & -1 & +1 \end{array}$$

$$|b_i|$$

$$\tilde{\Gamma}_{PR} = \tilde{\Gamma}_{BT} = \begin{bmatrix} 0 & 0 & a & 0 \\ 0 & 0 & 0 & b \\ c & 0 & 0 & 0 \\ 0 & d & 0 & 0 \end{bmatrix};$$

	<i>a</i>	<i>b</i>	<i>c</i>	<i>d</i>
$\tilde{\Gamma}_3$	$-i$	$-i$	$+i$	$+i$
$\tilde{\Gamma}_5$	$+1$	-1	$+1$	-1
$\tilde{\Gamma}_9$	$+1$	$+1$	$+1$	$+1$
$\tilde{\Gamma}_{11}$	$+i$	$-i$	$-i$	$+i$

ϕ_{13}, ϕ_{24}

$$\tilde{\Gamma}_{AD} = \tilde{\Gamma}_{BR} = \begin{bmatrix} 0 & 0 & 0 & a \\ 0 & 0 & b & 0 \\ 0 & c & 0 & 0 \\ d & 0 & 0 & 0 \end{bmatrix};$$

	<i>a</i>	<i>b</i>	<i>c</i>	<i>d</i>
$\tilde{\Gamma}_{14}$	-1	+1	+1	-1
$\tilde{\Gamma}_7$	- <i>i</i>	- <i>i</i>	+ <i>i</i>	+ <i>i</i>
$\tilde{\Gamma}_{16}$	+ <i>i</i>	- <i>i</i>	+ <i>i</i>	- <i>i</i>
$\tilde{\Gamma}_2$	+1	+1	+1	+1

ϕ_{14}, ϕ_{23}

$$\tilde{\Gamma}_{PL} = \tilde{\Gamma}_{TR} = \begin{bmatrix} 0 & a & 0 & 0 \\ b & 0 & 0 & 0 \\ 0 & 0 & 0 & c \\ 0 & 0 & d & 0 \end{bmatrix};$$

	<i>a</i>	<i>b</i>	<i>c</i>	<i>d</i>
$\tilde{\Gamma}_6$	-1	-1	+1	+1
$\tilde{\Gamma}_{13}$	+i	-i	-i	+i
$\tilde{\Gamma}_8$	-i	+i	-i	+i
$\tilde{\Gamma}_{15}$	-1	-1	-1	-1

ϕ_{12}, ϕ_{34}

Discrete Ambiguities in Helicity Basis

G. Keaton & R. Workman Phys.Rev. C53, 1434 (1996) gave discrete ambiguity relations associated with transformations of helicity amplitudes:

Ambiguity I: $H_1 \longleftrightarrow H_4$ $H_2 \longleftrightarrow -H_3$

$$\begin{bmatrix} H'_1 \\ H'_2 \\ H'_3 \\ H'_4 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 & +1 \\ 0 & 0 & -1 & 0 \\ 0 & -1 & 0 & 0 \\ +1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} H_1 \\ H_2 \\ H_3 \\ H_4 \end{bmatrix} = \Gamma^4 \begin{bmatrix} H_1 \\ H_2 \\ H_3 \\ H_4 \end{bmatrix}.$$

Ambiguity II:

$$H_1 \longrightarrow H_2 \quad H_2 \longrightarrow -H_1 \quad H_3 \longrightarrow H_4 \quad H_4 \longrightarrow -H_3$$

$$\begin{bmatrix} H'_1 \\ H'_2 \\ H'_3 \\ H'_4 \end{bmatrix} = \begin{bmatrix} 0 & +1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & +1 \\ 0 & 0 & -1 & 0 \end{bmatrix} \begin{bmatrix} H_1 \\ H_2 \\ H_3 \\ H_4 \end{bmatrix} = -i \Gamma^{10} \begin{bmatrix} H_1 \\ H_2 \\ H_3 \\ H_4 \end{bmatrix} .$$

Ambiguity III:

$$H_1 \longrightarrow H_3 \quad H_2 \longrightarrow H_4 \quad H_3 \longrightarrow -H_1 \quad H_4 \longrightarrow -H_2$$

$$\begin{bmatrix} H'_1 \\ H'_2 \\ H'_3 \\ H'_4 \end{bmatrix} = \begin{bmatrix} 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & +1 \\ -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{bmatrix} \begin{bmatrix} H_1 \\ H_2 \\ H_3 \\ H_4 \end{bmatrix} = i \Gamma^{12} \begin{bmatrix} H_1 \\ H_2 \\ H_3 \\ H_4 \end{bmatrix} .$$

