

Near light cone QCD

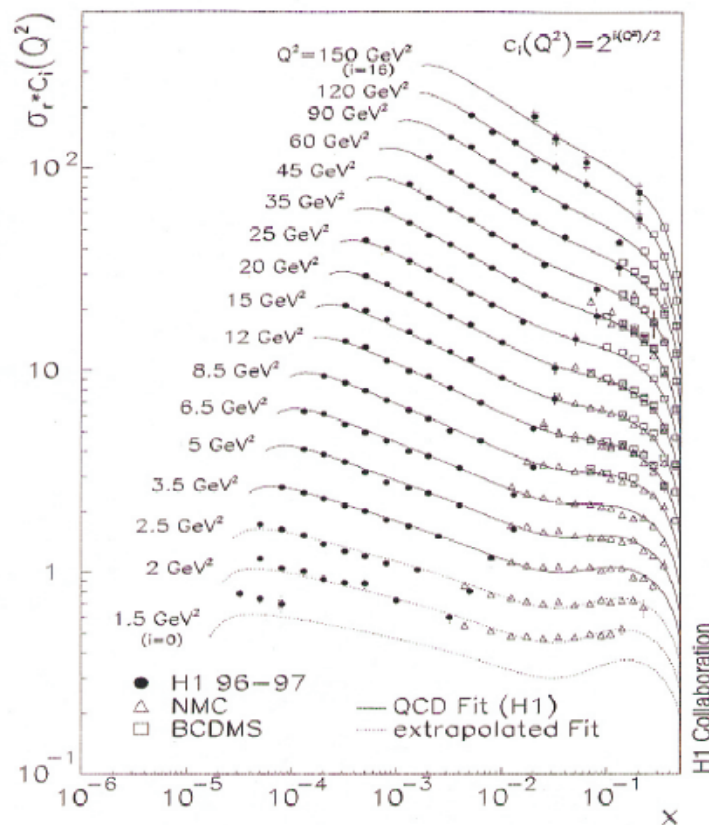
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Outline:

- Motivation: High energy total cross sections
- Near light cone Hamiltonian for Quark Antiquark QCD vacuum near the light cone
 - Nonperturbative Wilson Line correlations of the simplified vacuum Hamiltonian
 - Diffractive cross section
 - Lattice solution for the full Hamiltonian including transverse gauge fields

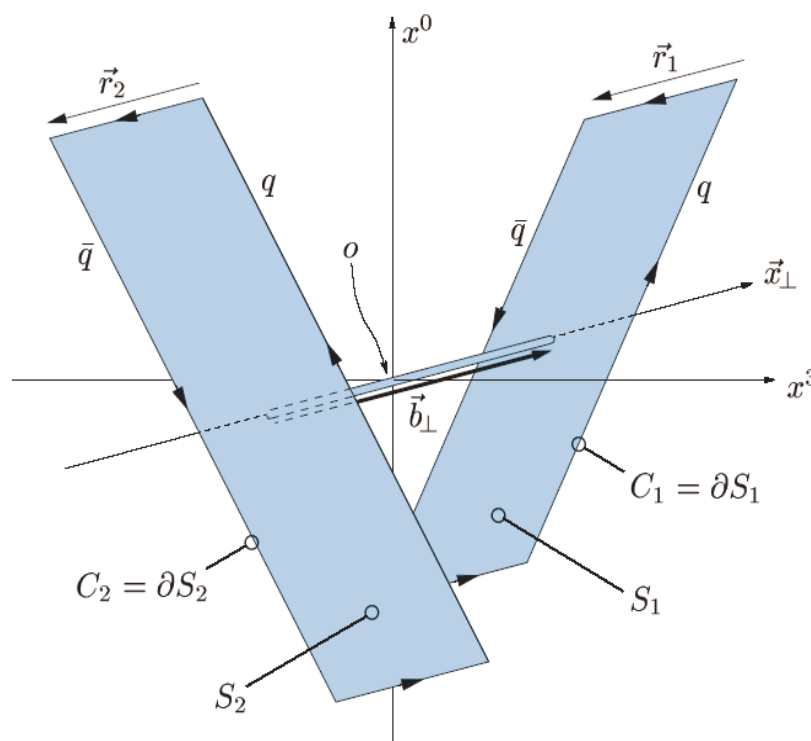
HERA: strong energy dependence of gamma - proton cross section



