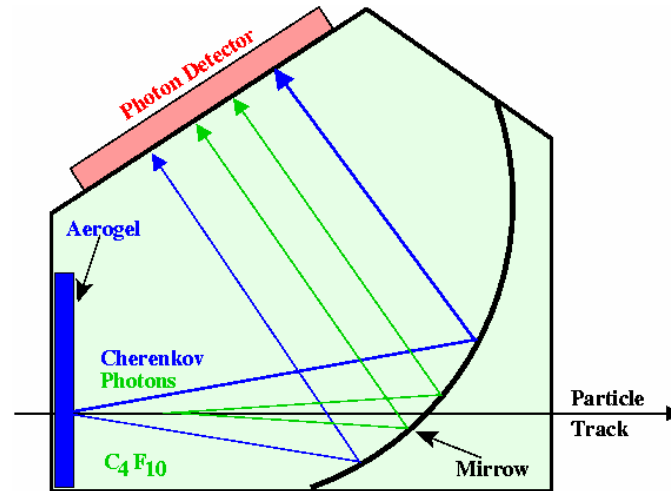


Strange Sea Contribution to the Nucleon Spin



Fatiha Benmokhtar

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with

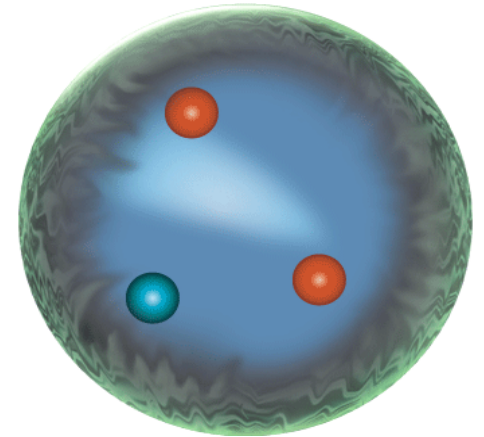
A. El Alaoui, (LPC Grenoble, France), **H. Avakian** (Jlab, USA),

K. Hafidi (ANL, USA)

Strange Sea Contribution to the Nucleon Spin

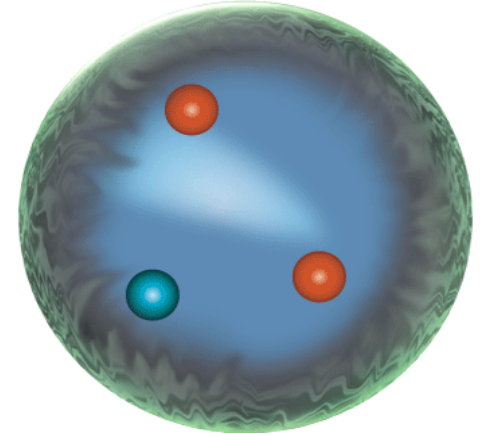
Outline

- Strange quark contributions to the nucleon properties (in general).
- Strange quark contribution to the nucleon spin.
- World data on ΔS
- Semi Inclusive DIS: - Five flavor tagging
- Isosclar Method
- **DeltaS with CLAS12?**



Nucleon constituents

$$P = \underset{\substack{\nearrow \\ \text{valence}}}{uud} + \underbrace{u\bar{u} + d\bar{d} + s\bar{s}}_{\text{« sea= virtual pairs »}} + g + \dots$$



- The sea contains all flavors, but
 - the **u** and **d** sea can't be distinguished from the valence
 - the heavier quarks (**c,b,t**) are too heavy to contribute much
 - **Strange quark** is the natural candidate to study the sea.

**With how much do virtual pairs contribute
to the structure of the nucleon ?
Mass? Momentum? Charge & Magnetisation
and Spin?**

Strange Quark Contribution to the Nucleon Properties

(in brief)

- Mass: Hyp \rightarrow $= \langle \bar{\psi} | \psi \rangle 130 \text{ MeV}$ \rightarrow 0 to 30 % with big theoretical uncertainties...
 π -N \rightarrow

- Longitudinal Momentum: -Study the spectral functions $q(x)$,
 -From unpol DI ν_{μ} -Nucleon scattering (NuTeV)

. For $x < 0.1$ $\int_0^1 x (s(x) + \bar{s}(x)) dx \sim 4\%$ \rightarrow Difficult to connect with ordinary observable

- Electromagnetic Form Factors: Parity violation experiments: G0, HAPPEX, PVA4.

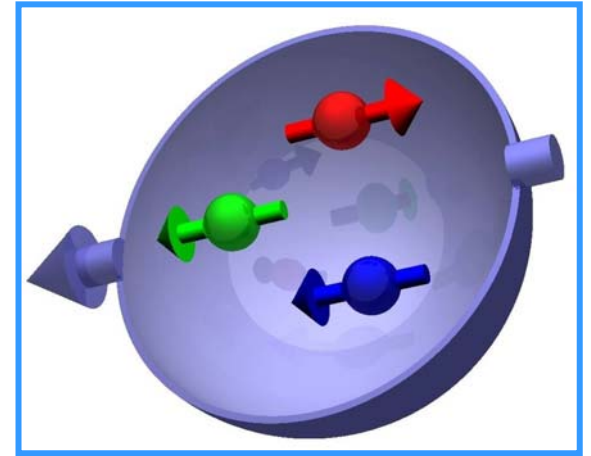
$$G_{E,M}^S = \left(1 - 4 \sin^2 \theta_W\right) G_{E,M}^{\gamma,p} - G_{E,M}^{\gamma,n} - G_{E,M}^{Z,p} \rightarrow \text{Underway} \text{ 😊}$$

- What about the contribution to the Nucleon Spin?

$$\langle N | \bar{s} \gamma^\mu \gamma^5 s | N \rangle = ???$$

Strangeness Contribution to the Nucleon Spin

$$s = \frac{1}{2} \Sigma, \quad \Sigma = \quad +$$
$$= \frac{4}{3} \quad d = \frac{1}{3}$$



BUT

In 1989 EMC measured

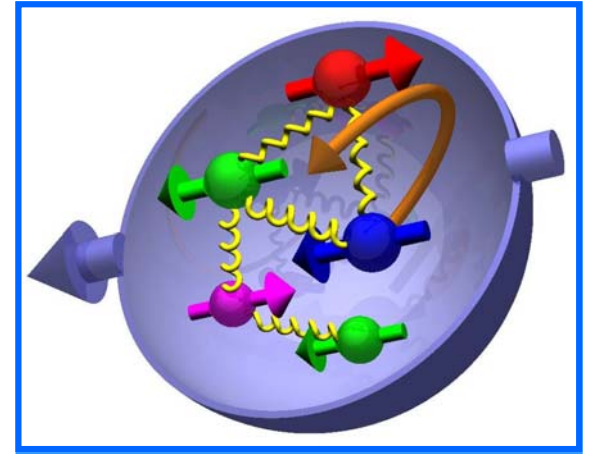
$$\Sigma = 0.120 \pm 0.094 \mp 0.138 \quad \rightarrow \quad \sim 20 \% \text{ of the nucleon spin!!!}$$

→ Spin Crisis!

