

Extracting Nucleon Polarisabilities from Proton and Deuteron Compton Scattering



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- 1 Why Polarisabilities?
- 2 Analysis
- 3 Consequences
- 4 Concluding Questions



How do constituents of the nucleon react to external fields?
How to reliably extract **neutron** and **spin** polarisabilities?



Comprehensive Theory Effort:

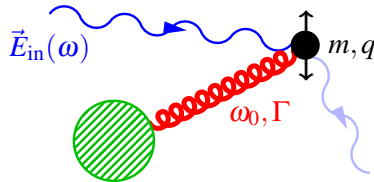
hg, J. McGovern (Manchester), D. R. Phillips (Ohio U)

+ G. Feldman (GW): **Prog. Part. Nucl. Phys.** 2012

Precursors: R. Hildebrandt/T. R. Hemmert/B. Pasquini/hg...2000-05, ...,
Beane/Malheiro/McGovern/Phillips/van Kolck 1999-2005; Choudhury Shukla/Phillips 2005-08
Friar 1975, Arenhövel/Weyrauch 1980-83, Karakowski/Miller 1999, Levchuk/L'vov 1994-2000

1. Why Polarisabilities?

Example: induced electric dipole radiation from harmonically bound charge, damping Γ Lorentz/Drude 1900/1905



$$\vec{d}_{\text{ind}}(\omega) = \frac{q^2}{m} \frac{1}{\omega_0^2 - \omega^2 - i\Gamma\omega} \vec{E}_{\text{in}}(\omega)$$

$$=: 4\pi \alpha_{E1}(\omega)$$

$$\mathcal{L}_{\text{pol}} = 2\pi \left[\underbrace{\alpha_{E1}(\omega)\vec{E}^2 + \beta_{M1}(\omega)\vec{B}^2}_{\text{electric, magnetic scalar dipole}} + \dots \right]$$

Energy- (ω)-dependent multipole-decomposition dis-entangles scales, symmetries & mechanisms of interactions with & among constituents.

Clean, perturbative probe explores: – spontaneously broken **chiral symmetry** of pion cloud

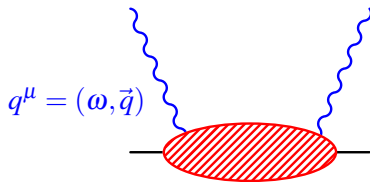
– $\Delta(1232)$ **resonance** properties

– constituents of the **nucleon spin**

– **proton** \leftrightarrow **neutron** iso-spin breaking:

\implies elmag. p-n self-energy difference from $\beta_{M1}^p - \beta_{M1}^n$ Walker-Loud/...2012

– 2γ contribution to Lamb shift in muonic H (β_{M1})

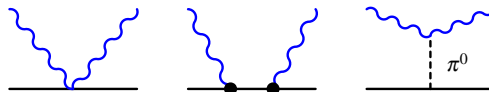


2. Analysis

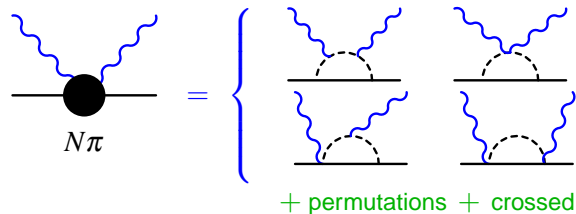
(a) Microscopic processes: χ EFT with Δ

Bernard/Kaiser/Meißner 1994, ...
 hg/Hemmer/Hildebrandt/Pasquini 2003-
 hg/McGovern/Phillips 2012

- **Powell**: point-like spin- $\frac{1}{2}$ with anom. mag. moment



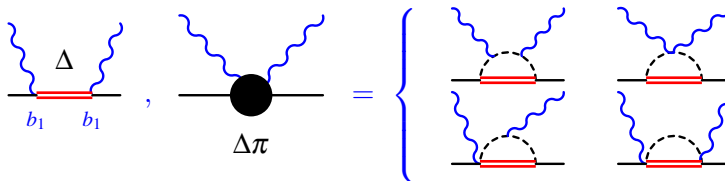
- **Chiral Dynamics**:
Cusp at $1-\pi$ production threshold.



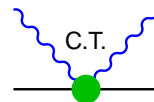
- **Large Δ para-magnetism**
 from N -to- Δ $M1$ transition.

$$\delta \bar{\beta}_\Delta \approx +10 \times 10^{-4} \text{ fm}^3$$

$$\Rightarrow \bar{\beta}^p \approx 2 \text{ "fine-tuned"}$$



- **Core Contribution** only for α_{E1} , β_{M1} : ω -indep. dia-magnetism.




Different scales and mechanisms dis-entangled by ω -dependence.

(b) Including the $\Delta(1232)$ in χ EFT

$\Delta(1232)$ lowest hadronic excitation **above** 1-pion threshold m_π , **below** χ EFT breakdown scale $\Lambda_\chi \approx 1000$ MeV

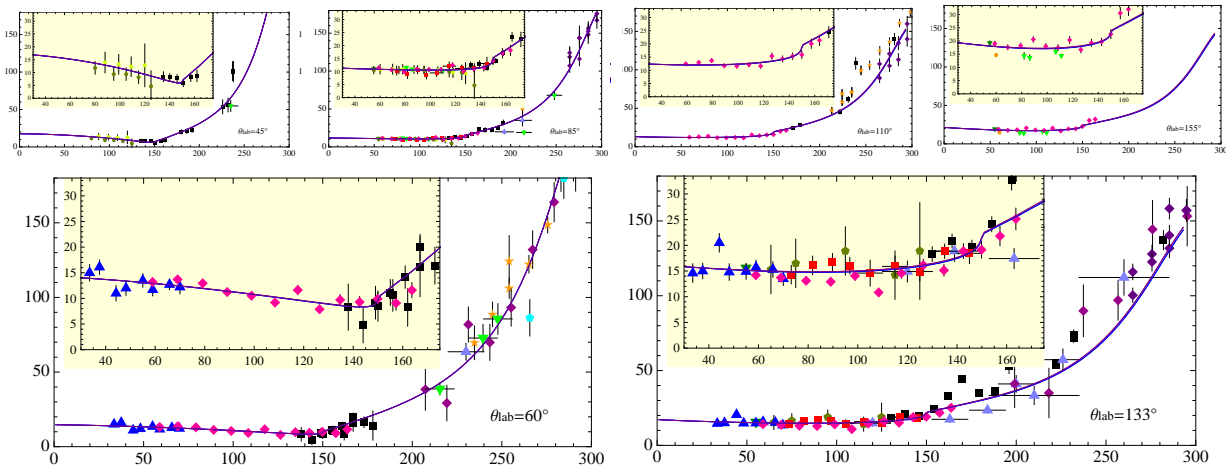
$$\Rightarrow \text{Expand in } \delta = \frac{M_\Delta - M_N}{\Lambda_\chi} \approx \sqrt{\frac{m_\pi}{\Lambda_\chi}} = \frac{p_{\text{typ}}}{\Lambda_\chi} \ll 1 \text{ (numerical fact) } \text{Pascalutsa/Phillips 2002}$$

Higher Energy: Δ propagator  $\propto \frac{1}{E - (M_\Delta - M_N)}$ enhanced as $E \rightarrow M_\Delta - M_N$

$$\Rightarrow \text{re-order \& re-sum} \quad \text{====} \rightarrow \text{====} + \text{====} \overset{\text{dashed loop}}{\text{====}} + \dots = \frac{1}{E - (M_\Delta - M_N) - \overset{\text{dashed loop}}{\text{====}}} + \text{relativity}$$

Complete at $\mathcal{O}(e^2 \delta^3)$ (N²LO). Probe non-zero Δ width, $M1$ and $E2$ transition strengths.

(c) Determine Proton Polarisabilities from All Low-Energy Data



Iterate: > 200 MeV: Δ parameters ($b_1 = 3.7$; $b_2/b_1 = -0.34$ fixed) \implies < 170 MeV: polarisabilities

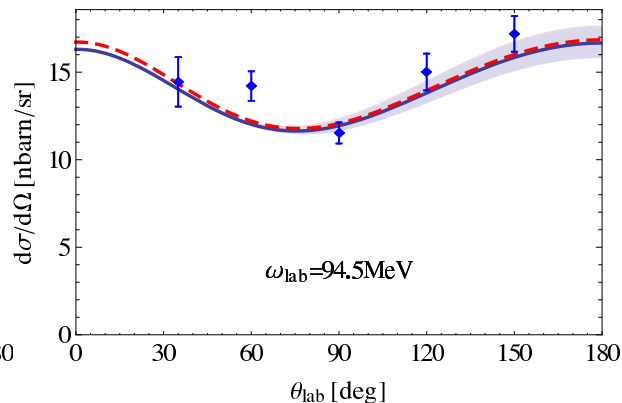
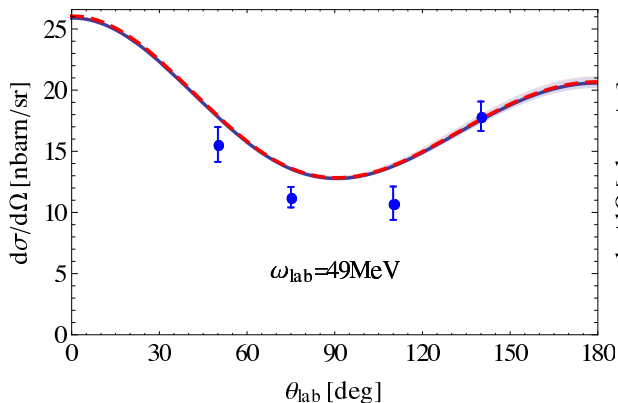
	$\bar{\alpha}_{E1}^p [10^{-4} \text{ fm}^3]$	$\bar{\beta}_{M1}^p [10^{-4} \text{ fm}^3]$	$\chi^2/\text{d.o.f.}$
LO, parameter-free Bernard/Kaiser/Meißner 1994	12.5	1.25	no fit
NLO, free fit	$10.5 \pm 0.5_{\text{stat}} \pm 0.8_{\text{theory}}$	$2.7 \pm 0.5_{\text{stat}} \pm 0.8_{\text{theory}}$	106.1/124
NLO, Baldin constrained $\bar{\alpha}_{E1}^s + \bar{\beta}_{M1}^s = 13.8 \pm 0.4$	$10.7 \pm 0.3_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$	$3.1 \mp 0.3_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$	106.5/125
Olmos de Leon MAMI/global 2001	$12.1 \pm 1.2_{\text{stat+model}} \pm 0.4_{\Sigma}$	$1.6 \mp 1.2_{\text{stat+model}} \pm 0.4_{\Sigma}$	