Ambiguity IV:

$$H_1 \longrightarrow -H_1^* \quad H_2 \longrightarrow H_2^* \quad H_3 \longrightarrow H_3^* \quad H_4 \longrightarrow -H_4^*$$

$$\begin{bmatrix} H'_1 \\ H'_2 \\ H'_3 \\ H'_4 \end{bmatrix} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} H_1^* \\ H_2^* \\ H_3^* \\ H_4^* \end{bmatrix} = \Gamma^{15} \begin{bmatrix} H_1^* \\ H_2^* \\ H_3^* \\ H_4^* \end{bmatrix} .$$

Ambiguities in the transversity basis

Ambiguity I:

$$b_1 \longrightarrow +b_1 \quad b_2 \longrightarrow +b_2 \quad b_3 \longrightarrow -b_3 \quad b_4 \longrightarrow -b_4$$

$$\begin{bmatrix} b'_1 \\ b'_2 \\ b'_3 \\ b'_4 \end{bmatrix} = \begin{bmatrix} +1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \\ b_4 \end{bmatrix} = \tilde{\Gamma}^4 \begin{bmatrix} b_1 \\ b_2 \\ b_3 \\ b_4 \end{bmatrix} .$$

Ambiguity II:

$$b_1 \longrightarrow -b_1 \quad b_2 \longrightarrow +b_2 \quad b_3 \longrightarrow +b_3 \quad b_4 \longrightarrow -b_4$$

$$\begin{bmatrix} b'_1 \\ b'_2 \\ b'_3 \\ b'_4 \end{bmatrix} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \\ b_4 \end{bmatrix} = -i \tilde{\Gamma}^{10} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \\ b_4 \end{bmatrix} .$$

Ambiguity III:

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$$\begin{bmatrix} b'_1 \\ b'_2 \\ b'_3 \\ b'_4 \end{bmatrix} = \begin{bmatrix} -1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & +1 \end{bmatrix} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \\ b_4 \end{bmatrix} = i\tilde{\Gamma}^{12} \begin{bmatrix} b_1 \\ b_2 \\ b_3 \\ b_4 \end{bmatrix} .$$

Ambiguity IV:

$$b_1 \longrightarrow -b_2^* \quad b_2 \longrightarrow -b_1^* \quad b_3 \longrightarrow -b_4^* \quad b_4 \longrightarrow -b_3^*$$

$$\begin{bmatrix} b'_1 \\ b'_2 \\ b'_3 \\ b'_4 \end{bmatrix} = \begin{bmatrix} 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \end{bmatrix} \begin{bmatrix} b_1^* \\ b_2^* \\ b_3^* \\ b_4^* \end{bmatrix} = \tilde{\Gamma}^{15} \begin{bmatrix} b_1^* \\ b_2^* \\ b_3^* \\ b_4^* \end{bmatrix} . \quad (1)$$

Linear ambiguity I, II and III $L = \tilde{\Gamma}^4, \tilde{\Gamma}^{10},$ and $\tilde{\Gamma}^{12}$

Antilinear ambiguity IV ambiguity $A = \tilde{\Gamma}^{15}K_0.$

Other three antilinear ambiguities $A = \tilde{\Gamma}^6K_0, \tilde{\Gamma}^{13}K_0,$ and $\tilde{\Gamma}^8K_0,$ can be constructed by Ambiguity IV and the three linear ambiguities I to III. I

$$\begin{aligned}\tilde{\Gamma}^6 &= \tilde{\Gamma}^4 \tilde{\Gamma}^{15} \\ \tilde{\Gamma}^{13} &= j \tilde{\Gamma}^{10} \tilde{\Gamma}^{15} \\ \tilde{\Gamma}^8 &= -j \tilde{\Gamma}^{12} \tilde{\Gamma}^{15} .\end{aligned}$$

See KW (Ref. [?]).

