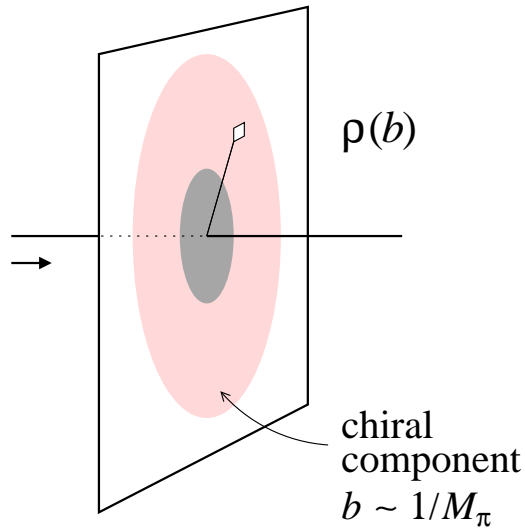


# Chiral dynamics in transverse nucleon structure

C. Weiss (JLab), Chiral Dynamics 2012, 08–Aug–12



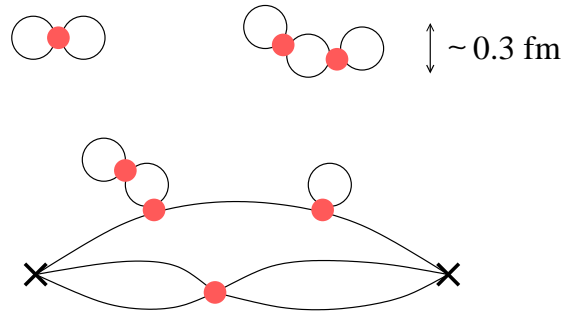
$b \sim 1/M_\pi$  new regime for  $\chi$ PT  
universal, model-independent

Chiral and non-chiral components  
identified by spatial size

Connection with GPDs,  
peripheral  $ep/pp$  processes

- Parton picture of nucleon structure
  - Wave function description
  - Transverse densities from elastic FFs
  - Connection with GPDs
- Chiral component of transverse density
  - Invariant  $\chi$ PT: Spectral function  $\text{Im } F_{1,2}(t)$   
Strikman, CW, PRC82 (2010) 042201
  - Light-front  $\chi$ PT:  $\pi N$  wave functions  
C. Granados, CW; in progress
  - Chiral vs. non-chiral component
  - Spectral analysis: Vector mesons  
Miller, Strikman, CW, PRC84 (2011) 045205
  - Effect on low- $t$  elastic FFs  
Simple estimates only!
- Chiral dynamics in parton densities
  - $x$ -dependent transverse size of nucleon
  - Peripheral high-energy  $ep/pp$  scattering

# Nucleon structure: Parton picture



- QCD vacuum not empty

Strong non-perturbative gluon fields  
of size  $\ll 1$  fm ← Lattice QCD, analytic models

Chiral symmetry breaking:  $\bar{q}q$  pair condensate,  
 $\pi$  as collective excitation

- Nucleon at rest

$\langle N | J_\mu | N \rangle$  from Euclidean correlation functns

No concept of particle content!

Cannot separate “constituents” from vacuum fluctuations

- Fast-moving nucleon  $P \gg \mu_{\text{vac}}$

Closed system: Wave function, Gribov, Feynman  
variable particle number,  $x_i, \mathbf{k}_{Ti}$

Current operators “count” particle nr  $\times$  charge

Physical properties:

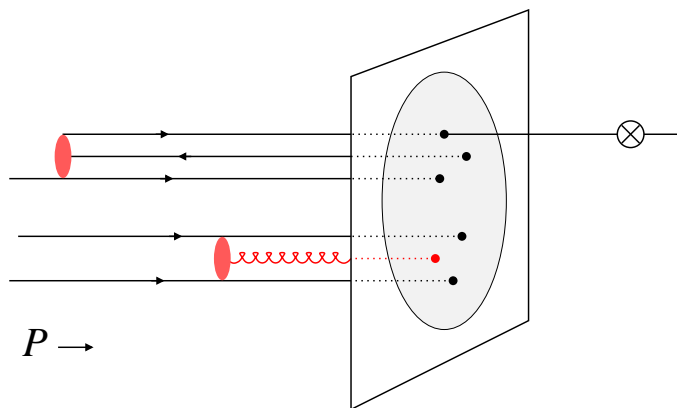
Longitudinal momentum densities PDFs

Transverse distributions → Form factors, GPDs

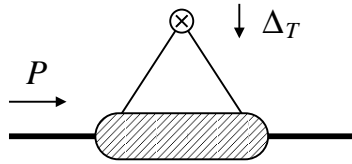
Orbital motion TMDs

Alt. view: Observer moves with velocity  $v \rightarrow 1$ .

Light-front quantization, frame-independent Brodsky et al.