(d) Determine Neutron Polarisabilities from all Deuteron Data

Illinois \circ , Lund \star , Saskatoon \blacklozenge

--- $N\pi + \Delta$ free fit

— $N\pi + \Delta$ + stat. error, Baldin constrained



$$\bar{\alpha}_{E1}^s [10^{-4} \text{ fm}^3]$$

$$\bar{\beta}_{M1}^s [10^{-4} \text{ fm}^3]$$

$$\chi^2/\text{d.o.f.}$$

free fit

$$10.5 \pm 2.0_{\text{stat}} \pm 0.8_{\text{theory}}$$

$$3.6 \pm 1.0_{\text{stat}} \pm 0.8_{\text{theory}}$$

24.3/24

Baldin constrained

$$\bar{\alpha}_{E1}^s + \bar{\beta}_{M1}^s = 14.5 \pm 0.3$$

$$10.9 \pm 0.9_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$$

$$3.6 \mp 0.9_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$$

24.4/25

proton (Baldin) hg/...2012

$$10.7 \pm 0.3_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$$

$$3.1 \mp 0.3_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$$

106.5/125

previous ranges:

$$[6 \dots 18]$$

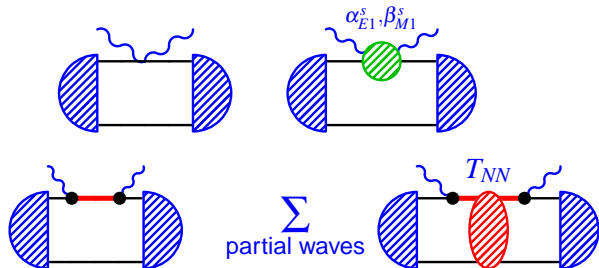
$$[-4 \dots 9]$$

⇒ neutron \approx proton polarisabilities

(e) Deuteron Compton Scattering at $\omega = 0 \dots 200$ MeV

One-body: electric, magnetic moment couplings

$$\omega \sim \frac{Q^2}{M} \approx 20 \text{ MeV} \quad \omega \sim Q \approx 100 \text{ MeV}$$



LO, N³LO

LO, ↗ NLO

LO

↘ NLO, N³LO

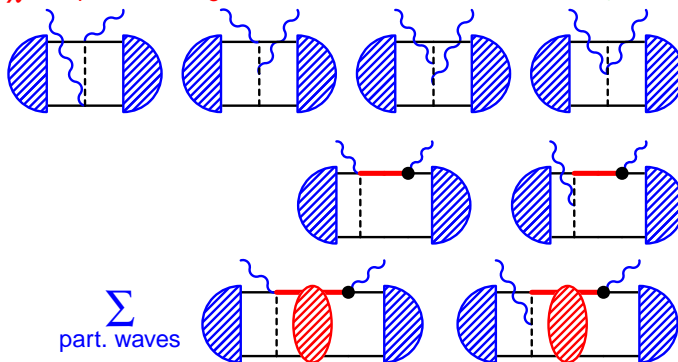
2N correlated

$$\frac{i}{B_d \pm \omega - \frac{q^2}{M}}$$

uncorrelated

χEFT pion-exchange currents:

Beane et al. 1999-2005; hg/...2005



NLO

→ NLO

NLO

↘ N²LO

NLO

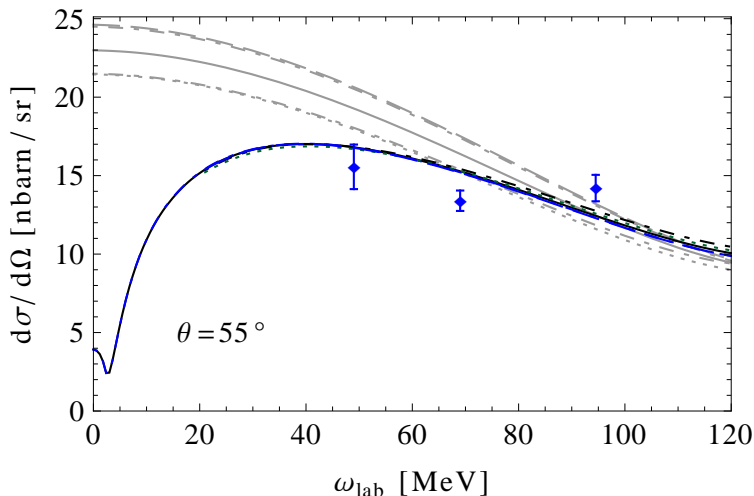
↘ N³LO, pert.

Full LO T_{NN} pivotal for current conservation. Arenhövel 1980

Low-Energy Theorem: Thomson limit $\mathcal{A}(\omega = 0) = -\frac{e^2}{M_d} \vec{\varepsilon} \cdot \vec{\varepsilon}'$.

Thirring 1950, Friar 1975, Arenhövel 1980: Thomson limit \iff current conservation \iff gauge invariance.

Exact Theorem \implies At each χ EFT order \implies Checks numerics.



Significantly reduces cross section for $\omega \lesssim 70$ MeV.

Urbana, Lund data

Numerically confirmed to $\lesssim 0.2\%$, irrespective of deuteron wave function & potential.

model-independence

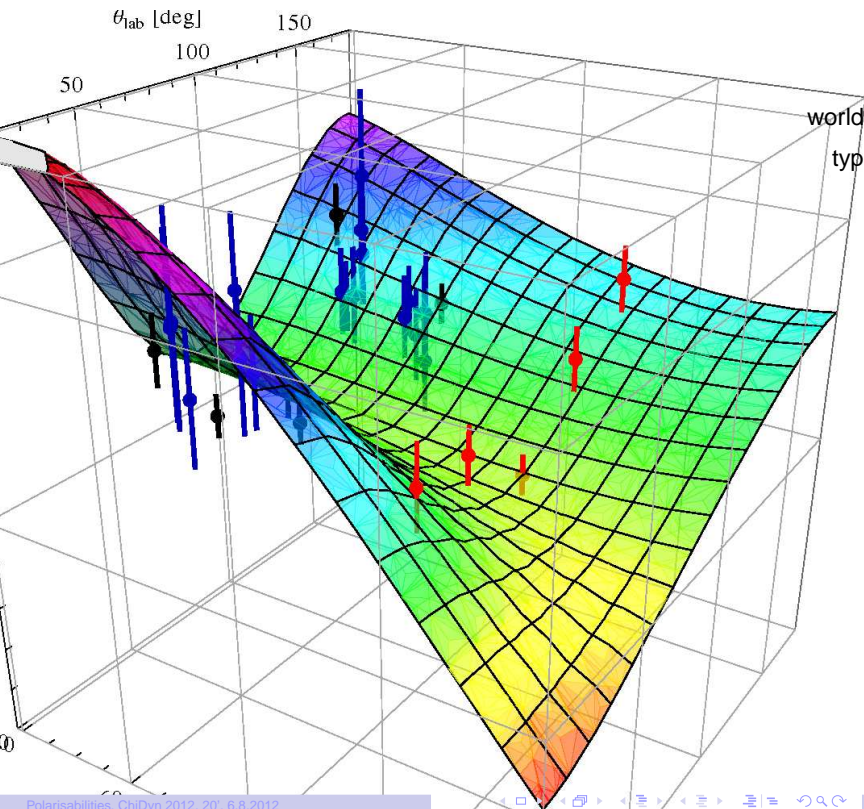
Wave function & potential dependence significantly reduced even as $\omega \rightarrow 150$ MeV \implies **gauge invariance.**

3. Consequences

(a) Un-Polarised Deuteron Compton Scattering

hg/...2005-2010

hg/McGovern/Phillips/Feldman PPNP 2012



$$\left. \frac{d\sigma}{d\Omega} \right|_{\text{lab}} \quad [\text{nbarn/sr}]$$

world data: **29 points** Illinois, Saskatoon, Lund

typ. stat.+uncorrel. sys.: $\pm[7 \dots 10]\%$

typ. correl. sys.: $\pm[3 \dots 5]\%$

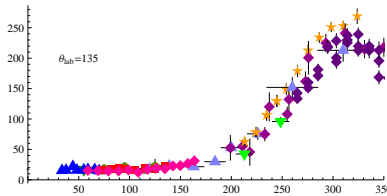
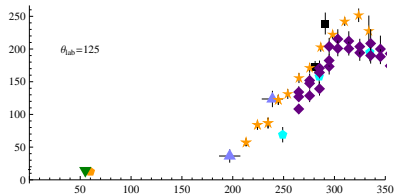
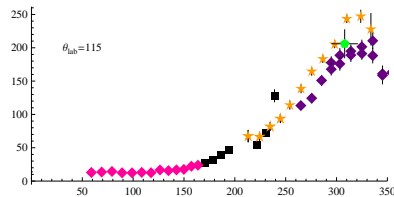
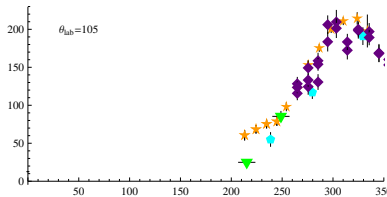
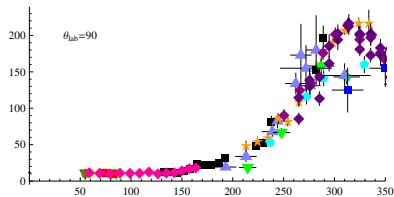
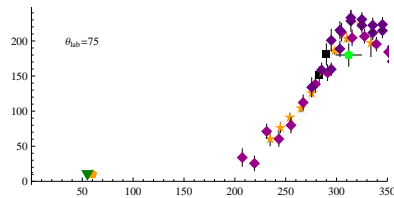
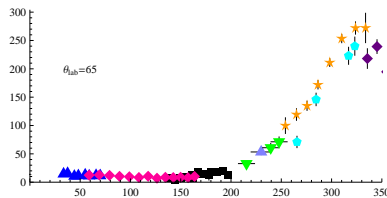
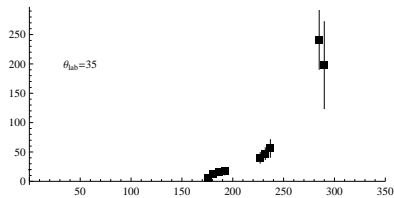
limited energy & angle coverage

Upcoming

MAXlab: [60; 200] MeV, data taken;

HIγS: [65; 100] MeV, approved

(b) Proton Compton Data



~ 300 data, mostly 1991-2001

New effort for better data:
MAMI, MAXlab, HIγS,...