Strangeness Contribution to the Nucleon Spin

$$S = \dots = - \sum_q \dots + \dots + \dots$$

$$= \dots + \dots - \dots + \dots - \dots + \dots - \dots$$



-> Gluon spin ΔG

-> Sea quark spin Δq

-> Orbital Angular Momentum L_G and L_q

Lattice Calculations predict:

$$\Delta s = -0.1$$

World Data on ΔS

- Polarized deep-inelastic inclusive scattering (SMC)

$$\Delta u + \Delta d + \Delta s = 0.20 \pm 0.10$$
$$\Delta s = -0.1 \pm 0.1$$

- ν -p elastic scattering (E734 BNL)

$$\Delta s = -0.15 \pm 0.09$$

- Polarized Semi-Inclusive DIS (HERMES)

* 5 flavor tagging \rightarrow published

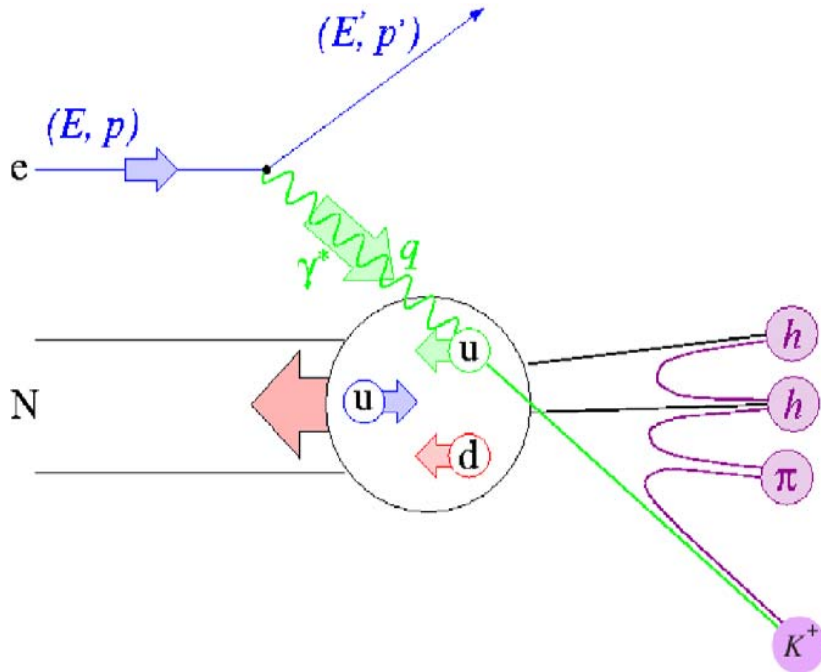
* Isoscalar method \rightarrow preliminary

$$\Delta s \sim$$

So far, the results vary widely and the uncertainties are big

$$\left\langle N \left| \bar{s} \gamma^\mu \gamma^5 s \right| N \right\rangle = ???$$

Semi Inclusive DIS



$$\sigma(e p \rightarrow e h X)$$

Sidis Selection:

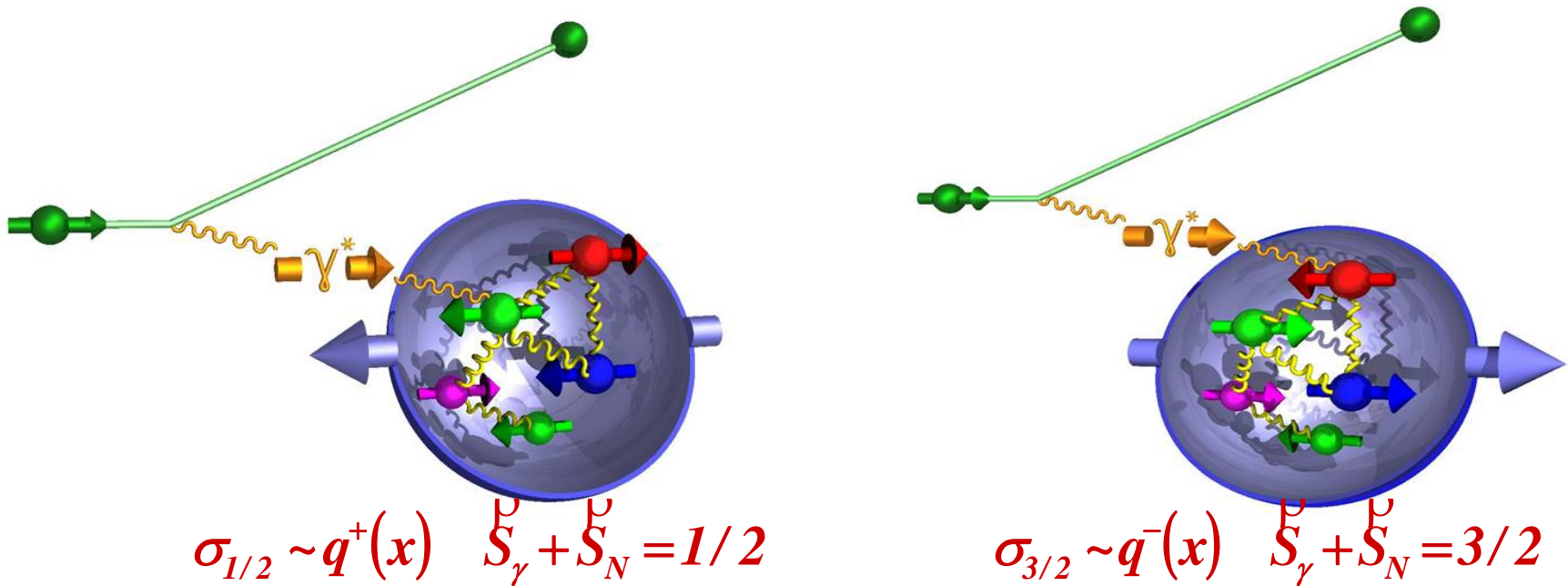
- $\nu = E - E'$
- $x = Q^2 / (2M \nu)$
- $y = \nu / E$
- $W^2 = M^2 - 2M \nu - Q^2$
- $Q^2 = 4EE' \sin^2(\theta/2)$

Hadron Selection:

- $z = E_h / \nu$
- $x_F = 2 |p_{||}| / W$

$$\frac{d^2 \sigma}{d\Omega dE'} = \frac{\alpha^2}{MQ^4} \frac{E}{E'} L_{\mu\nu} W^{\mu\nu}$$

Spin Asymmetries



• Select $q^-(x)$ or $q^+(x)$ by changing the orientation of target nucleon spin or helicity of incident lepton beam

Spin Independent Structure Function F_1

$$\sigma_{1/2} + \sigma_{3/2} \propto F_1(x) = \frac{1}{2} \sum_f e^2 (q^+(x) + q^-(x))$$

Spin Dependent Structure Function g_1

$$\sigma_{1/2} - \sigma_{3/2} \propto g_1(x) = \frac{1}{2} \sum_f e^2 \underbrace{(q^+(x) - q^-(x))}_{\Delta q}$$

Measurement of Δq in Semi-Inclusive DIS

$$A_1^{e(h)}(x, Q^2) = \frac{\sigma_{1/2}^{e(h)} - \sigma_{3/2}^{e(h)}}{\sigma_{1/2}^{e(h)} + \sigma_{3/2}^{e(h)}} = \frac{1}{(1 + \eta\gamma)D} \frac{1}{\langle P_B P_T \rangle} \frac{(N^{e(h)}/L)^{\bar{\leftarrow}} - (N^{e(h)}/L)^{\bar{\rightarrow}}}{(N^{e(h)}/L)^{\bar{\leftarrow}} + (N^{e(h)}/L)^{\bar{\rightarrow}}}$$

$$\sim \sum_q \frac{e_q^2 q(x) \int dz D_q^h(z)}{\sum_{q'} e_{q'}^2 q'(x) \int dz D_{q'}^h(z)} \frac{\Delta q(x)}{q(x)}$$

$$P_q^h(x, z)$$

- P_q^h Purity is a conditional probability that a hadron of type h observed in the final state is originated from a struck quark of flavor q in case of unpolarized beam/target.