# Nucleon structure: Transverse densities



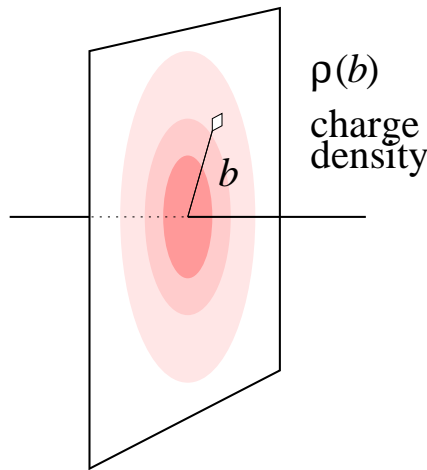
- Current matrix element parametrized by invariant form factors

$$\langle N' | J_\mu | N \rangle \rightarrow F_1(t), F_2(t) \quad \text{Dirac, Pauli}$$

- Transverse charge density  $t = -\Delta_T^2$

$$F_1(t) = \int d^2b e^{i\Delta_T \cdot b} \rho(b) \quad \text{2D Fourier}$$

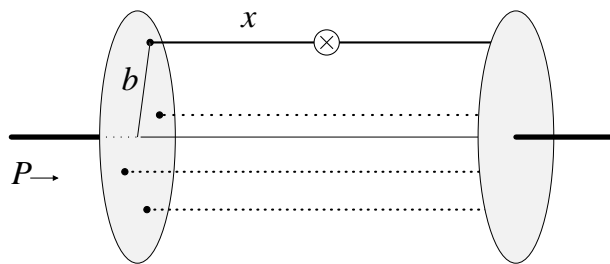
Transverse density of charge in fast-moving nucleon  
 $\mathbf{b}$  displacement from transverse C.M.



- Proper density for relativistic system

$$\rho(b) = \sum_N \text{charge} \int dx \psi^*(x, \mathbf{b}/\bar{x}, \dots) \psi(x, \mathbf{b}/\bar{x}, \dots)$$

Cumulative charge of constituents at transv. position  $\mathbf{b}$   
 Breit frame distribution not density: Miller 07

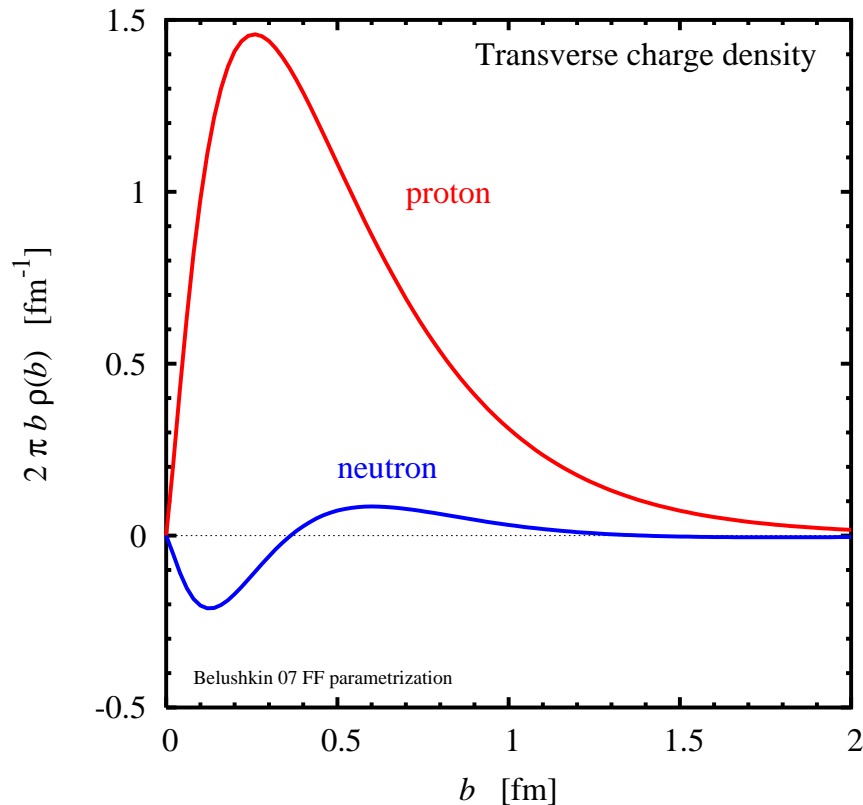


- Reduction of GPDs

$$\rho(b) = \int dx f_{q-\bar{q}}(x, \mathbf{b})$$

Transverse size in hard exclusive processes  
 universal, process-independent

# Nucleon structure: Densities from FF data



$$\rho(b) = \int_0^{\infty} \frac{d\Delta}{2\pi} \Delta J_0(\Delta b) F_1(t = -\Delta^2)$$

- Nucleon transverse charge density from spacelike form factor data

Experimental and incompleteness errors estimated [Venkat, Arrington, Miller, Zhan 10](#)

Recent low- $|t|$  data incorporated  
MAMI: [Vanderhaeghen, Walcher 10](#)

Neutron density positive at distances  $b \sim 0.5 - 1$  fm

[Miller 07](#). Different theoretical explanations

- Large-distance behavior?

$b \rightarrow \infty$  dominated by chiral dynamics: universal, model-independent

At what distances is the chiral component relevant?

Can it be seen experimentally?