Gaps for: $\omega \in [160; 250]$ MeV; $\theta \rightarrow 0^\circ$: Baldin check; $\theta \rightarrow 180^\circ$ for $\Delta(1232)$!

Small quoted systematics \implies tensions: MAMI vs. LEGS, but also others $\implies \frac{\chi^2}{\text{d.o.f.}} \approx 1$ cannot be achieved.

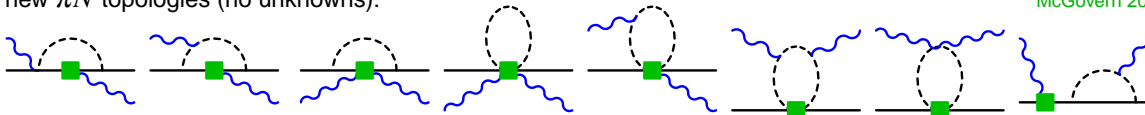
Not more, but more reliable data needed for unpolarised proton.

– higher-order πN interactions in same topologies from $\mathcal{L}_{\chi EFT}$

only unknowns: $\delta\alpha_{E1}, \delta\beta_{M1}$

– new πN topologies (no unknowns):

McGovern 2001



– No new $\Delta(1232)$ contributions.

– πN loops push intermed. angles too high \implies Fit γ_{M1M1} ?? Why only this? 6 parameters: flat directions.

– Preliminary: Small changes $N^2LO \rightarrow N^3LO \implies$ Extraction stable, converges.

(d) Understanding Energy Dependence

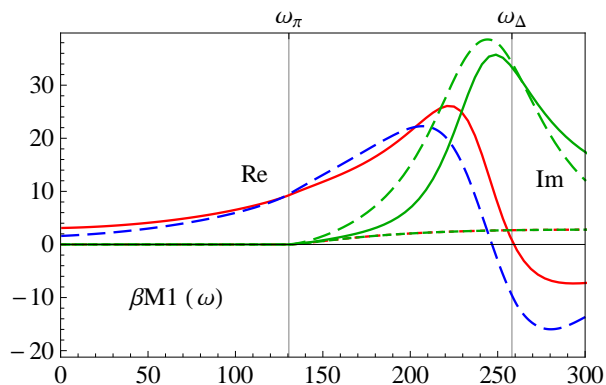
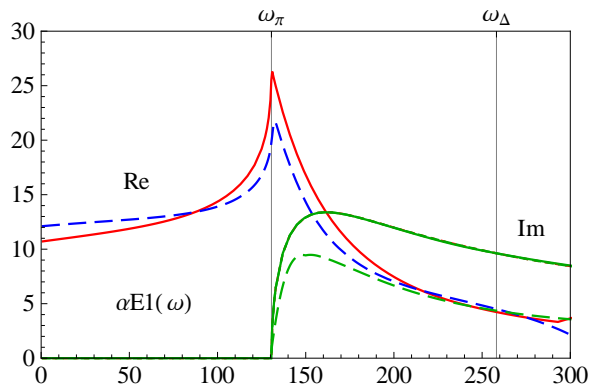
Dynamical Polarizabilities: Multipole decomposition of real Compton scattering at **fixed energy**.

$$2\pi \left[\alpha_{E1}(\omega) \vec{E}^2 + \beta_{M1}(\omega) \vec{B}^2 + \gamma_{E1E1}(\omega) \vec{\sigma} \cdot (\vec{E} \times \dot{\vec{E}}) + \gamma_{M1M1}(\omega) \vec{\sigma} \cdot (\vec{B} \times \dot{\vec{B}}) + \dots \right]$$

Neither more nor less information about **response** of constituents, but **more readily accessible**.

$\alpha_{E1}(\omega)$: Pion cusp well captured by single- $N\pi$.

$\beta_{M1}(\omega)$: para-magnetic N -to- Δ $M1$ -transition.



Re: refraction; Im: absorption

(e) What is the Physics of the Fit Parameters?

Pion & Δ dynamics capture ω -dependence.

\Rightarrow **Hypothesis:** $\delta\alpha_{E1}, \delta\beta_{M1}$ parameterise nucleon core.

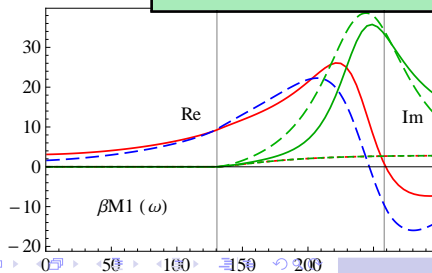
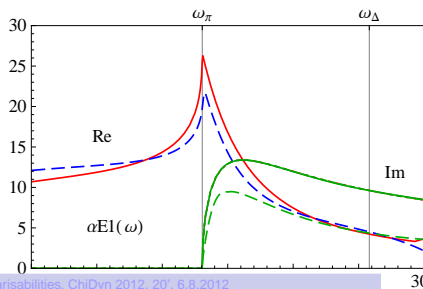
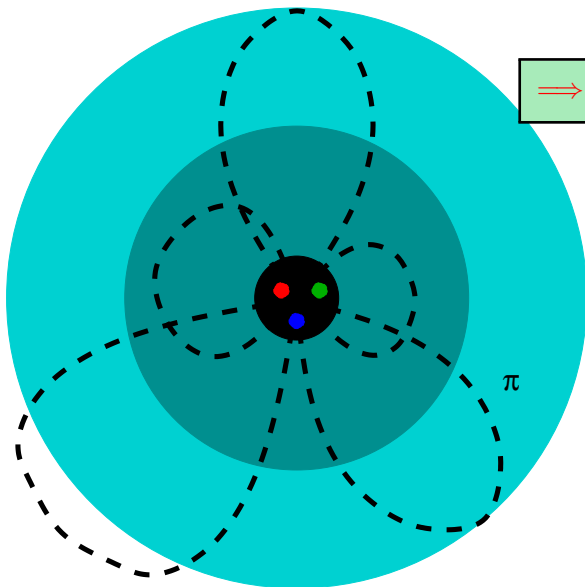
- $\chi^2/\text{d.o.f.} \approx 1 \Rightarrow \sim \omega$ -indep. $\lesssim 250$ MeV.
- **Naively NLO:** $|\delta\alpha_{E1}|, |\delta\beta_{M1}| \sim \frac{\alpha}{\Lambda_\chi^2 M} \approx 1$
- **Anomalously large:**

$$\delta\bar{\alpha}_{E1} = -5.4, \quad \delta\bar{\beta}_{M1} = -10.1$$
- **Cancel $\Delta, \Delta\pi$ effects** in static pols.
- **Iso-spin independent:** proton \approx neutron values.
- **Not necessary for spin-polarisabilities γ_i ?**

Speculation:

Guichon, Schumacher

Correlated 2π exchange in 0^{++} t -channel??



4. Concluding Questions

Dynamical polarisabilities: Energy-dependent multipole-decomposition dis-entangles scales, symmetries & mechanisms of interactions with & among constituents: χ iral symmetry of pion-cloud, iso-spin breaking, $\Delta(1232)$ properties, nucleon spin-constituents.

$\implies \chi$ EFT: unified frame-work off light nuclei: model-independent, systematic, reliable errors.

Scalar Dipole Polarisabilities from *all Compton data below 200 MeV*:

proton	$\bar{\alpha}^p = 10.7 \pm 0.3_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$	$\bar{\beta}^p = 3.1 \mp 0.3_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$
iso-scalar	$\bar{\alpha}^s = 10.9 \pm 0.9_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$	$\bar{\beta}^s = 3.6 \mp 0.9_{\text{stat}} \pm 0.2_{\Sigma} \pm 0.8_{\text{theory}}$
neutron	$\bar{\alpha}^n = 11.1 \pm 1.8_{\text{stat}} \pm 0.4_{\Sigma} \pm 0.8_{\text{theory}}$	$\bar{\beta}^n = 4.1 \mp 1.8_{\text{stat}} \pm 0.4_{\Sigma} \pm 0.8_{\text{theory}}$

Theory To-Do List: explore host of observables: expansion $\frac{p_{\text{typ}}}{\Lambda_{\chi}} \ll 1$ for credible error-bars.

math notebooks

- **One Nucleon:** χ EFT with $\Delta(1232)$ at δ^4 into resonance region
- **Deuteron:** extend beyond 1π -threshold into $\Delta(1232)$ region; break-up
- ^3He & heavier: Thomson limit; into $\Delta(1232)$ region; break-up

near-done

now focus of attention

near-term

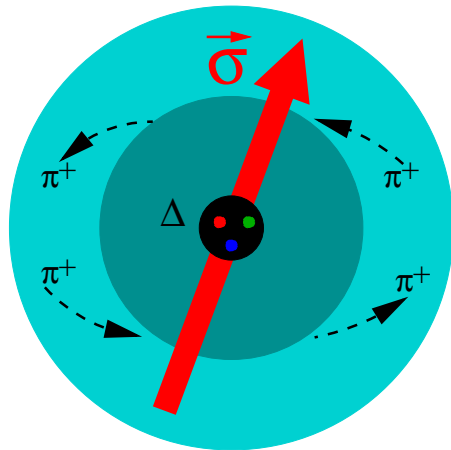
Data Needed: cross-sections & asymmetries – reliable systematics. \implies **spin-polarisabilities.**

Clean probe to explore the strong force at low energies.