Spin	Linear Transformation L			Antilinear Transformation A				
Observable	$\tilde{\Gamma}_4$	$\tilde{\Gamma}_{10}$	$\tilde{\Gamma}_{12}$	$\tilde{\Gamma}_6$	$\tilde{\Gamma}_8$	$\tilde{\Gamma}_{13}$	$\tilde{\Gamma}_{15}$	
$\sigma(\theta)$	+	+	+	+	+	+	+	\mathcal{S}
Σ	+	+	+	+	+	+	+	
T	+	+	+	+	+	+	+	
P	+	+	+	+	+	+	+	
G	-	-	+	+	-	+	-	\mathcal{BT}
H	-	-	+	-	+	-	+	
E	-	-	+	-	+	-	+	
F	-	-	+	+	-	+	-	

$$L: b_i \longrightarrow b'_i = L_{ij}b_j$$

$$A: b_i \longrightarrow b'_i = A_{ij}b_j^*$$

Spin	Linear Transformation L			Antilinear Transformation A			
Observable	$\tilde{\Gamma}_4$	$\tilde{\Gamma}_{10}$	$\tilde{\Gamma}_{12}$	$\tilde{\Gamma}_6$	$\tilde{\Gamma}_8$	$\tilde{\Gamma}_{13}$	$\tilde{\Gamma}_{15}$
O_x	-	+	-	-	-	+	+
O_z	-	+	-	+	+	-	-
C_x	-	+	-	+	+	-	-
C_z	-	+	-	-	-	+	+
T_x	+	-	-	+	-	-	+
T_z	+	-	-	-	+	+	-
L_x	+	-	-	-	+	+	-
L_z	+	-	-	+	-	-	+

$$L: b_i \longrightarrow b'_i = L_{ij}b_j$$

$$A: b_i \longrightarrow b'_i = A_{ij}b_j^*$$

Fierz for SU(4)

$$\Gamma_{ij}^{\alpha} \Gamma_{st}^{\beta} = \sum_{a,b=1}^{16} C_{a,b}^{\alpha,\beta} \Gamma_{it}^a \Gamma_{sj}^b \longrightarrow \Gamma^{\alpha} \Gamma^{\beta} = \sum_{a,b=1}^{16} C_{a,b}^{\alpha,\beta} \Gamma^a \Gamma^b$$

$$C_{a,b}^{\alpha,\beta} \equiv \frac{1}{16} \text{Tr}[\Gamma^a \Gamma^{\alpha} \Gamma^{\beta} \Gamma^b]$$

Fierzing Observables:

$$\Omega^{\alpha} \Omega^{\beta} = \sum_{a,b=1}^{16} C_{a,b}^{\alpha,\beta} \Omega^a \Omega^b$$

$16 \times 16/2 = 128 \longrightarrow x$ relations; for example:

$$\Gamma_{ij}^1 \Gamma_{st}^1 = \frac{1}{x} [\Gamma_{it}^1 \Gamma_{sj}^1 + \Gamma_{it}^2 \Gamma_{sj}^2 + \cdots + \Gamma_{it}^{16} \Gamma_{sj}^{16}]$$

Linear-Quadratic Relations (16):

$$\Omega_1 = 1 = \frac{1}{4} \sum_{\alpha=1}^{16} (\Omega_\alpha)^2$$

$$\Omega_4 = \Omega_{10}\Omega_{12} + \Omega_6\Omega_{15} - \Omega_8\Omega_{13}$$

$$\Omega_{10} = \Omega_4 \Omega_{12} + \Omega_2\Omega_{14} + \Omega_7\Omega_{16}$$

$$\Omega_{12} = \Omega_4 \Omega_{10} + \Omega_3\Omega_{11} - \Omega_5\Omega_9$$

$$\Omega_3 = +\Omega_{11}\Omega_{12} - \Omega_7\Omega_{15} + \Omega_{14}\Omega_8$$

$$\Omega_5 = -\Omega_9 \Omega_{12} + \Omega_7\Omega_{13} - \Omega_{14}\Omega_6$$

$$\Omega_9 = -\Omega_5 \Omega_{12} - \Omega_2\Omega_{15} - \Omega_{16}\Omega_8$$

$$\Omega_{11} = +\Omega_3 \Omega_{12} + \Omega_2\Omega_{13} + \Omega_{16}\Omega_6$$