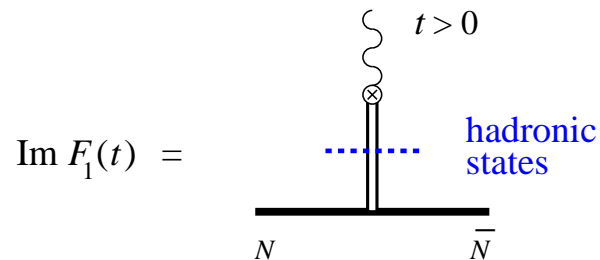
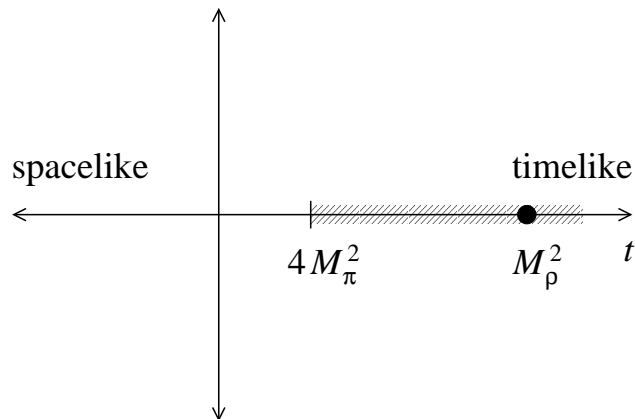
# Transverse density: Dispersion representation

- Dispersion representation of form factor

$$F_1(t) = \int_{4m_\pi^2}^{\infty} \frac{dt'}{t' - t - i0} \frac{\text{Im } F_1(t')}{\pi}$$

Spectral function  $\text{Im } F_1(t')$  describes “process”  
current  $\rightarrow$  hadronic states  $\rightarrow N\bar{N}$

$\text{Im } F_1(t')$  from form factor fits and theory:  
 $\chi$ PT near threshold, dispersion rels, pQCD  $t \rightarrow \infty$   
Höhler et al. 76; Belushkin, Hammer, Meissner 06



Isovector:  $\pi\pi, \rho, \rho', \dots$   
Isoscalar:  $\omega, \phi, K\bar{K}, \dots$

- Transverse density

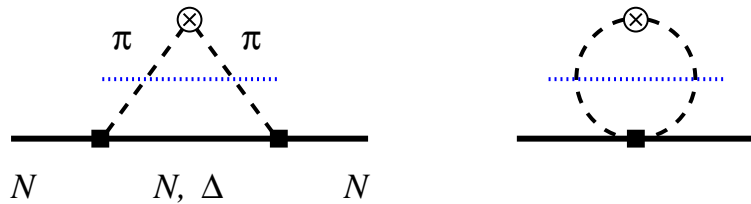
$$\rho(b) = \int_{4m_\pi^2}^{\infty} \frac{dt}{2\pi^2} K_0(\sqrt{t}b) \text{Im } F_1(t)$$

$K_0 \sim e^{-b\sqrt{t}}$  exponential suppression of large  $t$

Distance  $b$  selects masses  $\sqrt{t} \sim 1/b$ : “Filter”  
Cf. Borel transformation in QCD sum rules. Strikman, CW 10

Peripheral  $\rho(b) \longleftrightarrow$  low-mass hadronic states

# Chiral component: Spectral function



- Spectral function near threshold  
 $t - 4M_\pi^2 \sim \text{few } M_\pi^2$  from  $\chi$ PT Isovector

Relativistic chiral Lagrangian

Becher, Leutwyler 99

Heavy-baryon expansion not convergent due to subthreshold singularity

Becher, Leutwyler 99; Kaiser 03

- Asymptotic behavior for  $b \rightarrow \infty$

Strikman, CW 10

Yukawa tail with range  $2M_\pi$

Pre-exponential factor strongly varying

$$\rho^V(b) \propto \frac{\exp(-2M_\pi b)}{(2M_\pi b)^3} \times \left( 1 + \text{terms } \frac{M_N^2}{M_\pi^3 b} + \dots \right)$$

- Include intermediate  $\Delta$

Large coupling!

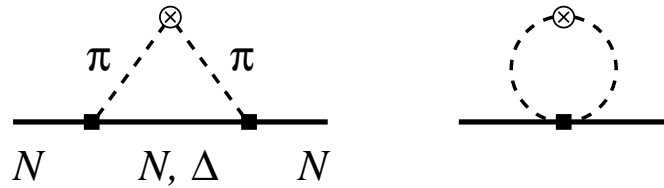
Large- $N_c$  limit of QCD:  $g_V \sim N_c^0$  requires cancellation between  $N$  and  $\Delta$  Cohen 95

$$M_N, M_\Delta \sim N_c$$

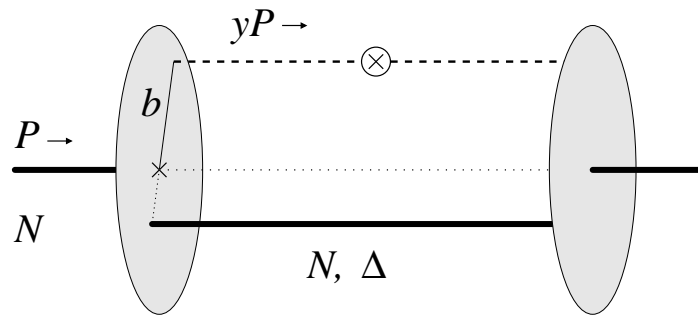
$$M_N - M_\Delta \sim 1/N_c \quad \text{degenerate}$$

Limits  $b \rightarrow \infty$  and  $M_\pi/M_N \rightarrow 0$  do not commute: Do not expand in  $M_\pi/M_N$

# Chiral component: Partonic formulation



Equivalent!