(a) Spin-Polarisabilities: Nucleonic Bi-Refringence and Faraday Effect

Response of **spin-degrees of freedom in nucleon** to real photon of definite multipolarity and non-zero energy ω .

⇒ **Multipole Analysis.**



$$\mathcal{L}_{\text{pol}} = 4\pi N^\dagger \times$$

$$\left\{ \frac{1}{2} \left[\alpha_{E1}(\omega) \vec{E}^2 + \beta_{M1}(\omega) \vec{B}^2 \right] \right. \quad \text{scalar dipole}$$

$$+ \frac{1}{2} \left[\gamma_{E1E1}(\omega) \vec{\sigma} \cdot (\vec{E} \times \dot{\vec{E}}) + \gamma_{M1M1}(\omega) \vec{\sigma} \cdot (\vec{B} \times \dot{\vec{B}}) \right] \quad \text{“pure” spin-dependent dipole}$$

$$- 2 \gamma_{M1E2}(\omega) \sigma_i B_j E_{ij} + 2 \gamma_{E1M2}(\omega) \sigma_i E_j B_{ij} \quad \text{“mixed” spin-dependent dipole}$$

$$+ \dots \left. \vphantom{\frac{1}{2}} \right\} N \quad \text{quadrupole etc.}$$

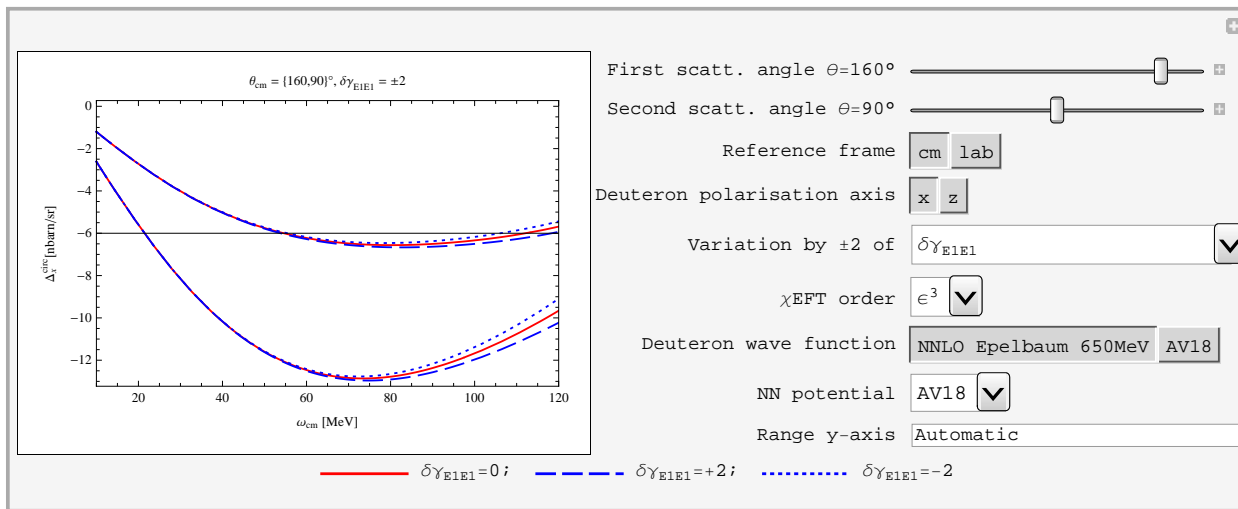
$$E_{ij} := \frac{1}{2} (\partial_i E_j + \partial_j E_i) \text{ etc.}$$

(b) Plethora of Polarisation Observables in Compton Scattering

2×6 observables, 6 polarisabilities, 3 kinemat. variables ω, θ, ϕ + additional constraints

⇒ Interactive *mathematica* 8.0 notebooks. [hg/...2010-](#)

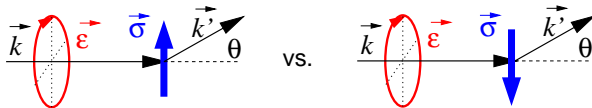
Example Screenshot



(c) Spin-Polarisabilities from Circular-Polarised Photon

Proton Best: Incoming γ circularly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, perpendicular to \vec{k} :

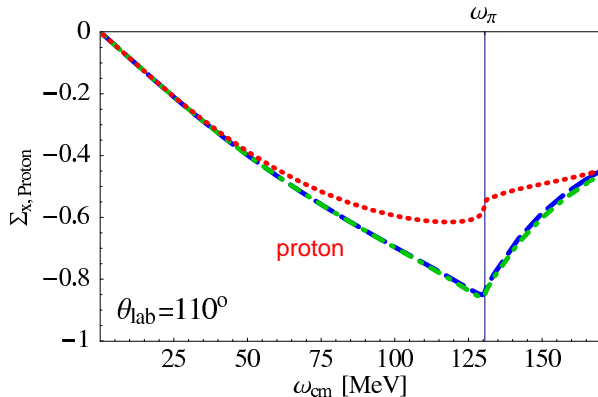
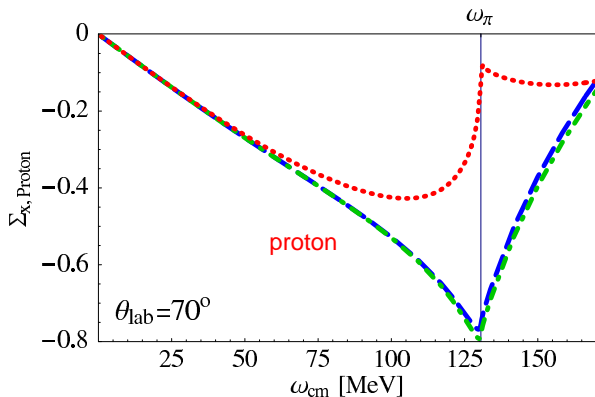
$$\Sigma_x = \frac{(\uparrow\rightarrow) - (\uparrow\leftarrow)}{(\uparrow\rightarrow) + (\uparrow\leftarrow)}$$



--- full $N\pi + \Delta$

..... no γ 's

--- no $l \geq 2$



- Dominated by structure

- Clear γ -dep.

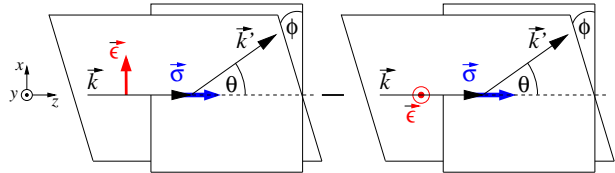
- Higher poles negligible

- HI γ S (?), MAMI(?)

Also good signal for **linear polarisations**.

Deuteron Best: Incoming γ linearly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, parallel to \vec{k} :

difference Δ_z^{lin} , asymmetry $\Sigma_z^{\text{lin}} = \frac{\Delta_z^{\text{lin}}}{\text{sum}}$



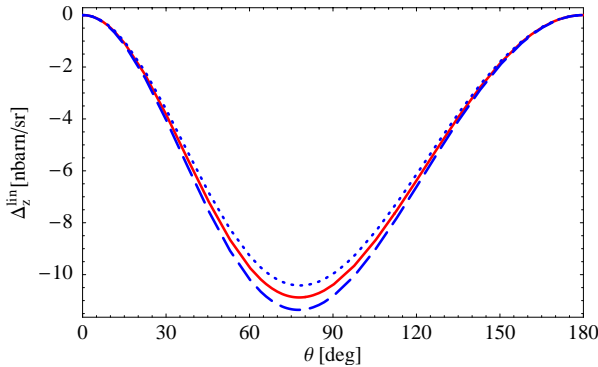
Sensitivity on neutron γ_{M1M1}

— 3.2; - - - - 3.2 + 2; ····· 3.2 - 2

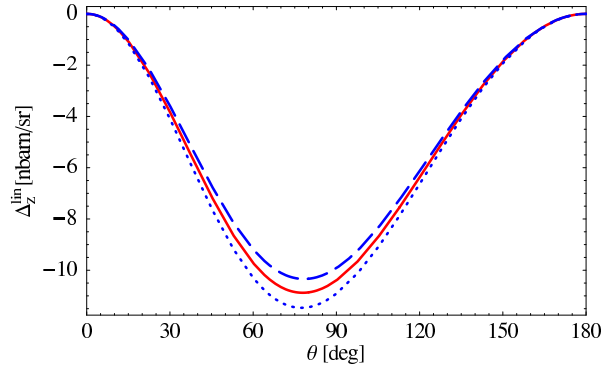
Sensitivity on neutron α_{E1}

— 11.3; - - - - 11.3 + 2; ····· 11.3 - 2

$\omega_{\text{lab}} = 125 \text{ MeV}, \delta\gamma_{M1M1} = \pm 2$



$\omega_{\text{lab}} = 125 \text{ MeV}, \delta\alpha_{E1} = \pm 2$



Sensitive to γ_{M1M1} , but must nail down α_{E1}, β_{M1} at lower energy.

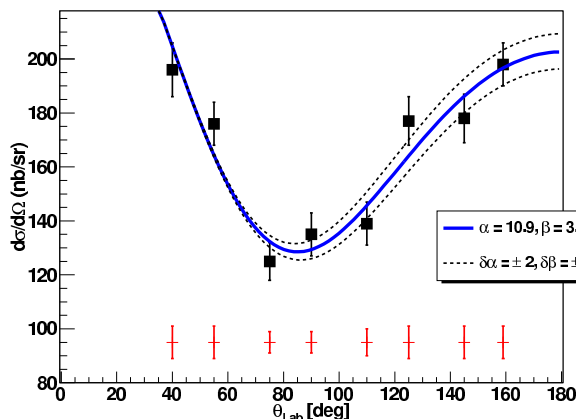
Experiment: coherent $\frac{d\sigma}{d\Omega} \propto (\text{target-charge})^{2 \text{ to } 1}$, more & easier targets \implies *heavier nuclei* for larger count-rates

Theory: To reliably extract nucleon pols., must accurately describe nuclear binding & levels

\implies *lighter nuclei* for simpler numerics & reliability

Find sweet-spot between competing forces: ^4He , ^6Li .