Linear-Quadratic Relations:

$$\Omega_{14} = \Omega_2 \Omega_{10} + \Omega_3 \Omega_8 - \Omega_5 \Omega_6$$

$$\Omega_7 = \Omega_{16} \Omega_{10} - \Omega_3 \Omega_{15} + \Omega_5 \Omega_{13}$$

$$\Omega_{16} = \Omega_7 \Omega_{10} - \Omega_9 \Omega_8 + \Omega_{11} \Omega_6$$

$$\Omega_2 = \Omega_{14} \Omega_{10} - \Omega_9 \Omega_{15} + \Omega_{11} \Omega_{13}$$

$$\Omega_6 = +\Omega_{15} \Omega_4 - \Omega_5 \Omega_{14} + \Omega_{11} \Omega_{16}$$

$$\Omega_{13} = -\Omega_8 \Omega_4 + \Omega_5 \Omega_7 + \Omega_{11} \Omega_2$$

$$\Omega_8 = -\Omega_{13} \Omega_4 + \Omega_3 \Omega_{14} - \Omega_9 \Omega_{16}$$

$$\Omega_{15} = +\Omega_6 \Omega_4 - \Omega_3 \Omega_7 - \Omega_9 \Omega_2$$

Linear-Quadratic Relations:

$$\mathcal{I} = 1 = \frac{1}{4} \sum_{\alpha=1}^{16} (\Omega_{\alpha})^2$$

$$\Sigma = -T P - T_x L_z + L_x T_z$$

$$T = -\Sigma P + C_z O_x - O_z C_x$$

$$P = -\Sigma T + G F - H E$$

$$\begin{pmatrix} 1 & 0 & 0 & -P \\ 0 & 1 & P & 0 \\ 0 & P & 1 & 0 \\ -P & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} G \\ H \\ E \\ F \end{pmatrix} = \begin{pmatrix} 0 & 0 & O_x & O_z \\ O_x & O_z & 0 & 0 \\ 0 & 0 & C_x & C_z \\ C_x & C_z & 0 & 0 \end{pmatrix} \begin{pmatrix} T_x \\ T_z \\ L_x \\ L_z \end{pmatrix}$$

$$S \quad \times \quad BT = \quad BR \quad \times \quad TR$$

$$\begin{pmatrix} 1 & 0 & 0 & -T \\ 0 & 1 & T & 0 \\ 0 & T & 1 & 0 \\ -T & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} O_x \\ O_z \\ C_x \\ C_z \end{pmatrix} = \begin{pmatrix} H & 0 & G & 0 \\ 0 & H & 0 & G \\ F & 0 & E & 0 \\ 0 & F & 0 & E \end{pmatrix} \begin{pmatrix} T_x \\ T_z \\ L_x \\ L_z \end{pmatrix}$$

$$S \quad \times \quad BR = \quad BT \quad \times \quad TR$$

$$\begin{pmatrix} 1 & 0 & 0 & \Sigma \\ 0 & 1 & -\Sigma & 0 \\ 0 & -\Sigma & 1 & 0 \\ \Sigma & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} T_x \\ T_z \\ L_x \\ L_z \end{pmatrix} = \begin{pmatrix} H & 0 & F & 0 \\ 0 & H & 0 & F \\ G & 0 & E & 0 \\ 0 & G & 0 & E \end{pmatrix} \begin{pmatrix} O_x \\ O_z \\ C_x \\ C_z \end{pmatrix}$$

$$S \quad \times \quad TR = \quad BT \quad \times \quad BR$$

Quadratic Relations (15):

$$\Omega_2 \Omega_7 - \Omega_{14}\Omega_{16} - \Omega_3 \Omega_9 - \Omega_5 \Omega_{11} = 0$$

$$\Omega_3 \Omega_5 + \Omega_9 \Omega_{11} + \Omega_6 \Omega_8 + \Omega_{13}\Omega_{15} = 0$$

$$\Omega_2 \Omega_{16} - \Omega_7 \Omega_{14} - \Omega_6 \Omega_{13} - \Omega_8 \Omega_{15} = 0$$

$$\Omega_4 \Omega_3 - \Omega_{10}\Omega_{11} + \Omega_7 \Omega_6 + \Omega_{14}\Omega_{13} = 0$$

$$\Omega_4 \Omega_5 + \Omega_{10}\Omega_9 + \Omega_7 \Omega_8 + \Omega_{14}\Omega_{15} = 0$$

$$\Omega_4 \Omega_9 + \Omega_{10}\Omega_5 + \Omega_2 \Omega_6 - \Omega_{16}\Omega_{13} = 0$$

$$\Omega_4 \Omega_{11} - \Omega_{10}\Omega_3 + \Omega_2 \Omega_8 - \Omega_{16}\Omega_{15} = 0$$

$$\Omega_4 \Omega_{14} - \Omega_{12}\Omega_2 + \Omega_3 \Omega_{13} + \Omega_5 \Omega_{15} = 0$$