Equivalence of invariant and partonic formulation of  $\chi$ PT

- Partonic representation of chiral component

$$\rho^V(b) = \int_0^1 dy \left[ \frac{2}{3} f_{\pi N}(y, b) - \frac{1}{3} f_{\pi \Delta}(y, b) \right]$$

$f_{\pi B}(y, b)$  pion longitudinal momentum and transverse coordinate density

Strikman, CW 09: Chiral contributions to parton densities

$y \sim M_\pi/M_N$  parametrically small

- Overlap of light-front wave functions

$$f_{\pi N}(y, b) = \sum_{L=0,1} \psi_{\pi N}^{L*}(y, b/\bar{y}) \psi_{\pi N}^L(y, b/\bar{y}) + \delta(y) \text{ term}$$

Orbital angular momentum of  $\pi N$  system!

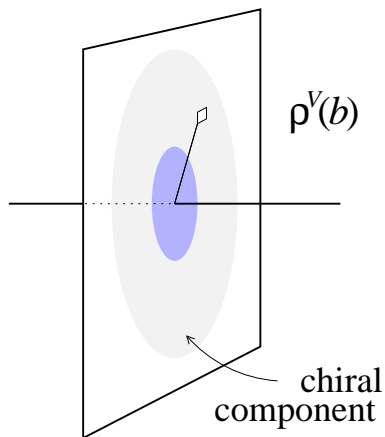
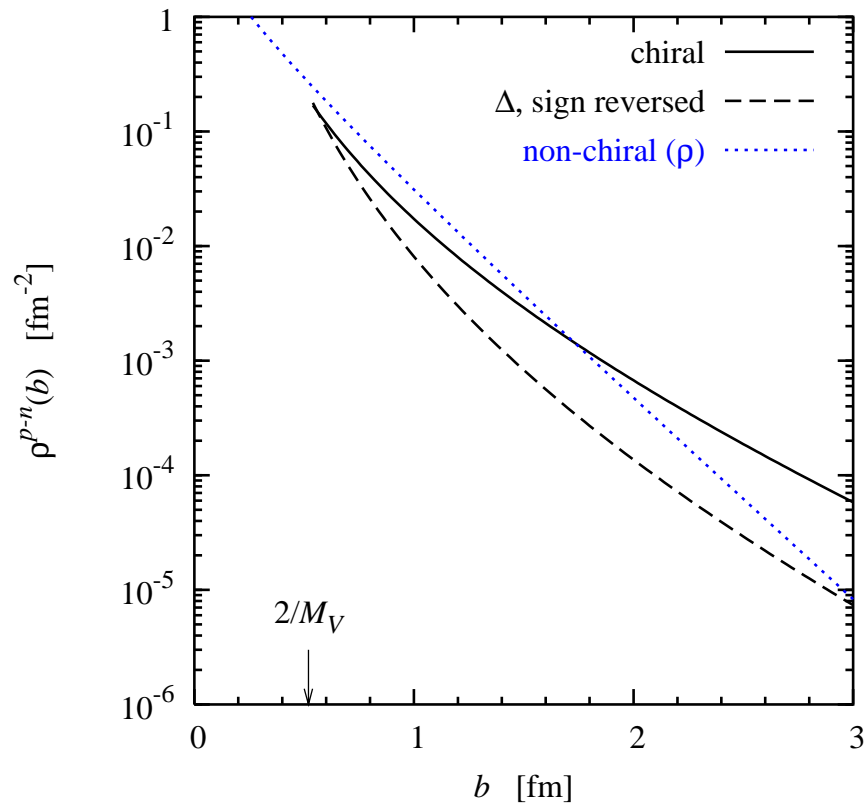
Also magnetic density:  $\Delta L = 1$ . Granados, CW 12

- Contact terms  $\delta(y)$  “Light-front zero modes”

Represent effect of high-mass intermediate states in time-ordered PT Granados, CW 12

Coefficient  $(1 - g_A^2)$ : Compositeness of nucleon

# Chiral component: Numerical evaluation



- Numerical evaluation

Contact term  $\delta(y)$  in  $\pi N$  contributes  $< 10\%$  at  $b > 1$  fm

$\pi\Delta$  negative, suppressed at large  $b$

- Compare with “non-chiral” density

Simple estimate:  $\rho$  meson pole

Chiral component dominates only at  $b > 2$  fm . . . surprisingly large!