^6Li experiment at HI γ S \implies Theory needs efficient method for $A > 3$ systems.



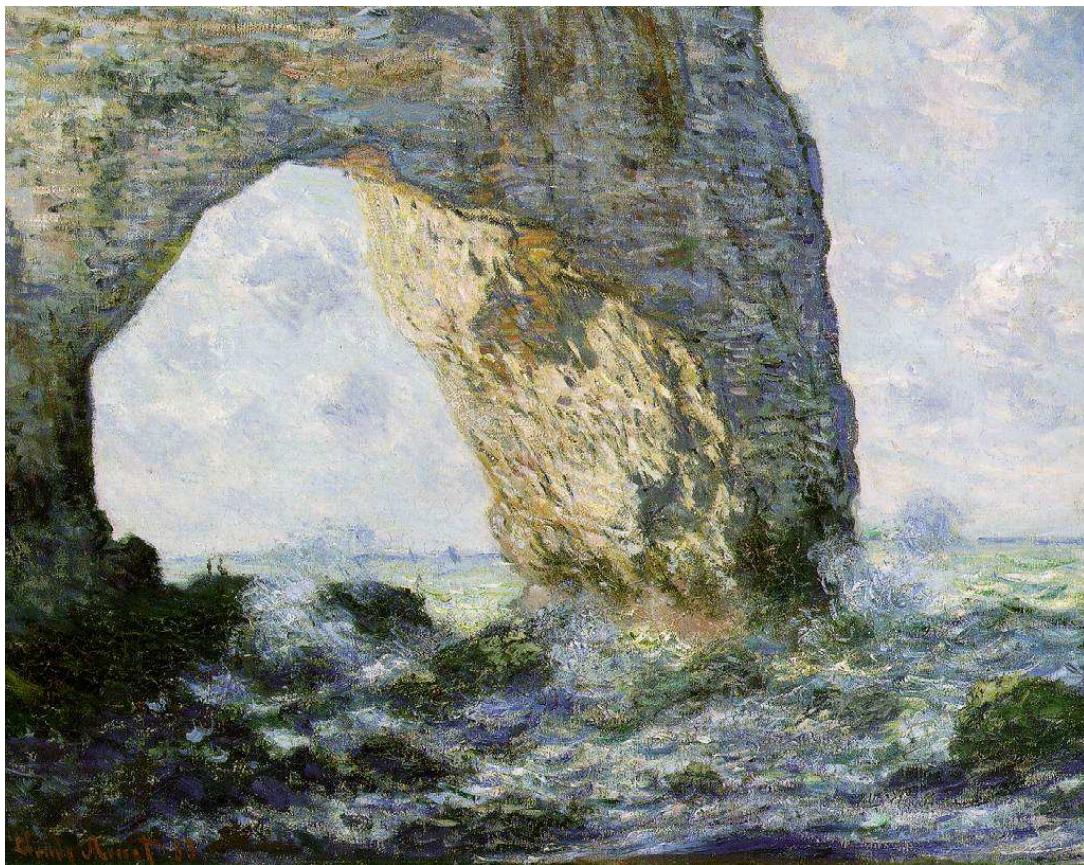
Myers/Feldman/... 2012 with shot-in-the-dark model

Promising:

Lorentz Integral Transformation method LIT

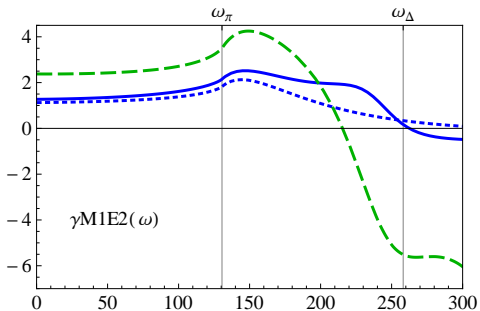
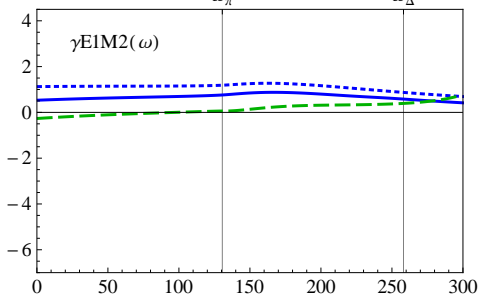
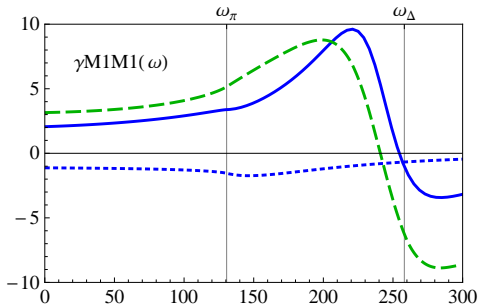
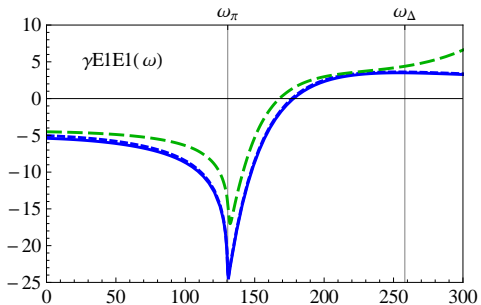
for rescattering contributions Efras/Leidemann/Orlandini 2000

Now tackling ^6Li . Bampa/Leidemann/Arenhövel/... in progress



Claude Monet: Rock Arch West of Etretat (The Manneport), 1883

Predicted in χ EFT: No N -core contributions \Rightarrow Spin-physics dominated by pion-cloud + Δ (γ_{M1M1} , γ_{M1E2}).



Pure polarisabilities

γ_{E1E1} , γ_{M1M1} agree.

— $N\pi + \Delta$

- - - $N\pi$

- - - Disp. Rel. .

Mixed polarisabilities

γ_{E1M2} , γ_{M1E2} small.

Uncertainties in DR?

Static values

units: [10^{-4} fm 4]

$$\bar{\gamma}_0 = -(\gamma_{E1E1} + \gamma_{M1M1} + \gamma_{E1M2} + \gamma_{M1E2})$$

$$\bar{\gamma}_\pi - (\pi^0\text{-pole}) = -\gamma_{E1E1} + \gamma_{M1M1} - \gamma_{E1M2} + \gamma_{M1E2}$$

χ EFT
iso-scalar

-0.7
8.6

DR
iso-scalar

-0.4
12

MAMI
proton

-1
~ 8

LEGS
proton

~ 18

MAMI
neutron

??

[12...16] $\pm 4_{\text{model}}$

(b) Consequences of NN-Rescattering: An Exact Low-Energy Theorem hg/Hemmert/ Hildebrandt 2010

Off-shell T_{NN} by **Green's function method**

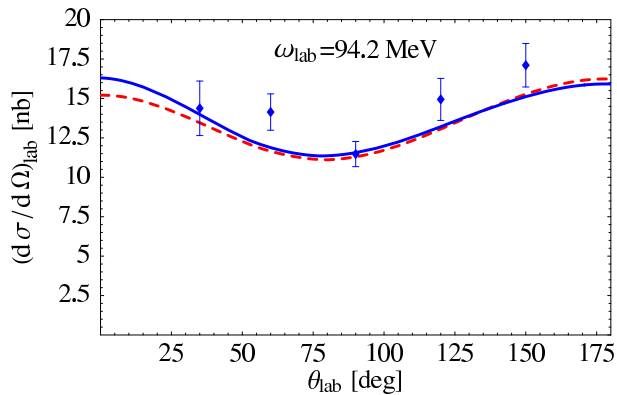
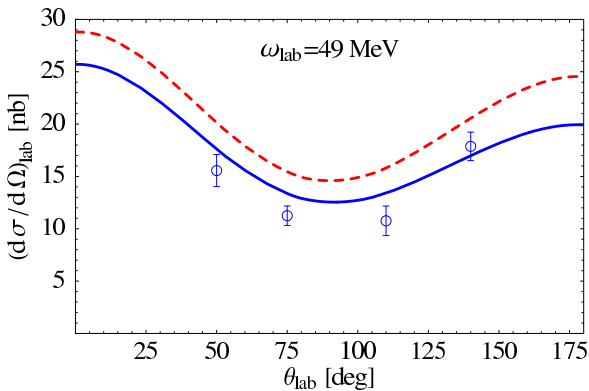
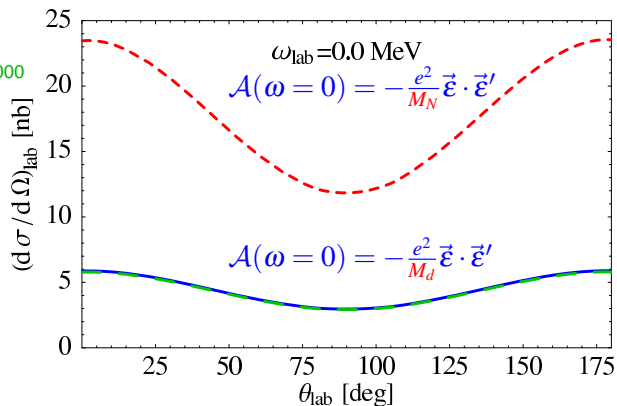
Arenhövel/Weyrauch 1980/83, Karakowski/Miller 1999, Levchuk/L'vov 2000

Thomson limit order-by-order in χ EFT.

Statistically significant only for $\omega \lesssim 70$ MeV.

Included higher-order diagrams indeed **small**.

---: without T_{NN} —: with T_{NN} - - -: predicted Thomson



(b) Consequences of NN-Rescattering: An Exact Low-Energy Theorem

Eliminate Dependence on Deuteron Wave-function also at high ω , for $\omega \rightarrow 0$ clear from Thomson.