- D_q^h is a measure of the probability that a quark of flavor q will fragment into a hadron of type h = fragmentation function

$$P A_1 = P \cdot Q$$

ΔS from SIDIS

$$\overset{P}{A}_1 = P \cdot \overset{P}{Q}$$

1- Five flavor decomposition (Δq)

2-Isoscalar extraction of ΔS

Five flavor decomposition (Δq with SIDIS)

$$\vec{A}_1 = \begin{pmatrix} A_{1,p} \\ A_{1,p}^{\pi^+} \\ A_{1,p}^{\pi^-} \\ A_{1,d} \\ \cdot \\ \cdot \\ \cdot \\ A_{1,d}^{K^-} \end{pmatrix} \quad P_q^h(x) = \frac{e_q^2 q(x) \int dz D_q^h(z)}{\sum_{q'} e_{q'}^2 q'(x) \int dz D_{q'}^h(z)} \quad \vec{Q} = \begin{pmatrix} \Delta u/u \\ \Delta d/d \\ \Delta \bar{u}/\bar{u} \\ \Delta \bar{d}/\bar{d} \\ \Delta s/s \end{pmatrix}$$

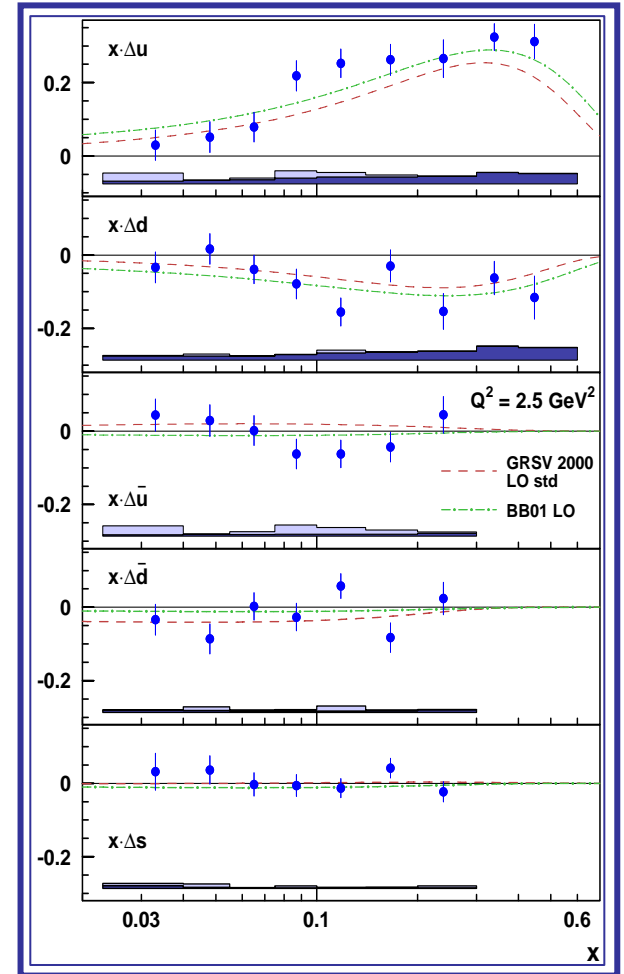
-HERMES: Purities calculated from MC simulation of the entire scattering process.

$$\Delta u = 0.601 \pm 0.039 \pm 0.049 \quad \Delta \bar{u} = -0.002 \pm 0.036 \pm 0.029$$

$$\Delta d = -0.226 \pm 0.039 \pm 0.050 \quad \Delta \bar{d} = -0.054 \pm 0.033 \pm 0.011$$

$$\Delta s = 0.028 \pm 0.033 \pm 0.009$$

Phys.Rev.D71(2005)012003



Isoscalar extraction of ΔS

$$K^+ = u\bar{s} \quad K^- = \bar{u}s$$

➤ Need a longitudinally polarized deuterium target

- strange quark in proton and neutron identical
- fragmentation simplifies

➤ Assumptions:

- isospin symmetry between proton and neutron
- charge-conjugation invariance in fragmentation

➤ Extraction from data of:

- inclusive $A_{1d}(x, Q^2)$ and kaon $A_{1d}^K(x, Q^2)$ double spin asymmetries
- kaon multiplicities

Isoscalar extraction of ΔS :

$$\begin{pmatrix} A_d(x) \\ A_d^K(x) \end{pmatrix} = \begin{pmatrix} P_Q(x) & P_S(x) \\ P_Q^K(x) & P_S^K(x) \end{pmatrix} \begin{pmatrix} \Delta Q(x)/Q(x) \\ \Delta S(x)/S(x) \end{pmatrix}$$

• Measure inclusive asymmetry $A_{1,d}$ and kaon asymmetries $A_{1(K^+ + K^-)}$

• Extract isoscalar combinations of $\Delta Q(x)$ and $\Delta S(x)$

$$\Delta S(x) \equiv \Delta s(x) + \Delta \bar{s}(x)$$

$$\Delta Q(x) \equiv \Delta u(x) + \Delta \bar{u}(x) + \Delta d(x) + \Delta \bar{d}(x)$$

• Inclusive purities from PDFs
(CTEQ6, MRST, GRV...)

$$P_Q(x) = \frac{5Q(x)}{5Q(x) + 2S(x)}, P_S(x) = \frac{2S(x)}{5Q(x) + 2S(x)}$$

• Kaon purities can be computed from the kaon multiplicities and the pdfs: **How?** →

Kaon Multiplicities

$$P_Q^K(x) = \frac{Q(x) \mathcal{D}_{\text{non strange}}^K}{Q(x) \mathcal{D}_{\text{non strange}}^K + 2S(x) \mathcal{D}_{\text{strange}}^K}$$

$$P_S^K(x) = \frac{S(x) \mathcal{D}_{\text{strange}}^K}{Q(x) \mathcal{D}_{\text{non strange}}^K + 2S(x) \mathcal{D}_{\text{strange}}^K}$$

- Using charge symmetry $D_q^{K^+ + K^-}(z) = D_{\bar{q}}^{K^+ + K^-}(z)$

$$\frac{dN^K(x)/dx}{dN^{\text{DIS}}(x)/dx} = \frac{Q(x) \mathcal{D}_{\text{non strange}}^K + 2S(x) \mathcal{D}_{\text{strange}}^K}{5Q(x) + 2S(x)}$$