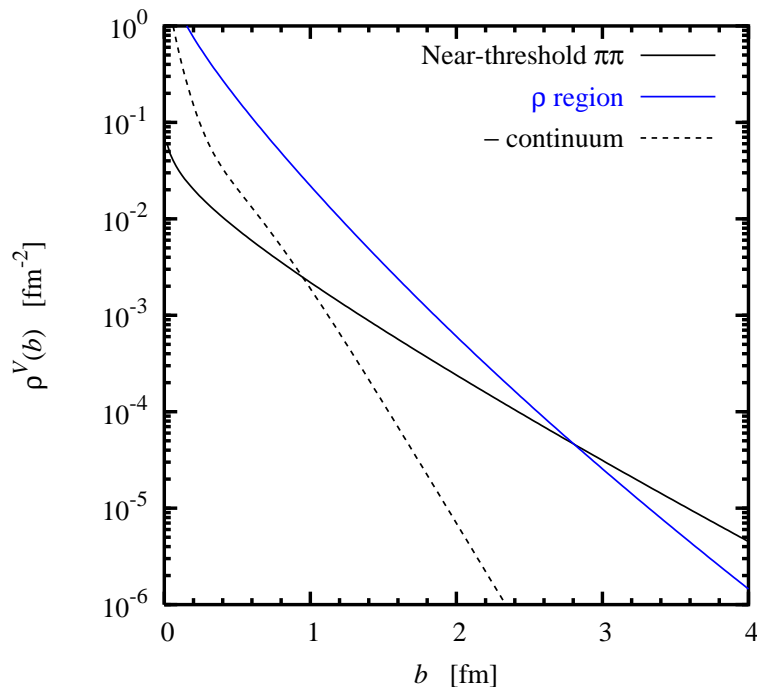
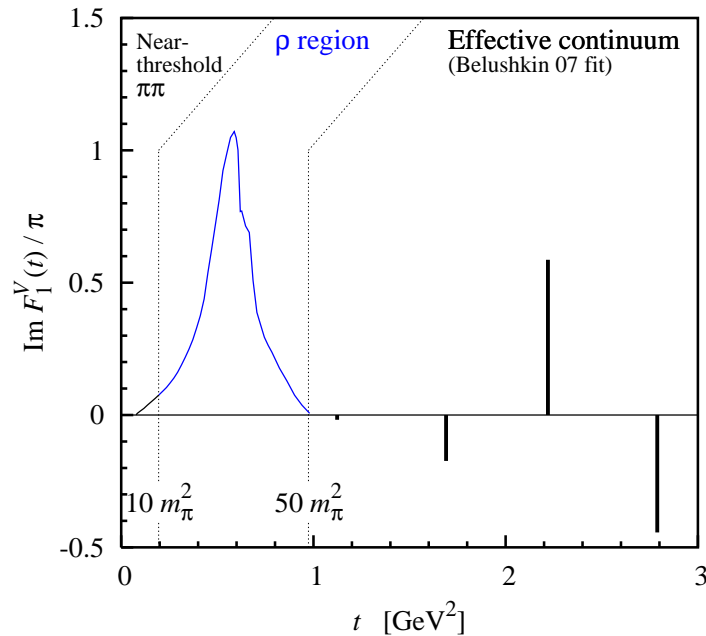
- Impact parameter  $b$  as new way of identifying chiral component

Model-independent!

$b$  observable in exclusive processes, objectively defined  $\leftrightarrow$  Breit frame radius



# Chiral component: Spectral analysis



- Empirical isovector spectral function

Near-threshold  $\pi\pi$  from chiral dynamics  
 $\rho$  region from  $\pi\pi$  phase shifts Höhler 76  
 High-mass continuum from form factor fits  
 Belushkin, Hammer, Meissner 07

- Spectral analysis of isovector density

Strikman, CW 10; Miller, Strikman, CW 11

Near-threshold  $\pi\pi$  relevant only at  $b > 2$  fm

Intermediate  $b = 0.5 - 1$  fm dominated by  $\rho$ ,  
 with  $\sim 10\%$  correction from higher masses  
 “Vector dominance” quantified