- Increase depends on Q^2
- Is much stronger than in hadronic collisions
- Can also be obtained by DGLAP evolving suitably parametrized structure functions
- Many theoretical models

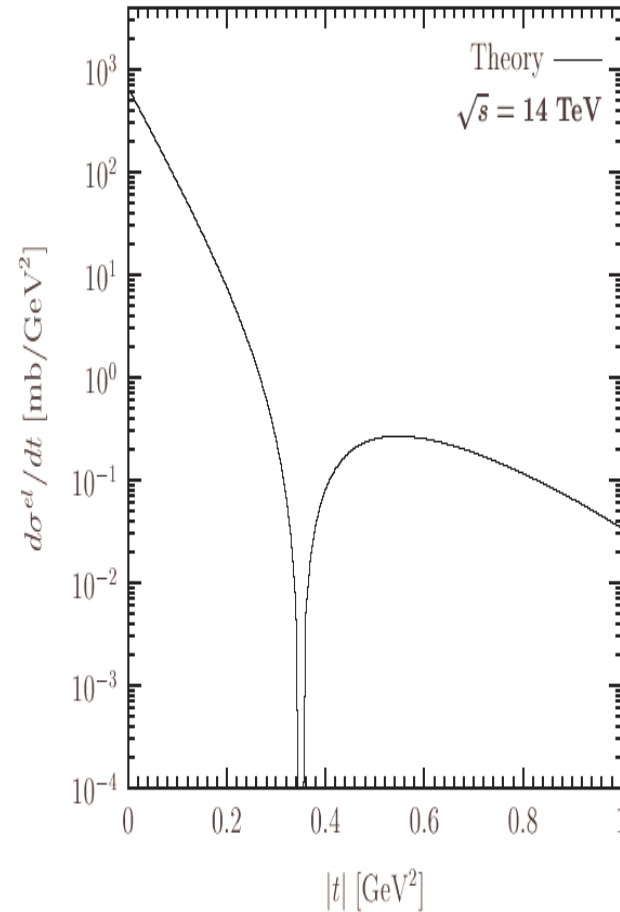
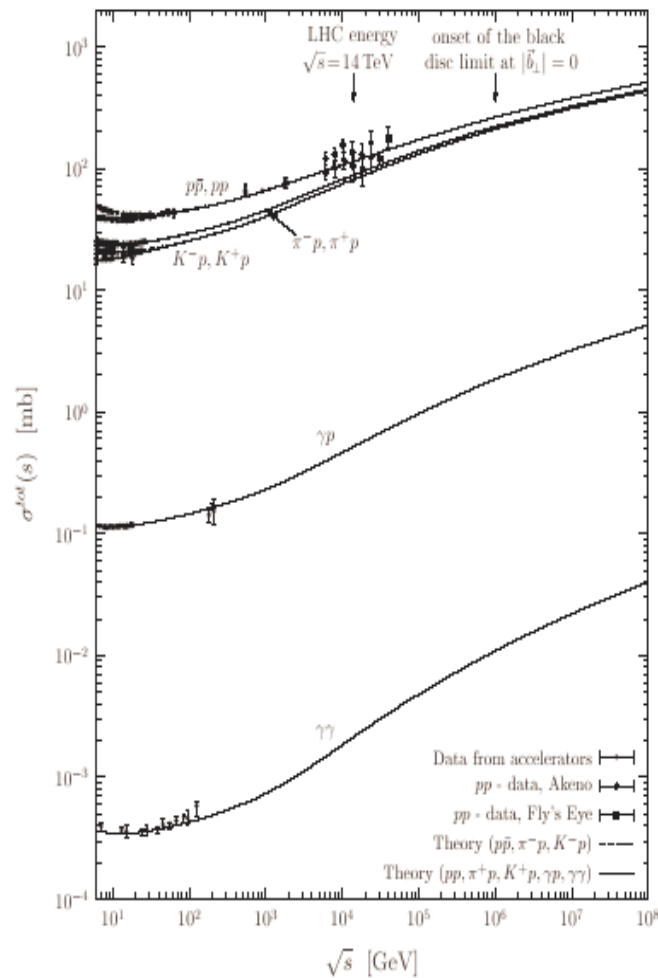
Loop Loop Correlation Model