Fit parameters

Measure

Multiplicities

HERMES Preliminary results (isoscalar)

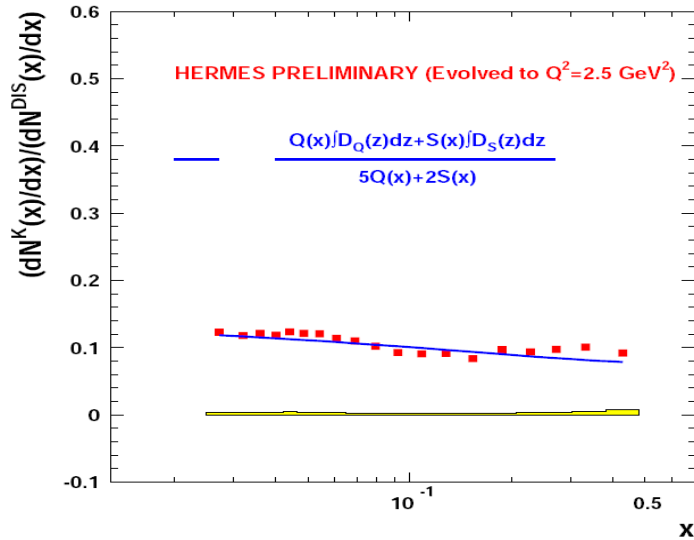


Fig. 1. Multiplicity in 4π of charged kaons in semi-inclusive DIS on a deuterium target as a function of Bjorken x . The statistical error bars are not visible, and the bands at the bottom represent the systematic uncertainties.

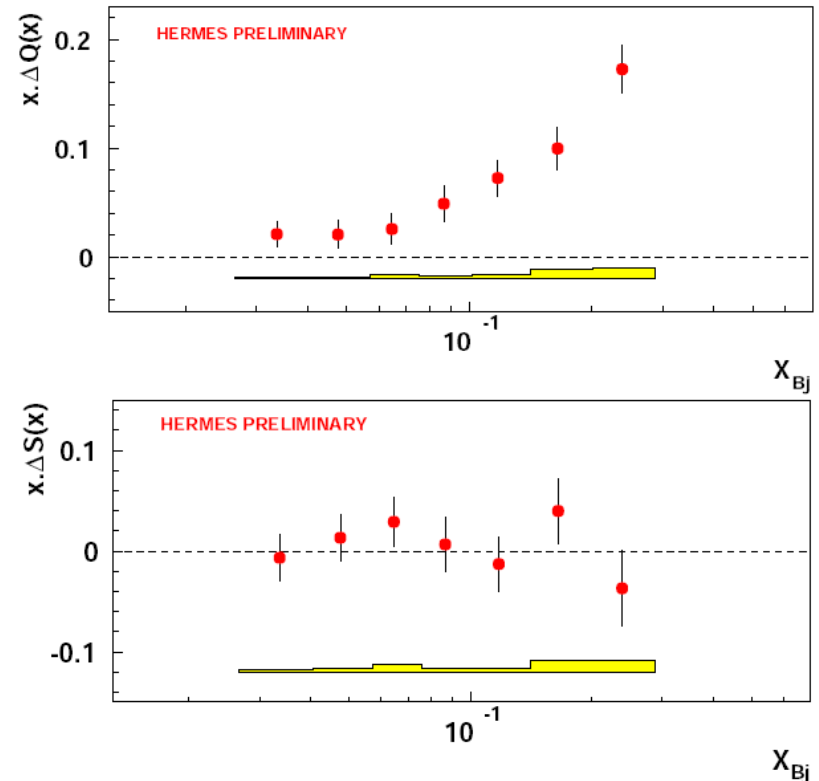
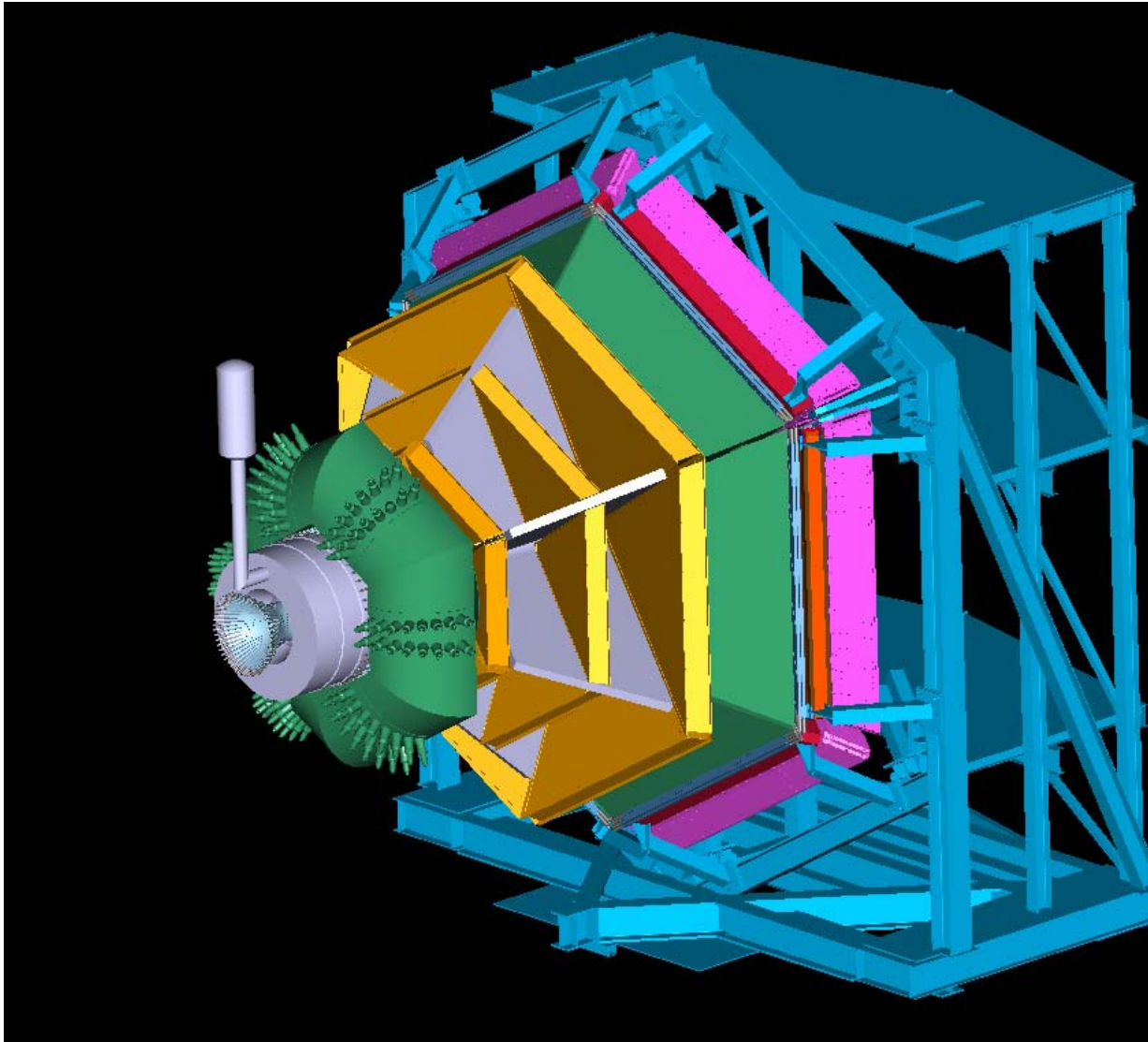


Fig. 2. Strange and non-strange quark helicity distributions at $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$ as a function of Bjorken x . The error bars are statistical, and the bands at the bottom represent the systematic uncertainties.