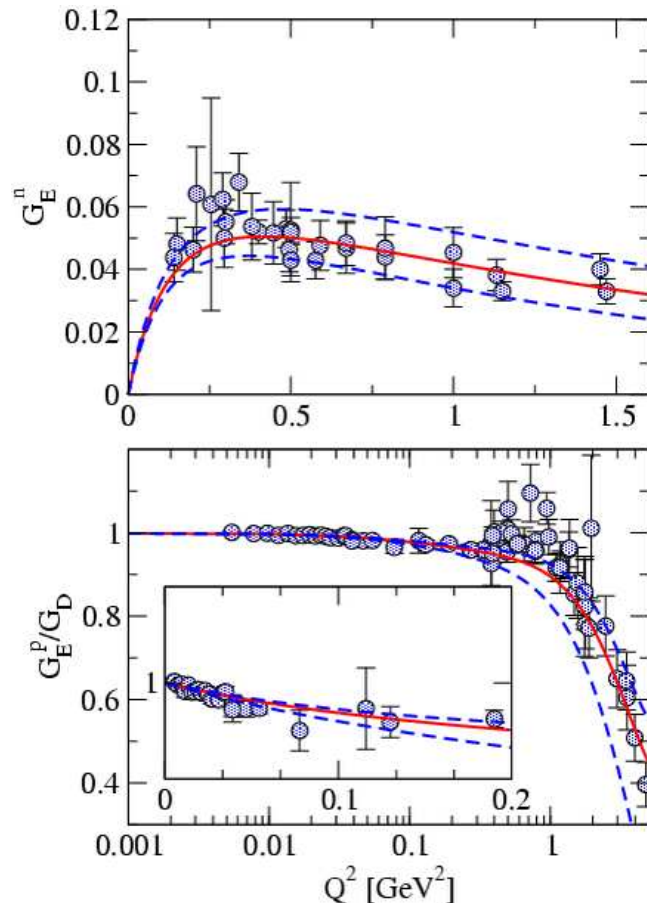
- Isoscalar density

$\omega$  dominates at  $b > 1.5$  fm.

Large cancellations between  $\omega$  and  
 higher-mass states at  $b = 0.5 - 1$  fm

Model-independent identification of chiral component, “vector dominance” in QCD

# Chiral component: Effect on form factors Simple estimates!



Dispersion fit Belushkin, Hammer, Meissner 07  
New data from Bates, MAMI, JLab

- Moments of transverse charge density

$$\begin{aligned}\langle b^2 \rangle &= \int d^2b \, b^2 \rho(b) = 4 F_1'(0) \\ \langle b^4 \rangle &= 32 F_1''(0)\end{aligned}$$

- Contribution of chiral component isovector

$$\langle b^2 \rangle_{\text{chiral}} \approx 0.2 \times \langle b^2 \rangle_{\text{fit}} \quad \text{small}$$

$$\langle b^4 \rangle_{\text{chiral}} \approx 1.5 \times \langle b^2 \rangle_{\text{fit}}^2 \quad \text{sizable}$$

Chiral component should be visible in “unnatural” second and higher derivatives of FF at  $Q^2 = 0$

Can we extract it?

- Analyticity of form factor fit is essential

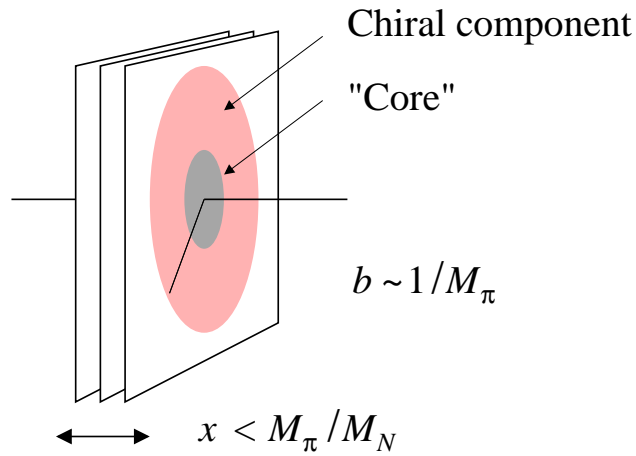
Needs dispersion analysis: Belushkin et al. 07

- Affects extrapolation to  $t \rightarrow 0$

CLAS/PRIMEX 12 GeV experiment at  $Q^2 = 10^{-4} - 10^{-2} \text{ GeV}^2$

PR12-11-106 Gasparian et al.

# Parton distributions: Chiral component

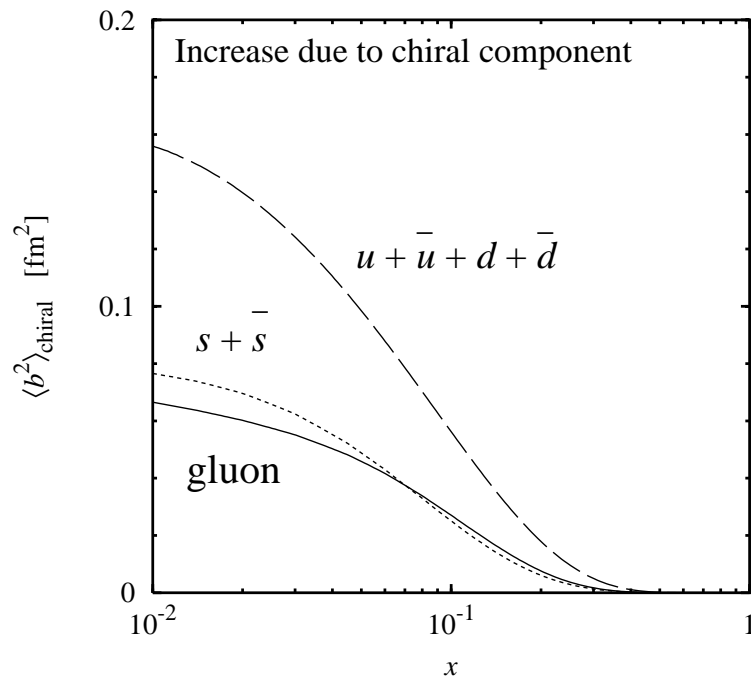


- Large-distance component of quark/gluon densities at  $b \sim M_\pi^{-1}$  and  $x < M_\pi/M_N$   
Strikman, CW PRD69:054012,2004; PRD80:114029,2009