- The loop loop correlation model gives naively no s -dependence, but one can calculate NLO in $\text{Ln}(1/x)$ contribution in the gluon distribution, then evolve with DGLAP



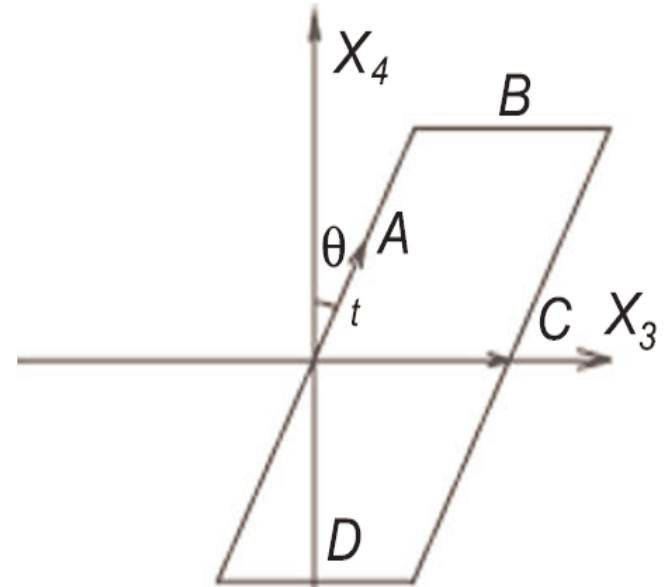
H.J.Pirner, A.Shoshi
G.Soyez Eur.Phys.
C33, 2004,63-74

Predictions: Total and elastic cross section



Calculation of near light cone Hamiltonian

- Start from Wilson loop near the light cone
- Go to Minkowski space
- Find $P^- = E - P_z$ from the phase factor



H.J. Pirner and N. Nurpeissov, **Phys.Lett.B595:379-386,2004**

Light cone Hamiltonian for q \bar{q}

$$u_\mu = (\gamma, 0_\perp, \gamma\beta).$$

Phasefactor= $e^{-ig \int d\tau (\gamma A^0 - \gamma\beta A^3)}$

$$\langle W_r[C] \rangle = e^{-i\gamma(P^0 - P^3)}$$

$$P^- = \frac{(\mu^2 + k_\perp^2)P^+}{2(1/4P^{+2} - k^{+2})} + \frac{1}{\sqrt{2}}\sigma\sqrt{x_3^2 + x_\perp^2/\gamma^2}.$$

Final Mass square Operator

$$M_c^2 = 2P^+ P^- = \frac{(\mu^2 + k_\perp^2)}{1/4 - \xi^2} + 2\sigma \sqrt{\rho^2 + M_c^2 x_\perp^2}$$

The final light cone mass operator has a potential term which confines the quark antiquark pair by the string tension. This Hamiltonian has been solved and gives typical mass 1 GeV

QCD Vacuum near the Light Cone

$$x^t = x^+ = \frac{1}{\sqrt{2}} \left\{ \left(1 + \frac{\eta^2}{2} \right) x^0 + \left(1 - \frac{\eta^2}{2} \right) x^3 \right\},$$
$$x^- = \frac{1}{\sqrt{2}} (x^0 - x^3).$$

η

- Quantization on a space like interval
- Formalism with η

**Always related
to equal time
Hamiltonian by a
boost**

:

H.Naus,H.J.P,T.Fields,J.P.Vary Phys. Rev. D 56,8062,1997

M.Ilgenfritz, Y.P.Ivanov, H.J.P. Phys. Rev D 62, 054006,2000

Lattice Hamiltonian

$$\begin{aligned}
 T_+^+ &= \frac{1}{2}g^2 \left(\frac{a_{\parallel}}{a_{\perp}} \right)^2 \Pi_-^{a^2} \\
 &\quad - 2 \frac{1}{g^2} \left(\frac{a_{\parallel}}{a_{\perp}} \right)^2 \text{Tr}(U_{\square}(1, 2) - 1) \\
 &\quad + \frac{1}{2}g^2 \frac{1}{\eta^2} \left[\Pi_i^a + i \frac{1}{g^2} \text{Tr}(\tau^a U_{\square}(-, i)) \right]^2
 \end{aligned}$$

Lattice constants a in parallel and transverse directions are differentiated. Lorentz factor for boosted transverse E and B fields is directly visible

Two Strategies for $\eta \rightarrow 0$

- Calculate only terms diverging on the light cone
- Multiply Hamiltonian with eta
- Look for fixed point in the effective coupling
- Solution does not have transverse A-fields
- Find a wave function which allows a solution of the full Hamiltonian for infinitesimal eta
- Try quantum evolution of fast and classical evolution of slow fields

I: Collective variables φ

- Wilson line integrals are the collective variables:
- Lattice Hamiltonian is defined with these variables

$$a_-^{c_0} = \frac{1}{L} \int_0^L dx^- A_-^{c_0}(\vec{x}_\perp, x^-)$$

$$\varphi^{c_0}(\vec{b}_\perp) = \frac{1}{2} g L a_-^{c_0}(b_\perp),$$

Critical behaviour of Z(3)

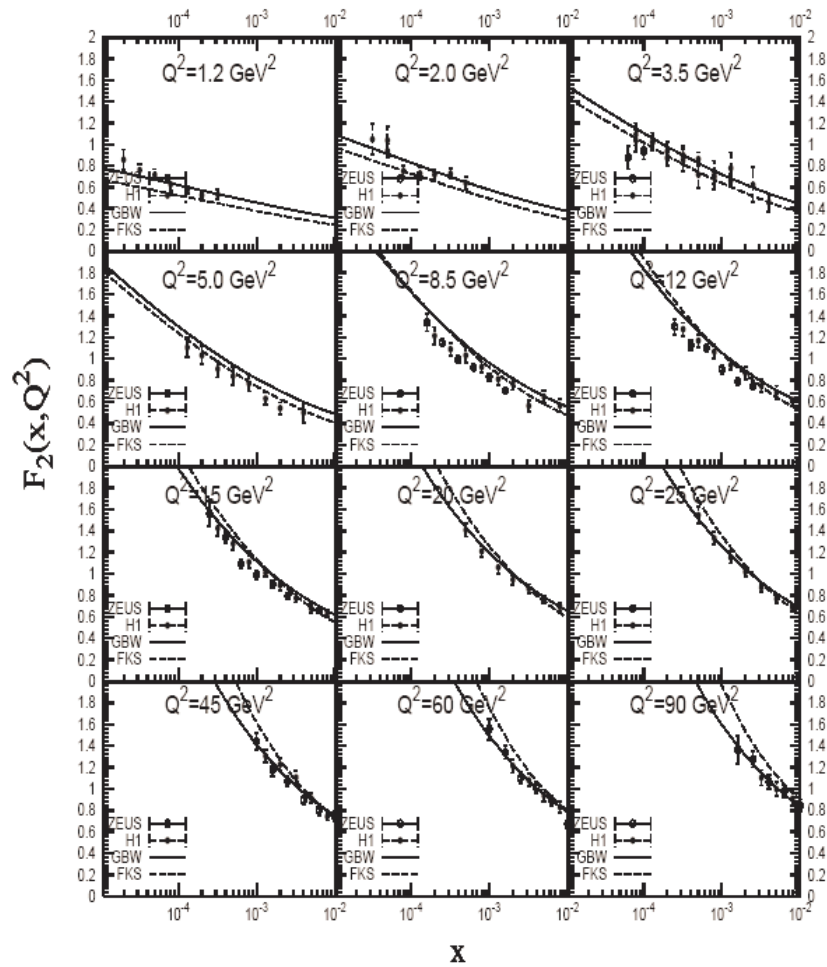
- The correlation length of Wilson lines grows with decreasing x