H. E. Jackson proceedings in EPJ 2006

What can we do with CLAS12?

CLAS12 @ 11GeV



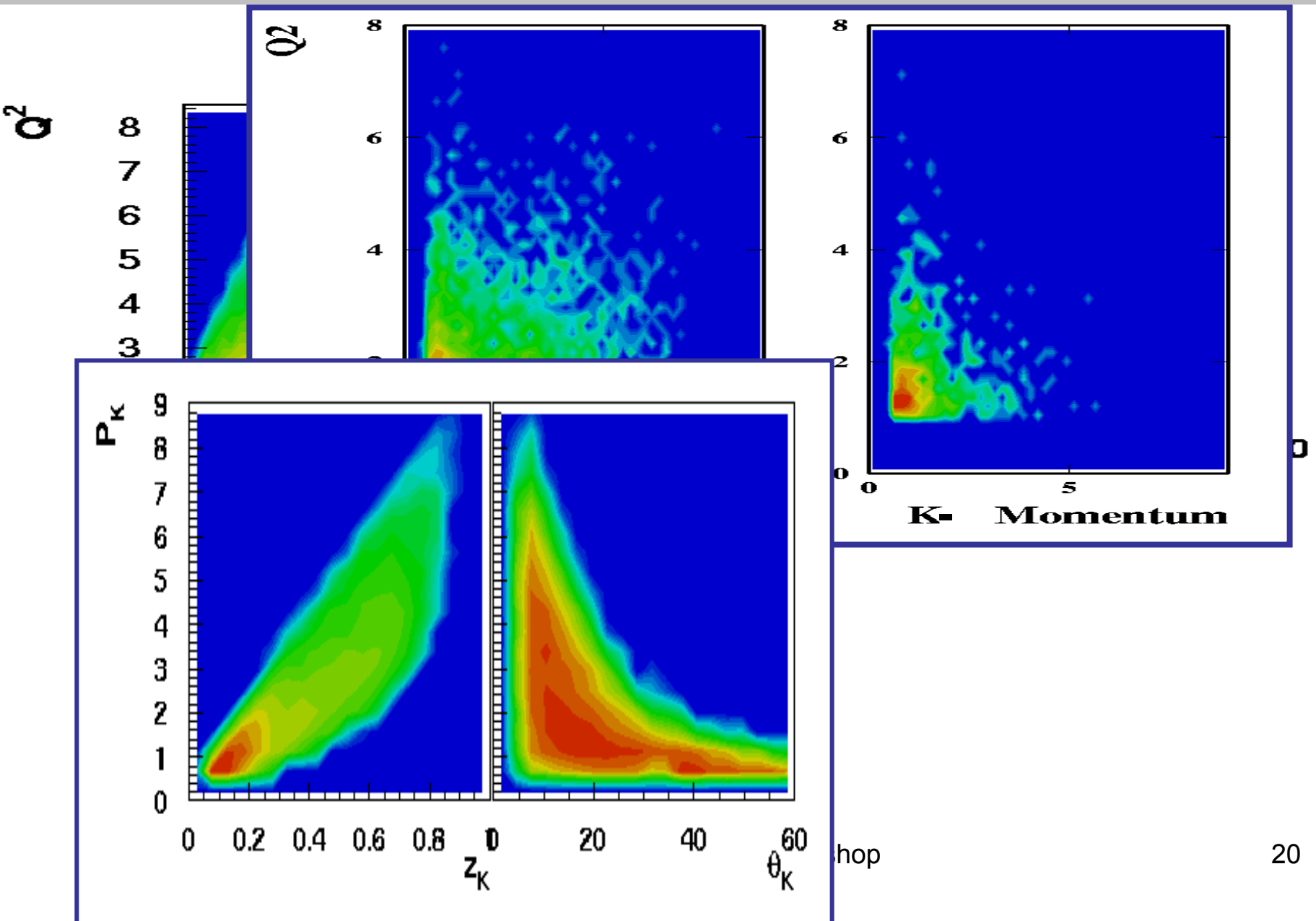
Wide detector and physics acceptance
(current/target fragmentation)
High beam polarization 85%
High target polarization 85%
NH₃, ND₃ targets

Track resolutions:

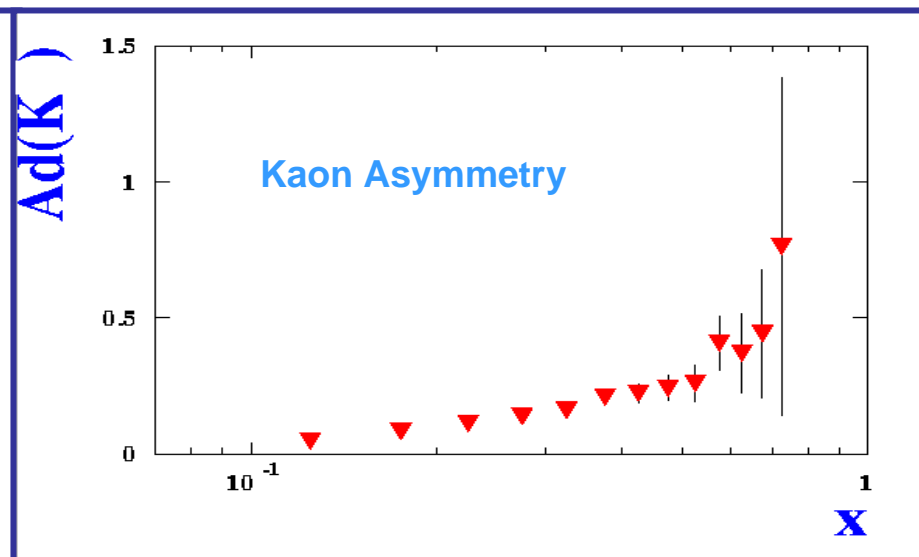
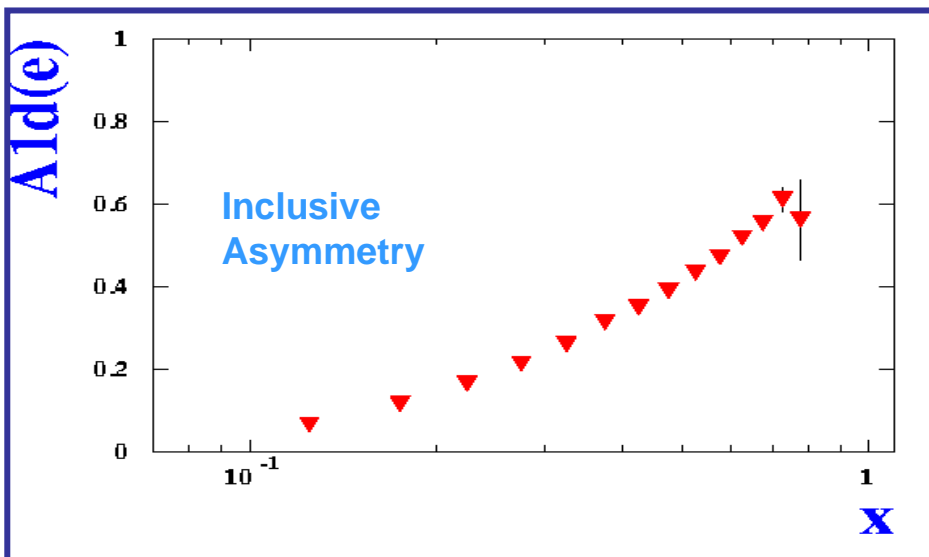
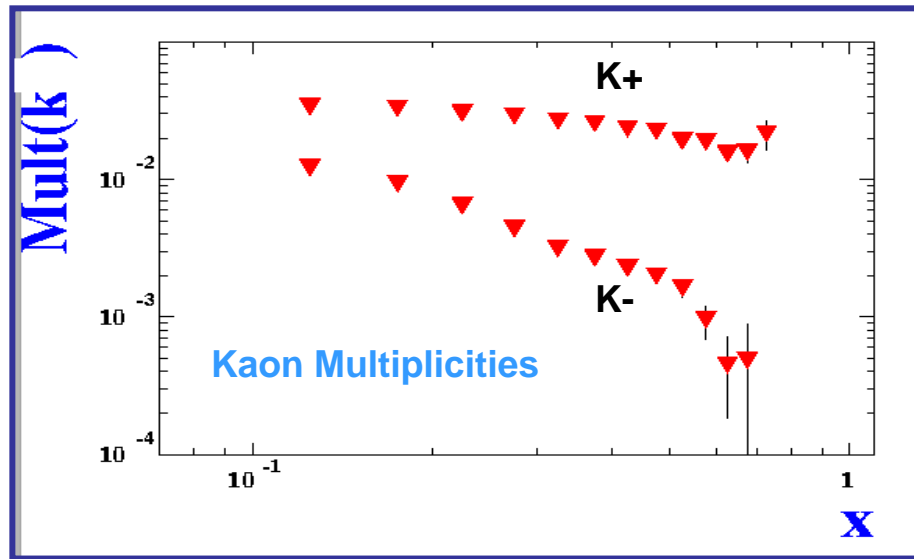
$$\begin{array}{ll} \delta p \text{ (GeV/c)} & 0.003p + 0.001p^2 \\ \delta\theta \text{ (mr)} & < 1 \\ \delta\phi \text{ (mr)} & < 3 \end{array}$$

$$\text{Lumi} > 10^{35} \text{cm}^{-2} \text{s}^{-1}$$

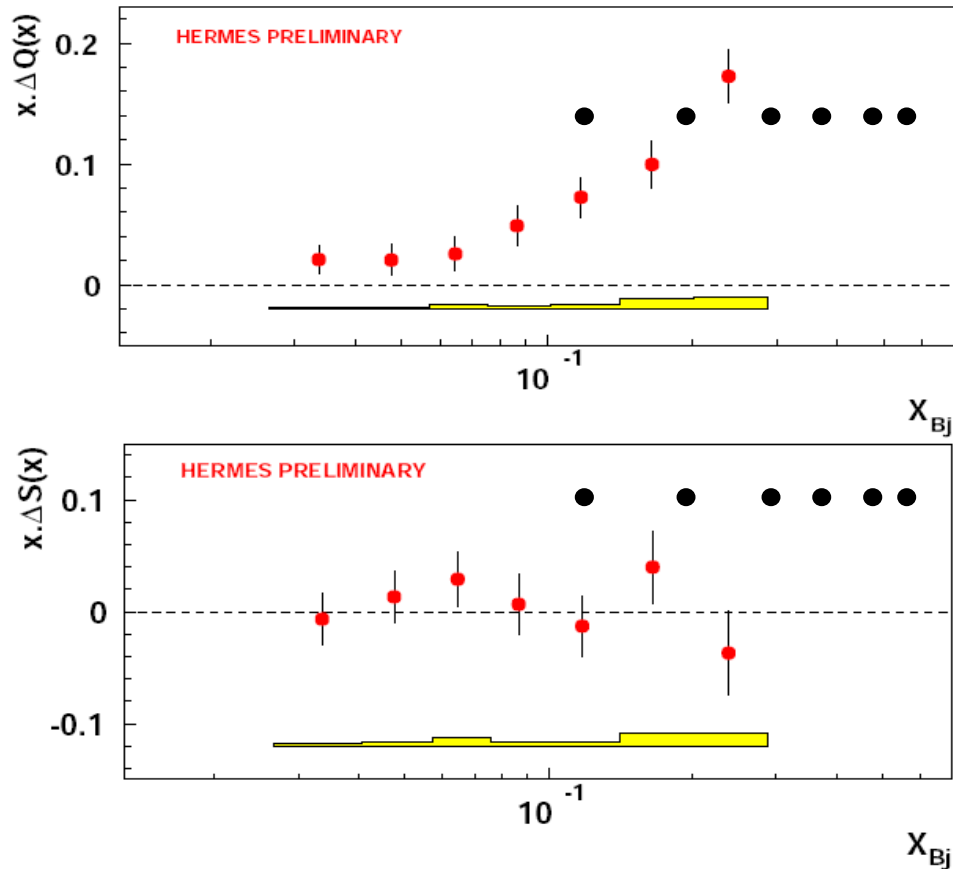
CLAS12 Kaon Distributions



CLAS12 kaon SIDIS Simulations



CLAS12 points



- Overlap with Hermes
- What can one do with the data at $x > 0.1$?

-Does it worth it to go ahead with a proposal? 😊

Support slides....

Beam time???

clasdisde.p1.e11.000.emn0.75tmn.12.xs**57.78nb**.dis92.dat

Data analyzed with 200 runs pol+, 200 runs pol-
& 450 runs unpol...> for what?

57.78nb per .dat ,

Number of events: 57780 per dat.

We have 400.dat polarized file in this simulation. --> 23.112.000 total events.

Lumi > $10^{35} \text{cm}^{-1} \text{s}^{-1}$

Neutrino-p technique

. ν -p elastic scattering (E734 BNL)