Model-independent, calculable: "Yukawa tail"

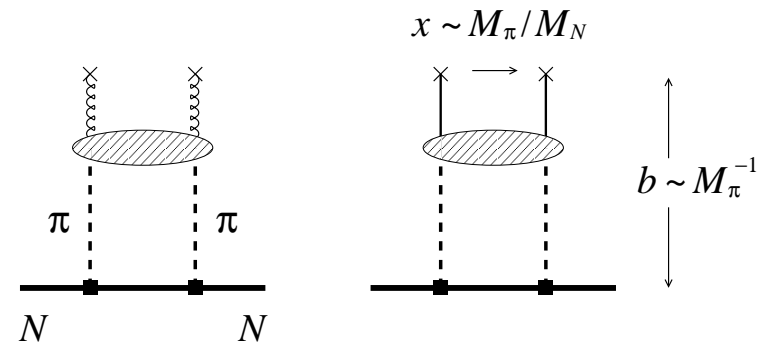
Strong in isoscalar quarks and gluons, suppressed in isovector  $\bar{d} - \bar{u}$

Small fraction of total number:  
Most partons sit at distances  $b \lesssim 0.5$  fm

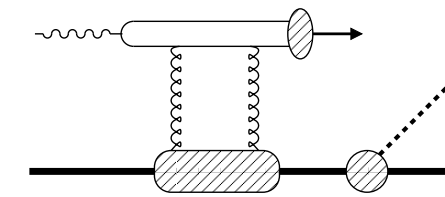
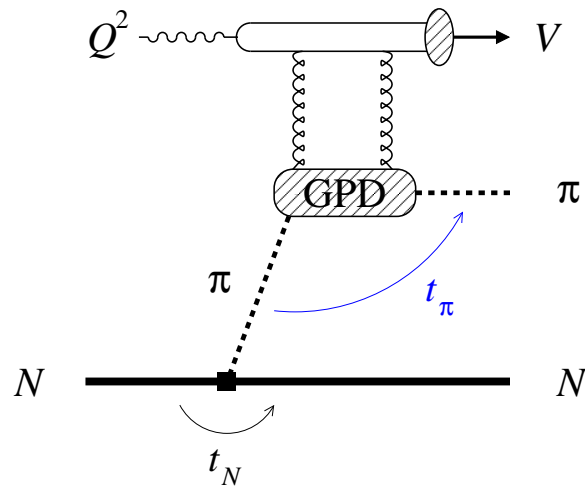


- Increase of nucleon's transverse radius below  $x \sim M_\pi/M_N$

Larger for quarks than for gluons!



# Parton distributions: Peripheral hard processes



suppressed!

- Hard exclusive process on peripheral pion  
Strikman, CW PRD69:054012,2004

$$k_\pi^2 \sim M_\pi^2 \text{ quasi-real}$$

Requires  $x \ll M_\pi/M_N \sim 0.1$

- Kinematics with  $p_T(\pi) \gg p_T(N)$  suppresses production on nucleon

$$F_{\pi NN}(t) \text{ softer than } \text{GPD}_\pi(t)$$

- Probe gluon GPD in pion at  $|t_\pi| \sim 1 \text{ GeV}^2$

Fundamental interest

Moments calculable in Lattice QCD

- Requires detection of forward nucleon and moderate- $p_T$  pion

Feasible with Electron-Ion Collider EIC

Direct probe of chiral component of nucleon's partonic structure!

# Summary

- Transverse densities relate elastic FFs to universal partonic structure

Impact parameter  $b$  “observable” unlike Breit frame radius

Dispersion representation connects partonic structure at given  $b$  with hadronic exchange mechanisms

- Chiral component of transverse densities at large  $b$

Chiral expansion justified by  $b \sim M_\pi^{-1}$ . . . new arena for  $\chi$ PT

Equivalence of invariant and light–front formulations of  $\chi$ PT

Chiral and non–chiral components identified by spatial size

Chiral component dominant only at large  $b \gtrsim 2$  fm

- Dispersion fits to FF data contain much interesting information

Correct analyticity crucial for studying peripheral densities

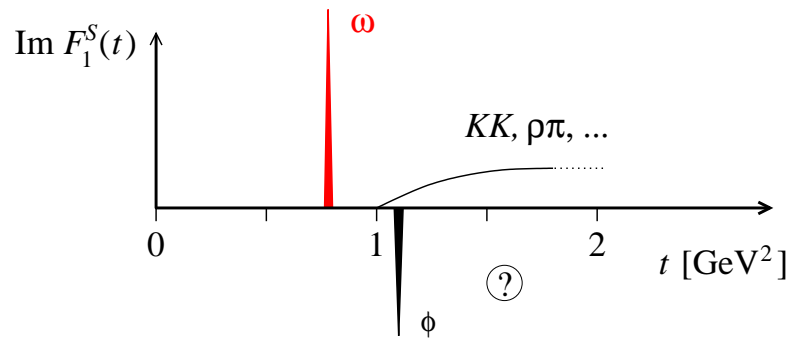
Should be updated with new data JLab 6/12 GeV