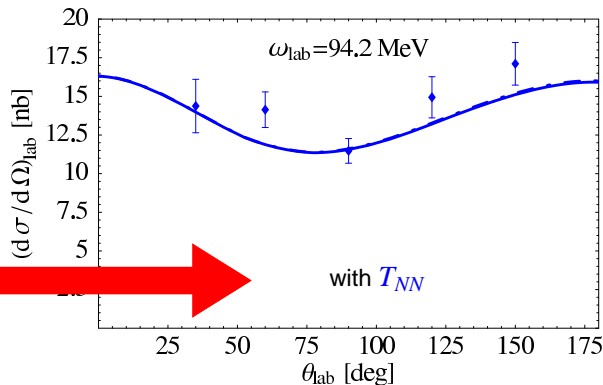
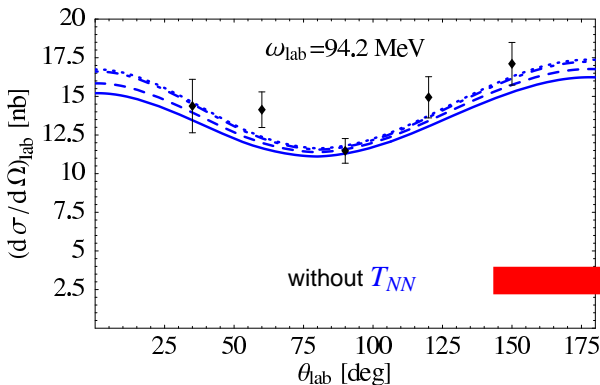
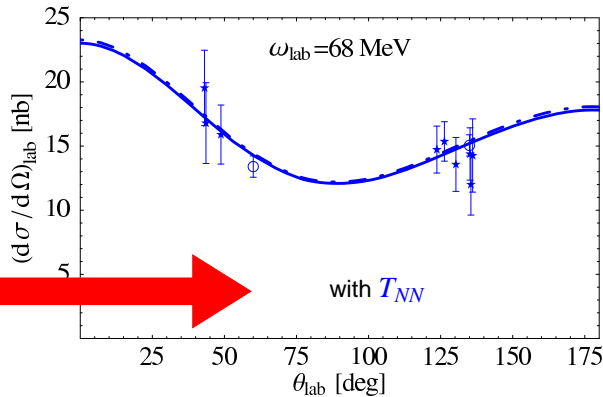
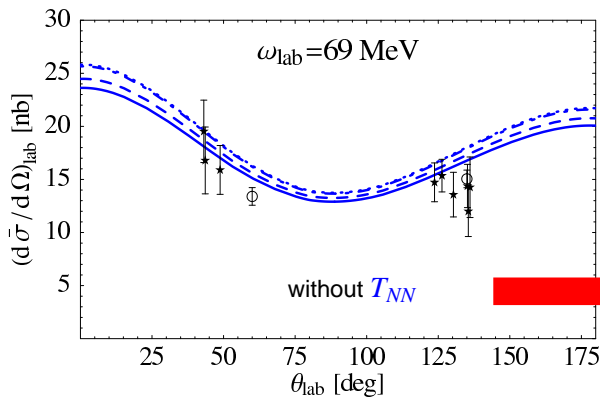
Illinois \circ , Lund \star , Saskatoon \blacklozenge

— “NNLO χ PT”

- - - AV18

- - - CD-Bonn

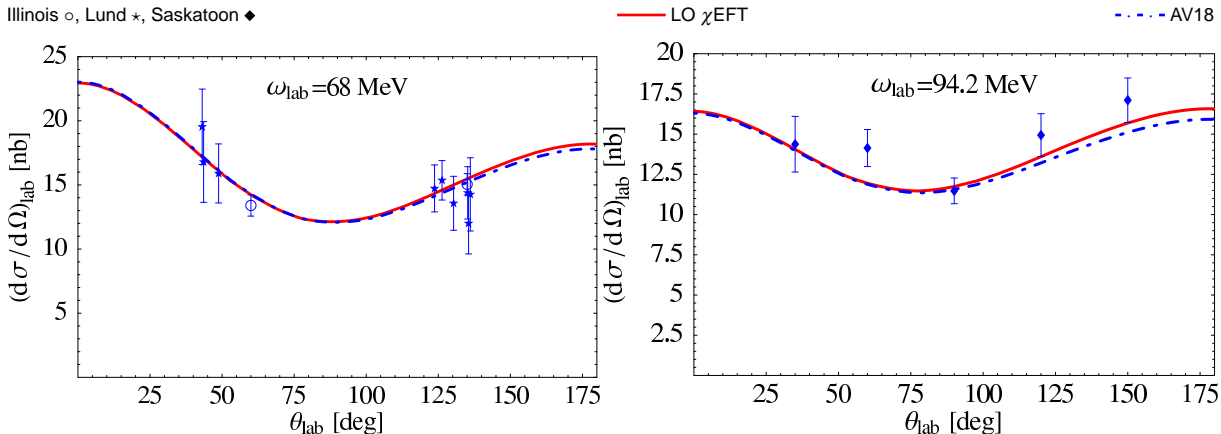
⋯⋯⋯ Nijmegen 93



(b) Consequences of NN -Rescattering: An Exact Low-Energy Theorem

Dependence of T_{NN} on NN -potential \cong short-distance, for $\omega \rightarrow 0$ clear from Thomson.

Illinois \circ , Lund \star , Saskatoon \blacklozenge

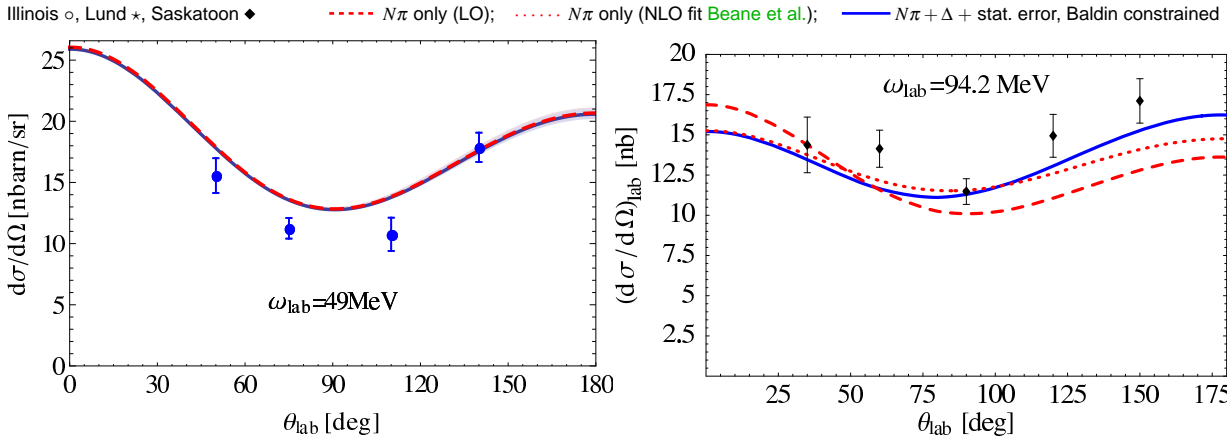


$$\text{LO } \chi\text{EFT-potential: } \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} + \begin{array}{c} C_{0,P} \\ \diagup \quad \diagdown \\ \diagdown \quad \diagup \end{array} \sim Q^{-1}$$

Consistent for Compton at NLO: $\mathcal{O}(Q^0)$ -correction of NN -potential presumed zero.

AV18 provides $< 3\%$ corrections \implies suggests higher-order indeed $Q^1 \approx \left(\frac{1}{7}\right)^2$.

(c) Determine Neutron Polarisabilities from all Deuteron Data



constrained:

$$\tilde{\alpha}^s = 11.0 \pm 0.9_{\text{stat}} \pm 0.3_{\Sigma} \pm 0.8_{\text{theory}}$$

$$\tilde{\beta}^s = 3.2 \mp 0.7_{\text{stat}} \pm 0.3_{\Sigma} \pm 0.8_{\text{theory}}$$

Proton hg/...2003:

$$\tilde{\alpha}^p = 11.0 \pm 1.4_{\text{stat}} \pm 0.4_{\Sigma} \pm 0.8_{\text{theory}}$$

$$\tilde{\beta}^p = 2.8 \mp 1.4_{\text{stat}} \pm 0.4_{\Sigma} \pm 0.8_{\text{theory}}$$

previous ranges:

$$[6 \dots 18]$$

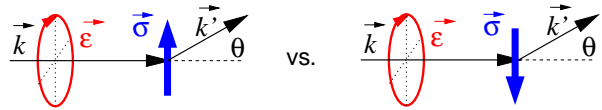
$$[-4 \dots 9]$$

estimate theory uncertainty ($\lesssim \pm 1$): higher-order $1N$; AV18 vs. LO χ EFT, d wave-fu., with vs. without T_{NN} .