.Cross section contains a form factor G_1^s

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·

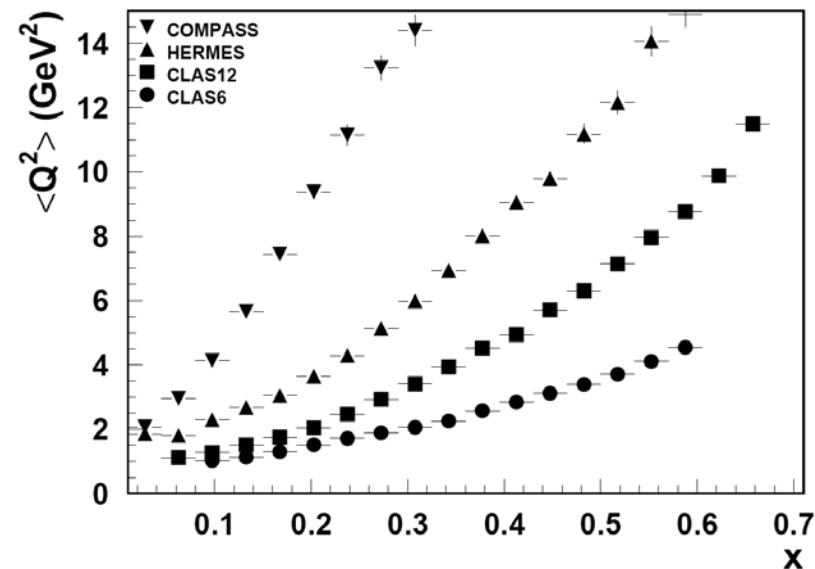
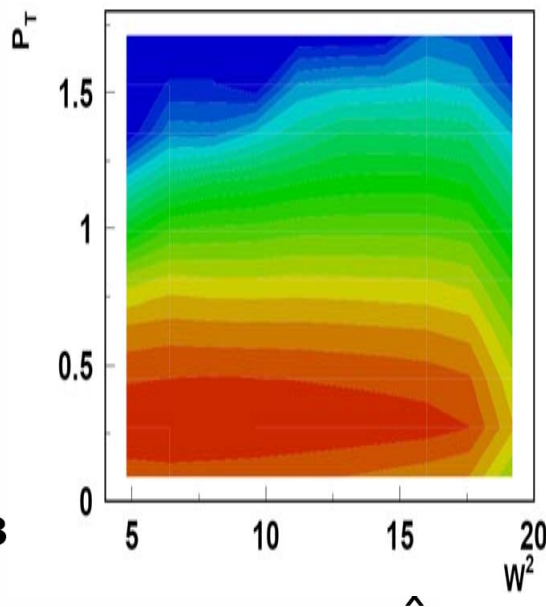
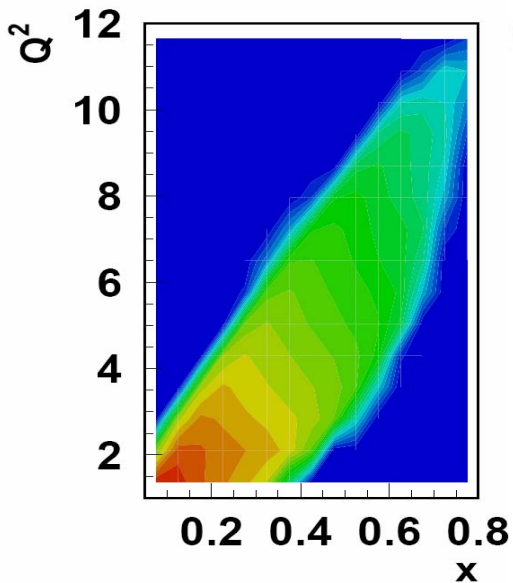
$$G_1^s(Q^2 \rightarrow 0) = \Delta s$$

$$\Delta s = -0.15 \pm 0.09$$

Parity Violation and DeltaS

- Neutral weak probes have sensitivity to D_s
 D_s contributes to parity violation electron scattering, however its sensitivity is suppressed due to the smallness of electron weak charge (at least at forward angles)
- . The contribution of D_s is not suppressed in elastic neutrino-nucleon scattering. It can be determined by fitting the cross section with nucleon EM form factors inputs. \rightarrow small and negative! Strange axial form factor behaves as
G Ae at $Q^2 \rightarrow 0$

CLAS12: Kinematical coverage



**SIDIS
kinematics**

$Q^2 > 1 \text{ GeV}^2$
 $W^2 > 4 \text{ GeV}^2 (10)$
 $y < 0.85$
 $M_X > 2 \text{ GeV}$

$x = 0.3 \rightarrow Q^2 \approx 2 \text{ GeV}^2$ (CLAS),
 $\sim 5 \text{ GeV}^2$ (HERMES)
 $\sim 15 \text{ GeV}^2$ (COMPASS)

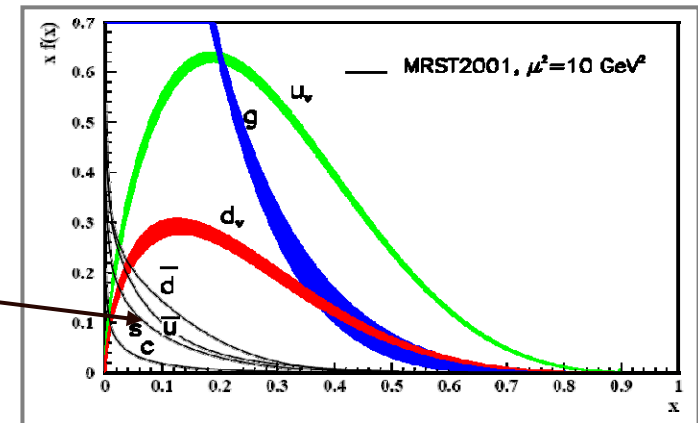
Large Q^2 accessible with CLAS12 are important for **separation of HT contributions...need to change this.**

- DIS kinematics, $Q^2 > 1 \text{ GeV}^2$, $W^2 > 4 \text{ GeV}^2$, $y < 0.85$
- $0.4 > z > 0.7$, $M_X^2 > 2 \text{ GeV}^2$

Strangeness Contribution to the Nucleon Longitudinal Momentum

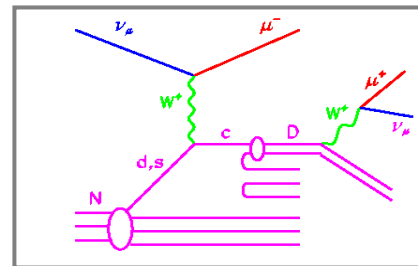
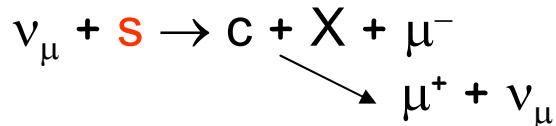
-Study the spectral functions $q(x)$, $= \frac{1}{v}$

-From unpolarized deep inelastic ν_μ -Nucleon scattering Experiments (NuTeV)



Momentum "fraction"

2 muon from charm quark production:



$$\int_0^1 x (s(x) + \bar{s}(x)) dx \sim 4\%$$

For $x < 0.1$ → Difficult to connect with ordinary observable

Strangeness Contribution to the Nucleon Mass

Mass of the Nucleon: $M_N = \langle N | H_{QCD} | N \rangle$

Quark mass term $H_{QCD} = \sum_i m_i \bar{q}_i q_i$

Chiral limit: $M_N = M_0 \neq 0$ \rightarrow gluon and $q\bar{q}$ condensate.