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$$\Omega_4 \Omega_7 - \Omega_{12}\Omega_{16} + \Omega_3 \Omega_6 + \Omega_5 \Omega_8 = 0$$

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$$\Omega_4 \Omega_2 - \Omega_{12}\Omega_{14} + \Omega_9 \Omega_6 + \Omega_{11}\Omega_8 = 0$$

$$\Omega_{10}\Omega_6 - \Omega_{12}\Omega_{15} + \Omega_5 \Omega_2 - \Omega_{11}\Omega_7 = 0$$

$$\Omega_{10}\Omega_{13} + \Omega_{12}\Omega_8 - \Omega_5 \Omega_{16} - \Omega_{11}\Omega_{14} = 0$$

$$\Omega_{10}\Omega_8 + \Omega_{12}\Omega_{13} - \Omega_3 \Omega_2 + \Omega_9 \Omega_7 = 0$$

$$\Omega_{10}\Omega_{15} - \Omega_{12}\Omega_6 + \Omega_3 \Omega_{16} + \Omega_9 \Omega_{14} = 0$$

Square Relations(6):

$$(\Omega_3)^2 + (\Omega_5)^2 + (\Omega_9)^2 + (\Omega_{11})^2 = (\Omega_1)^2 - (\Omega_4)^2 - (\Omega_{10})^2 + (\Omega_{12})^2$$

$$(\Omega_{14})^2 + (\Omega_7)^2 + (\Omega_{16})^2 + (\Omega_2)^2 = (\Omega_1)^2 - (\Omega_4)^2 + (\Omega_{10})^2 - (\Omega_{12})^2$$

$$(\Omega_6)^2 + (\Omega_{13})^2 + (\Omega_8)^2 + (\Omega_{15})^2 = (\Omega_1)^2 + (\Omega_4)^2 - (\Omega_{10})^2 - (\Omega_{12})^2$$

$$(\Omega_3)^2 + (\Omega_5)^2 - (\Omega_9)^2 - (\Omega_{11})^2 = (\Omega_{14})^2 + (\Omega_7)^2 - (\Omega_{16})^2 - (\Omega_2)^2$$

$$-(\Omega_3)^2 + (\Omega_5)^2 - (\Omega_9)^2 + (\Omega_{11})^2 = (\Omega_6)^2 + (\Omega_{13})^2 - (\Omega_8)^2 - (\Omega_{15})^2$$

$$(\Omega_{14})^2 - (\Omega_7)^2 + (\Omega_{16})^2 - (\Omega_2)^2 = (\Omega_6)^2 - (\Omega_{13})^2 + (\Omega_8)^2 - (\Omega_{15})^2$$

Bounds on measurements

From

$$\begin{aligned} & (\Omega_6)^2 + (\Omega_{13})^2 + (\Omega_8)^2 + (\Omega_{15})^2 \pm 2(\Omega_6 \Omega_{15} - \Omega_8 \Omega_{13}) \\ = & (\Omega_1)^2 + (\Omega_4)^2 - (\Omega_{10})^2 - (\Omega_{12})^2 \pm 2(\Omega_1 \Omega_4 - \Omega_{10} \Omega_{12}) \end{aligned}$$

we obtain

$$(\Omega_6 \pm \Omega_{15})^2 + (\Omega_8 \mp \Omega_{13})^2 = (\Omega_1 \pm \Omega_4)^2 - (\Omega_{10} \pm \Omega_{12})^2 .$$

The left hand side of the equation is positive, so is the right hand side. Therefore

$$\Omega_1 \pm \Omega_4 \geq |\Omega_{10} \pm \Omega_{12}| \quad \text{or} \quad 1 \pm \Sigma \geq |T \pm P| .$$

Other bounds, within the set \mathcal{S} , can be derived in the same way:

$$1 \pm T \geq |P \pm \Sigma| , \quad 1 \pm P \geq |\Sigma \pm T| .$$

We can deduce the bounds ¹

$$1 + \Sigma^2 \geq P^2 + T^2$$

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The BDS rule:

Barker, Donnachie, & Storrow, *Nucl. Phys.* **B95**, 347 (1975)

In BDS, the following rule was promulgated:

In order to determine all amplitudes without discrete ambiguities, one has to measure five double spin observables along with the four type S measurements, provided no four double spin observables are selected from the same set of BT , BR and TR .