\Rightarrow **neutron \approx proton polarisabilities**

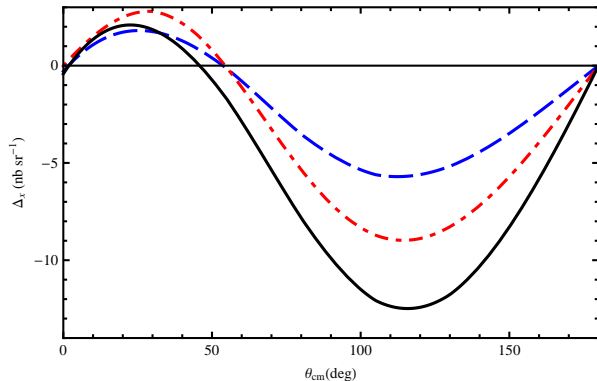
Deuteron Best: Incoming γ circularly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, perpendicular to \vec{k} :

difference Δ_x^{circ} , asymmetry $\Sigma_x^{\text{circ}} = \frac{\Delta_x^{\text{circ}}}{\text{sum}}$



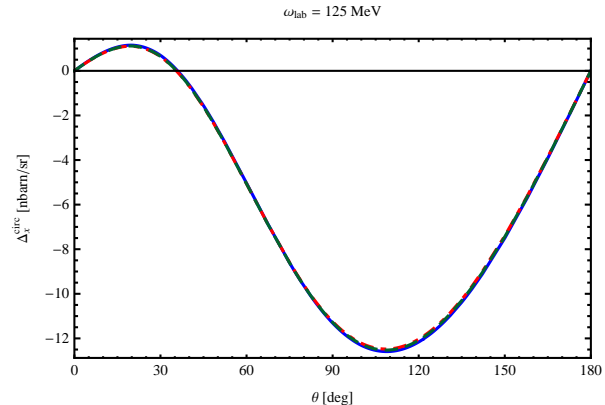
Sensitivity on Δ & NN -rescattering:

--- $N\pi$, no NN ; - - - $N\pi + \Delta$, no NN ; — $N\pi + \Delta + NN$



Sensitivity on wave-function:

NNLO Epelbaum 650 MeV, AV18, Nijmegen 93

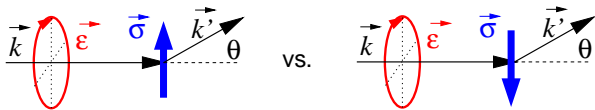


- More pronounced by explicit $\Delta(1232)$
- Thomson (NN rescatt.) important even at high $\omega = 125$ MeV

- No residual deuteron wave-function dependence
- Higher poles negligible

Deuteron Best: Incoming γ circularly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, perpendicular to \vec{k} :

difference Δ_x^{circ} , asymmetry $\Sigma_x^{\text{circ}} = \frac{\Delta_x^{\text{circ}}}{\text{sum}}$



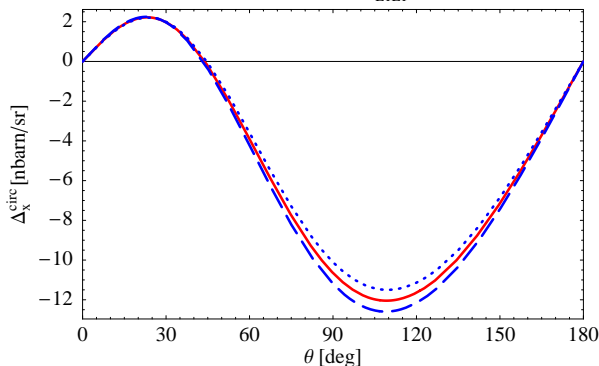
Sensitivity on neutron γ_{E1E1}

— -5.2 ; - - - $-5.2 + 2$; ····· $-5.2 - 2$

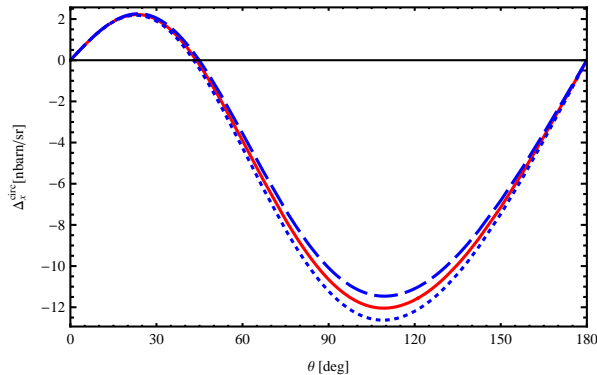
Sensitivity on neutron $\alpha_{E1} - \beta_{M1}$; Baldin- Σ fixed

— 8.2 ; - - - $8.2 + 2$; ····· $8.2 - 2$

$\omega_{\text{lab}} = 125 \text{ MeV}, \delta\gamma_{E1E1} = \pm 2$



$\omega_{\text{lab}} = 125 \text{ MeV}, \delta(\alpha_{E1} - \beta_{M1}) = \pm 2, \text{ Baldin fixed}$

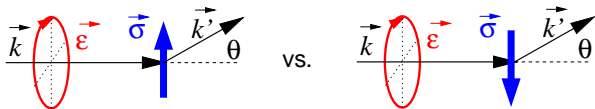


Sensitive to γ_{E1E1} , but must nail down α_{E1}, β_{M1} at lower energy.

Similarly good signal for linear polarisation Δ_x^{lin} .

Deuteron Best: Incoming γ circularly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, perpendicular to \vec{k} :

difference Δ_x^{circ} , asymmetry $\Sigma_x^{\text{circ}} = \frac{\Delta_x^{\text{circ}}}{\text{sum}}$



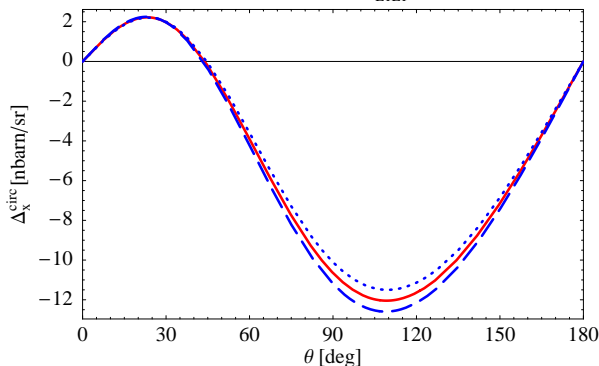
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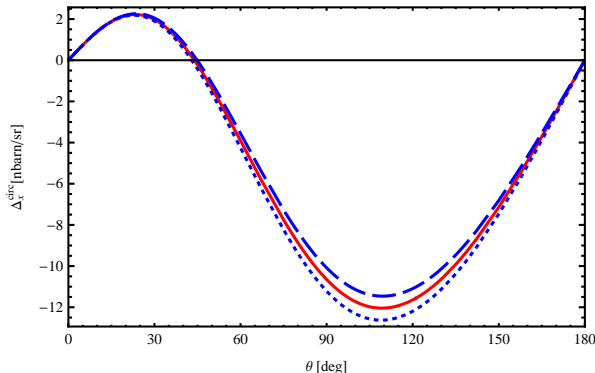
Sensitivity on neutron $\alpha_{E1} - \beta_{M1}$; Baldin- Σ fixed

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$\omega_{\text{lab}} = 125 \text{ MeV}, \delta\gamma_{E1E1} = \pm 2$



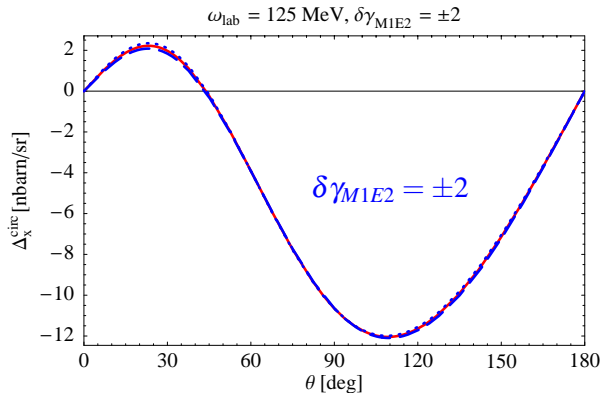
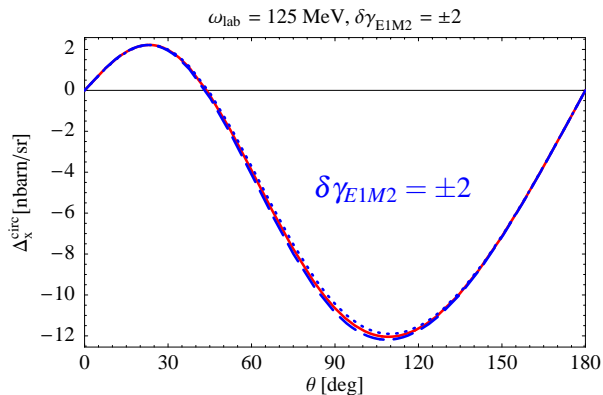
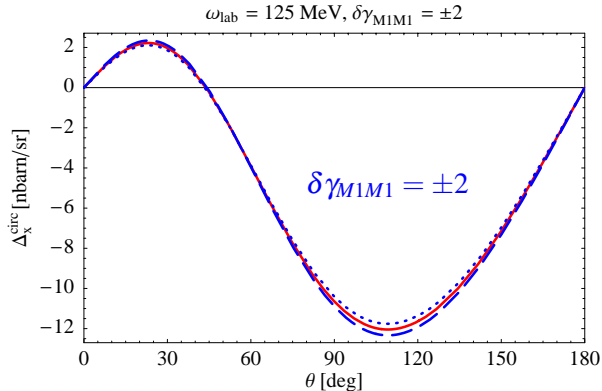
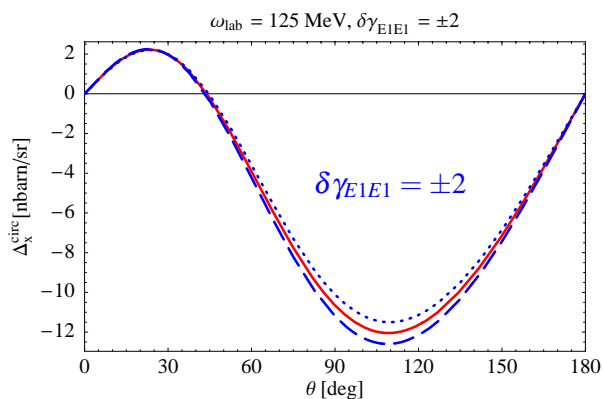
$\omega_{\text{lab}} = 125 \text{ MeV}, \delta(\alpha_{E1} - \beta_{M1}) = \pm 2, \text{ Baldin fixed}$



Sensitive to γ_{E1E1} , but must nail down α_{E1}, β_{M1} at lower energy.

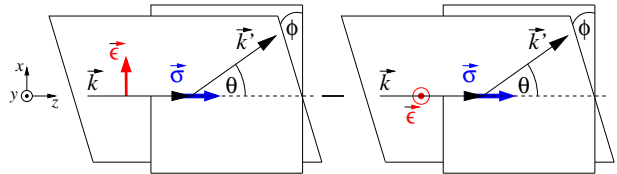
$\Delta(1232)$ and re-scattering increase signal.