Turn on the quark masses: $M_N = M_0 + \mathcal{O} + \sigma_s + \text{heavy quark term}$

\mathcal{O} and σ_s : scalar FF. = $f(Q^2)$

$$\mathcal{O} = m \langle N | \bar{u}u + \bar{d}d | N \rangle, \sigma_s = ms \langle N | \bar{s}s | N \rangle \quad m = \frac{1}{2}(m_u + m_d)$$

-First constraint:

Hyperon mass splitting due the SU(3) flavor symmetry breaking effect.

$$\frac{1}{3} \left(1 - \frac{m_s}{m}\right) (1 - y) \hat{\sigma} = M_\Lambda - M_\Xi$$

$$y = \frac{2 \langle N | \bar{s}s | N \rangle}{\langle N | \bar{u}u + \bar{d}d | N \rangle}$$

strange content of the nucleon

- canonical ratio: $m_s/m \approx 26$
- assume $y = 0$
- Higher order chiral corrections

$$\hat{\sigma} = 35 \text{ MeV}$$

-Second constraint:

$\pi - N$ sigma term

$$\Sigma_{\pi N} = F_\pi^2 \bar{D}^+(s = M_N^2, t = 2m_\pi^2)$$

Which is related to the isospin even $\pi - N$ scattering amplitude .

s = invariant mass of the $\pi - N$ system

t = four-momentum transfer

Lowest order PQCD low energy theorem states: $\Sigma_{\pi N} \approx \hat{\sigma}(t = 2m_\pi^2)$

($t = 0$) **Extrapolation**

F. Benmokhtar, RICH work

$$\hat{\sigma} = 45 \text{ MeV}$$

Strangeness Contribution to the Nucleon Mass

$$\begin{array}{l}
 \text{Hyp} \rightarrow \hat{\sigma} = 35 \text{ MeV} \\
 \pi\text{-N} \rightarrow \hat{\sigma} = 45 \text{ MeV}
 \end{array}
 \left. \vphantom{\begin{array}{l} \text{Hyp} \\ \pi\text{-N} \end{array}} \right\} y = 2\% \Rightarrow \langle |^- | \rangle = 130 \text{ MeV}$$

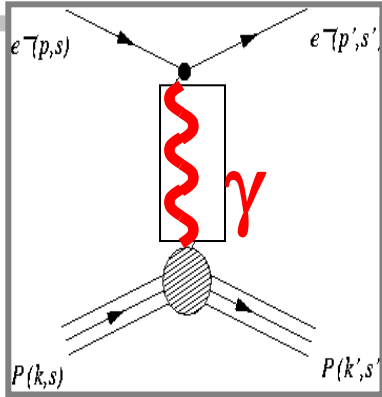
Sea quark contribute to a sizable amount to the nucleon mass!

However: uncertainty quite large due to many theoretical or experimental factors.

- . Other lattice calculation lead to: $\sigma \sim 53 \text{ MeV}$ and $y = 0.36$
 - . Experimental input suggests $\sigma \sim 90 \pm 8 \text{ MeV}$
 - . Analysis based on dispersion sum rules: $\sigma \sim 70 \pm 9 \text{ MeV}$.
- Etc..

Conclusion Mass : 0 to 30 % with big uncertainties...

Strange Electromagnetic Form Factors



Define vector (EM) form factors:

$$\langle N | J_{EM}^\mu | N \rangle \Rightarrow G_{E,M}^\gamma$$

Distribution of nucleon's charge and magnetization.

$$G_{E,M}^\gamma = \frac{2}{3} G_{E,M}^u - \frac{1}{3} G_{E,M}^d - \frac{1}{3} G_{E,M}^s$$

$$J_{EM}^\mu = \sum e_q \bar{q} \gamma^\mu q$$

▪ Charge symmetry

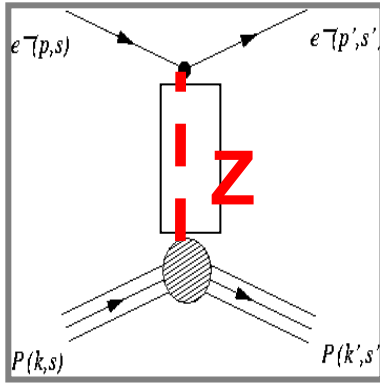
$$G^{u,p} = G^{d,n}; G^{d,p} = G^{u,n}; G^{s,p} = G^{s,n}$$

$$G_{E,M}^{\gamma,p} = \frac{2}{3} G_{E,M}^u - \frac{1}{3} G_{E,M}^d - \frac{1}{3} G_{E,M}^s$$

$$G_{E,M}^{\gamma,n} = \frac{2}{3} G_{E,M}^d - \frac{1}{3} G_{E,M}^u - \frac{1}{3} G_{E,M}^s$$

Need one more constraint ...

Neutralweak Form Factors — Additional Constraint



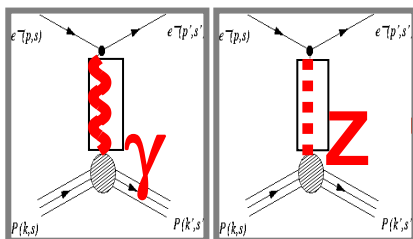
$$\langle N | J_{NC}^{\mu,V} | N \rangle \Rightarrow G_{E,M}^Z$$

Flavor decomposition of proton neutral weak form factor:

$$G_{E,M}^{Z,p} = \left(1 - \frac{8}{3} \sin^2 \theta_W\right) G_{E,M}^u + \left(-1 + \frac{4}{3} \sin^2 \theta_W\right) G_{E,M}^d + \left(-1 + \frac{4}{3} \sin^2 \theta_W\right) G_{E,M}^s$$

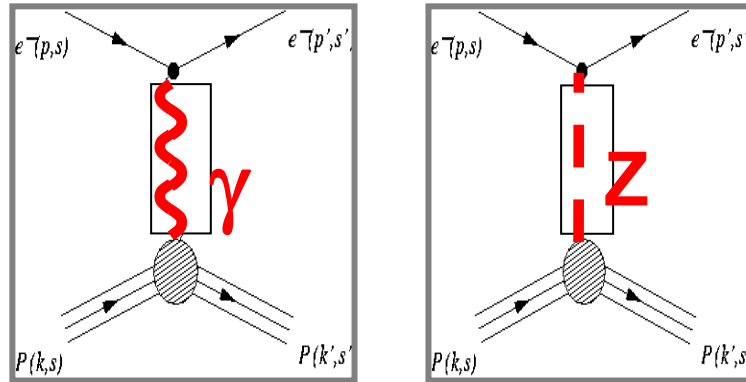
Quark weak charges in the unified electroweak theory in SM

$$\sin^2 \theta_W = 0.2312 \pm 0.00015$$



$$G_{E,M}^S = \left(1 - 4 \sin^2 \theta_W\right) G_{E,M}^{\gamma,p} - G_{E,M}^{\gamma,n} - G_{E,M}^{Z,p}$$

Parity Violation Asymmetry



Interference: $\sigma \sim |M^{EM}|^2 + |M^{NC}|^2 + 2\text{Re}(M^{EM*})M^{NC}$, $\frac{M_\gamma}{M_Z} \approx 10^5$

- Scatter polarized electrons off unpolarized target
- Tiny ($\sim 10^{-6}$) cross section asymmetry isolates weak interaction

$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim \frac{|M_{PV}^{NC}|}{|M^{EM}|} \sim \frac{Q^2}{(M_Z)^2}$$

$$A_{PV} = \left[\frac{-G_F Q^2}{4\pi\sqrt{2}} \right] \frac{A_E + A_M + A_A}{\sigma_p}$$

F. Benmokhtar, RICH workshop

Electric Magnetic Axial