Thus, they say **nine** experiments are required.

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Chiang & Tabakin *Phys. Rev. C* **55**, 2054 (1997)

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Pion Nucleon Elastic Case

Amplitude

$$T = f_1 + i \vec{\sigma} \cdot \hat{n} f_2$$

Transversity amplitudes

$$b_1 = \frac{f_1 + i f_2}{\sqrt{2}} \quad b_2 = \frac{f_1 - i f_2}{\sqrt{2}}$$

Observables

$$\mathcal{I}(\theta) = |b_1|^2 + |b_2|^2 = \langle b | \sigma^0 | b \rangle$$

$$\mathcal{I}(\theta)P = \hat{P} = |b_1|^2 - |b_2|^2 = A_N^2 = \langle b | \sigma^z | b \rangle$$

$$\mathcal{I}(\theta)A = \hat{A} = 2 \Re(b_1^* b_2) = 2|b_1||b_2| \cos(\theta) = \langle b | \sigma^x | b \rangle$$

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Fierzing Pion Nucleon Elastic Case

Fierz for SU(2)

$$\sigma_{ij}^{\alpha} \sigma_{st}^{\beta} = \sum_{a,b=0}^3 C_{a,b}^{\alpha,\beta} \sigma_{it}^a \sigma_{sj}^b \longrightarrow \sigma^{\alpha} \sigma^{\beta} = \sum_{a,b=0}^3 C_{a,b}^{\alpha,\beta} \sigma^a \sigma^b$$

$$4 \times 4/2 = 8 \longrightarrow 1$$

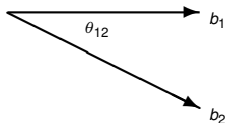
$$\sigma_{ij}^0 \sigma_{st}^0 = \sigma_{it}^1 \sigma_{sj}^1 + \sigma_{it}^2 \sigma_{sj}^2 + \sigma_{it}^3 \sigma_{sj}^3$$

$$P^2 + A^2 + R^2 = 1$$

Thus deduce bounds: $|P| \leq 1$ and $0 \leq A^2 + R^2 \leq 1$

Complete set: $\mathcal{I}, P, A = 3$ plus 1 for sign of R, for example. **The**

Two Complex Amplitudes In Transversity Space:



Fierzing Pion Nucleon Elastic Case

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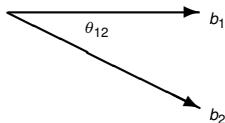
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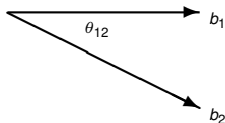
$$\sigma_{ij}^0 \sigma_{st}^0 = \sigma_{it}^1 \sigma_{sj}^1 + \sigma_{it}^2 \sigma_{sj}^2 + \sigma_{it}^3 \sigma_{sj}^3$$

$$P^2 + A^2 + R^2 = 1$$

Thus deduce bounds: $|P| \leq 1$ and $0 \leq A^2 + R^2 \leq 1$

Complete set: $\mathcal{I}, P, A=3$ plus 1 for sign of R, for example. **The**

Two Complex Amplitudes In Transversity Space:



Fierzing Pion Nucleon Elastic Case

Fierz for SU(2)

$$\sigma_{ij}^{\alpha} \sigma_{st}^{\beta} = \sum_{a,b=0}^3 C_{a,b}^{\alpha,\beta} \sigma_{it}^a \sigma_{sj}^b \longrightarrow \sigma^{\alpha} \sigma^{\beta} = \sum_{a,b=0}^3 C_{a,b}^{\alpha,\beta} \sigma^a \sigma^b$$

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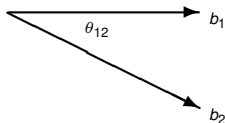
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Two Complex Amplitudes In Transversity Space:



Complete Double spin measurements

'X's' → 3 selected measurements,

'O's' → possible 4th observable to resolve ambiguities.