Deuteron Best: Incoming γ circularly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, perpendicular to \vec{k} :



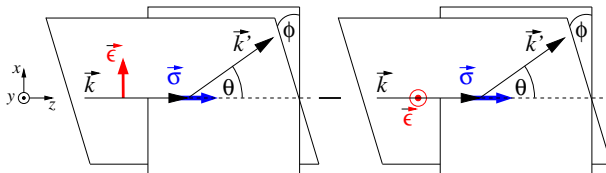
Deuteron Best: Incoming γ linearly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, parallel to \vec{k} :

difference Δ_z^{lin} , asymmetry $\Sigma_z^{\text{lin}} = \frac{\Delta_z^{\text{lin}}}{\text{sum}}$



Deuteron Best: Incoming γ linearly polarised, sum over final states. N -spin in (\vec{k}, \vec{k}') -plane, parallel to \vec{k} :

difference Δ_z^{lin} , asymmetry $\Sigma_z^{\text{lin}} = \frac{\Delta_z^{\text{lin}}}{\text{sum}}$



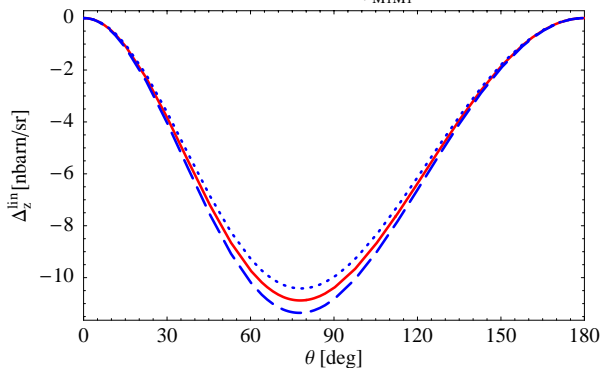
Sensitivity on neutron γ_{M1M1}

— 3.2; - - - - 3.2 + 2; ····· 3.2 - 2

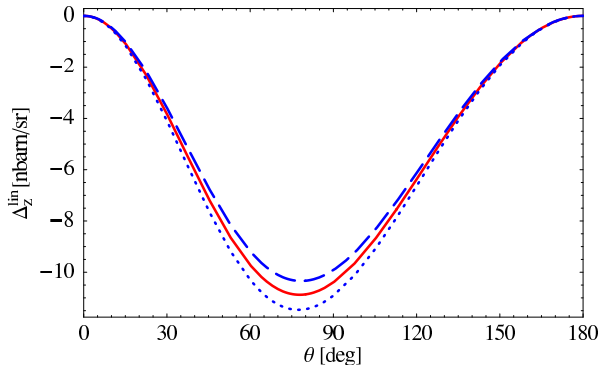
Sensitivity on neutron α_{E1}

— 11.3; - - - - 11.3 + 2; ····· 11.3 - 2

$\omega_{\text{lab}} = 125 \text{ MeV}, \delta\gamma_{M1M1} = \pm 2$



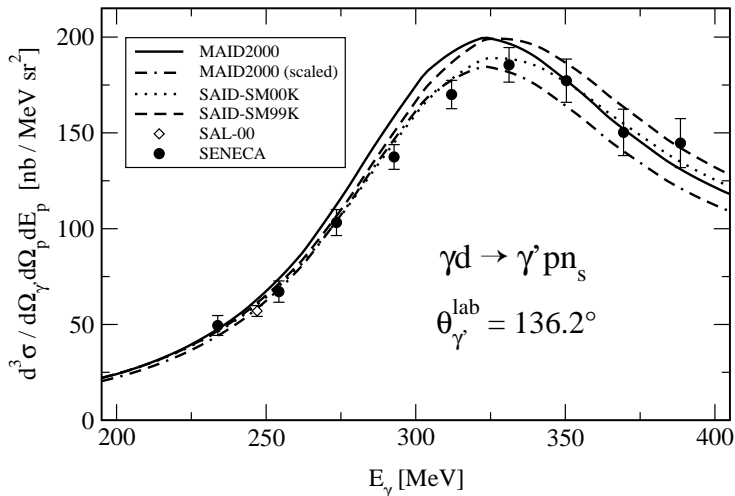
$\omega_{\text{lab}} = 125 \text{ MeV}, \delta\alpha_{E1} = \pm 2$



Sensitive to γ_{M1M1} , but must nail down α_{E1}, β_{M1} at lower energy.

Nucleon polarisabilities from centre of quasi-inelastic peak in $A(\gamma, \gamma A')N$

Only 10 data for $d(\gamma, \gamma p)n$ at $\omega \in [230; 400]$ MeV.



Kossert et al. 2003 found $\alpha_{E1}^n = 12.5 \pm 1.8(\text{stat})_{-0.6}^{+1.1}(\text{syst}) \pm 1.1(\text{model})$, β_{M1} from Baldin

sys. & model-error *under-estimated?*:
 π production, SAID/MAID-2000 amplitudes,
 π -exchange currents not chirally consistent,
 ...

To Do: Theory starting up

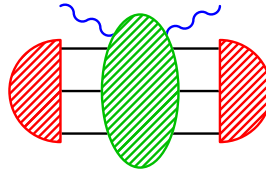
Analyse elastic & inelastic in **unified** χ EFT frame, test quasi-free hypothesis.
 Enhancement by $\Delta(1232)$ peak \implies accurate Δ theory.

Sensitivity of single-/double-polarised observables, breakup asymmetries.

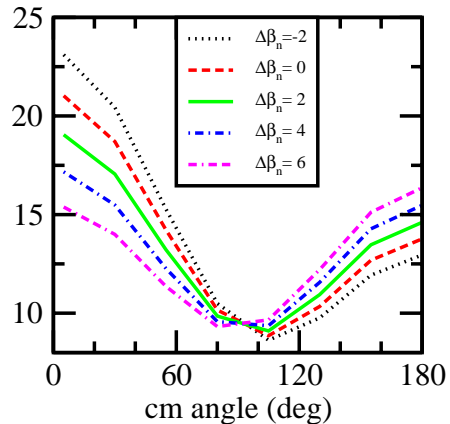
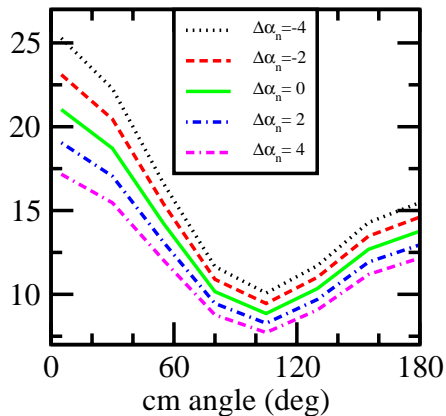
To Do: Experiment

Better data.
 Lower energies.

Alternative targets: ^3He , ^4He , ^6Li , ... , also for proton?



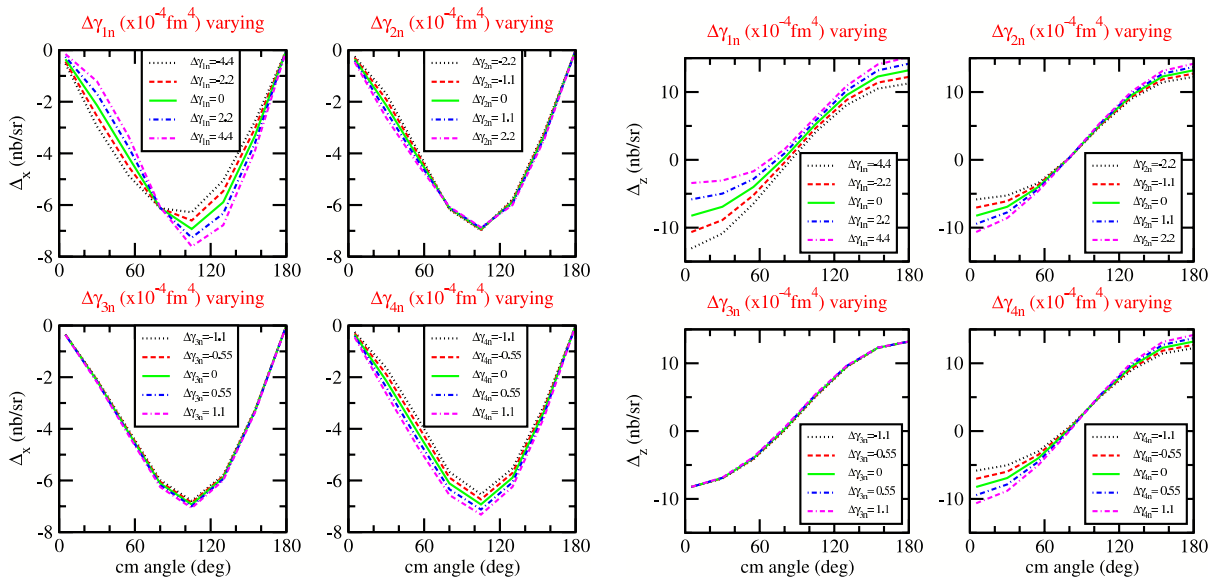
Example: Sensitivity of unpolarised cross-section at $\omega_{\text{lab}} = 120$ MeV on α^n, β^n



- First Compton cross section on ^3He . - ^3He as effective neutron spin target.
- Extend beyond $\omega \in [80; 120]$ MeV: re-scattering (Thomson, T_{NN}), explicit $\Delta(1232)$, threshold corrections

\Rightarrow Effects will become more pronounced.

Example: Sensitivity of polarised cross-section at $\omega_{\text{lab}} = 120 \text{ MeV}$ on γ_i^i 's



– First Compton cross section on ^3He .

– ^3He as effective neutron spin target.

– Extend beyond $\omega \in [80; 120] \text{ MeV}$: re-scattering (Thomson, T_{NN}), explicit $\Delta(1232)$, threshold corrections

⇒ Effects will become more pronounced.