$$\xi \propto \left(\frac{x}{x_0} \right)^{-\frac{1}{2\lambda_2}} f_h(0).$$

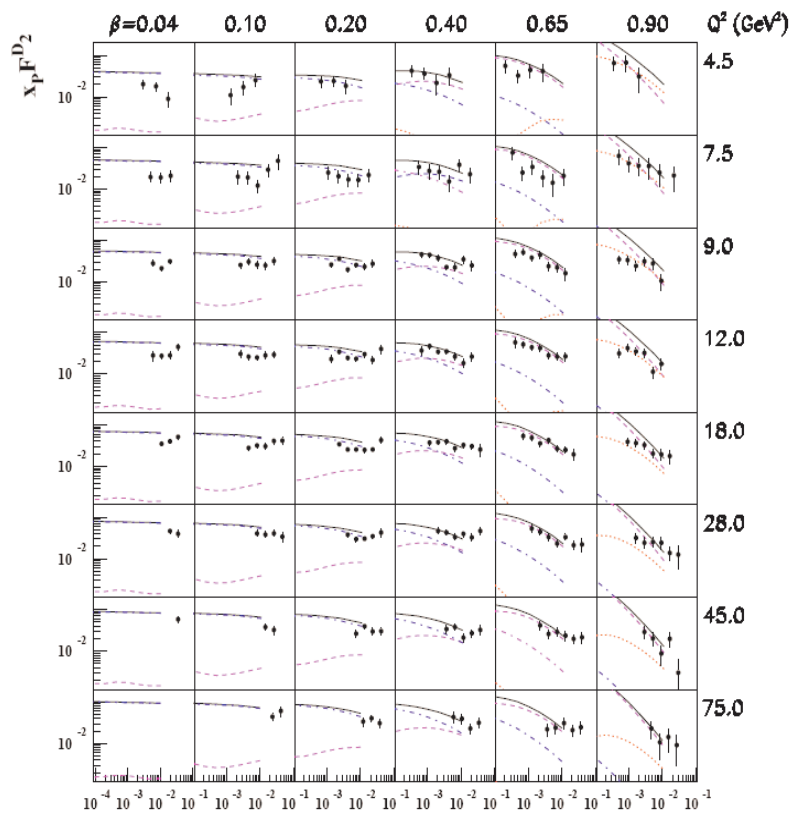
- The correlation length becoming bigger than $1/Q$ signals the dipole gas to dipole liquid transition

Results:

H.J. Pirner and
Yuan Feng
**Phys.Rev.D66:
034020,2002**



Diffractive Structure Function



With S.Munier unpublished

- Contributions from transverse photons via $q \text{ anti}q$
- Contributions from longitudinal photons
- Triple Pomeron contribution from quark antiquark gluon states

II: Exact Wave Function for the diverging terms

- For $A_- = 0$ one can find a starting solution for the leading term in the Hamiltonian
- This solution contains only positive frequency terms of the (fast) A_i fields
- In the classical evolution step the P_- -fields are updated with Gauss Law
- The A_- (slow) fields are updated with the classical equations of motion

M. Ilgenfritz, H.J. Pirner, D. Grünewald and E. Prokhlatilov
work in progress

Conclusions

- Near light cone valence wave quark Hamiltonian can be generated from tilted Euclidean Wilson loops
- Simplified Lattice Hamiltonian gives Wilson line correlations which give a consistent nonperturbative picture of F2 and F3
- Full Lattice Hamiltonian is solvable near the light cone including including zero mode fields and transverse gluon fields