<i>G</i>	X	X	X	X	X	X	X	X	X	X	X	X	X	X
<i>H</i>	X	X	X	X	X	X	X	X						
<i>E</i>									X	X	X	X	X	X
<i>F</i>														
<i>O_x</i>	X		O		O	O	O	O	X	O	O	O	O	O
<i>O_z</i>		X		O	O	O	O	O	O	X	O	O	O	O
<i>C_x</i>	O		X		O	O	O	O	O	O	X	O	O	O
<i>C_z</i>		O		X	O	O	O	O	O	O	O	X	O	O
<i>T_x</i>	O	O	O	O	X	O	O	O	O	O	O	O	X	
<i>T_z</i>	O	O	O	O	O	X	O	O	O	O	O	O		X
<i>L_x</i>	O	O	O	O	O	O	X	O	O	O	O	O	O	
<i>L_z</i>	O	O	O	O	O	O	O	X	O	O	O	O		O

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<i>E</i>	X	X	X	X	X	X	X	X						
<i>F</i>									X	X	X	X	X	X
<i>O_x</i>	X		O			O	O		X	O	O	O	O	O
<i>O_z</i>		X		O	O			O	O	X	O	O	O	O
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<i>T_x</i>		O	O		X		O		O	O	O	O	X	
<i>T_z</i>	O			O		X		O	O	O	O	O		X
<i>L_x</i>	O			O	O		X		O	O	O	O	O	
<i>L_z</i>		O	O			O		X	O	O	O	O		O

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F		O		X	O	O	O	O	O	O	O	X	O	O
O_x	X	X	X	X	X	X	X	X	X	X	X	X	X	X
O_z	X	X	X	X	X	X	X	X						
C_x									X	X	X	X	X	X
C_z														
T_x	O	O	O	O	X	O	O	O	O	O	O	O	X	O
T_z	O	O	O	O	O	X	O	O	O	O	O	O	O	X
L_x	O	O	O	O	O	O	X	O	O	O	O	O		
L_z	O	O	O	O	O	O	O	X	O	O	O	O		

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F		O		X	O			O	O	O	O	X	O	O
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C_x	X	X	X	X	X	X	X	X						
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T_x	O			O	X	O			O	O	O	O	X	O
T_z		O	O		O	X			O	O	O	O	O	X
L_x		O	O				X	O	O	O	O	O		
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F			O	X	O	O	O	O	O	O	O	X	O	O
O_x	O	O	O	O	X	O	O	O	O	O	O	O	X	O
O_z	O	O	O	O	O	X	O	O	O	O	O	O	O	X
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C_z	O	O	O	O	O	O	O	X	O	O	O	O		
T_x	X	X	X	X	X	X	X	X	X	X	X	X	X	X
T_z	X	X	X	X	X	X	X	X						
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F			O	X	O			O	O	O	O	X	O	O
O_x	O			O	X	O			O	O	O	O	X	O
O_z		O	O		O	X			O	O	O	O	O	X
C_x		O	O				X	O	O	O	O	O		
C_z	O			O			O	X	O	O	O	O		
T_x														
T_z	X	X	X	X	X	X	X	X	X	X	X	X	X	X
L_x	X	X	X	X	X	X	X	X						
L_z									X	X	X	X	X	X

Partial Wave Extraction is made difficult by:

- Unknown Energy and Angle dependent Overall Phase
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- Error bars

Some Observations:

Determination of unique Underlining dynamics is almost always limited in QM. Examples: Phase equivalent potentials & Redundant Inverse scattering solutions

We live with that and instead of theorem we surround the problem.

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- The 16 Measurements are related by a measurement algebra
- That algebra is equivalent to that of 16 Hermitian 4×4 SU(4) group
- More than 7 experiments are needed due to discrete ambiguities, which are also described using SU(4) algebra.
- Constraints and Bounds can be deduced from that algebra
- Constrains allow for one less experiment needed and is why CT say 8 and BDS say 9 needed for complete set.
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Fierz \rightarrow explicit and rigorous relationships between observables. Of course, such relationships can be derived from the bilinear structure of the observables, with much effort. That effort is now replaced by simply invoking the well-known Fierz rules as a general property. That allows us to avoid much algebra and to find all relations in one step. There are direct physical consequences of these relations. If double spin observables in a type set are known, then the fourth member of that type is uniquely determined. The fourth measurement is thus redundant.

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Thanks for including me!