- Many interesting extensions

Current vs. charge densities,  $x$ –dependent structures

Supplementary material

# Spectral analysis: Isoscalar density

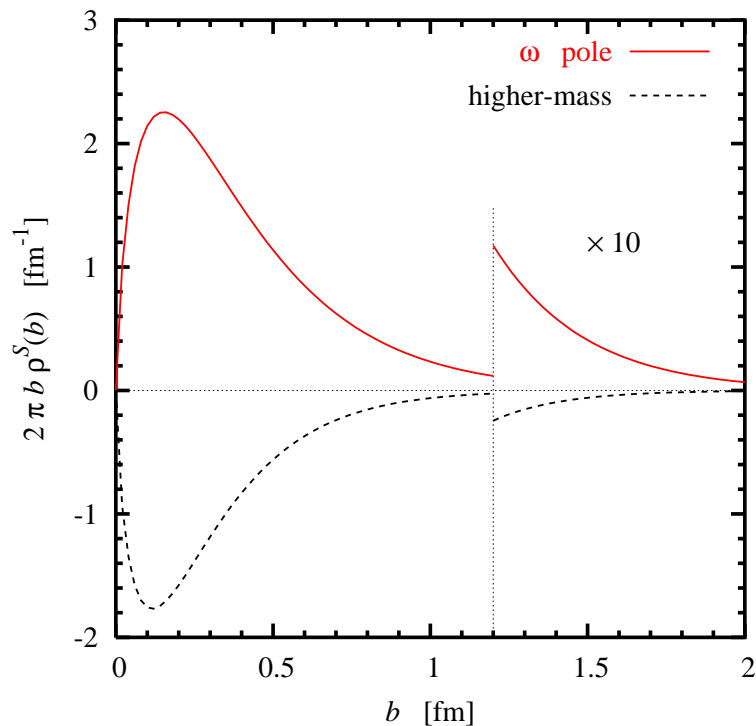


- Isoscalar spectral function

$\omega$  exhausts strength below  $1 \text{ GeV}^2$   
 Non-resonant  $3\pi$  negligible

Large negative strength above  $1 \text{ GeV}^2$ ,  
 dynamical origin unclear  
 $\phi NN$  coupling  $\leftrightarrow$   $s\bar{s}$  content of nucleon

High-mass continuum from form factor fits  
 Belushkin, Hammer, Meissner 07



- Spectral analysis of isoscalar density  
 Miller, Strikman, CW 11

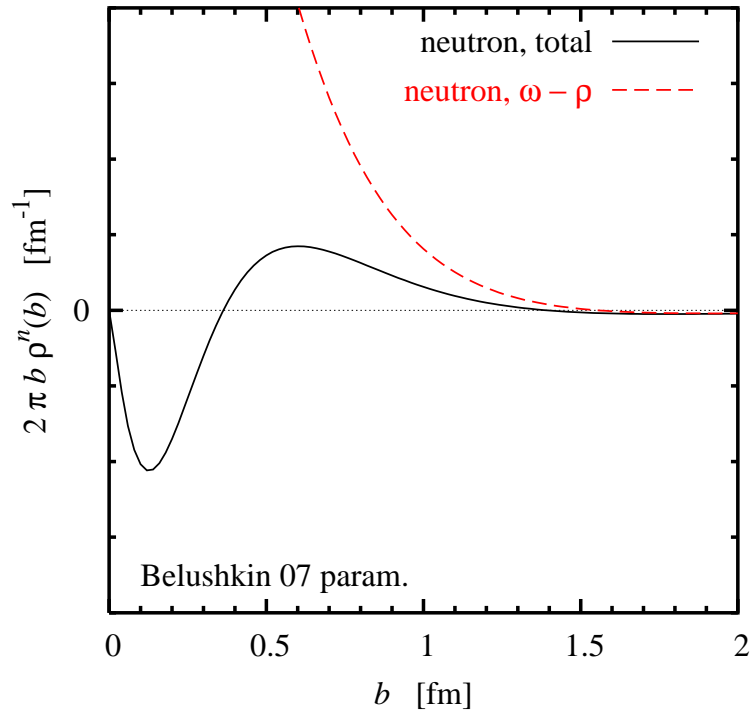
$\omega$  dominates at  $b > 1.5 \text{ fm}$   
 Fit uncertainty in  $\omega NN$  coupling  $\pm 15\%$

Large cancellations between  $\omega$  and  
 higher-mass states at  $b = 0.5 - 1 \text{ fm}$

- Impact of future form factor data

Sensitivity to  $\omega NN$  coupling broadly  
 distributed at spacelike  $|t| \lesssim 1 \text{ GeV}^2$   
 Does not require measurements at extremely small  $|t|$

# Spectral analysis: Neutron density



- Spectral analysis of neutron density

$\omega - \rho$  alone gives large positive density!

Substantially reduced by higher-mass states in isoscalar spectral function

Neutron form factor measurements can help to determine isoscalar spectral function  
→  $\phi NN$  coupling,  $s\bar{s}$  in nucleon